

# Conformal magnetic effect in scalar QED

Maxim Chernodub

Institut Denis Poisson, Tours, France

Pacific Quantum Center, Vladivostok, Russia

in collaboration with

Vladimir Goy, Alexander Molochkov

Pacific Quantum Center, Vladivostok, Russia

Partially supported by grant No. 0657-2020-0015 of the Ministry of Science and Higher Education of Russia.



# Massless Dirac fermions

A generic system in particle physics, cosmology, solid state ...

Covariant formulation (quantum field theory)

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \bar{\psi}i\not{D}\psi$$

Dirac semimetals (solid state):

$$\bar{\psi} \left[ i\gamma^0 \hbar \frac{\partial}{\partial t} + v_F \boldsymbol{\gamma} (i\hbar \boldsymbol{\nabla} - e\mathbf{A}) \right] \psi$$

$$\not{D} = \gamma^\mu D_\mu \quad D_\mu = \partial_\mu + ieA_\mu$$

$$F_{\mu\nu} = \partial_\mu A_\nu - \partial_\nu A_\mu$$

# Classical symmetries

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \bar{\psi}i\not{D}\psi$$

## Vector

$$\psi \rightarrow e^{i\omega_V} \psi$$

local/gauge symmetry

vector current is classically conserved

$$j_V^\mu = \bar{\psi}\gamma^\mu\psi \quad \partial_\mu j_V^\mu = 0$$

## Axial

$$\psi \rightarrow e^{i\omega_5\gamma^5} \psi$$

global symmetry (no axial gauge field)

axial current is classically conserved

$$j_A^\mu = \bar{\psi}\gamma^5\gamma^\mu\psi \quad \partial_\mu j_A^\mu = 0$$

## Conformal

global scale transformations

$$x \rightarrow \lambda^{-1}x, \quad A_\mu \rightarrow \lambda A_\mu, \quad \psi \rightarrow \lambda^{3/2}\psi$$

Dilatation current is classically conserved

$$j_D^\mu = T^{\mu\nu}x_\nu \quad \partial_\mu j_D^\mu \equiv T^\mu_\mu \equiv 0$$

$$(T^\mu_\mu)_{\text{cl}} \equiv 0$$

Energy-Momentum tensor

$$T^{\mu\nu} = -F^{\mu\alpha}F^\nu_\alpha + \frac{1}{4}\eta^{\mu\nu}F_{\alpha\beta}F^{\alpha\beta}$$

$$+ \frac{i}{2}\bar{\psi}(\gamma^\mu D^\nu + \gamma^\nu D^\mu)\psi - \eta^{\mu\nu}\bar{\psi}i\not{D}\psi$$

# Conformal anomaly and the beta function

Massless Dirac fermions

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \bar{\psi}i\not{D}\psi$$

are (classically) invariant under the global (scale) transformations:

$$x \rightarrow \lambda^{-1}x, \quad A_\mu \rightarrow \lambda A_\mu, \quad \psi \rightarrow \lambda^{3/2}\psi$$

The quantum theory generates an intrinsic scale due to a renormalization (in this particular case) of the electric charge:

$$\beta(e) = \frac{de(\mu)}{d \ln \mu}$$

In QED (for one Dirac fermion)

$$\beta_{\text{QED}}^{1\text{-loop}} = \frac{e^3}{12\pi^2}$$

renormalization scale



→ conformal symmetry is broken at the quantum level

# Quantum anomaly → anomalous transport

**Axial anomaly**  $\partial_\mu j_A^\mu = \frac{e^2}{16\pi^2} F_{\mu\nu} \tilde{F}^{\mu\nu}$

**Transport**  $j_A = \frac{\mu_V}{2\pi^2} eB, \quad j_V = \frac{\mu_A}{2\pi^2} eB$

Chiral separation and chiral magnetic effects

topological = exact in one loop = interesting!

**Mixed axial-gravitational anomaly**  $\partial_\mu j_A^\mu = -\frac{1}{384\pi^2} R_{\mu\nu\alpha\beta} \tilde{R}^{\mu\nu\alpha\beta}$

**Transport**  $j_V = \frac{\mu_V \mu_A}{\pi^2} \Omega, \quad j_A = \left( \frac{T^2}{6} + \frac{\mu_V^2 + \mu_A^2}{2\pi^2} \right) \Omega$

Thermal contribution to chiral vortical effects

topological = exact in one loop = interesting!

**Conformal anomaly**

$$\partial_\mu j_D^\mu = T^\alpha_\alpha \qquad \langle T^\mu_\mu \rangle = \frac{\beta(e)}{2e} F_{\mu\nu} F^{\mu\nu}$$

**Transport?**

not topological ... not one-loop exact ... not interesting?

# Conformal anomaly →

## Scale Electric Effect (SEE) and Scale Magnetic Effect (SME)

(Conformal Magnetic Effect = CME → interferes with Chiral Magnetic Effect ... already taken, too late)

### Oversimplified picture

Gravitational background: Weyl-transformed flat space

$$g_{\mu\nu}(x) = e^{2\tau(x)} \eta_{\mu\nu}$$

flat (Minkowski) metric

scale factor (arbitrary function of coordinates)

The conformal anomaly  
leads to scale electromagnetic effects:

$$\langle T^\mu_\mu \rangle = \frac{\beta(e)}{2e} F_{\mu\nu} F^{\mu\nu}$$

$$j^\mu \equiv \langle j^\mu_V \rangle = -\frac{2\beta(e)}{e} F^{\mu\nu} \partial_\nu \tau$$

“Generation of an electric current in a background of electromagnetic and gravitational fields”

**Fine print: 1) natural in a linear-response theory**

**2) very unnatural (would-be-wrong) in general relativity**

**→ to be refined (and made less simple) later**

# Scale electric effect (SEE)

Time-dependent background:  $\tau = \tau(t)$

**Metric:**  $g_{\mu\nu}(x) = e^{2\tau(x)}\eta_{\mu\nu}$

**Scale Electric Effect:**

$$\langle \mathbf{j}(t, \mathbf{x}) \rangle_{\text{scale}} = \sigma(t) \mathbf{E}(t, \mathbf{x}) \quad \text{for } \nabla \tau = 0$$

**Conformal conductivity:**

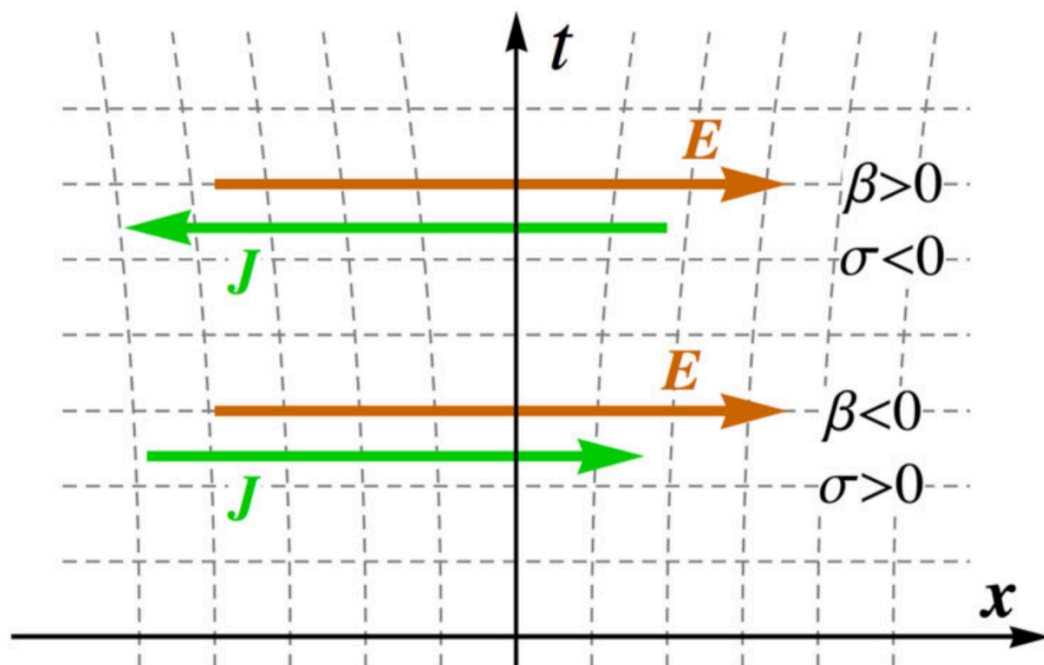
$$\sigma(t, \mathbf{x}) = -\frac{2\beta(e)}{e} \frac{\partial \tau(t, \mathbf{x})}{\partial t}$$

**Negative conductivity  
in an expanding space-time!**

**Cosmology:**

Independently obtained in the de Sitter spacetime (a version of the Schwinger effect, both for fermions and bosons)

- T. Hayashinaka, T. Fujita, and J. Yokoyama, Fermionic Schwinger effect and induced current in de Sitter space, *J. Cosmol. Astropart. Phys.* 07 (2016) 010; T. Hayashinaka and J. Yokoyama, Point splitting renormalization of Schwinger induced current in de Sitter spacetime, *J. Cosmol. Astropart. Phys.* 07 (2016) 012.
- T. Kobayashi and N. Afshordi, Schwinger effect in 4D de Sitter space and constraints on magnetogenesis in the early Universe, *J. High Energy Phys.* 10 (2014) 166.





# Scale magnetic effect (SME)

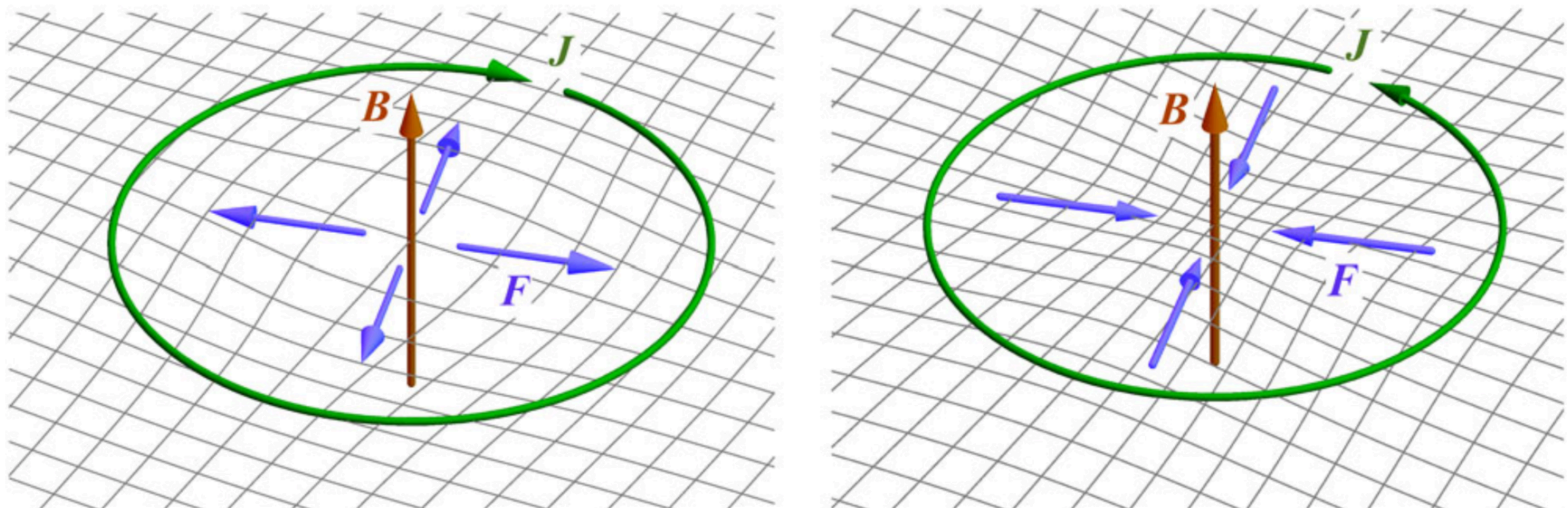
Space-dependent background:  $\tau = \tau(\mathbf{x})$

**Scale Magnetic Effect:**

$$\langle \mathbf{j}(t, \mathbf{x}) \rangle_{\text{scale}} = \mathbf{F}(\mathbf{x}) \times \mathbf{B}(t, \mathbf{x}) \quad \text{for } \partial_t \tau = 0$$

**Gravitational deformation vector:**

$$\mathbf{F}(t, \mathbf{x}) = \frac{2\beta(e)}{e} \nabla \tau(t, \mathbf{x})$$



**Distantly similar to Hall effects (BUT: no electric field, no matter: the vacuum effect).**



# Conformal anomaly and transport effects at the edge

What is about the boundaries? Consider a system with a reflective boundary.

electric current      beta function      normal to the boundary

$$J^\mu = -\frac{2c\beta_e}{e\hbar} \frac{F^{\mu\nu} n_\nu}{x}$$

distance to the reflective boundary

D.M. McAvity, H. Osborn, Class. Quantum Gravity 8, 603 (1991).

C.-S. Chu and R.-X. Miao, JHEP 07, 005 (2018), PRL 121, 251602 (2018).

## Scale (=Conformal) Magnetic Effect at the Edge (SMEE):

Electric current along the edge due to tangential magnetic field

$$\mathbf{j}(\mathbf{x}) = -f(\mathbf{x}) \mathbf{n} \times \mathbf{B}$$

$$f(\mathbf{x}) = \frac{2\beta(e)}{e^2} \frac{1}{x_\perp}$$

- Effect due to conformal anomaly
- No topology (Berry, Chern etc)

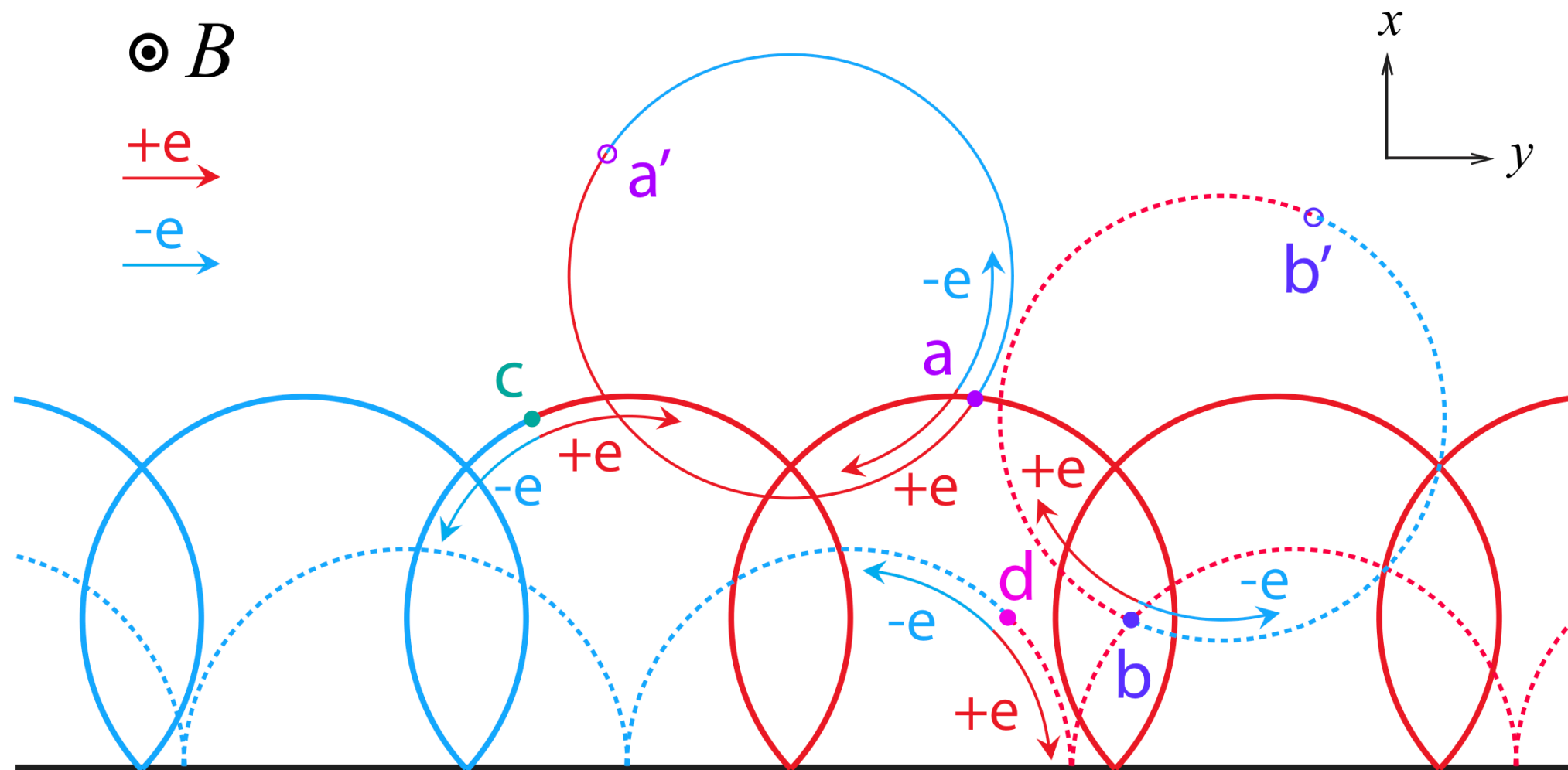
diverges at the boundary!

# Scale (=conformal) Magnetic Effect at the Edge

## A physical picture

C.-S. Chu and R.-X. Miao, JHEP 07, 005 (2018), PRL 121, 251602 (2018).

Ingredients: vacuum, edge and magnetic field



Skipping orbits (like in the Hall effect, but now in the vacuum)  
Absent: No Fermi surface, no temperature.

Works for fermions and bosons (= for usual and scalar QED)

# Conformal magnetic effect on the lattice

Generation of electric current at the boundary?

**Scalar electrodynamics at a conformal point (= no mass):**

$$\mathcal{L}_{\text{sQED}} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + [(\partial_\mu - ieA_\mu)\phi]^* (\partial^\mu - ieA^\mu)\phi$$

Massless one-component electrically-charged scalar field

---

**Lattice action:**

$$S = \beta_{\text{latt}} \sum_x \sum_{\mu < \nu=1}^4 (1 - \cos \theta_{x,\mu\nu}) \\ + \sum_x \sum_{\mu=1}^4 \left| \phi_x - e^{i(\theta_{x\mu} + \theta_{x\mu}^B)} \phi_{x+\hat{\mu}} \right|^2 \\ + \sum_x \left( -\kappa |\phi_x|^2 + \lambda |\phi_x|^4 \right),$$

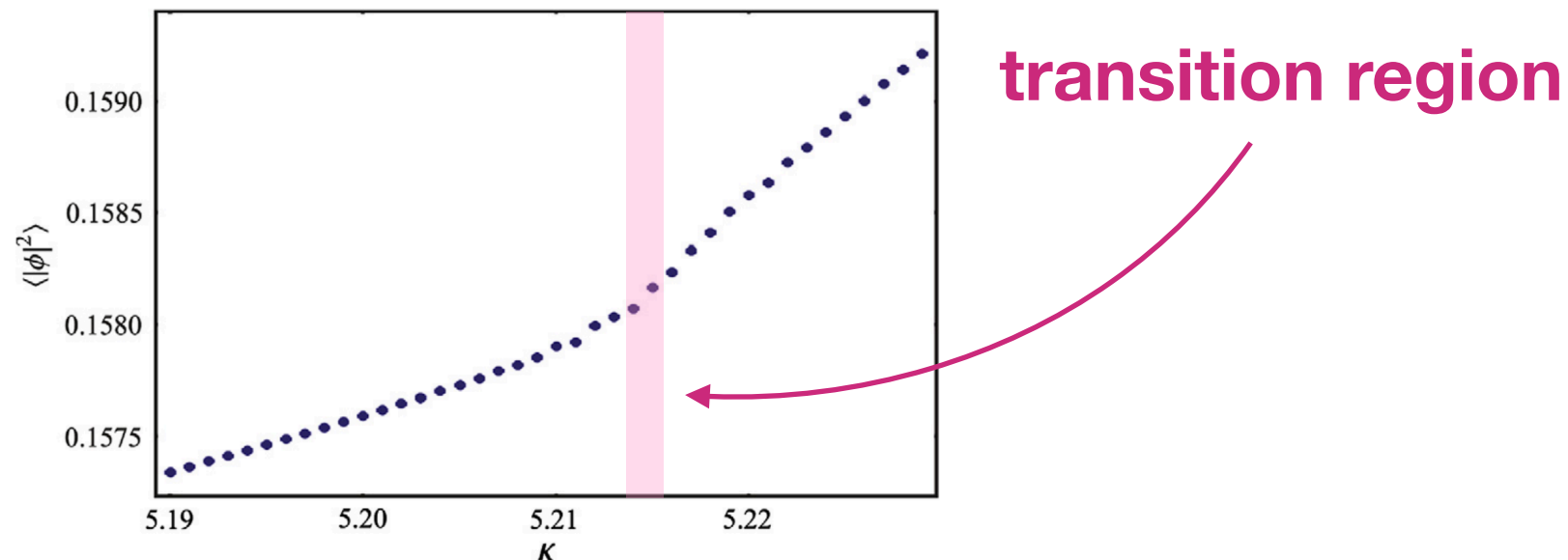
**We fix the couplings numerically in order to reach a conformal point in the vicinity of a second-order phase transition**

$$\beta_{\text{latt}} = 4 \\ \lambda = 10$$

**the coupling  $\kappa$  is fine-tuned**

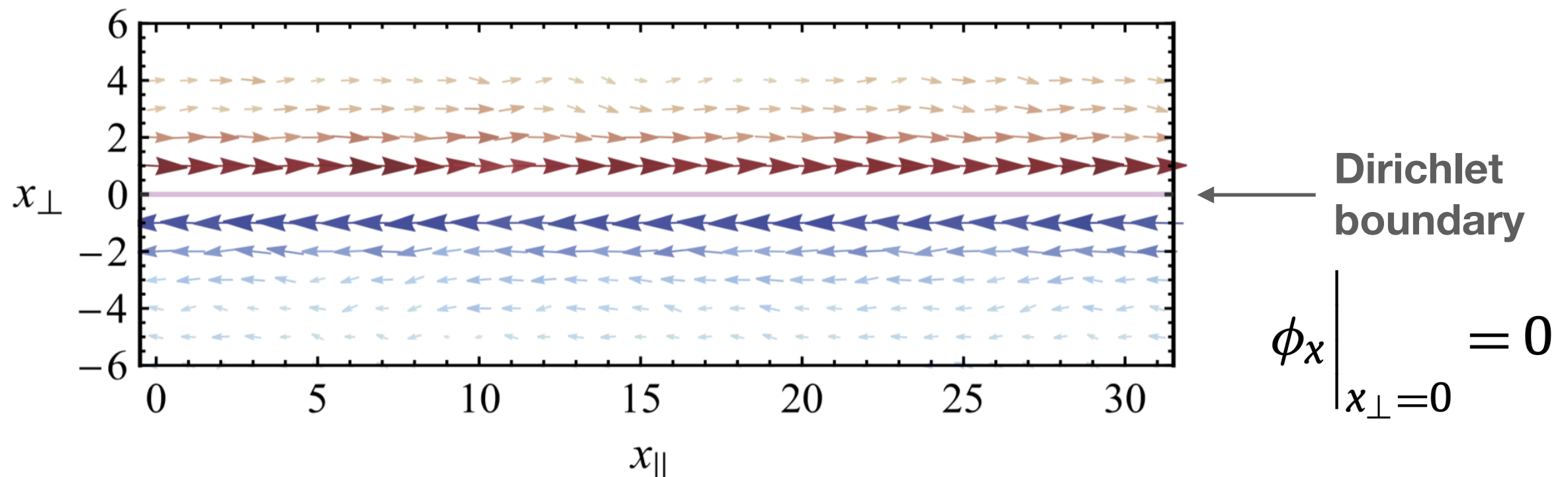
# Conformal magnetic effect on the lattice

## Approaching the conformal point



$$\beta_{\text{latt}} = 4$$
$$\lambda = 10$$

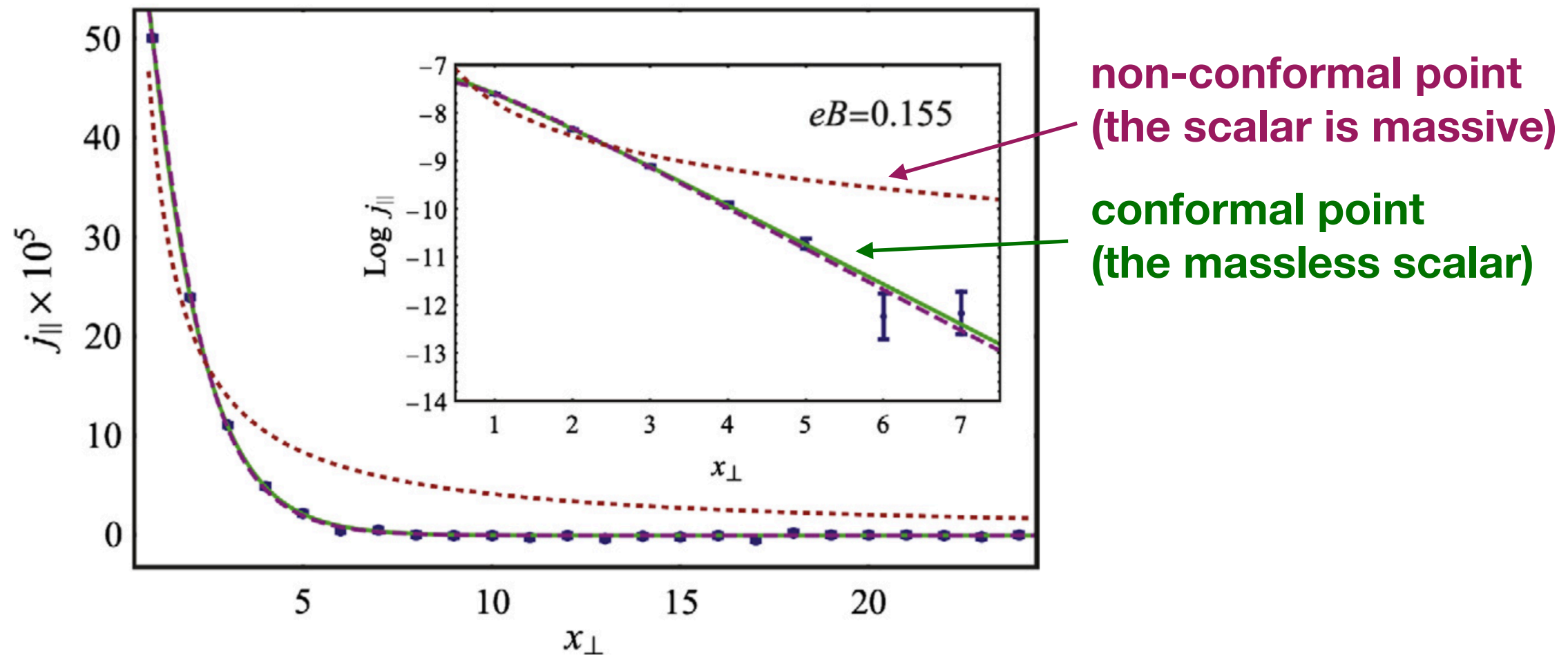
**We see the generated electric current!**



# Conformal magnetic effect: numerical check

Diverges at the boundary?

We see the  $1/x$  divergence of the current at the boundary



We see the correct coefficient and recover the beta function!

$$\beta_{\text{sQED}}^{1\text{-loop}} = \frac{e^3}{48\pi^2}$$

(Notice that the beta function of the scalar QED is four times smaller than the beta function in the usual QED)

# Scale magnetic effect at the edge and superconductivity

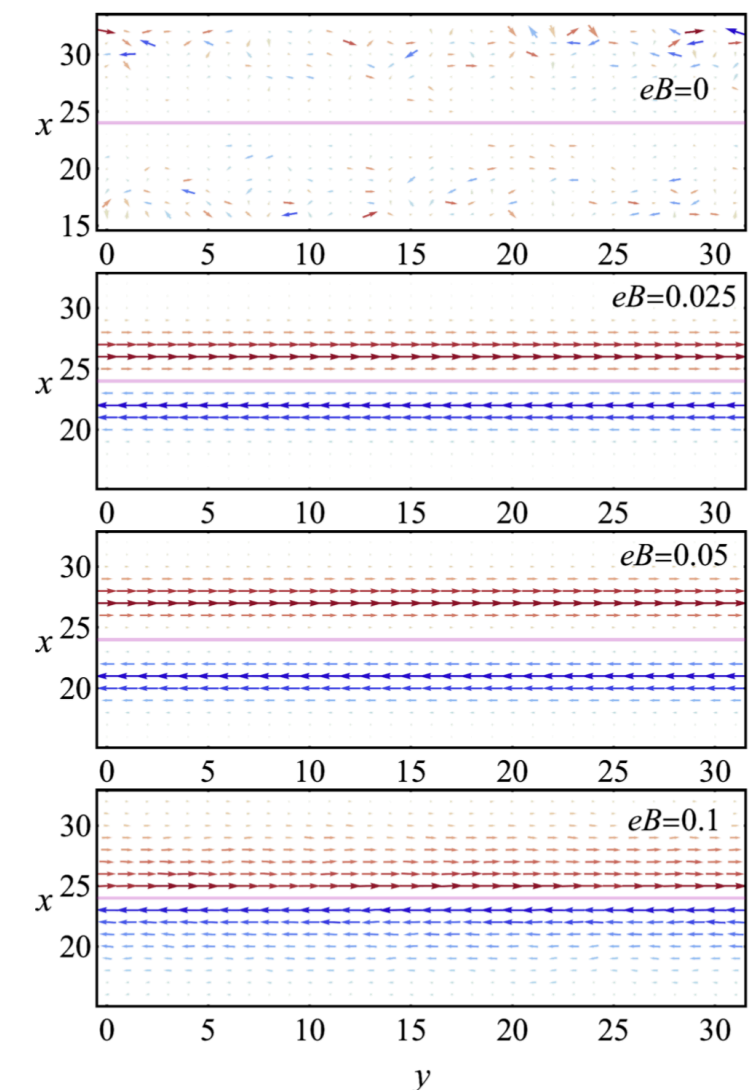
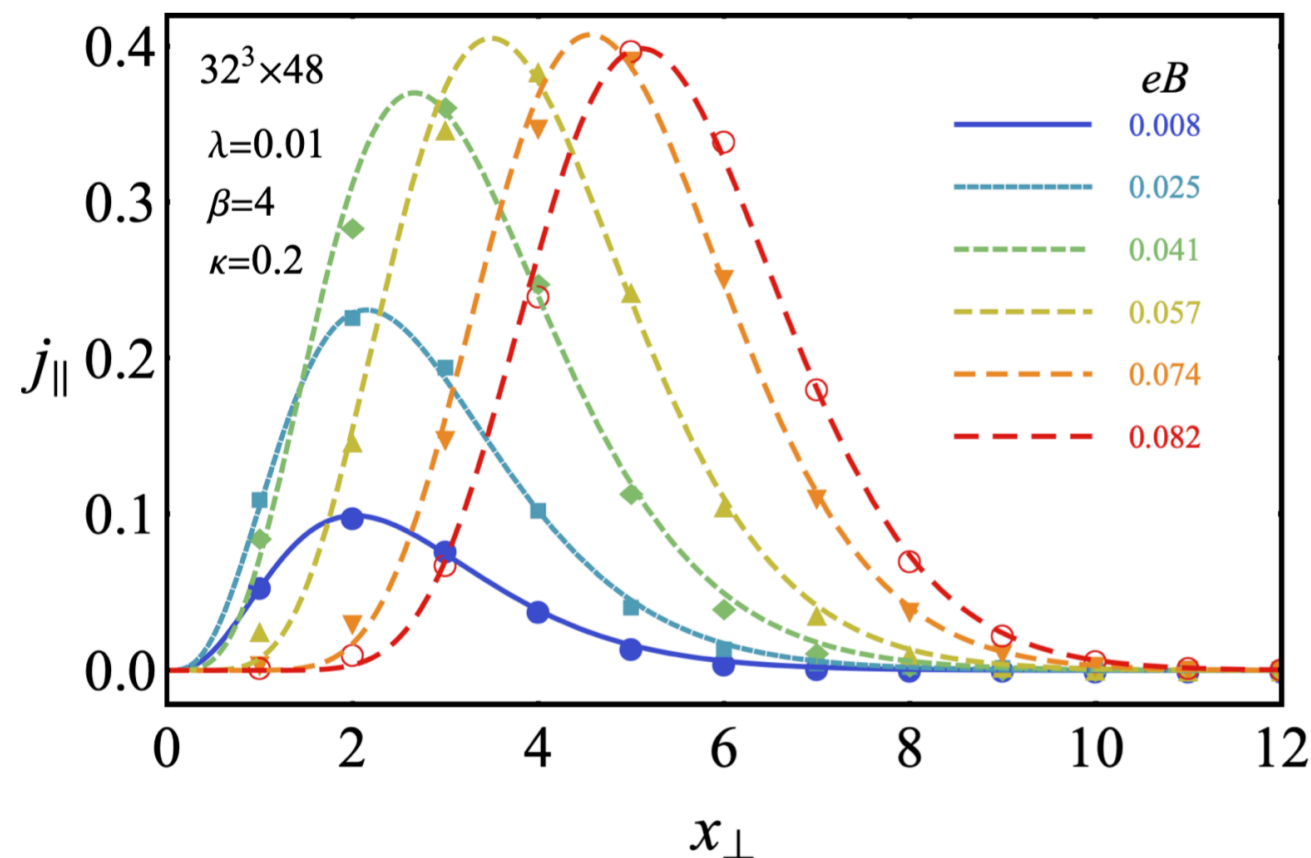
Scalar electrodynamics at a conformal point (massless scalar):

$$\mathcal{L}_{\text{sQED}} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + [(\partial_\mu - ieA_\mu)\phi]^* (\partial^\mu - ieA^\mu)\phi$$

Charged fields may condense!

$$\mathcal{L} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + (D_\mu\varphi)^* D^\mu\varphi - \overset{\text{potential}}{V(\varphi)}$$

What happens to the electric current in the superconducting phase?



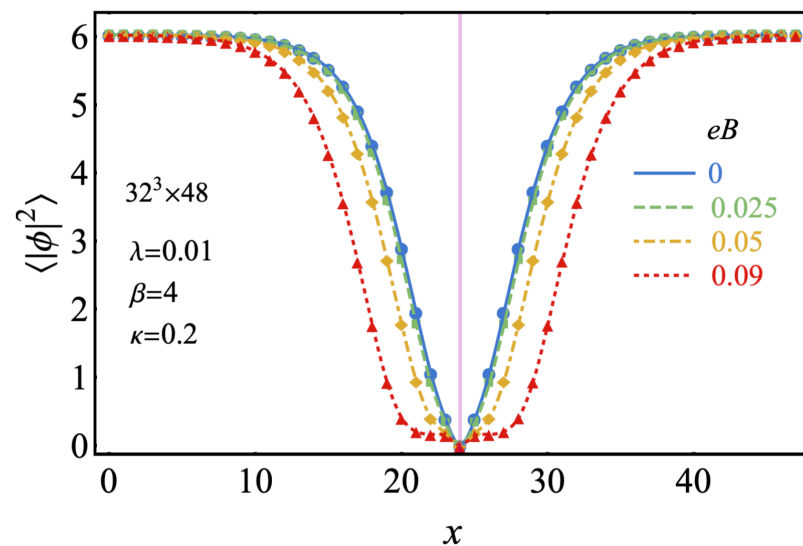


# Scale magnetic effect at the edge and superconductivity

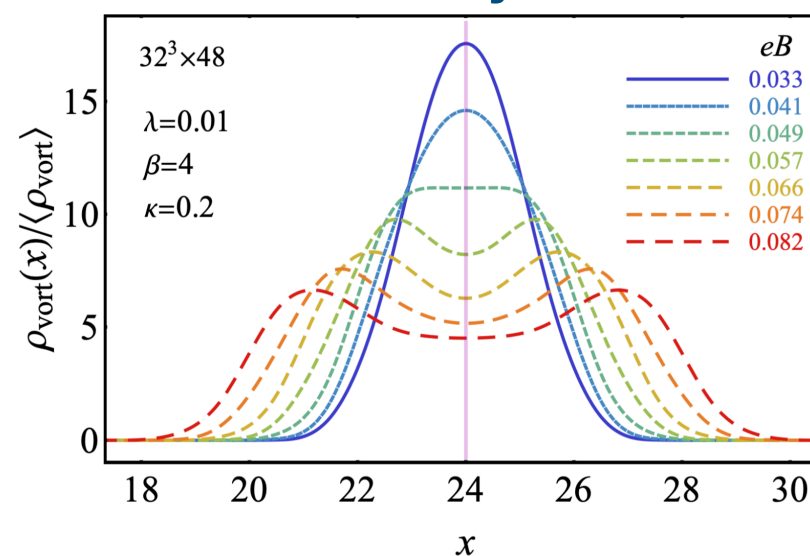
What happens to the SEEE boundary current in superconducting phase?

**It becomes the Meissner current!**

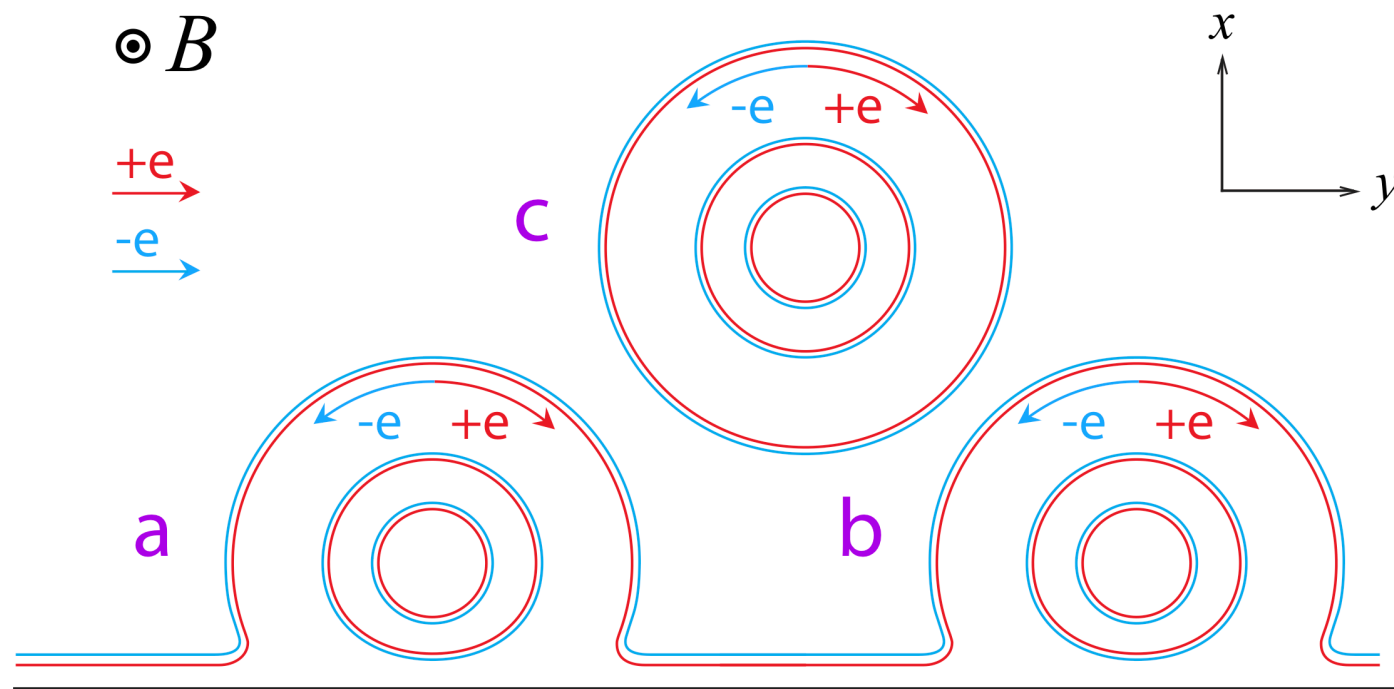
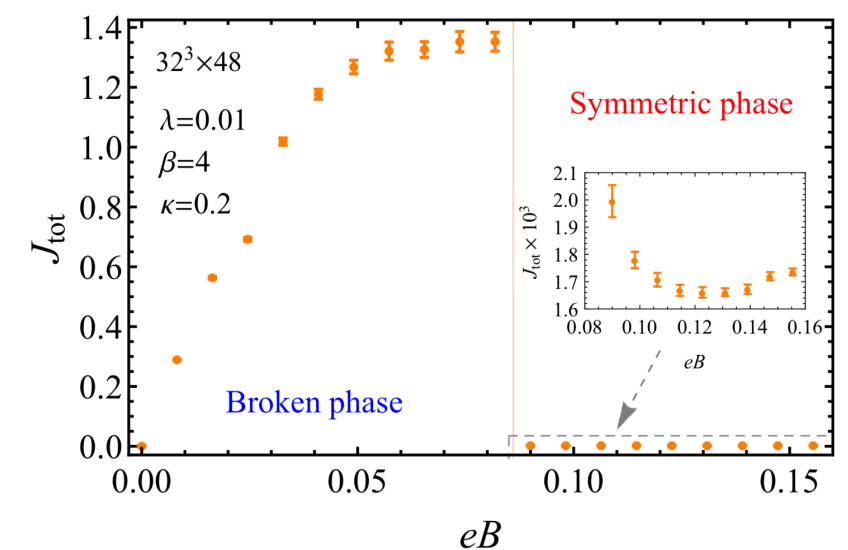
Superconducting condensate



Vortex density



Total (integrated) electric current



**Meissner:** total current is large, but it vanishes at the boundary

**Conformal:** total current is small, but it diverges at the boundary

# Summary

Conformal anomaly leads to a number of new transport effects:

- in the bulk (unbounded systems)
- at reflective boundaries (edges) of bounded systems

Generated electric currents are proportional to the beta function.  
(Accessible experimentally in Dirac and Weyl semimetals)

Scale (conformal) magnetic effect:

generates edge (boundary) currents in the absence of matter

We have shown numerically, that

1. the boundary electric current does exist in the scalar QED at the conformal point;
2. if the scalar field condenses, the boundary current gets supplemented by the usual Meissner current.