

Complex Langevin analysis of four-dimensional SU(2) gauge theory with a theta term

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Gauge theory with a θ term

☆ θ term: **topological** property of the gauge theory, **nonperturbative**

$$S_\theta = -i\theta Q = -\frac{i\theta}{32\pi^2} \int d^4x \epsilon_{\mu\nu\rho\sigma} \text{Tr} [F_{\mu\nu} F_{\rho\sigma}]$$

- **strong CP problem** of QCD

The experimental bound of θ is extremely small: $|\theta| < 10^{-10}$

→ no reason for it theoretically

- phase structure of 4D SU(N) YM around $\theta = \pi$

interesting prediction by the 't Hooft **anomaly matching**

Phase structure at $\theta = \pi$

☆ 't Hooft anomaly matching of 4D $SU(2)$ YM

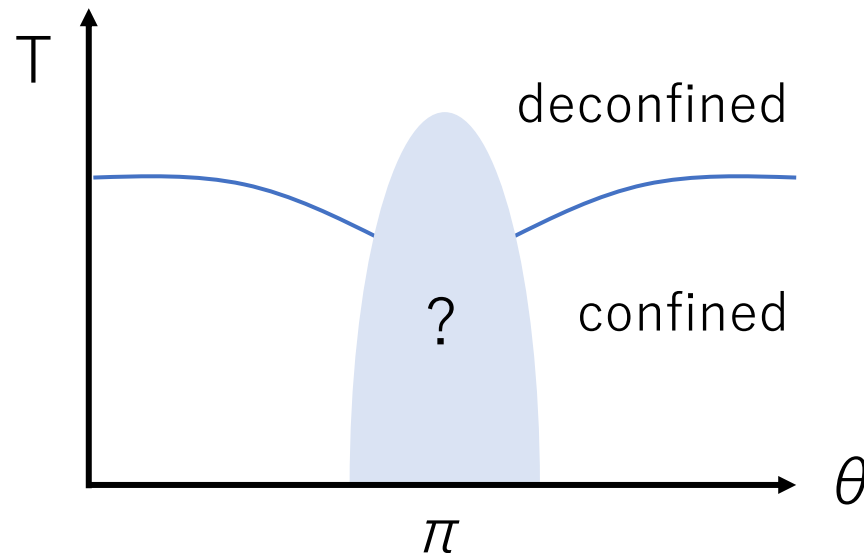
→ constrain the phase structure at $\theta = \pi$

mixed 't Hooft anomaly between
CP symmetry & Z_2 1-form center symmetry at $\theta = \pi$



- SSB of CP
- SSB of $Z_2^{(1)}$
- gapless
- topological QFT

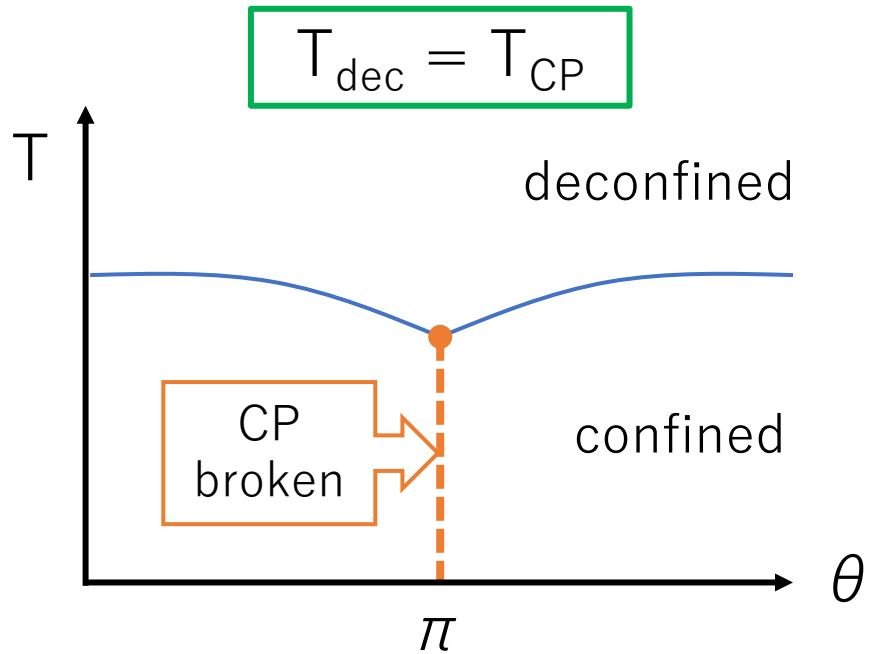
[D. Gaiotto, A. Kapustin, Z. Komargodski, N. Seiberg (2017)]



T_{dec} VS T_{CP}

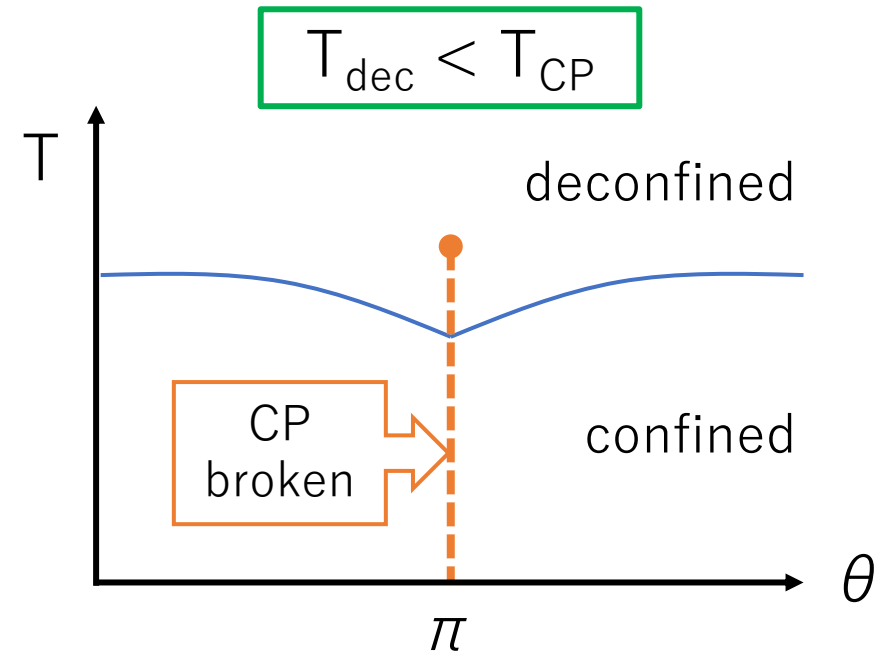
☆ anomaly matching $\rightarrow T_{\text{dec}} \leq T_{\text{CP}}$

• example of possible (θ, T) phase diagram



holography for large N supports

[F. Bigazzi, A. L. Cotrone, R. Sisca (2015)]



soft SUSY breaking of SYM supports

[S. Chen, K. Fukushima, H. Nishimura, Y. Tanizaki (2020)]

Numerical study of the θ term

- ☆ Monte Carlo simulation of the lattice gauge theory with a θ term
 - θ term is purely imaginary \rightarrow the action “S” is complex
 - impossible to interpret Boltzmann weight “ e^{-S} ” as a probability
 - \rightarrow sign problem
 - It arises in various cases, not only the θ term
 - finite density QCD, chiral fermion, real time dynamics, ...
- Many approaches
 - Lefschetz thimble, density of states, tensor renormalization group, ...
 - this work \rightarrow complex Langevin method

Complex Langevin method

complex Langevin method (CLM)

[G. Parisi (1983)] [J. R. Klauder (1983)]

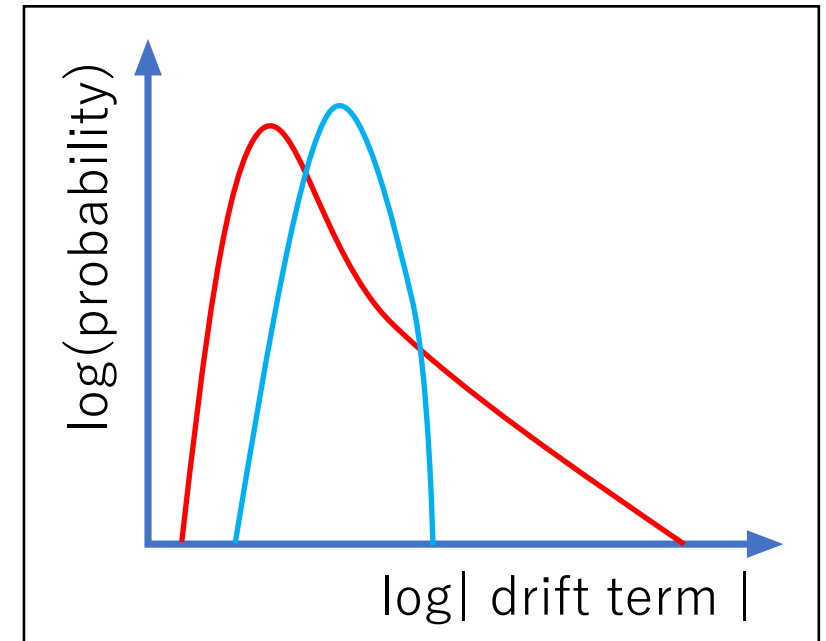
- Langevin equation: fictitious time evolution of dynamical variables
- real variable \rightarrow complex variable

$$\frac{dz(t)}{dt} = -\frac{\partial S(t)}{\partial z} + \eta(t) \quad x \mapsto z = x + iy$$

Diagram illustrating the Langevin equation components:

- The term $-\frac{\partial S(t)}{\partial z}$ is labeled as the "drift term".
- The term $\eta(t)$ is labeled as "Gaussian noise".

- do not use “probability” \rightarrow ~~sign problem~~
- condition required to be satisfied



The distribution of the drift term falls off exponentially or faster.

[K. Nagata, J. Nishimura, S. Shimasaki (2016)]

4D SU(2) gauge theory with a theta term

- simple example of a gauge theory with a θ term in 4D
- nevertheless it has a nontrivial phase structure at $\theta = \pi$

$$S = S_g + S_\theta$$

$$S_g = \frac{1}{2g^2} \int d^4x \text{Tr} [F_{\mu\nu} F_{\mu\nu}] \quad S_\theta = -i\theta Q$$

- topological charge

$$Q = \frac{1}{32\pi^2} \int d^4x \epsilon_{\mu\nu\rho\sigma} \text{Tr} [F_{\mu\nu} F_{\rho\sigma}]$$

integer value on a compact manifold

Lattice regularization

- **kinetic term:** standard Wilson action

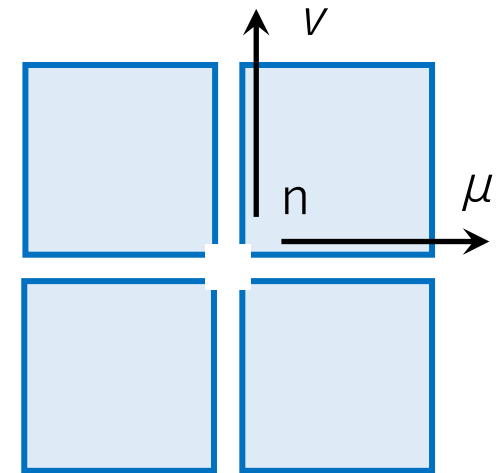
$$S_\beta = -\frac{\beta}{4} \sum_n \sum_{\mu \neq \nu} \text{Tr} [P_n^{\mu\nu}] \quad P_n^{\mu\nu} : \text{plaquette}$$

- **θ term:** clover leaf (symmetrized “figure 8”)

[P. Di Vecchia, K. Fabricius, G. C. Rossi, G. Veneziano (1981)]

$$Q_{\text{clov}} = -\frac{1}{32\pi^2} \sum_n \frac{1}{16} \sum_{\mu, \nu, \rho, \sigma=1}^4 \epsilon_{\mu\nu\rho\sigma} \text{Tr} [\bar{P}_n^{\mu\nu} \bar{P}_n^{\rho\sigma}]$$

$$\bar{P}_n^{\mu\nu} = P_n^{\mu\nu} - P_n^{-\mu\nu} - P_n^{\mu-\nu} + P_n^{-\mu-\nu} : \text{clover leaf}$$



Application of CLM

- discretized **complex Langevin equation** for the link variable $U_{n,\mu}$

$$U_{n,\mu}(t + \epsilon) = U_{n,\mu}(t) \exp(-i\epsilon D_{n,\mu} S(t) + i\sqrt{\epsilon} \eta_{n,\mu}(t))$$

$$U_{n,\mu} \in \text{SL}(2, \mathbb{C})$$

drift term

- gauge group is extended: $\text{SU}(2) \rightarrow \text{SL}(2, \mathbb{C})$

$$U_{n,\mu}^\dagger \rightarrow U_{n,\mu}^{-1}$$

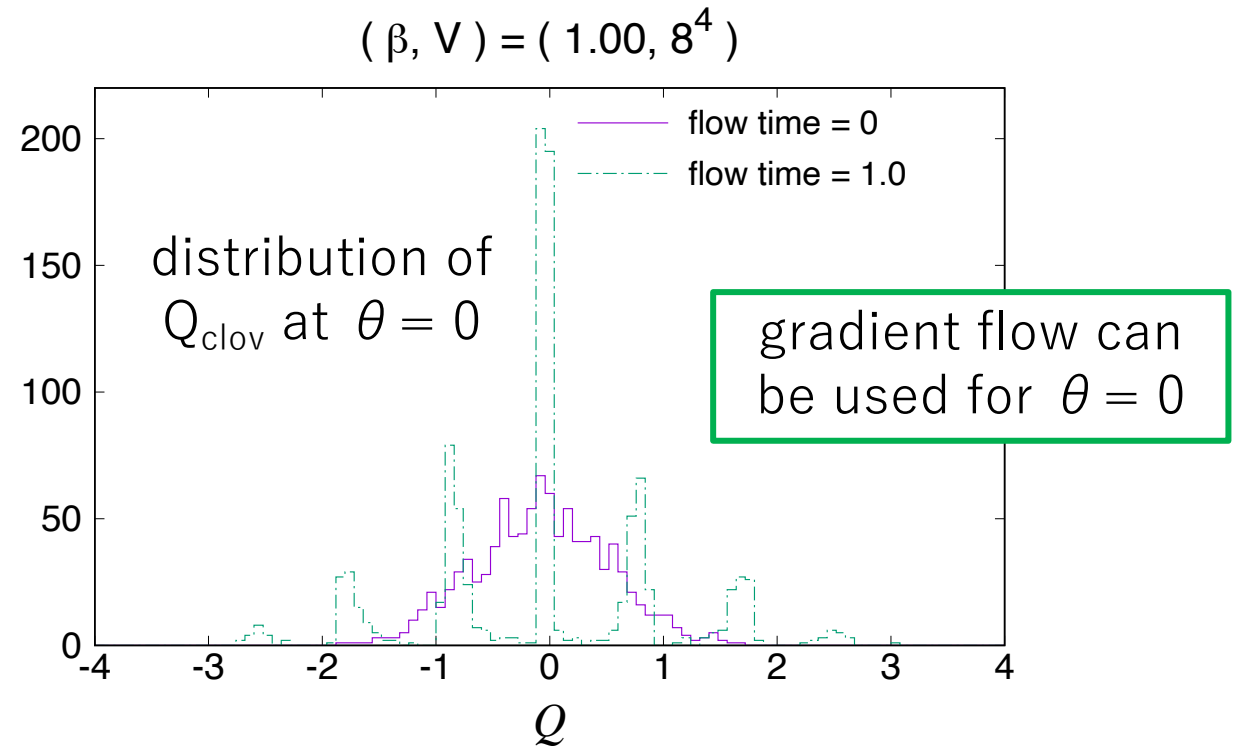
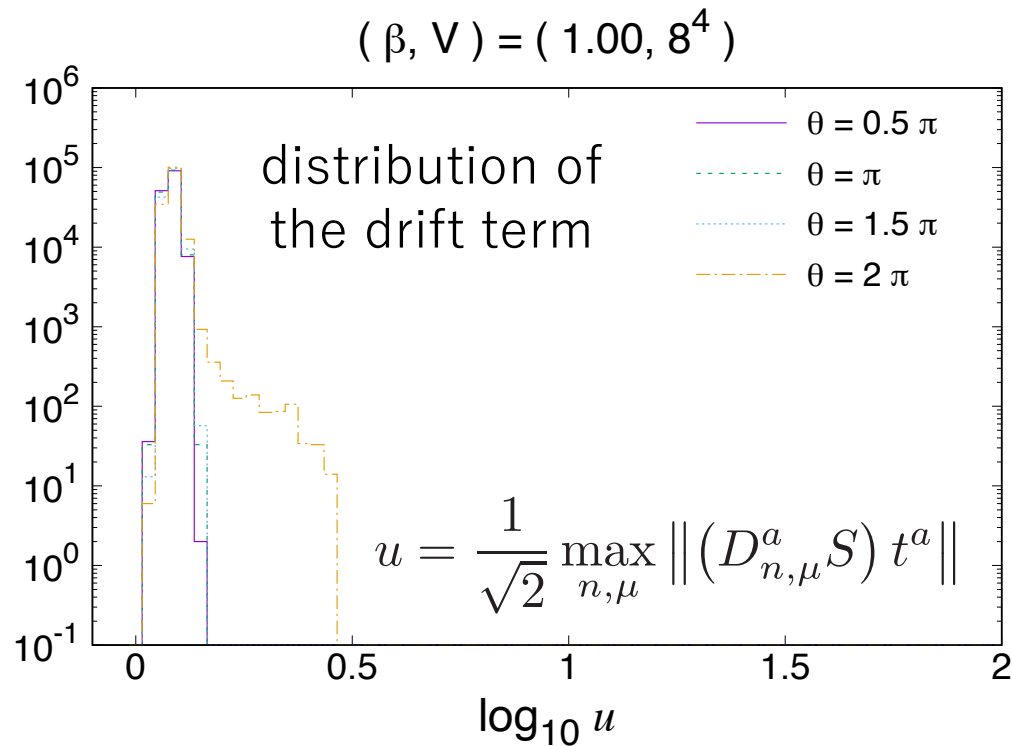
- drift term and observables have to respect **holomorphicity**
- control the non-unitarity by **gauge cooling**
 - gauge transformation to keep the link variable close to unitary
 - not affect gauge invariant observables

[K. Nagata, J. Nishimura, S. Shimasaki (2016)]

Validity test of CLM

- The condition for the correct convergence is satisfied without topology freezing !

☆ different from 2D U(1) gauge theory



Observable

☆ topological charge

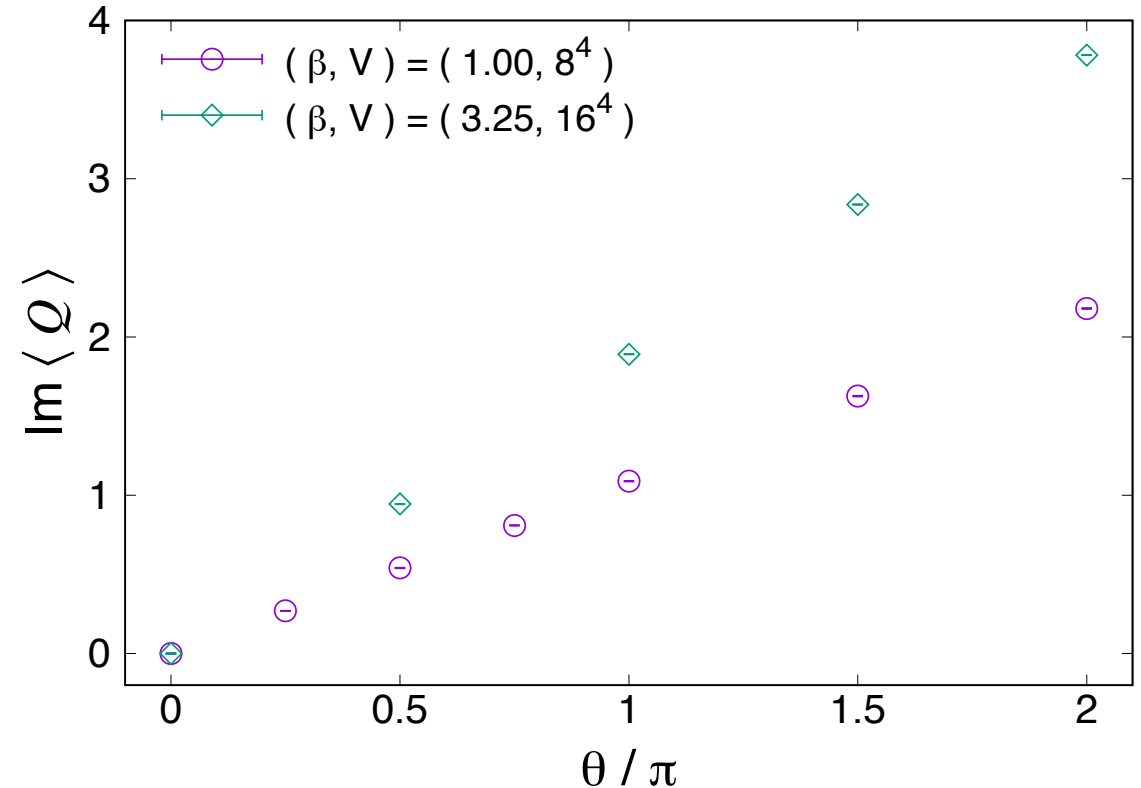
$$\langle Q \rangle = -i \frac{1}{Z} \frac{\partial Z}{\partial \theta}$$

- linear dependence on θ

2π -periodicity is absent



Q_{clov} is not an integer
due to short-range fluctuation

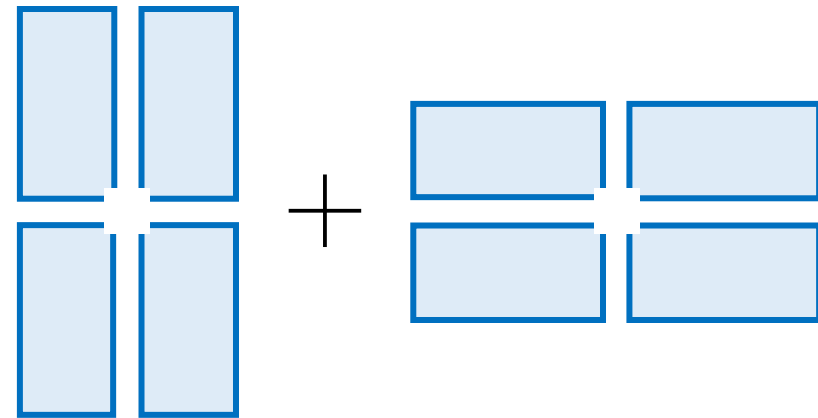


Strategy to recover “topology on the lattice”

- gradient flow / cooling
→ may not be justified for CLM ($\theta \neq 0$)
[L. Bongiovanni, G. Aarts, E. Seiler, D. Sexty (2014)]
- approach the continuum limit (increase V and β)
→ We increased V to 32^4 but Q_{clov} is still not close to an integer.

- improve the action by introducing 1×2 and 2×1 Wilson loops

☆ ongoing work



[Y. Iwasaki (1983)] [P. Weisz (1983)]

Summary

- The recent work on 't Hooft anomaly matching for 4D SU(2) YM predicted a nontrivial phase structure at $\theta = \pi$.
- We use the complex Langevin method to simulate the theory with the θ term, avoiding the sign problem.
- CLM for 4D SU(2) on the torus works without topology freezing unlike CLM for 2D U(1).
- However, further improvement is necessary to recover the topological property of the θ term on the lattice.

Future prospect

- We are now trying to improve the action by introducing rectangular Wilson loops.
- The 2π -periodicity of θ will be recovered if the topological charge close to an integer.
→ However, it is possible that CLM does not work in that case.
- We expect that introducing a puncture makes the convergence of CLM better.

Thank you!

Approach to complex action systems

➤ Reweighting method

- treat the phase of e^{-S} as an observable
- does not work if the phase oscillates rapidly

➤ Lefschetz thimble method

- reduce the phase oscillation by deforming the integral path from the real axis to the complex plane

➤ Complex Langevin method

- low computational cost
- has to meet a condition to justify the result

➤ Tensor renormalization group, Density of state, ...