

# Complex Langevin analysis of 2D U(1) gauge theory on a torus with a $\theta$ term

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# Theta term and the sign problem

$$S = S_G - i\theta Q$$

- The  $\theta$  term, which has many physical significances, is purely **non-perturbative**, and is the source of the **sign problem** in Monte Carlo simulations.
- Many approaches can be considered to overcome such problem, including **Lefschetz thimble**, **density of states**, **tensor renormalization group**, **complex Langevin method**, among others.
- Complex Langevin method (CLM) is known to be of **low computational costs** and is **straightforward to implement** – our approach in this work.

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# Complex Langevin method (CLM)

(G. Parisi, 1984; J.R.Klauder, 1983)

Degrees of freedom:

$$x \in \mathbb{R} \longrightarrow z \in \mathbb{C}$$

Observables:

$$O(x) \longrightarrow \mathcal{O}(z)$$

holomorphic  
function

Expectation values:

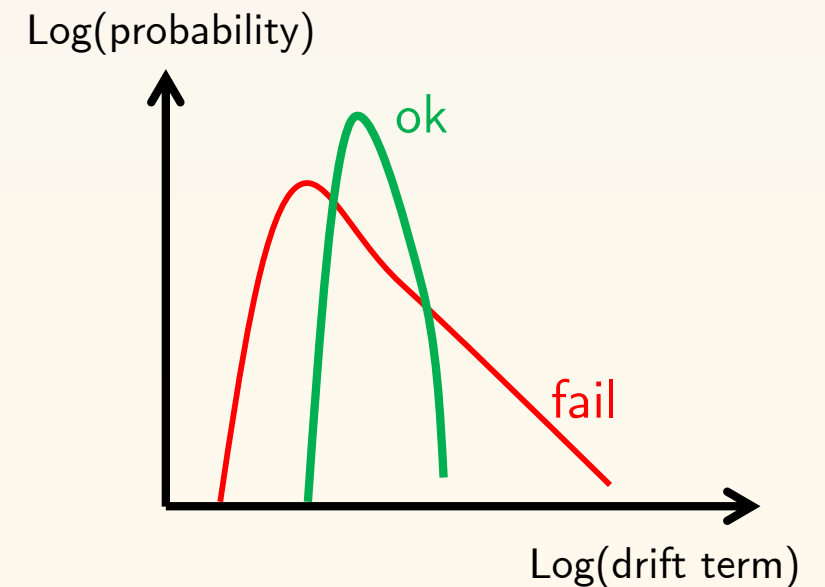
$$\langle O \rangle = \frac{\int dx O(x) P(x)}{\int dx P(x)} \longrightarrow \langle \mathcal{O} \rangle = \frac{1}{\Delta T} \int_{T_0}^{T_0 + \Delta T} dt \mathcal{O}(z(t))$$

Complex Langevin equation

$$\frac{dz(t)}{dt} = - \underbrace{\frac{\partial S(t)}{\partial z}}_{\text{Drift term}} + \eta(t)$$

Gaussian noise

Criterion of correct convergence



(Nagata, Nishimura, Shimasaki, 2016)

# 2D U(1) gauge theory with a theta term

This model is analytically solvable = good testing ground

$$S = S_G - i\theta Q$$

Lattice regularization

$$\vec{x} \rightarrow a\vec{n}$$

$$U_{n,\mu} = e^{iaA_{n,\mu}}$$

$$P_n = U_{n,1}U_{n+1,2}U_{n+2,1}^{-1}U_{n,2}^{-1} \simeq e^{ia^2F_{n,12}}$$

$$S_G = -\frac{\beta}{2} \sum_n (P_n + P_n^{-1})$$

Topological charge

“Log” definition  $Q_{\log} := \frac{1}{2\pi i} \sum_n \log P_n$

“Sine” definition  $Q_{\sin} := \frac{1}{4\pi i} \sum_n (P_n - P_n^{-1})$

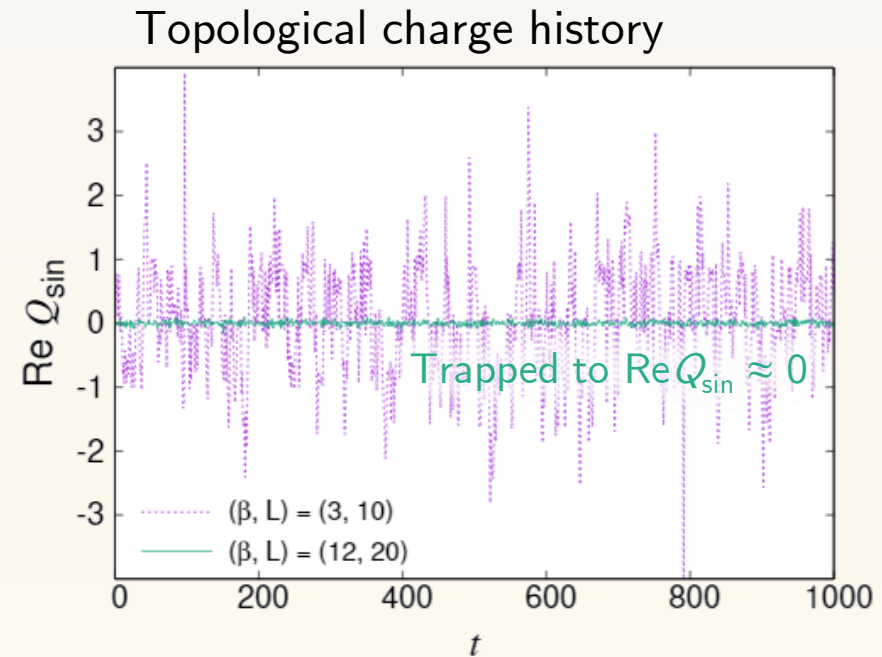
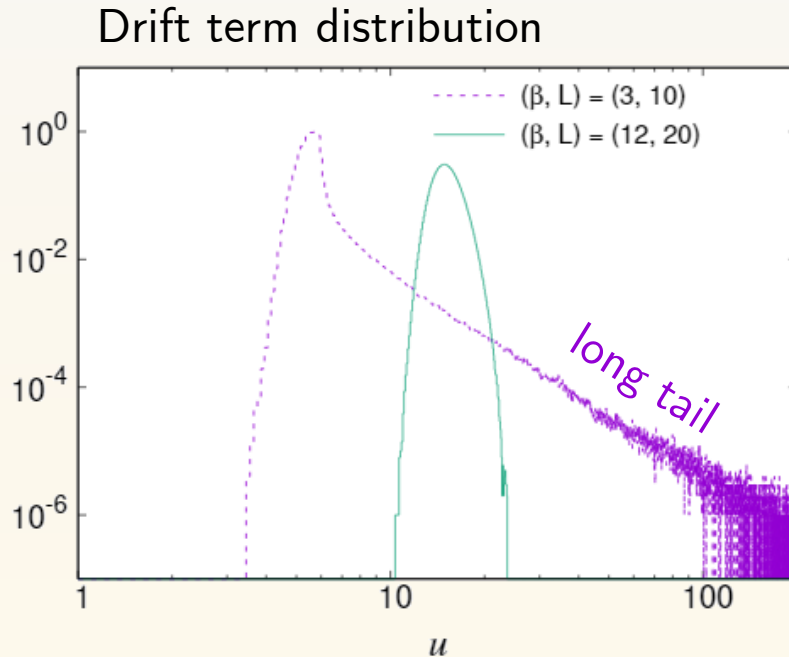
Langevin time evolution

$$U_{n,\mu}(t + \epsilon) = U_{n,\mu}(t) \exp(-i\epsilon D_{n,\mu} S(t) + i\sqrt{\epsilon} \eta_{n,\mu}(t))$$

# Result of the naive implementation (with sine definition)

Two cases: **small  $\beta$**  and **large  $\beta$**  (with fixed physical volume)

$$\theta = \pi$$



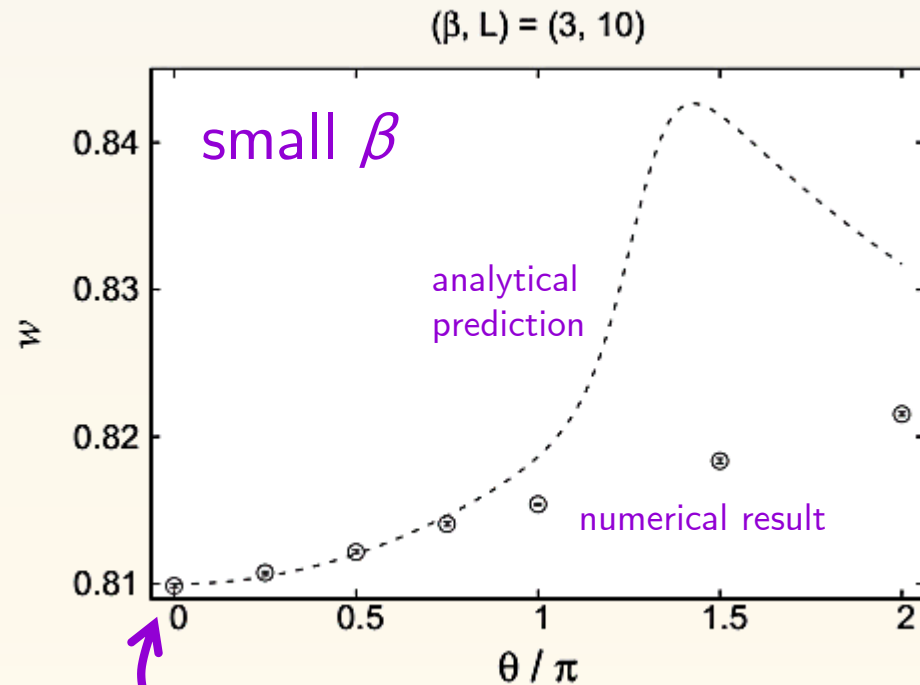
For **small  $\beta$**  : **incorrect convergence** (long tail in the drift term histogram)

For **large  $\beta$**  : **topological freezing** (trapped in a single topological sector)

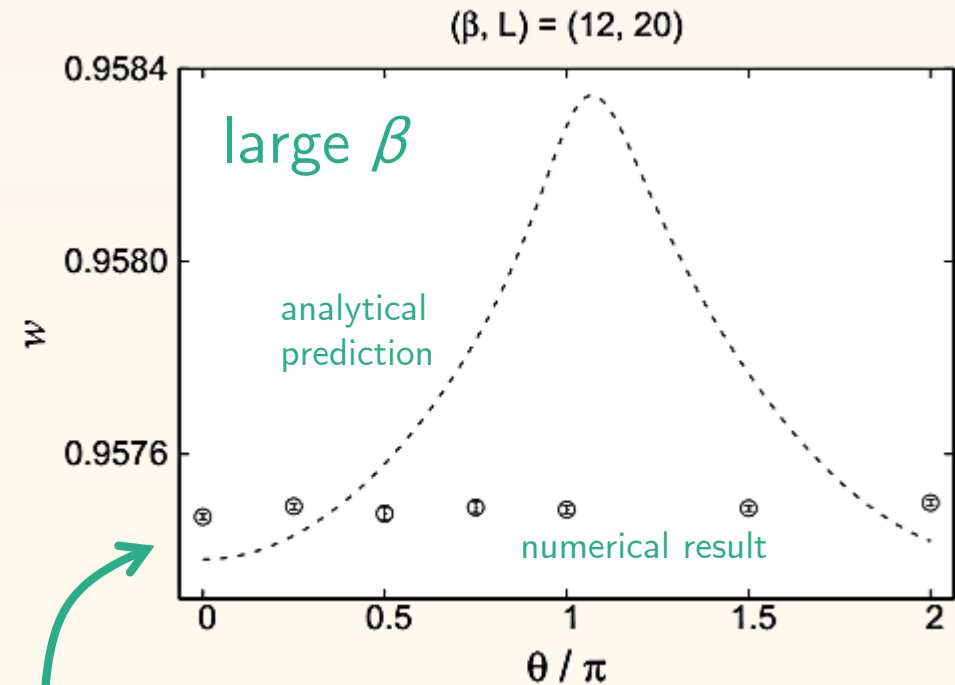
There is no intermediate  $\beta$  where both problems disappear!

# Observable: average plaquette

$$w = \frac{1}{V} \frac{\partial}{\partial \beta} \log Z = -\frac{1}{\beta V} \langle S_G \rangle$$



no singular drift problem here



incorrect even at  $\theta = 0$  because of the freezing problem

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# Punctured model

- We can introduce a puncture on the lattice to avoid at least the freezing problem
- A puncture allows  $Q$  to change freely = no topological freezing

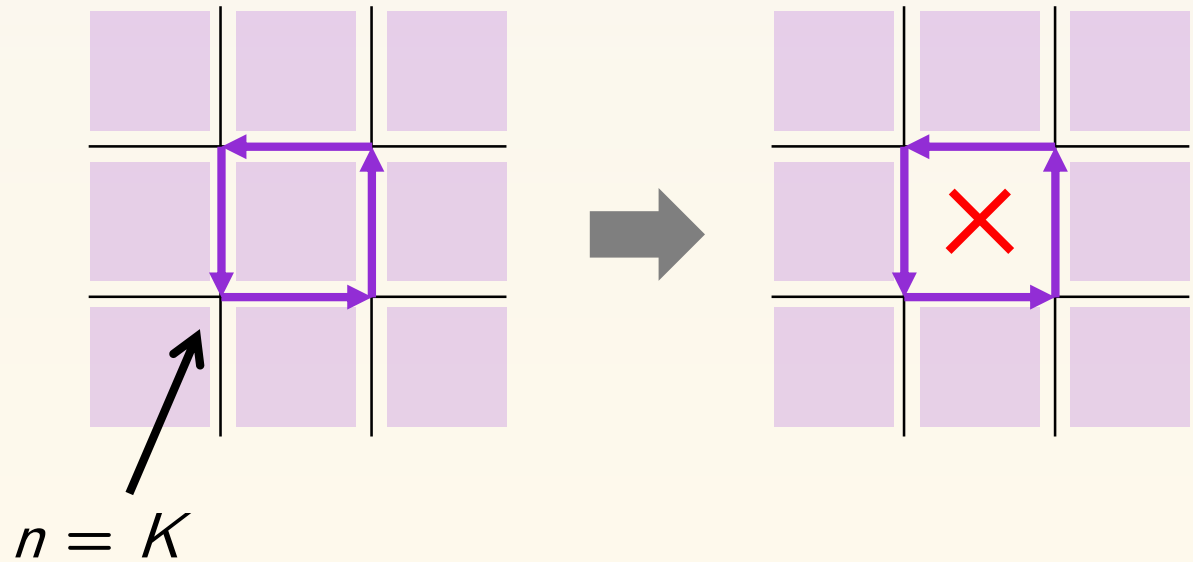
adding a **puncture** =  
removing a **specific plaquette**

Ex:

$$S_G = -\frac{\beta}{2} \sum_{n \neq K} (P_n + P_n^{-1})$$

$$Q_{\log} = \frac{1}{2\pi i} \sum_{n \neq K} \log P_n$$

charge is now a real number  
in any definition

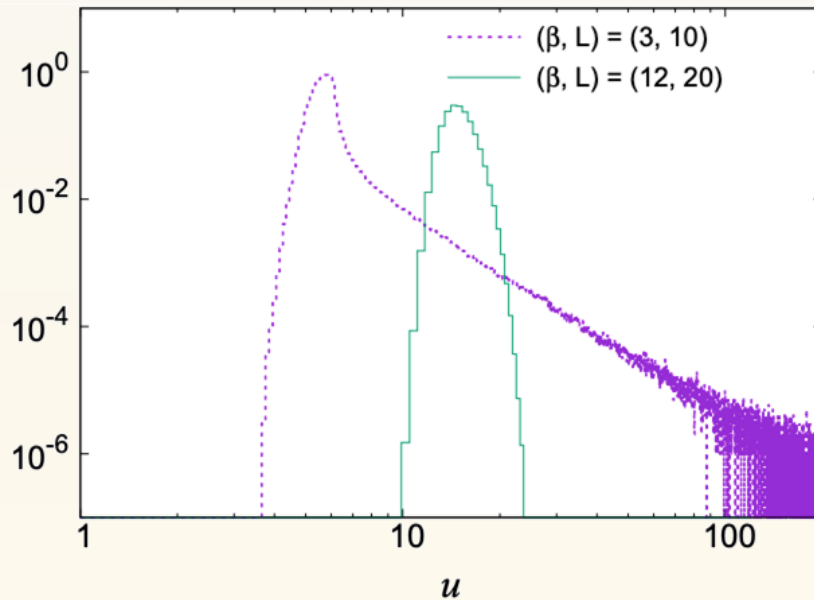


# Numerical results of the punctured model

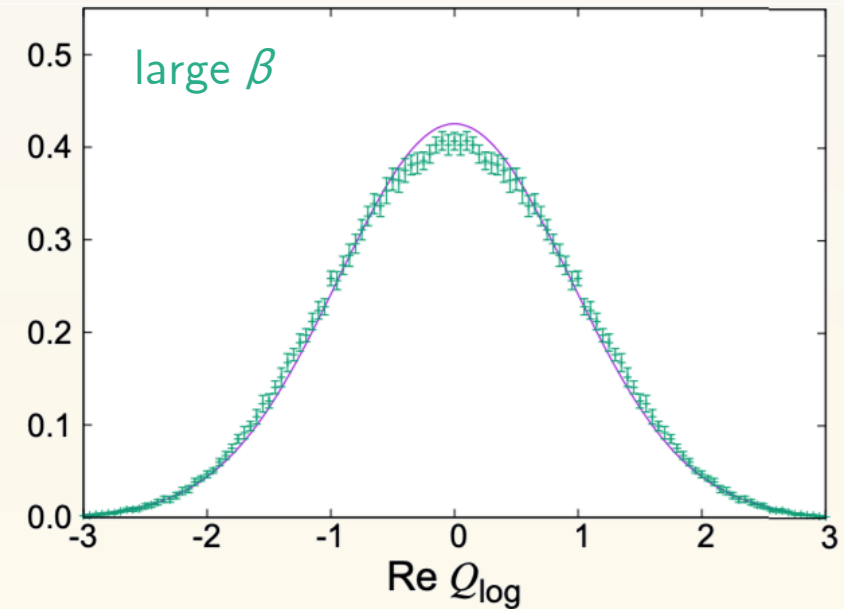
Two cases: **small  $\beta$**  and **large  $\beta$**  (with fixed physical volume)

$$\theta = \pi$$

Drift term distribution



Topological charge distribution

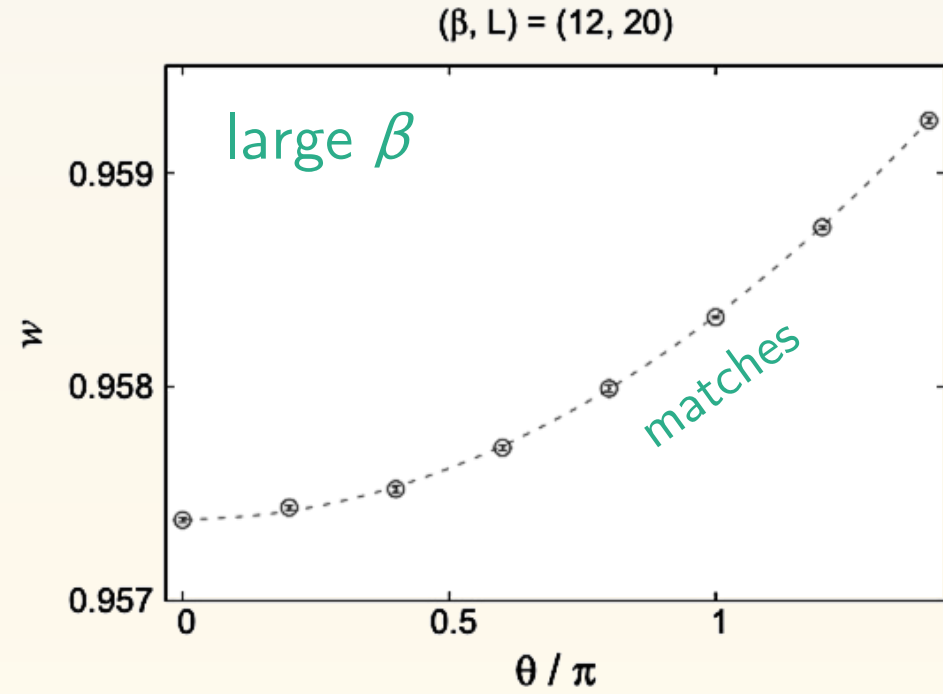
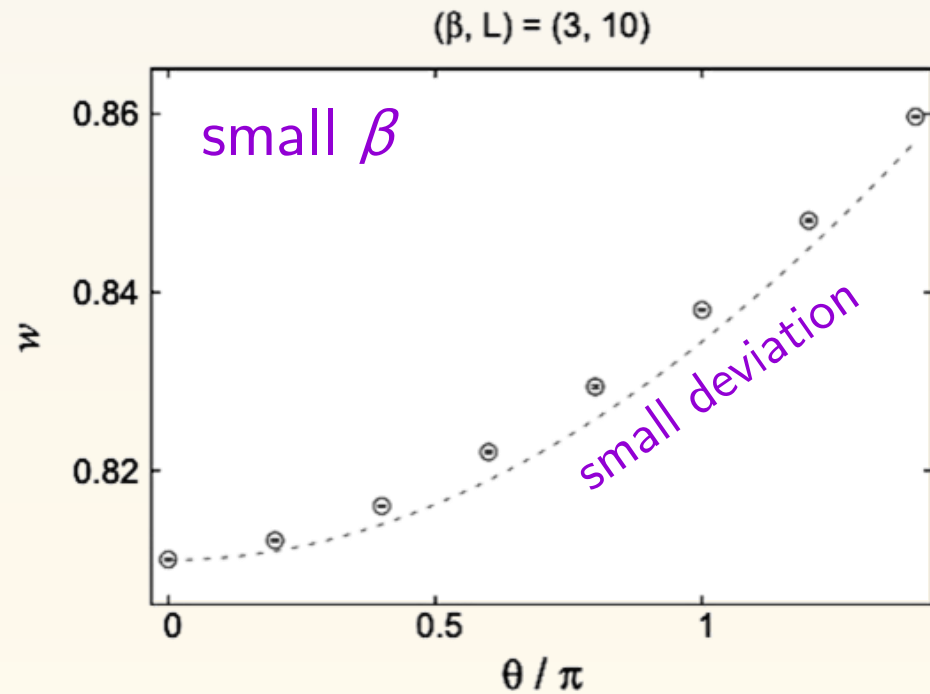


For **small  $\beta$**  : still have singular drift problem (see next slide)  
For **large  $\beta$**  : no more topological freezing **CLM works perfectly!**

← phenomenologically more interesting

# Observable: average plaquette

$$w = \frac{1}{V} \frac{\partial}{\partial \beta} \log Z = -\frac{1}{\beta V} \langle S_G \rangle$$



# Summary

- Monte Carlo simulation of gauge theory with a  $\theta$  term is difficult due to the **sign problem**
  - We use the **complex Langevin method** to study 2D U(1) gauge theory with a  $\theta$  term
- Naive implementation fails due to either the **incorrect convergence** or the **freezing problem**
- We introduce a **puncture** to avoid topological freezing
  - Both problems are resolved, giving correct results
  - The punctured model is equivalent to the infinite-volume limit of the original model for  $|\theta| < \pi$

# Ongoing work

## 4D SU(2) gauge theory with a $\theta$ term (next talk)

(K. Hatakeyama, M. Hirasawa, M. Honda, Y. Ito, A. Matsumoto, J. Nishimura, A.Y.)

Goal: confirming the nature of phase transition at  $\theta = \pi$  (SSB or gapless?)

(Gaiotto, Kapustin, Komargodski, Seiberg; 2017)

So far:

- Unlike 2D, CLM works well for small  $\beta$
- We now attempt to approach continuum limit (first without a puncture)

See next talk for more detail

Thank you!