

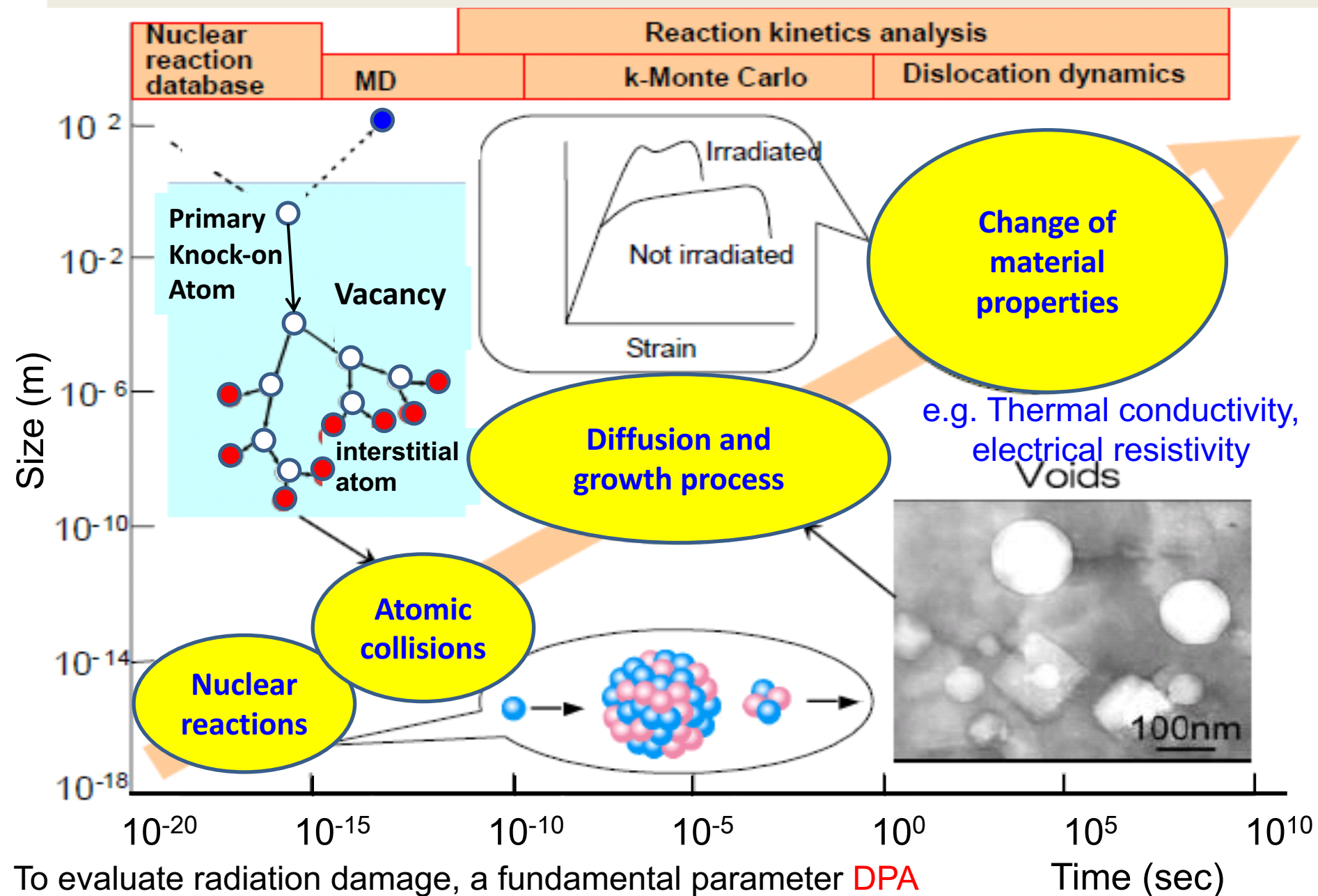
# Introduction of radiation damage calculation in PHITS for high-energy region

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Japan Atomic Energy Agency

- Outline
- Introduction
  - Displacement per atom (DPA) calculation method
  - Comparison between PHITS and other codes
  - RaDIATE examples with PHITS
  - Summary

# Scale of irradiation effect



To evaluate radiation damage, a fundamental parameter **DPA** that characterizes lattice displacement events is required.

# Microscopic effects on material

DPA: average number of displaced atoms per atom of a material

$$\text{DPA} = \int \sigma_{disp}(E) \phi(E) dE$$

$\sigma_{disp}$  : **displacement cross-section**  
 $\phi$  : irradiation fluence (particles/cm<sup>2</sup>)



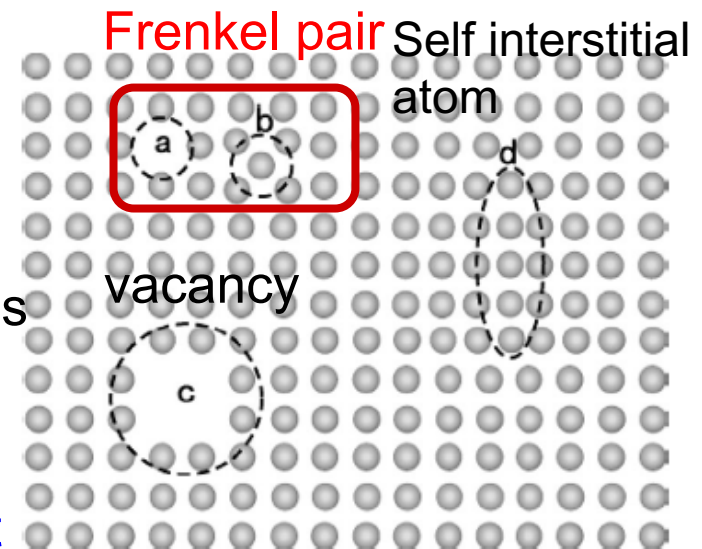
Related to the number of **Frenkel pairs**  $N_F$ :

$$\text{DPA} = \sum N_i N_F^i$$

$N_i$  = number of particles

a scale of radiation damage intensity.

DPA is used as a damage-based exposure unit  
and used to compare radiation damage by different radiation sources.



MC codes (PHITS, MARS, FLUKA, MCNP...) calculate DPA values:

1. Using database of  $\sigma_{disp}$  → limitation for kind of incident particles and materials
2. Calculating  $\sigma_{disp}$  for all particles with physics models event by event

This presentation shows radiation damage calculation with physics models in PHITS

# Physical Processes included in PHITS

T. Sato et al., JNST 50 (2013) 913.



**Transport**  
between  
collisions

**External Field**  
and Optical devices

**Ionization process**  
for charge particle

**Collision**  
with  
nucleus

**Low-energy Neutrons**  
**Photons, Electrons**

**High-energy nucleons**

**Heavy Ions**

- **Magnetic Field**
- **Gravity**
- **Super mirror (reflection)**
- **Mechanical devices, T0 chopper**

• **dE/dx : SPAR, ATIMA code**  
*Continuous-slowing-down  
Approximation (CSDA)*

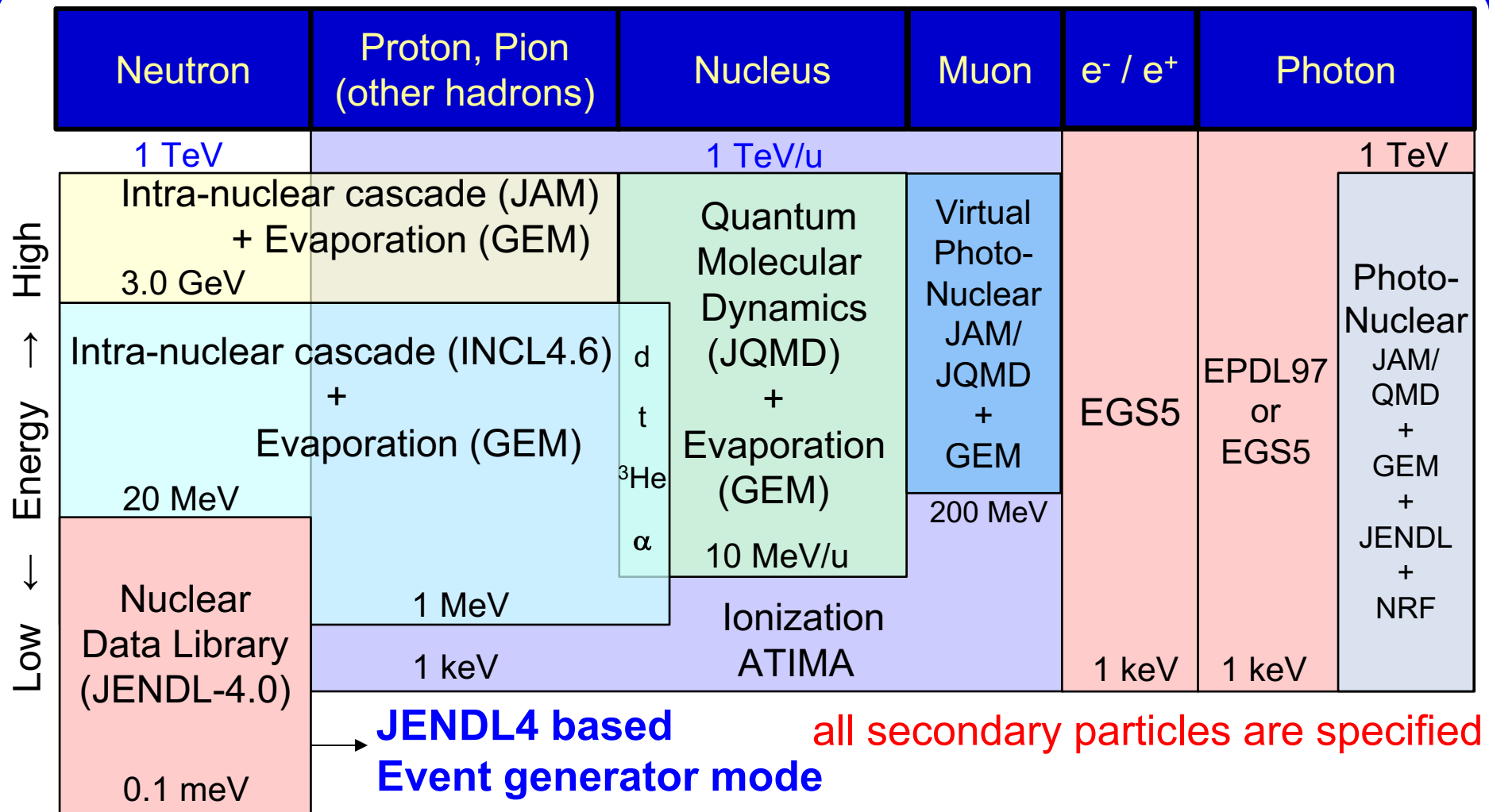
- **$\delta$ -ray generation**
- **Microdosimetric function**
- **Track-structure simulation**

**Nuclear Data based (JENDL-4.0 etc.)**  
**Event Generator Mode**  
**Electromagnetic cascade with EGS5**  
**Intra-Nuclear Cascade**

**Evaporation**

**Quantum Molecular Dynamics**

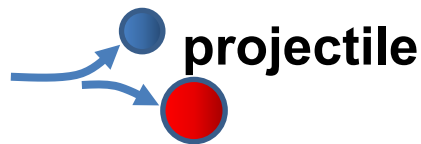
# Map of Models in PHITS



Next slide: DPA calculation method

# DPA calculation method in PHITS

## (1) Energy transfer with Coulomb scat.



target PKA

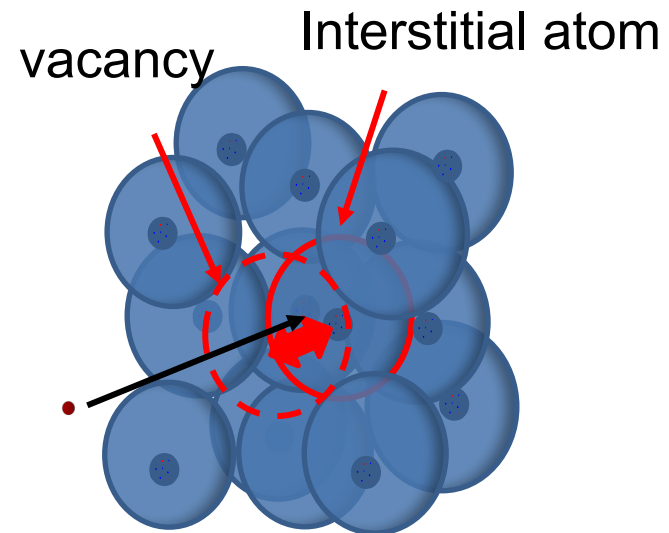
secondary



Nuclear  
reaction

target PKA

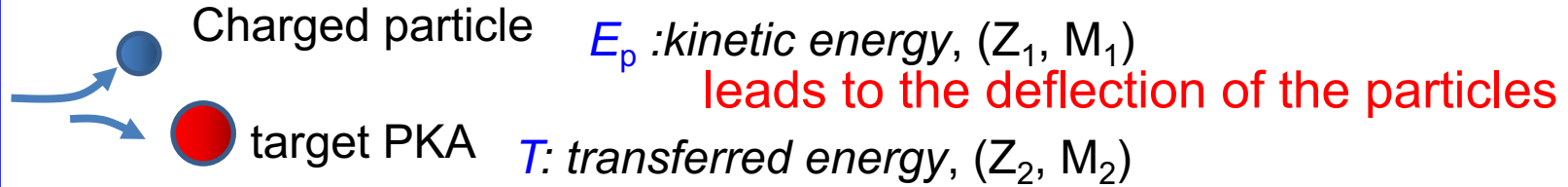
## (2) Cascade damage approximation



$$\sigma_{disp} = \int_{t_d}^{t_{max}} \underbrace{\frac{d\sigma_{Coul}(t)}{dt}}_{\text{Coulomb scattering}} \underbrace{\eta \frac{0.8}{2 \cdot T_d} T_{dam}}_{\text{Number of displacement atoms}} dt$$

# (1)Energy transfer with Coulomb scattering

M. Nastasi et al., "Ion-Solid Interaction: Fundamentals and Applications "



Coulomb scat. cross section: one parameter

$$d\sigma_{coul}(t) = \frac{\pi a_{TF}^2}{2} \frac{f(t^{1/2})}{t^{3/2}} dt$$

dimensionless collision parameter:

$$t \equiv \varepsilon^2 \frac{T}{T_{max}} = \varepsilon^2 \sin^2\left(\frac{\theta}{2}\right)$$

$T$ : Transferred energy to target atom

$T_{max}$ : maximum transferred energy

$$= \frac{4M_1M_2E_p}{(M_1 + M_2)^2}$$

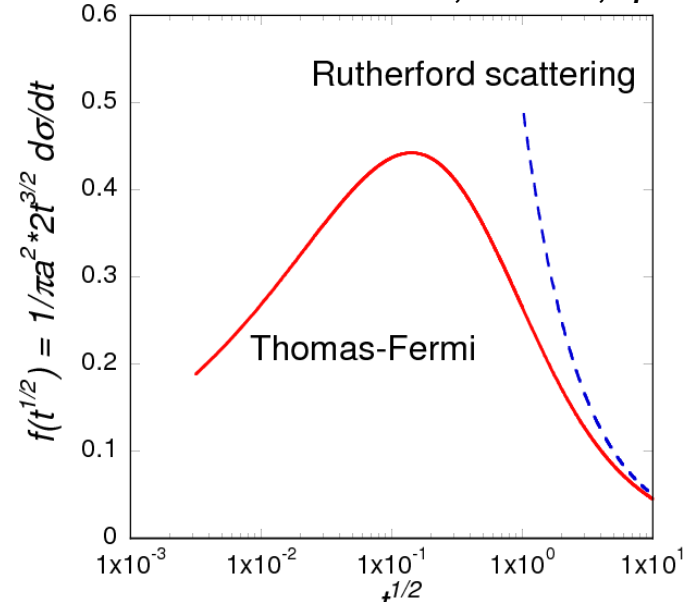
$\varepsilon$ : dimensionless energy

$$= \frac{Ea_{TF}M_2}{Z_1Z_2e^2(M_1 + M_2)}$$

Screening functions:

$$f(t^{1/2}) = \lambda t^{1/2-m} \left[ 1 + (2\lambda t^{1-m})^q \right]^{-1/q}$$

Thomas-Fermi  $\lambda=1.309$ ,  $m=1/3$ ,  $q=2/3$



Large  $t \rightarrow$  large  $T$  in close collisions

Small  $t \rightarrow$  small  $T$  in distance collisions

Next slide: cascade damage approximation

## (2) Cascade damage approximation

$$\sigma_{disp} = \int_{t_d}^{t_{max}} \frac{d\sigma_{Coul}(t)}{dt} \eta \frac{0.8}{2 \cdot T_d} T_{dam} dt$$

Defect production efficiency

$$T_{dam} = \xi(\varepsilon)T(\varepsilon) = \frac{1}{1 + k_{cas} \cdot g(\varepsilon)} T(\varepsilon)$$

Damage energy: transferred to lattice atoms reduced by the losses for electronic stopping atoms in the displacement cascade

$$g(\varepsilon) = \varepsilon + 0.40244 \cdot \varepsilon^{3/4} + 3.4008 \cdot \varepsilon^{1/6}$$

$$k_{cascade} = 0.1337 Z_{target}^{\frac{1}{6}} (Z_{target}/A_{target})^{1/2}$$

number of displaced atoms using phenomenological approach:  $N_{NRT}$

NRT(Norgett, Robinson, Torrens: 1975)

- 0.8: displacement efficiency derived from MD simulation of Robinson, Torrens 1972
- $T_d$ : threshold displacement energy. Bonds should be broken to displace an atom.  
e.g. set to 40 eV in Ti but varies 15 – 90 eV in other atom



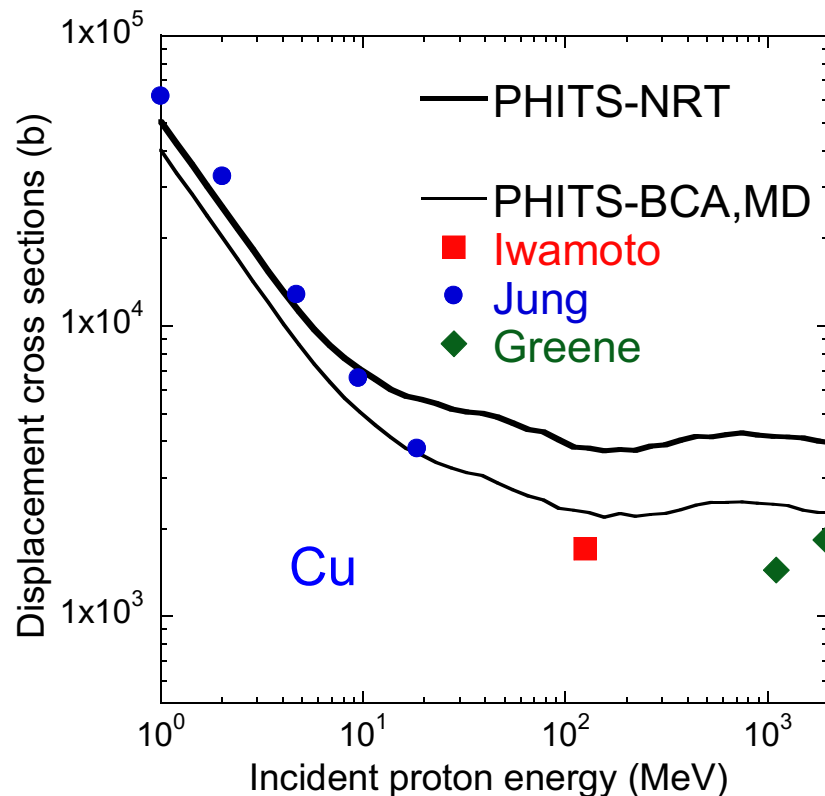
# Efficiency of the defect production in material

$$\eta = \frac{N_D}{N_{NRT}} = 0.7066 T_{dam}^{-0.437} + 2.28 \times 10^{-3} T_{dam} \quad \text{for Cu}$$

M.J. Caturla et al., J. Nucl. Mater. 296 (2001) 90.

$N_D$ : number of stable displacements at the end of collision cascade ← MD

$N_{NRT}$ : number of defects calculated by NRT model      Account for atom recombination in elastic cascading



$\sigma_{disp}$  with  $\eta$  reproduces experimental data in the high-energy region up to GeV for Cu.

Y. Iwamoto et al., JNM 458 (2015) 369.

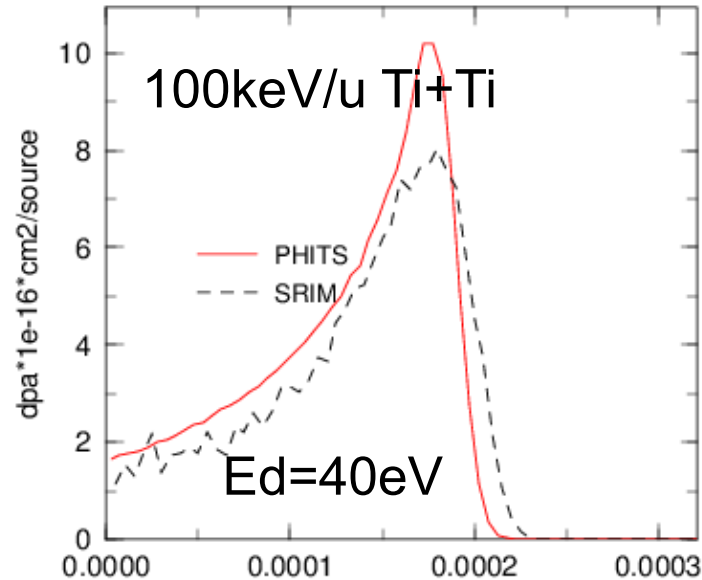
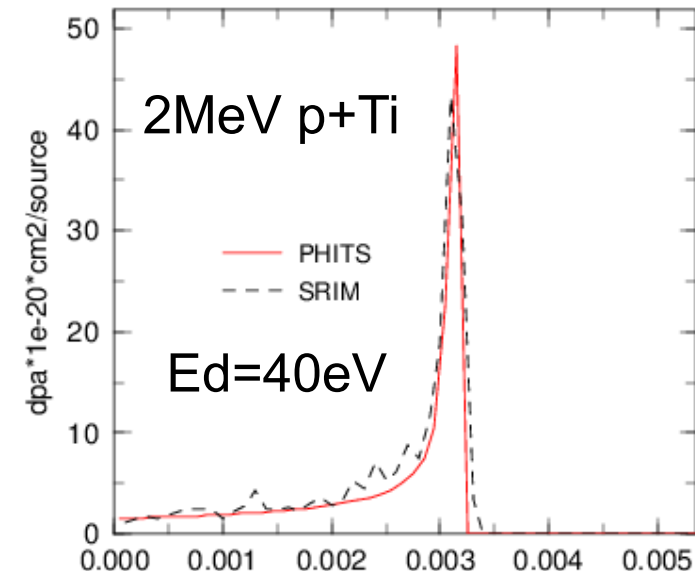
P. Jung, JNM 117 (1983) 70.

G.A. Greene et al., Proc. of AccApp'03 (2004) 881.

Note: there are new efficiencies proposed by Stoller and Nordlund.  
PHITS will include them soon.

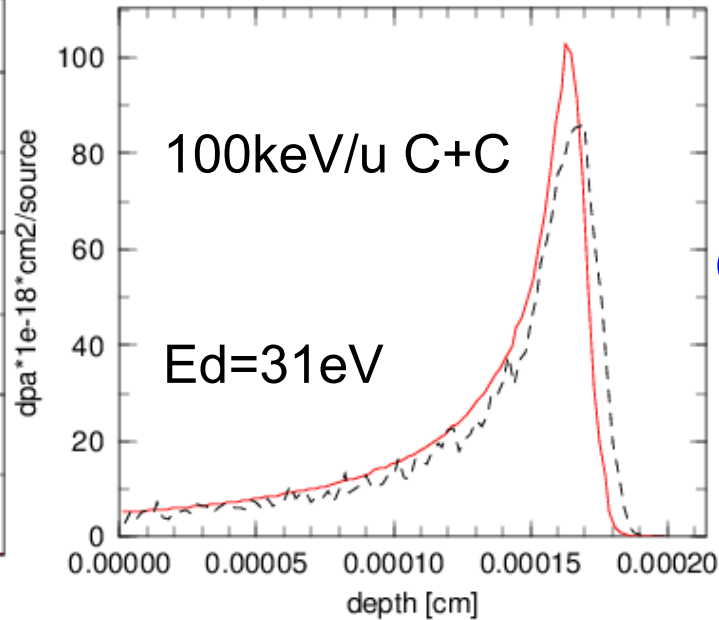
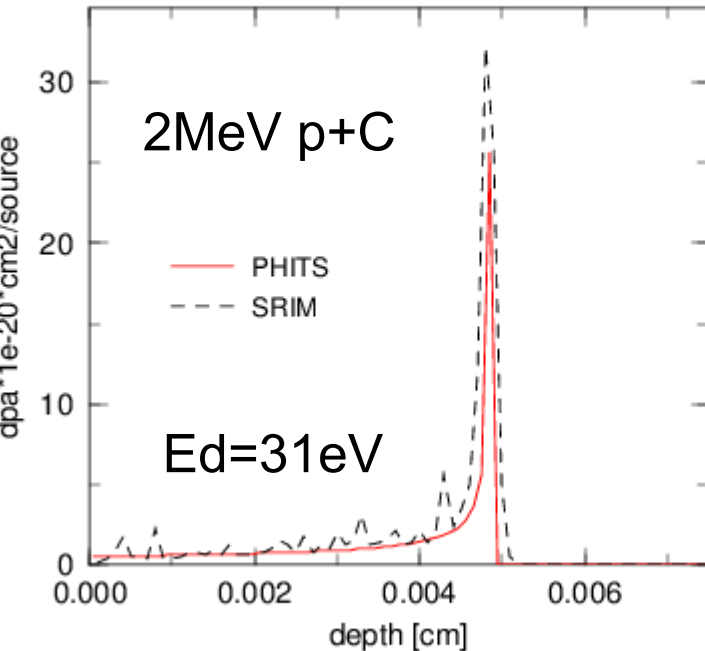
In this presentation, NRT-DPAs are reported.

# Comparison between PHITS and SRIM codes



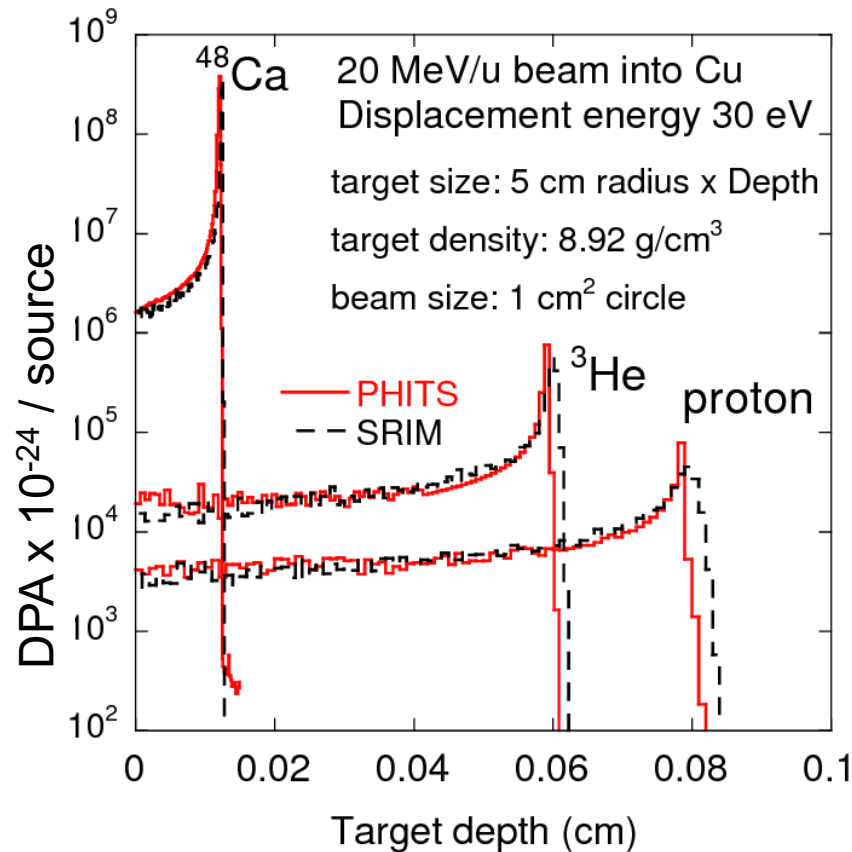
SRIM can simulate the transport of ions in matter without nuclear reaction.

<http://www.srim.org/>

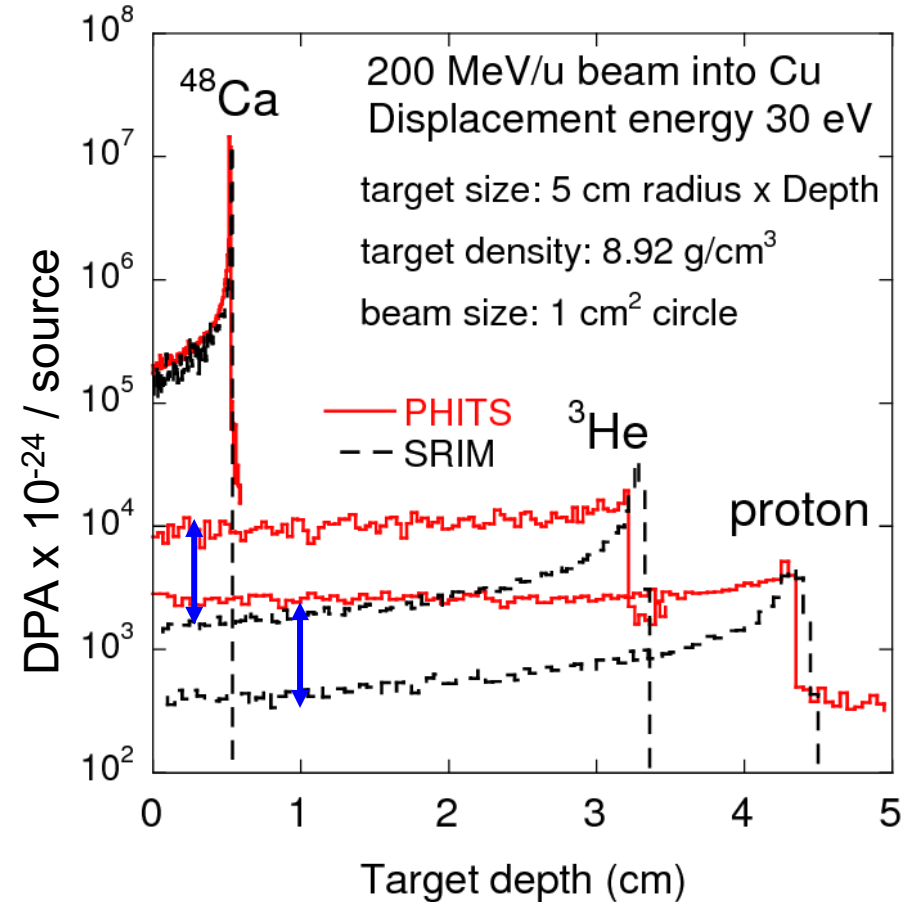


Good agreement.

# Comparison between PHITS and SRIM codes



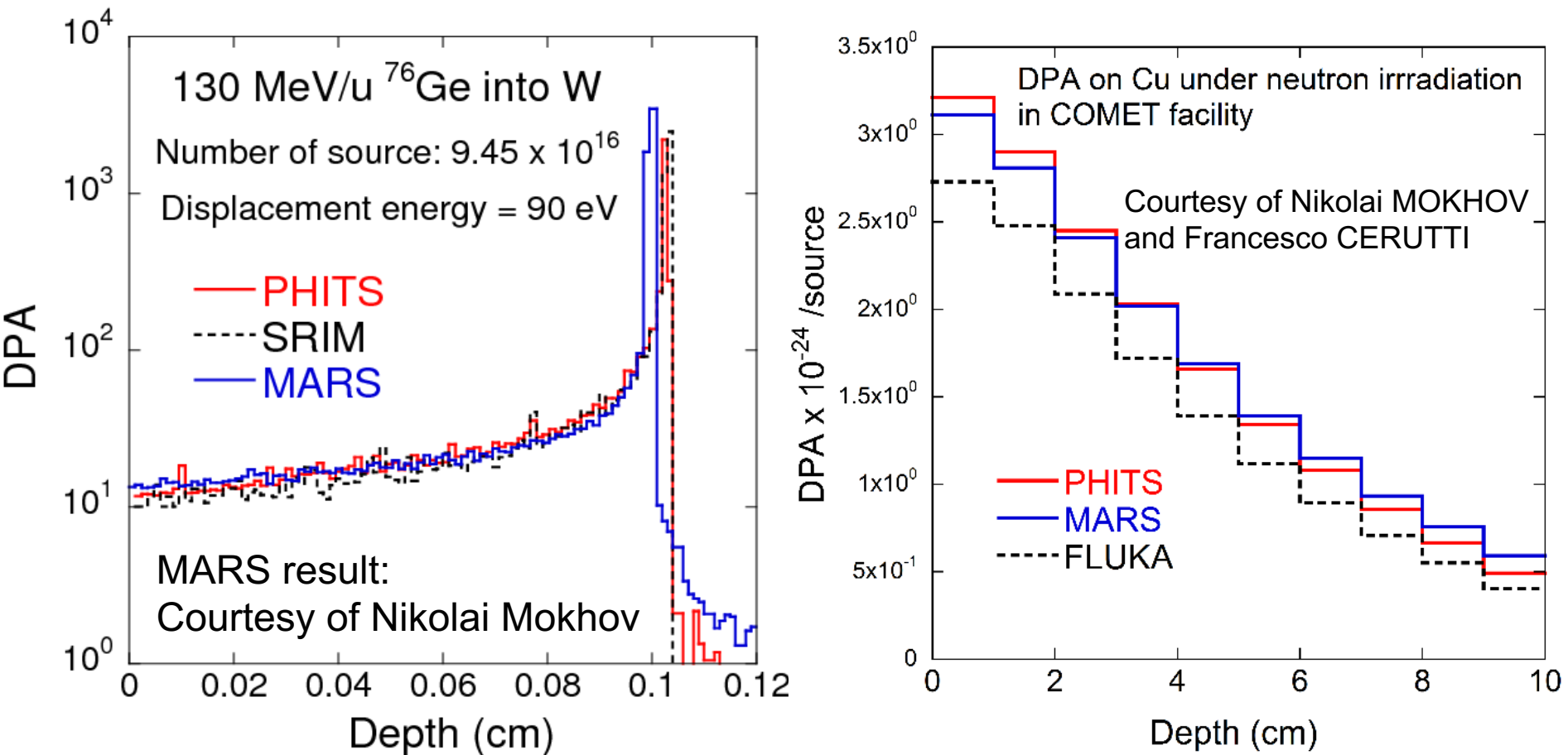
✓ PHITS results are in good agreement with SRIM results.



✓ differences between two results:

→ Secondary particles created from sequential nuclear reactions

# Comparison between PHITS, SRIM and MARS codes



Good agreement.

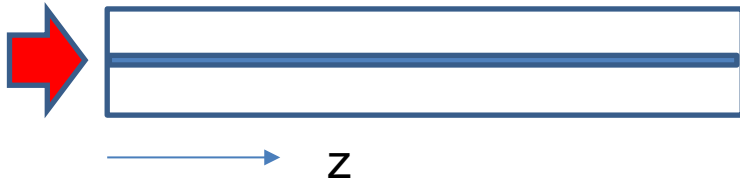
# RaDIATE example with PHITS

Proton gaussian beam

Energy: 0.18, 0.8, 3, 30, 120, 400 GeV  
(PHITS is not available at 7 TeV.)

proton

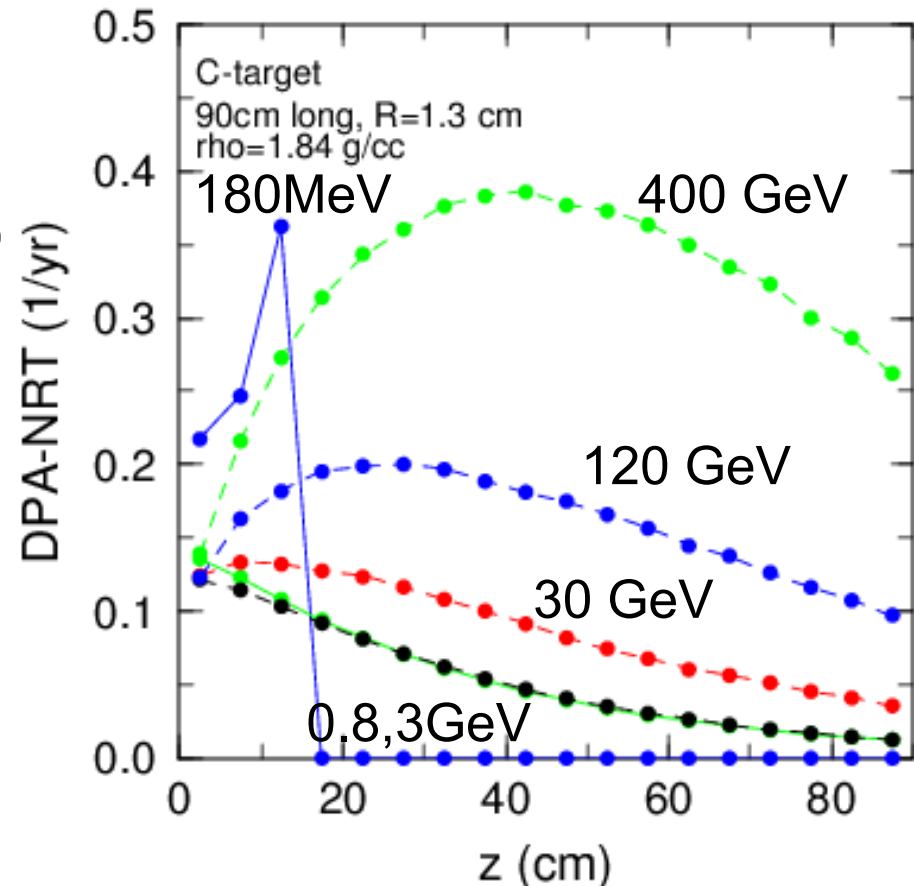
$$\sigma_x = \sigma_y = 0.43 \text{ cm}$$



90 cm long,  $R=1.3$  cm

C-target, density= $1.84 \text{ g/cm}^3$

Tally region: 90 cm long,  $r=0.2 \text{ cm}$



DPA depth distribution in C-target

180 MeV: Bragg peak of protons contributes DPA.

0.8-30 GeV: Damage may occur at surface.

120 GeV, 400 GeV: DPA increases with depth.

# Summary

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The displacement calculation method in PHITS is available for all incident particles, wide energy range (eV-TeV), and all materials.

Agreements between PHITS and other codes are good.

In the high energy region ( $> \sim 20$  MeV) for proton and neutron, DPA created by secondaries increase due to nuclear reactions.

## Future works

New defect production efficiency (Nordlund) will be implemented.  
(just implemented!)

Benchmark experiments will be performed at RCNP (100-400MeV) and J-PARC (400MeV-30GeV, see Meigo-san's talk).

# Experimental displacement cross section

Displacement cross section could be experimentally validated in Irradiation on metal at cryogenic temperature.

Recombination of Frenkel pairs by thermal motion is well suppressed.

Experimental displacement cross section

**Damage rate**

$$\sigma_{\text{exp}} = \frac{1}{\rho_{FP}} \frac{\Delta\rho_{\text{metal}}}{\phi}$$

J. Nucl. Mater. 49 (1973/74) 161.

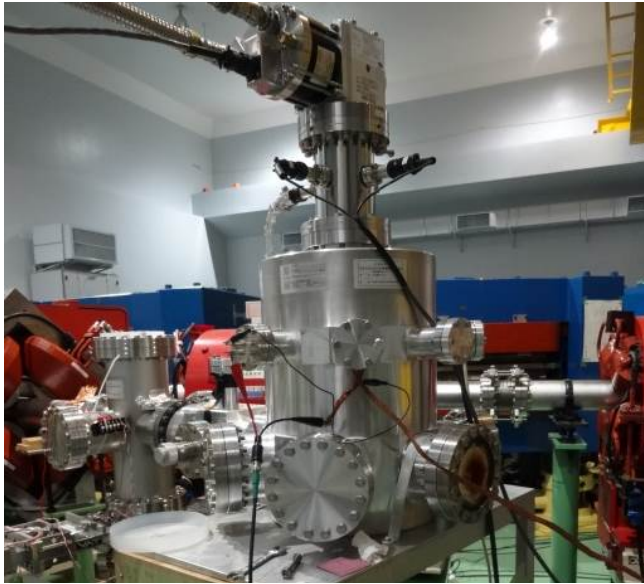
$\Delta\rho_{\text{metal}}$ : Electrical resistivity change( $\Omega\text{m}$ )

$\Phi$ : Beam fluence( $1/\text{m}^2$ )

$\rho_{FP}$ : Frenkel-pair resistivity ( $\Omega\text{m}$ )

Resistivity increase is the sum of resistivity per Frenkel pair

# Development of cryogen-free cooling system at FFAG/KURRI

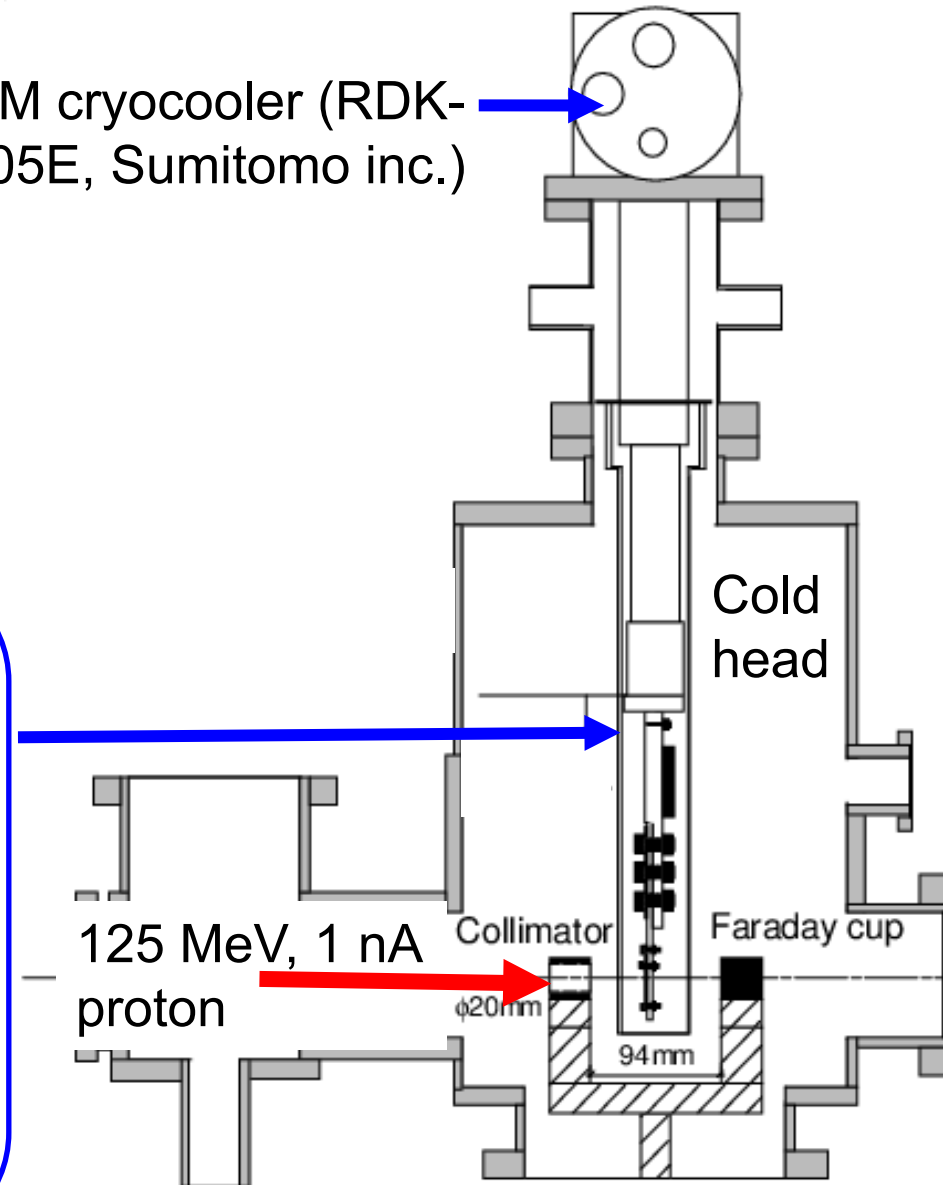


GM cryocooler (RDK-205E, Sumitomo inc.)



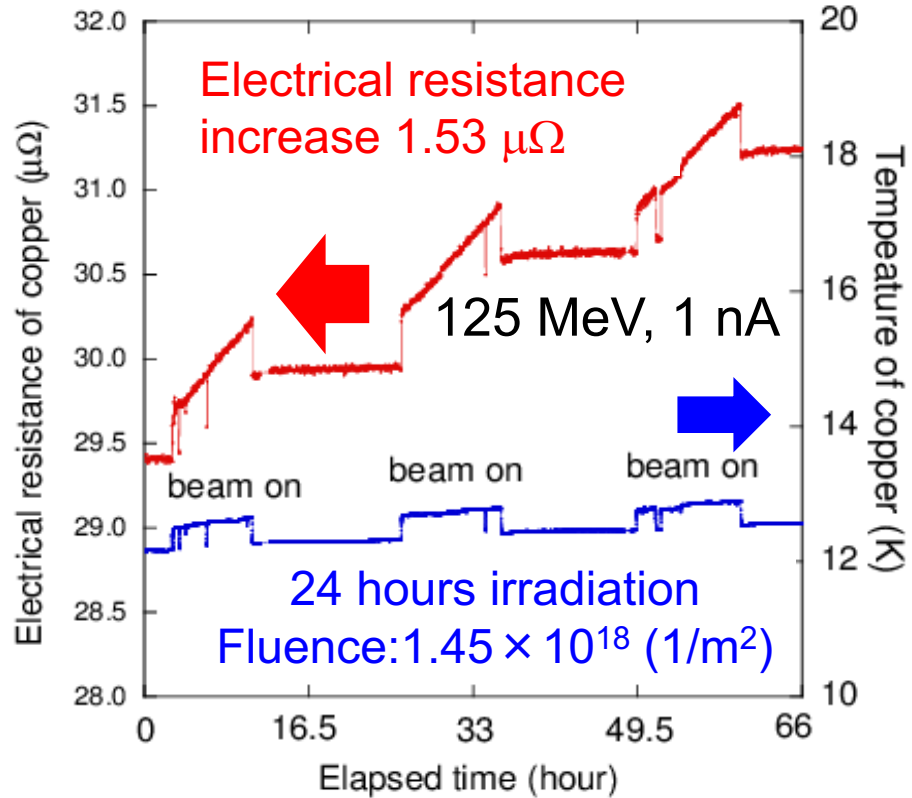
Sample was cooled by conduction coolant via Al and oxygen-free high-conductivity copper (OFHC).

Target assembly

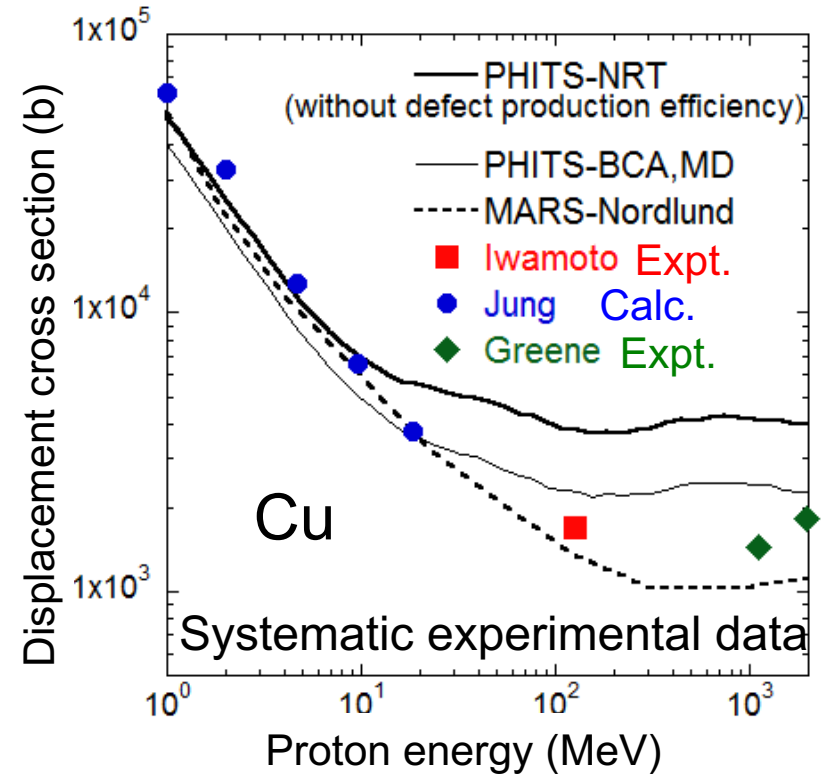




# Experimental results of copper



Electrical resistivity changes of copper



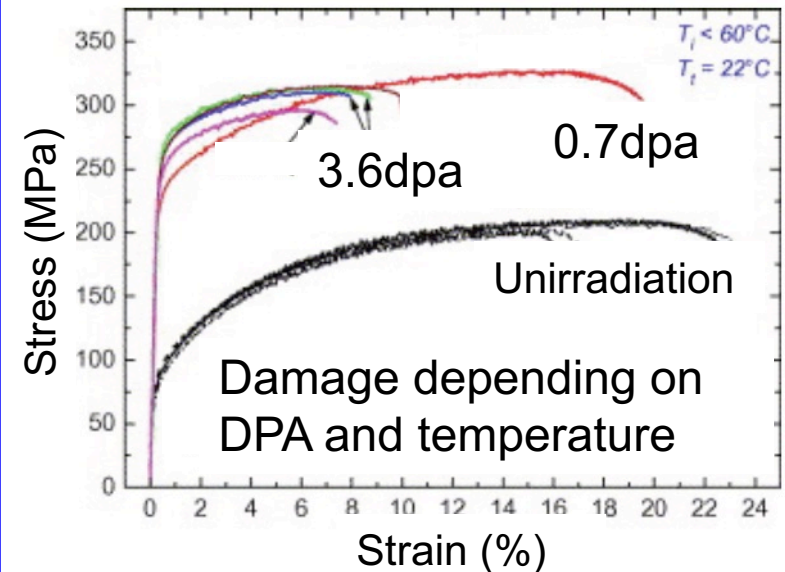
Displacement cross section of copper

# Evaluation of radiation damage

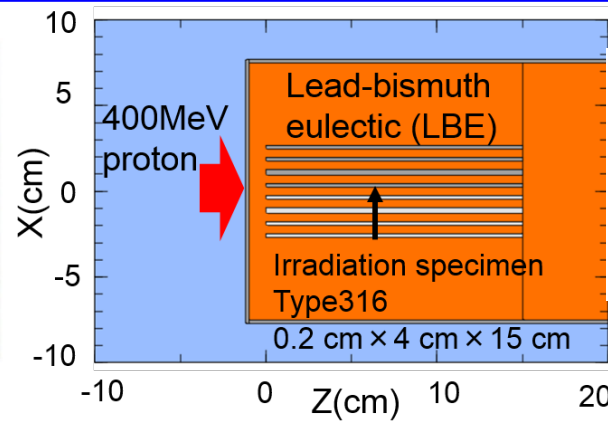
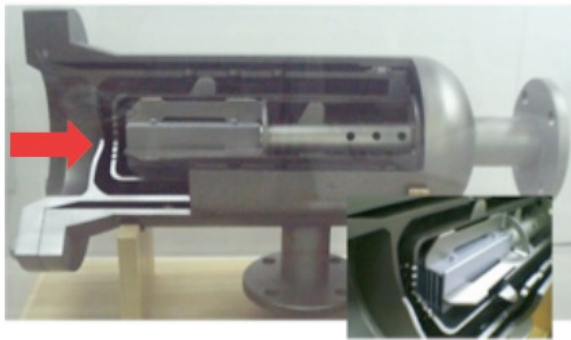
The average number of Displaced atoms Per Atom of a material (DPA) is used in evaluation of reactor and accelerator as a damage-based exposure unit.

$$\text{DPA} = \int \sigma_{\text{disp.}}(E) \phi(E) dE$$

Displacement cross section      Fluence

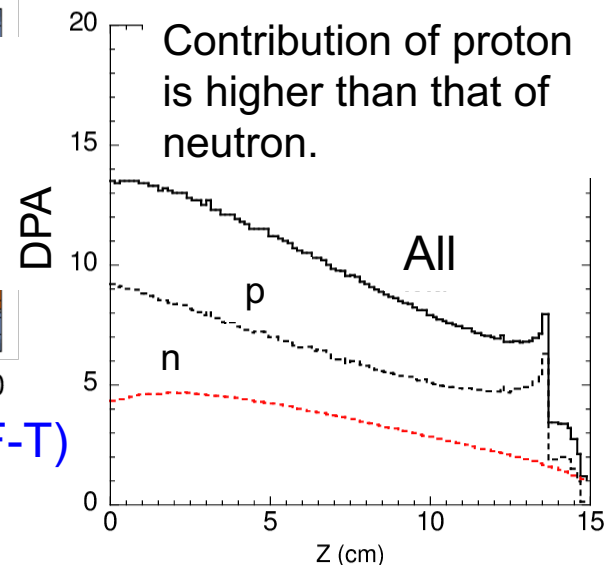


Damage depending on DPA and temperature  
Irradiation effect on AlMg3 for 600 MeV proton irradiation

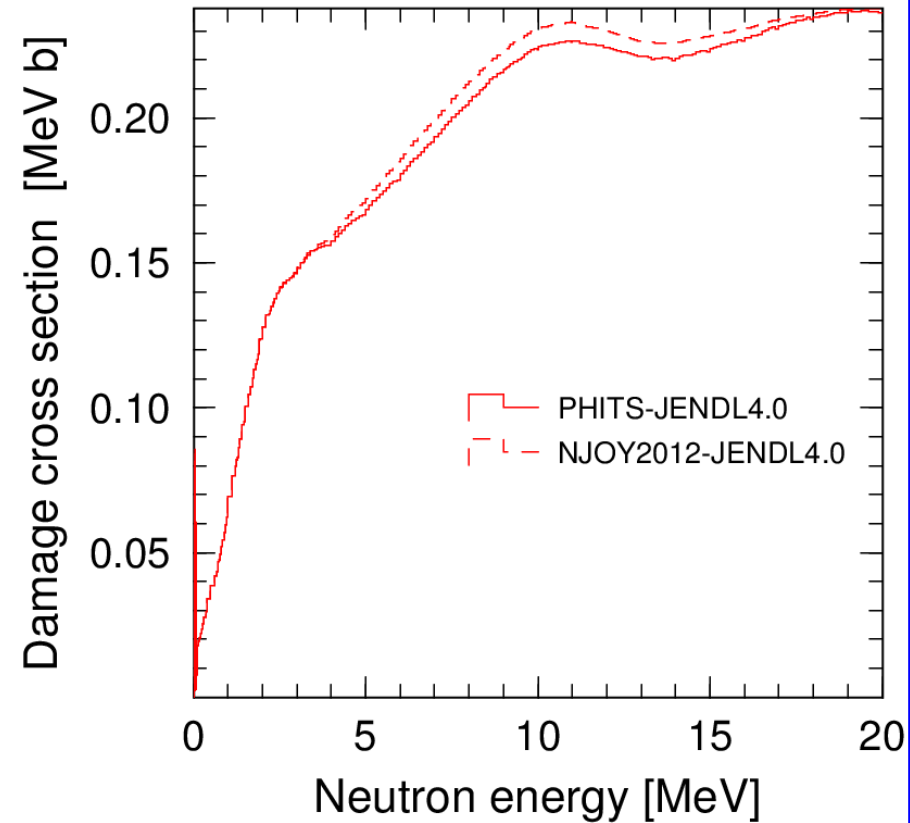
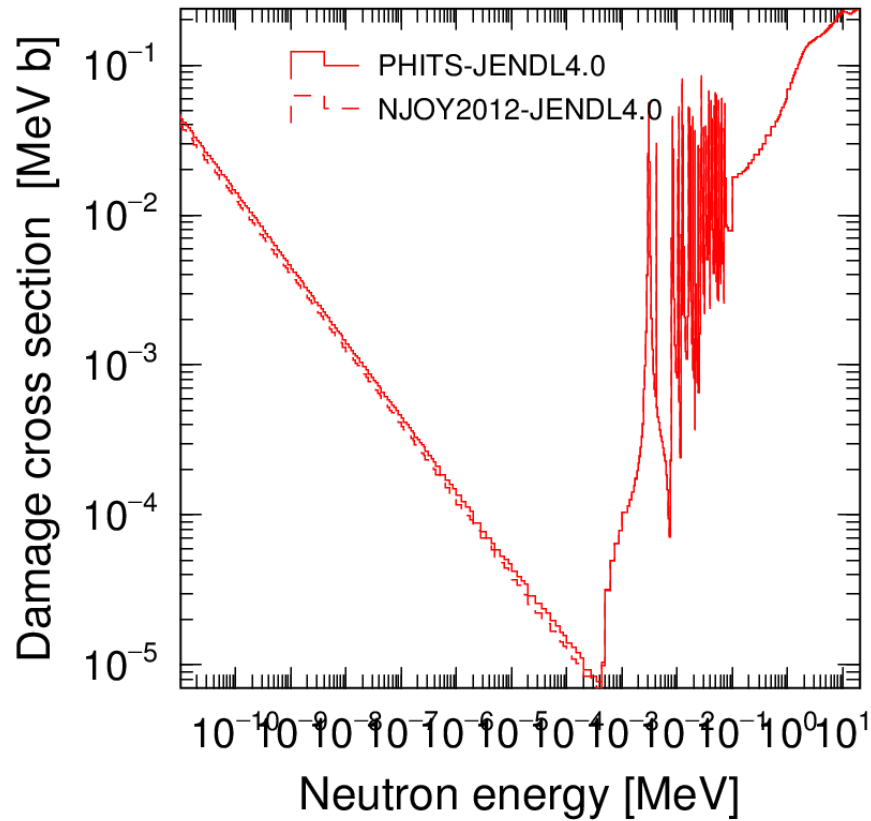


DPA for target at J-PARC ADS Target Test Facility (TEF-T) using Monte Carlo particle transport code PHITS

Y. Iwamoto, et al., J. Nucl. Sci. Technol 51 (2014) 98-107.



# ① JENDL-4.0を用いたPHITSとNJOY2012 の損傷エネルギー断面積の比較



$^{47}\text{Ti}$

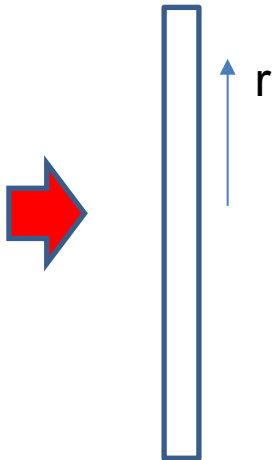
# RaDIATE example with PHITS

Proton gaussian beam

Energy: 2, 30 MeV

0.18, 0.8, 3, 30, 120, 400GeV

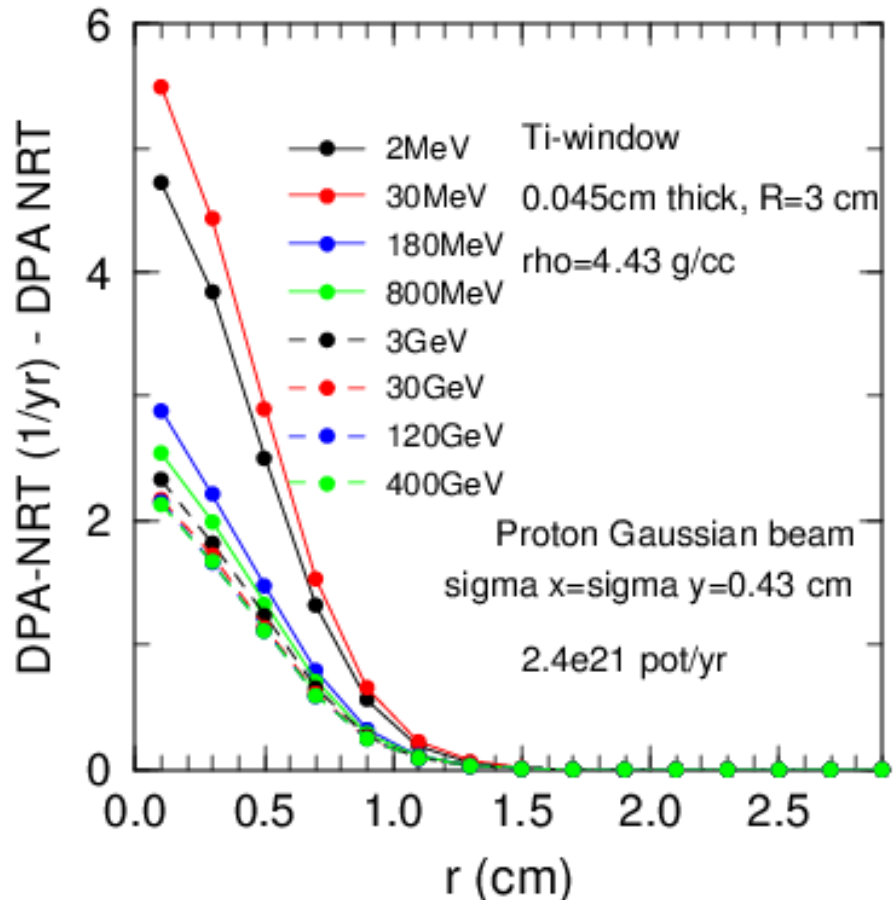
$$\sigma_x = \sigma_y = 0.43 \text{ cm}$$



0.045cm thick, R=1.3 cm

C-target, density=1.84g/cm<sup>3</sup>

Tally region: r=0 – 1.3 divided by 10



DPA radial distribution in Ti-target

# Two dimensional DPA distribution

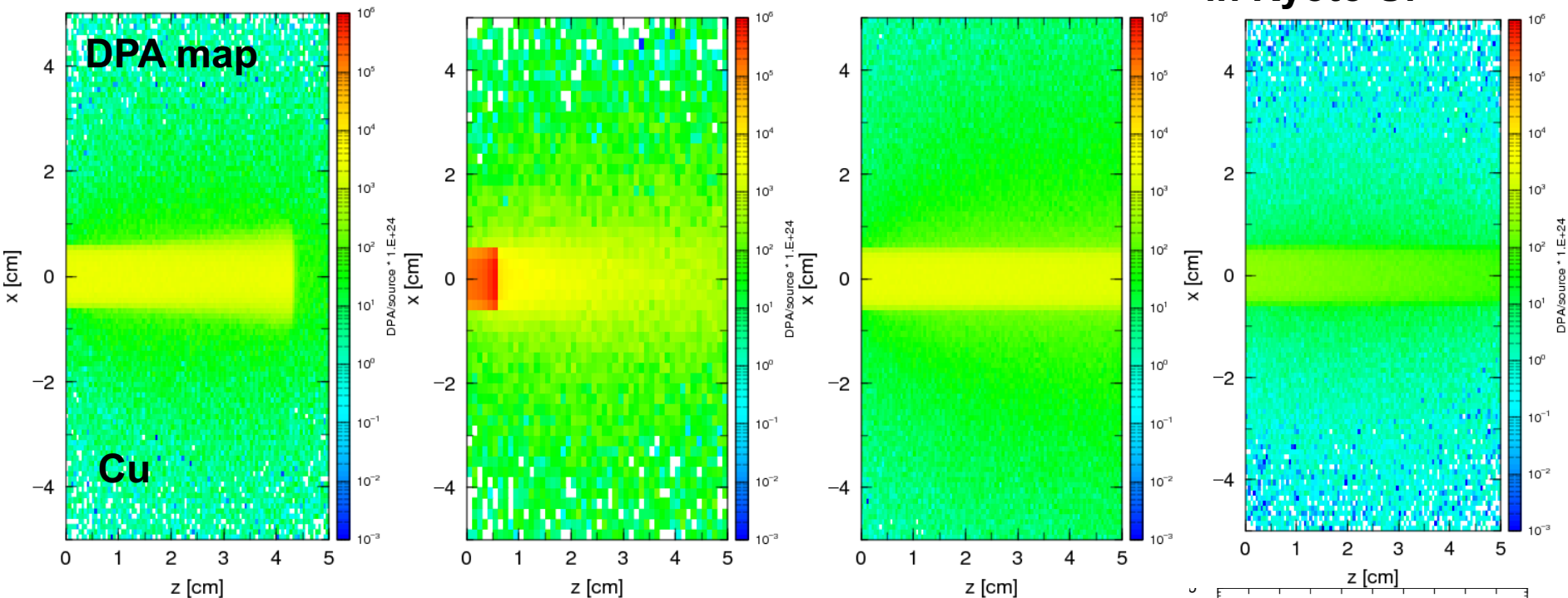
## Calculation condition

(1) 200 MeV proton

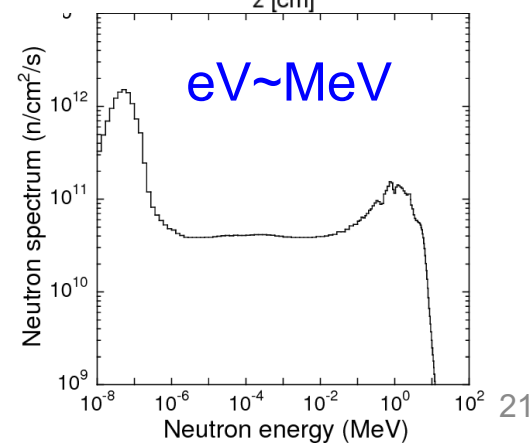
(2) 200 MeV/u  $^{48}\text{Ca}$

(3) 200 MeV neutron

(4) Reactor neutron in Kyoto U.



- ✓ Beam size:  $1\text{cm}^2$
- ✓ Target: 5 cm radius x depth Cu
- ✓ Displacement energy: 30 eV



# Introduction

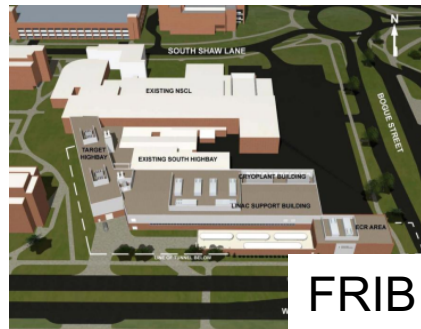
Prediction of structural damage to materials under irradiation is essential for the design.

Spallation source  
J-PARC, SNS, ESS



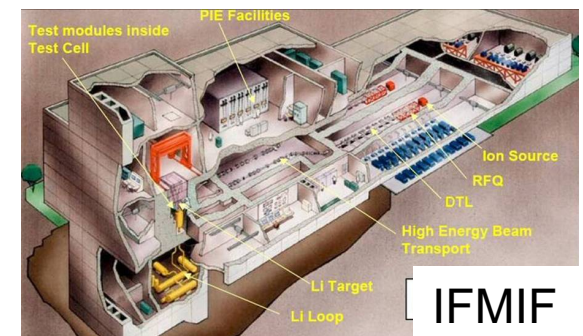
J-PARC  
proton, neutron  
thermal-TeV

Heavy ion facility  
FRIB, RIBF, GSI



FRIB  
heavy-ion  
MeV-GeV/nucleon

International Fusion  
Materials Irradiation Facility



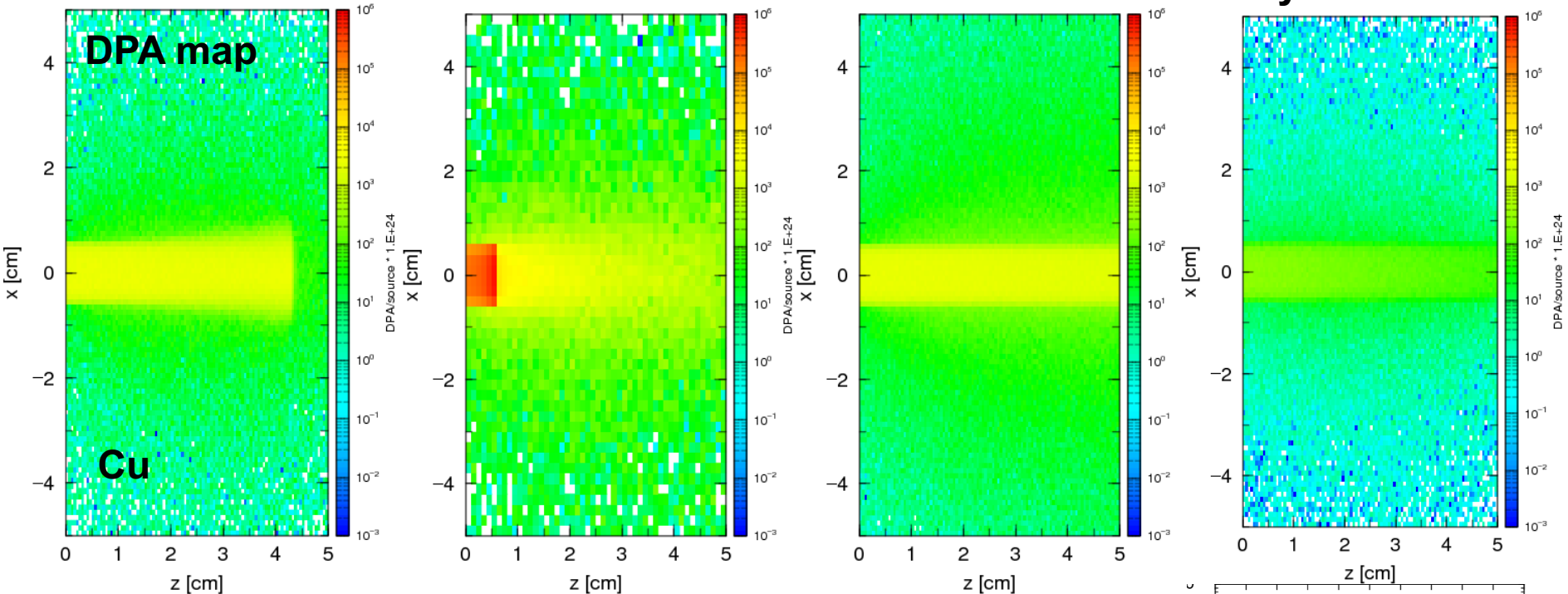
IFMIF  
neutron, deuteron  
~14MeV

To evaluate radiation damage, a fundamental parameter that characterizes lattice displacement events is required.

# Effect of nuclear reaction and elastic scattering

## Calculation condition

- (1) 200 MeV proton    (2) 200 MeV/u  $^{48}\text{Ca}$     (3) 200 MeV neutron    (4) Reactor neutron in Kyoto U.



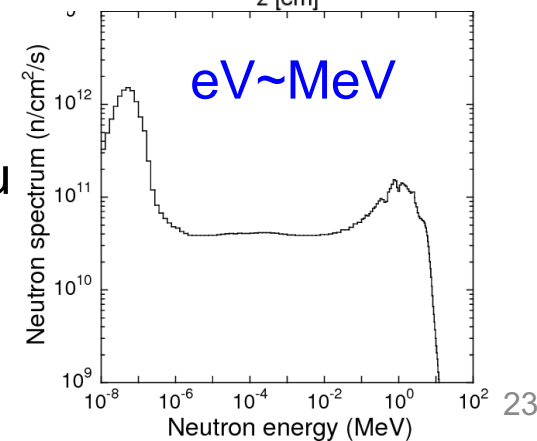
(5) 14 MeV proton

(6) 14 MeV neutron

✓ Beam size: 1 cm<sup>2</sup>

✓ Target: 5 cm radius x depth Cu

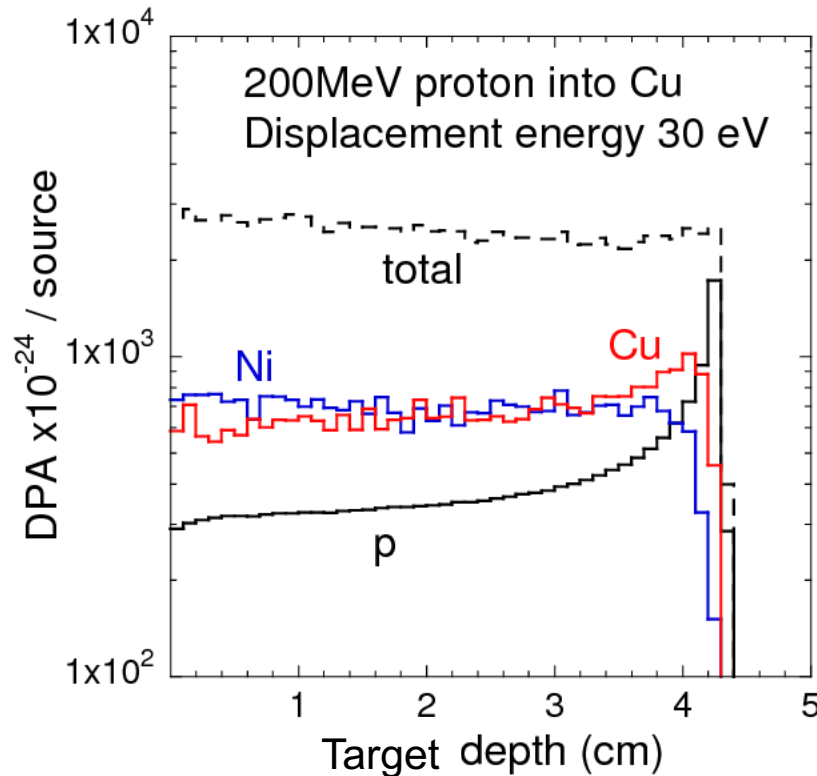
✓ Displacement energy: 30 eV



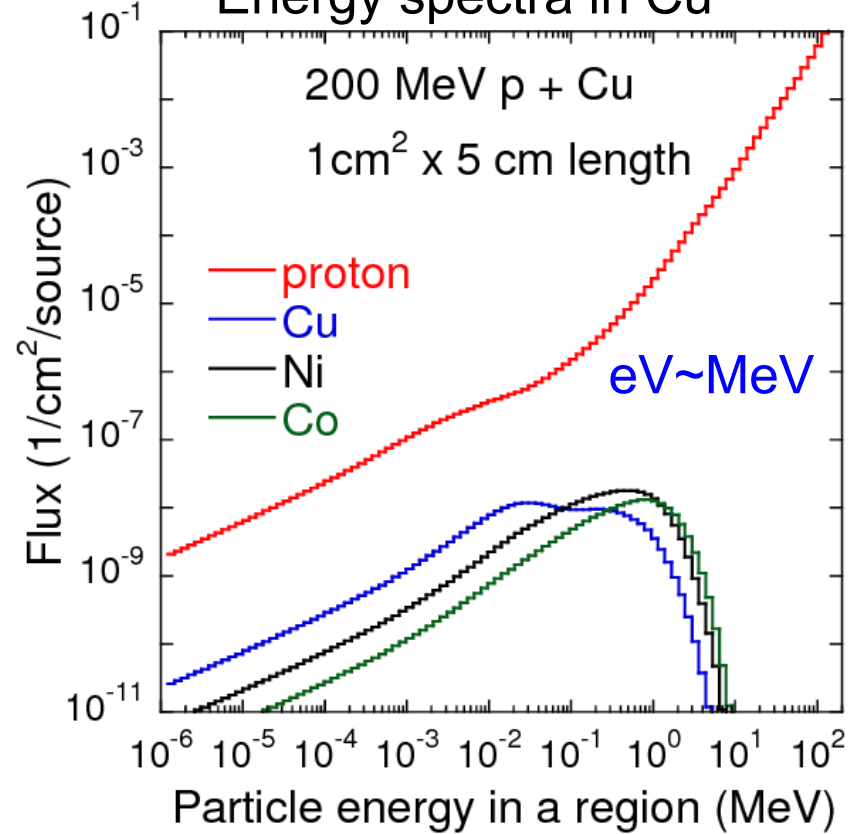


# (1) 200 MeV proton into Cu

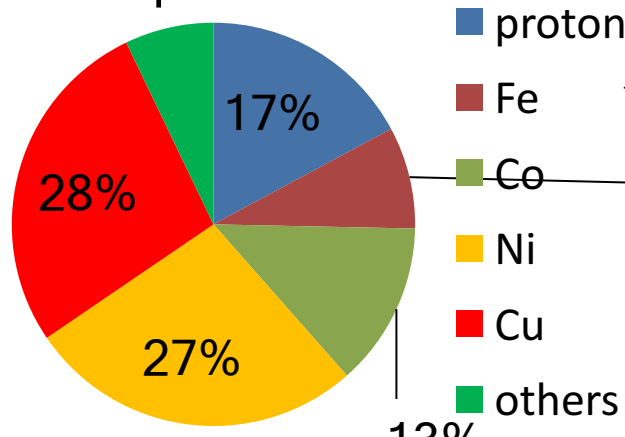
## DPA distribution



## Energy spectra in Cu



## Ratio of partial DPA to total



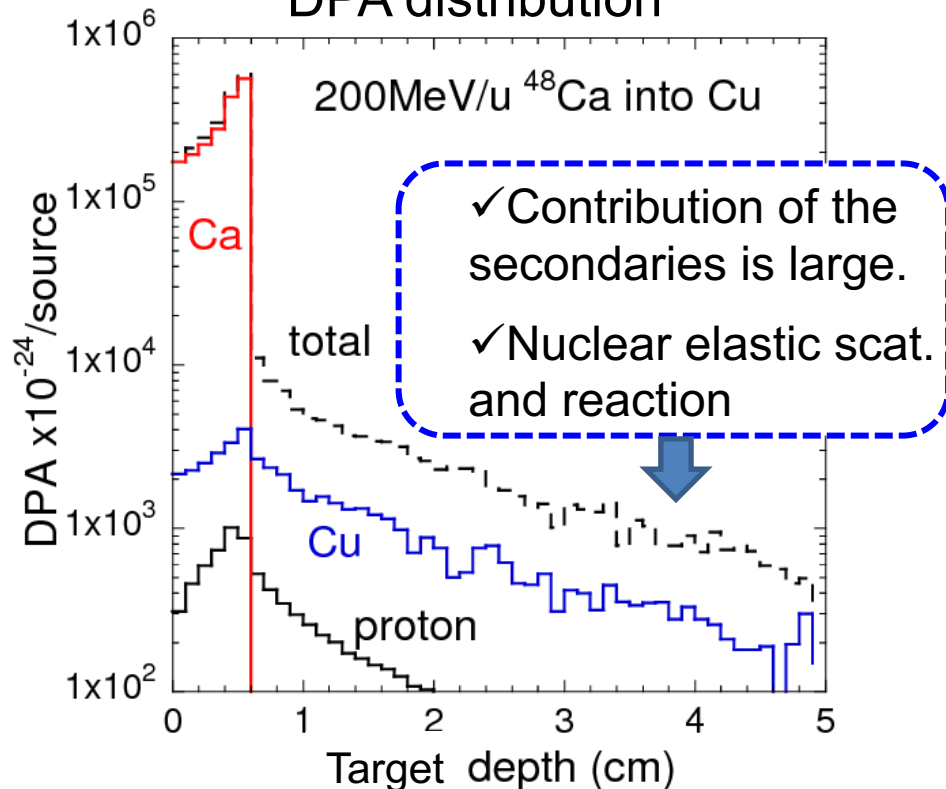
✓ Types of Particles around Cu increase due to nuclear reactions and these particles contribute to total DPA .

✓ Proton DPA is smaller than for heavy-ions because Coulomb scattering cross section of proton is much smaller than that of heavy ions.

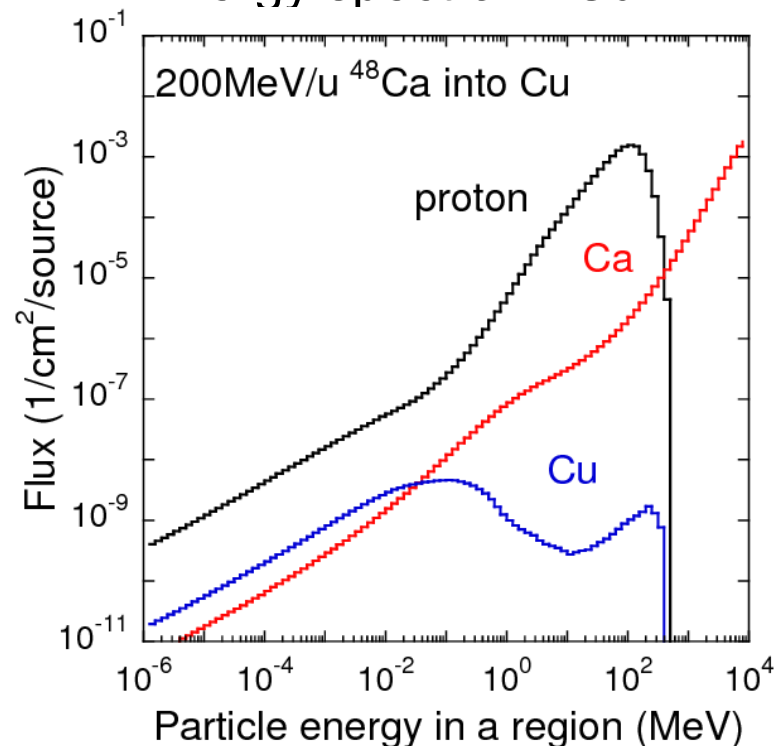


# (2) 200 MeV/nucleon $^{48}\text{Ca}$ into Cu

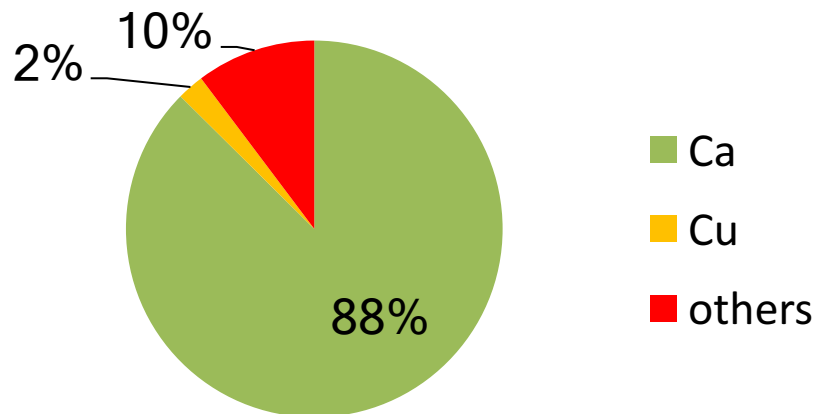
DPA distribution



Energy spectra in Cu



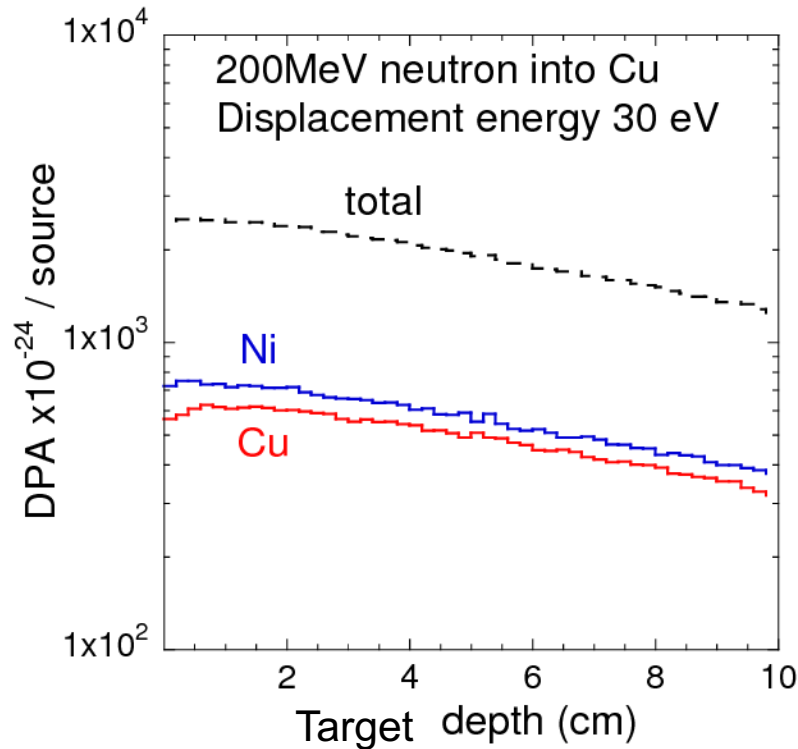
Ratio of partial DPA to total



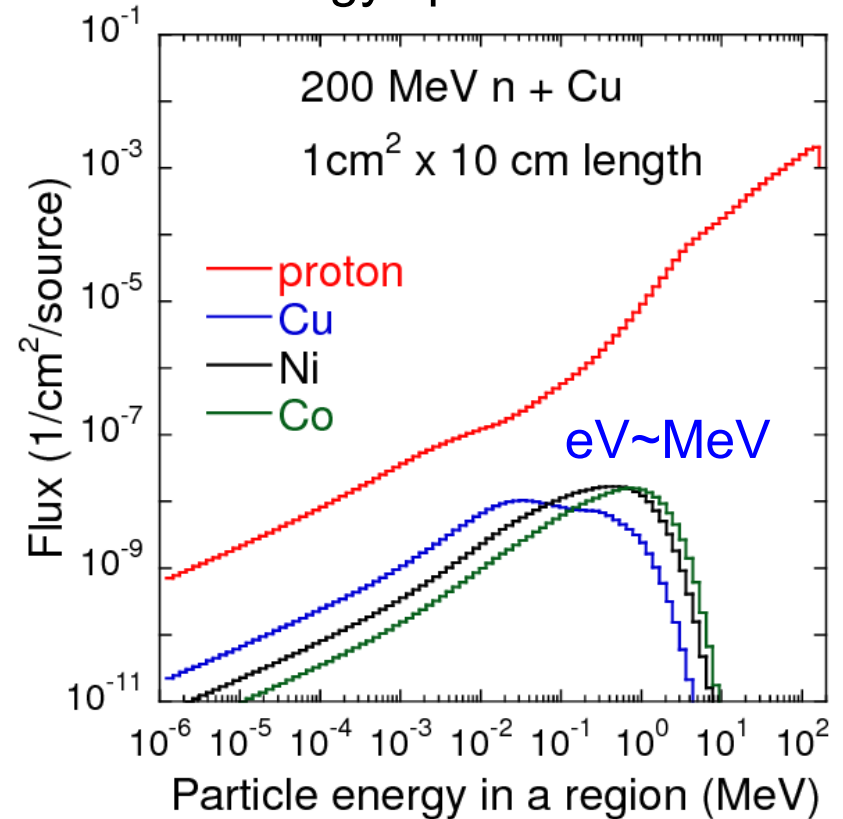
DPA produced by the primary beam is much larger than DPA produced by other contributors.

# (3) 200 MeV neutron into Cu

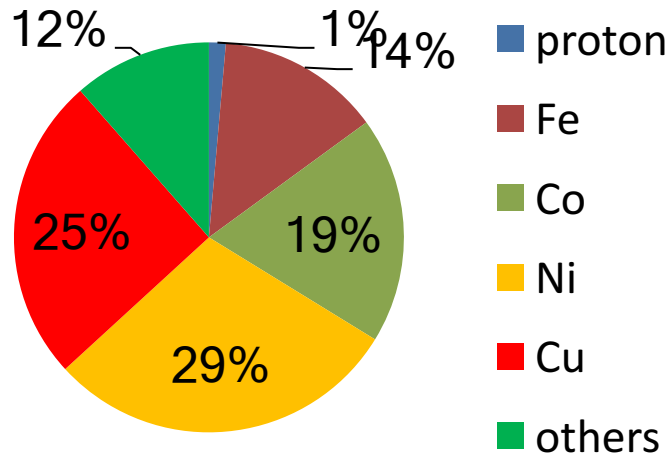
## DPA distribution



## Energy spectra in Cu



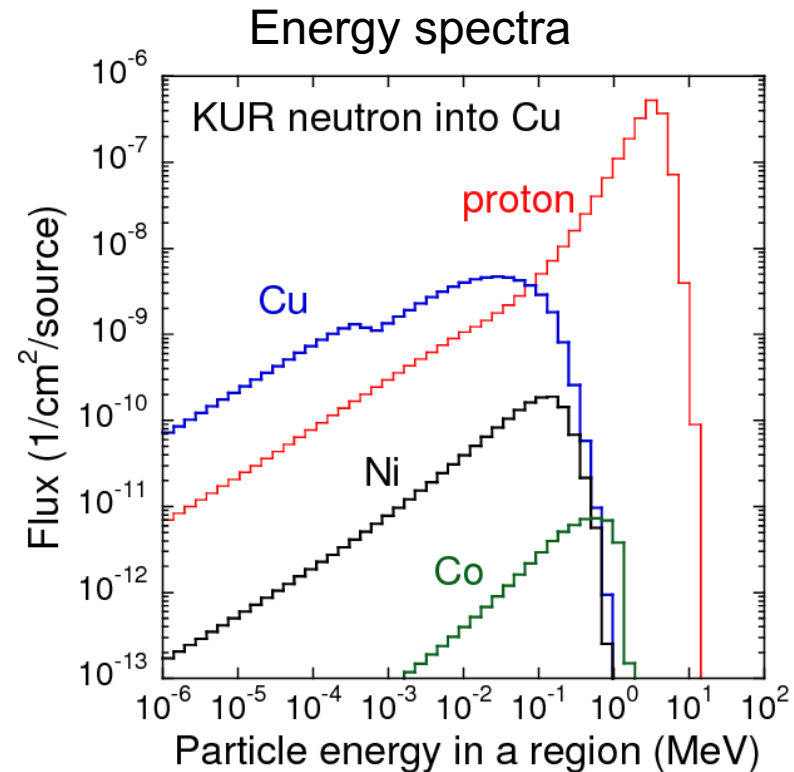
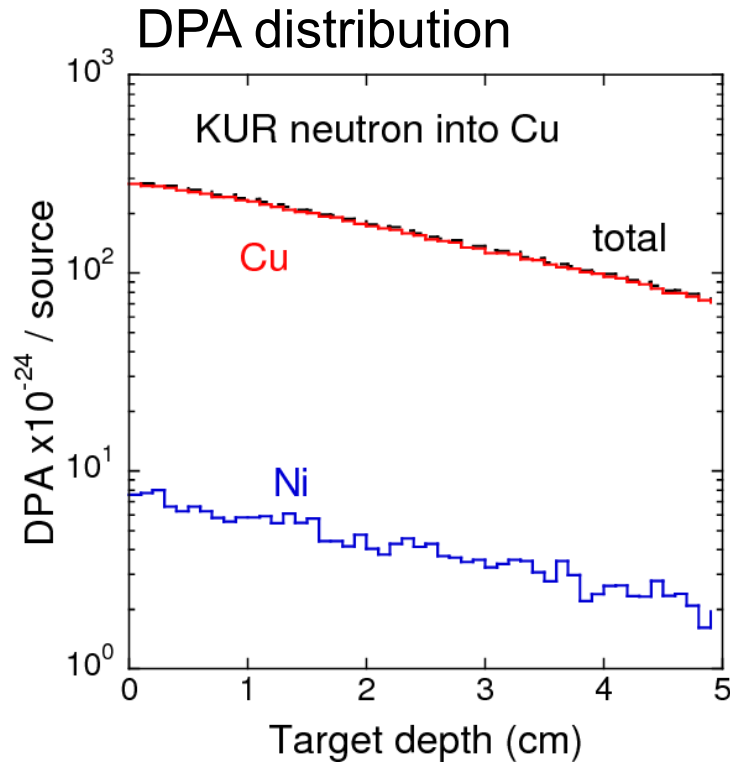
## Ratio of partial DPA to total



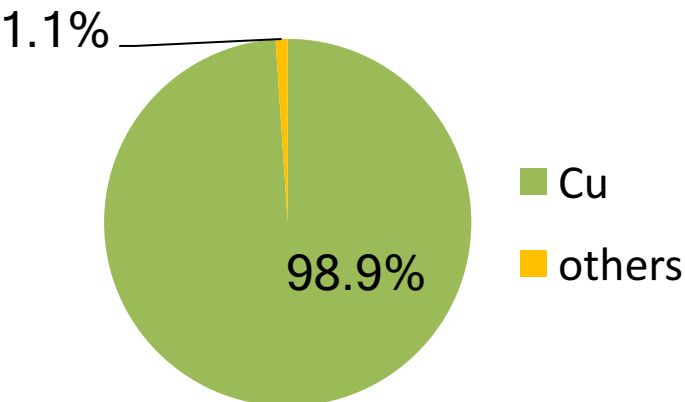
✓ Contributions to total DPA by various particles around Cu **increase due to nuclear reactions**.

✓ Secondary particle distributions for neutron are similar with that for protons.

## (4) Reactor neutron into Cu



Ratio of partial DPA to total



For the low-energy neutron incidence, the target atom is scattered by incident neutron elastic scattering and it contributes to the DPA value

# Summary of effect of nuclear reactions

5 cm radius and depth Cu target

	ratio of partial DPA to total (%)						
	proton	<sup>48</sup> Ca	Fe	Co	Ni	Cu	others
14 MeV proton	89	-	-	-	2	6	3
200 MeV proton	17	-	8	13	27	28	7
14MeV/nucleon <sup>48</sup> Ca	-	99.8	-	-	-	-	0.2
200MeV/nucleon <sup>48</sup> Ca	-	88	-	-	-	2	10
Reactor neutron in Kyoto U.	-	-	-	-	-	99	1
14 MeV neutron	-	-	-	1	31	68	-
200 MeV neutron	1	-	14	19	29	25	12

Proton: DPA value created by projectile decreased with energy.

DPA created by secondary (Cu, Ni) increase with energy.

Neutrons: reactor: n-Cu elastic scattering produce Cu and contribute to DPA.

Secondary particles produced by nuclear reactions increase with neutron energy.