

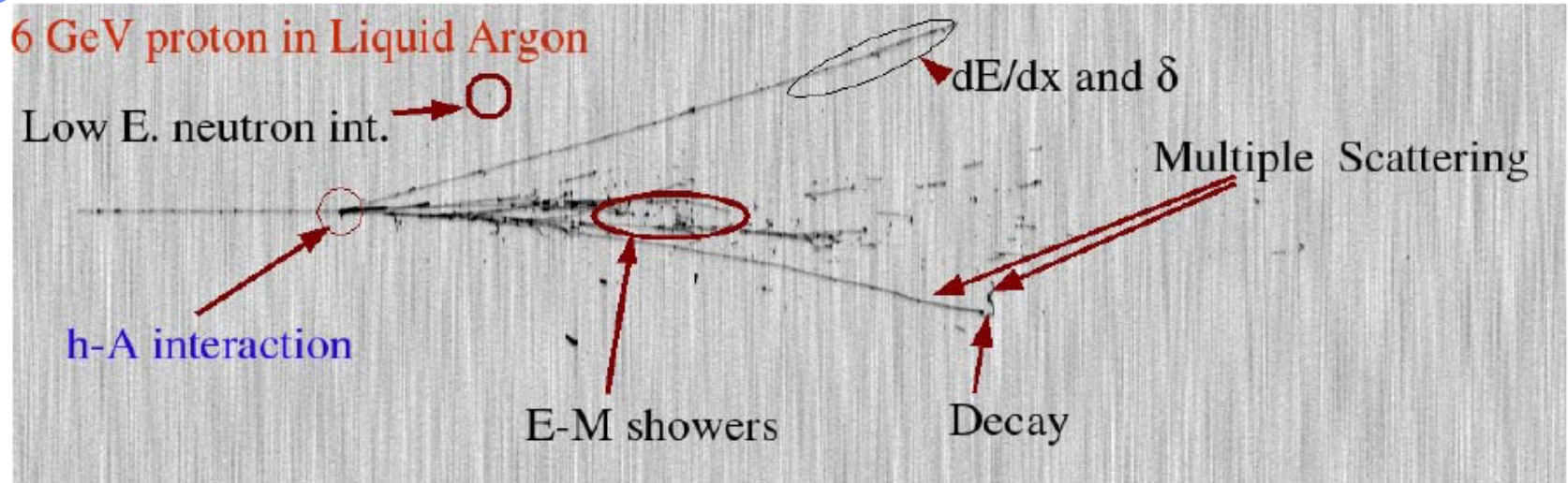


# Modeling of Radiation Induced Damage in FLUKA

RADIATE 4<sup>th</sup> Collaboration Meeting  
Tokai-Mura, Ibaraki, Japan, Sep. 18-22, 2017

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# Interaction and Transport Monte Carlo Code



- Hadron-nucleus interactions
- Nucleus-Nucleus interactions
- Electron interactions
- Photon interactions
- Muon interactions (inc. photonuclear)
- Neutrino interactions
- Decay
- Low energy neutrons
- Ionization
- Multiple scattering
- Combinatorial geometry
- Voxel geometry
- Magnetic field
- Analogue or biased
- On-line buildup and evolution of induced radioactivity and dose
- User-friendly GUI thanks to *Flair*

# Different kind of damage

*from Radiation-Matter interaction*

## Precious materials (healthy/tragic damage)

energy (dose) deposition radioisotope production and decay & positron annihilation and photon pair detection

## Oxidation

by generation of chemically active radicals (e.g. PVC dehydrochlorination by X and g-rays, radiolysis,...)

## Accidents

energy (power) deposition

## Degradation

energy (dose) deposition, particle fluence, DPA

## Gas production

residual nuclei production

## Electronics

high energy hadron fluence, neutron fluence, energy (dose) deposition

## Activation

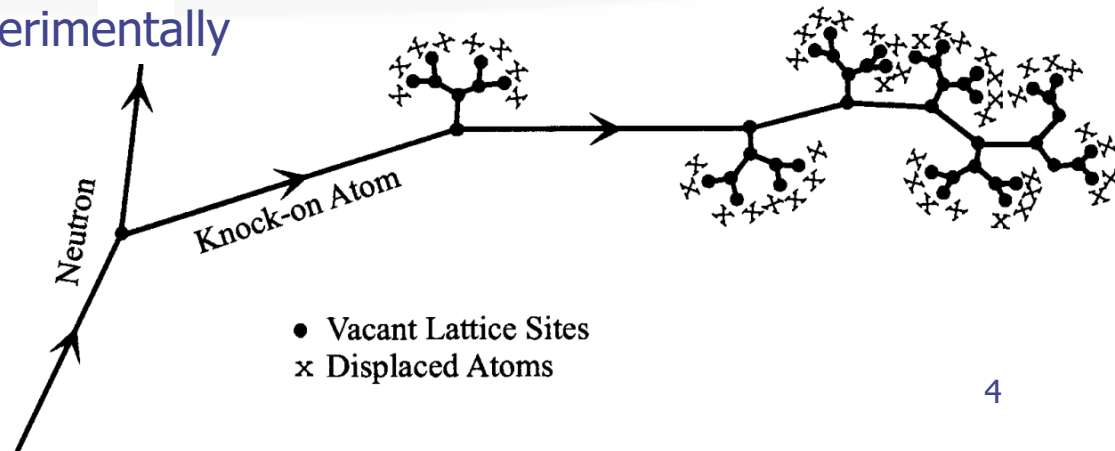
residual activity and dose rate

# DPA

- The unit that is commonly used to link the “*radiation damage effects*” with “*macroscopic structural damage*” is the **displacement per atoms**
- It is a “measure” of the amount of radiation damage in irradiated materials

*3 dpa means each atom in the material has been displaced from its site within the structural lattice an average of 3 times*

- a quantity directly linked to the Non Ionizing Energy Losses (**NIEL**) but restricted in energy
- dpa is a strong function of **projectile type, energy and charge** as well as material properties and can be induced by all particles in the cascade
- However dpa for the moment is a “**mathematical**” quantity that cannot be directly measured experimentally



# Frenkel pairs

- Frenkel pairs  $N_F$  (defect or disorder), is a compound crystallographic defect in which an **interstitial** lies near the **vacancy**. A Frenkel defect forms when an atom or ion leaves its place in the lattice (leaving a vacancy), and lodges nearby in the crystal (becoming an interstitial)

$$N_{NRT} \equiv N_F = \kappa \frac{\xi(T)T}{2E_{th}}$$

$N_{NRT}$   
 $\kappa=0.8$   
 $T$

Defects by Norgert, Robinson and Torrens  
is the displacement efficiency  
kinetic energy of the primary  
knock-on atom (PKA)

$\xi(T)$

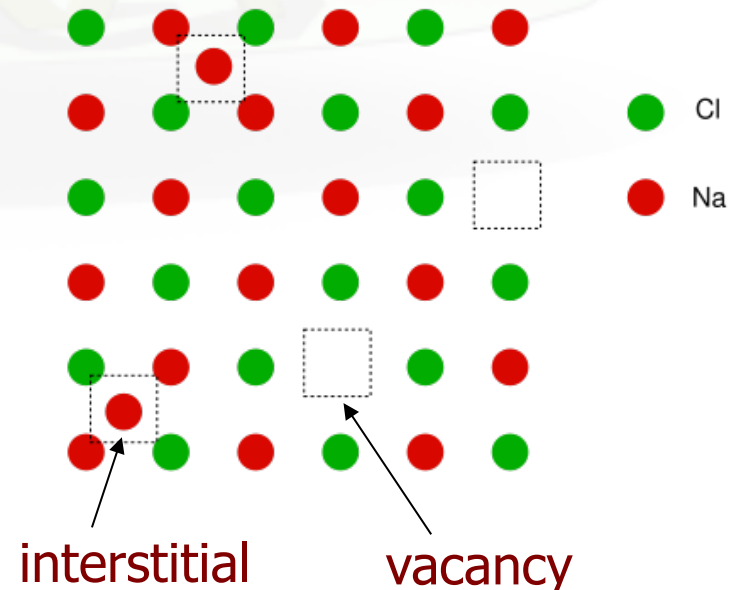
partition function (LSS theory)

$\xi(T) T$

directly related to the **NIEL**  
(non ionizing energy loss)

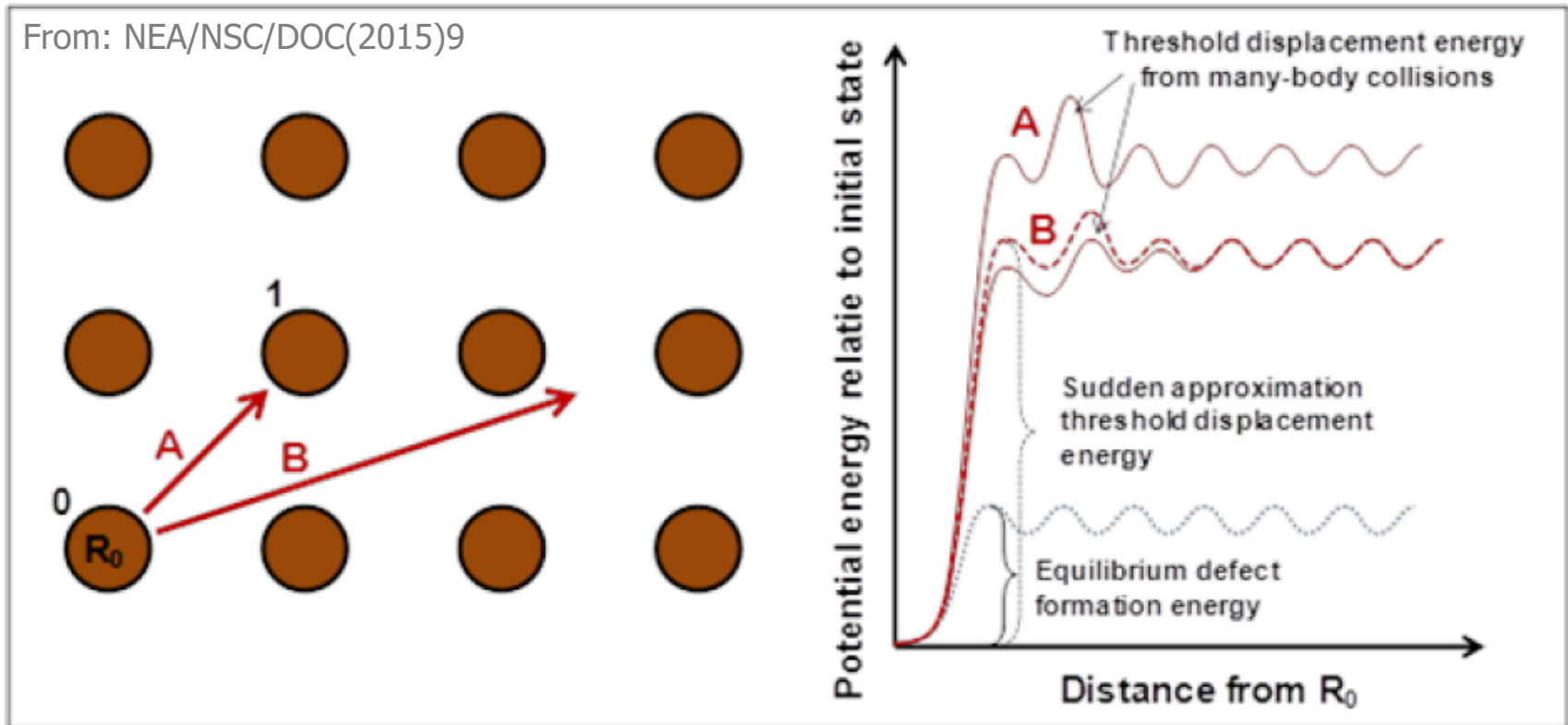
$E_{th}$

damage threshold energy



# Damage Threshold

From: NEA/NSC/DOC(2015)9



- Damage threshold depends on the direction of the recoil in the crystal lattice.
- FLUKA Use: the “average” threshold over all crystallographic directions (user defined)
- Typically of the order of **10-50 eV**

# Damage Threshold in Compounds

- NJOY (MT=444) sums up the cross section multiplied by the damage energies, which is the damage production cross section representing the effective kinetic energy of recoiled atom for reaction types  $i$  at neutron energy  $E_n$

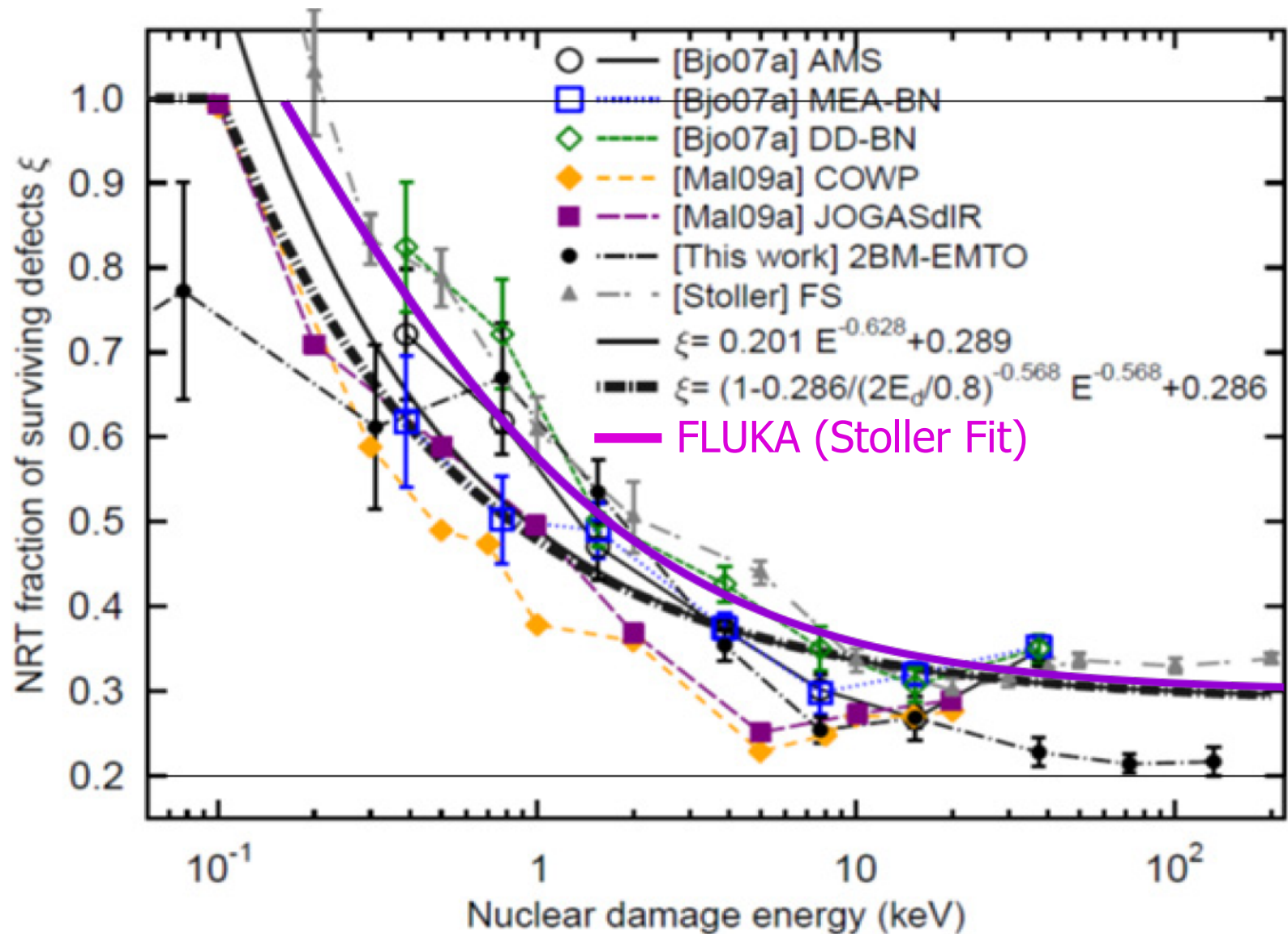
$$(E\sigma)_{DPA} = \sum_i E_{th,i} \sigma_i(E_n)$$

## Problematic:

- Damage threshold depends on the lattice structure.
- Damage threshold can be quite different for each combination for the specific compound  
e.g. NaCl:  $E_{th}(\text{Na-Na})$ ,  $E_{th}(\text{Na-Cl})$ ,  $E_{th}(\text{Cl-Na})$ ,  $E_{th}(\text{Cl-Cl})$
- Simple weighting with the atom/mass fraction doesn't work
- FLUKA's approximation is using a unique average damage threshold  $E_{th}$  for the compounds as well

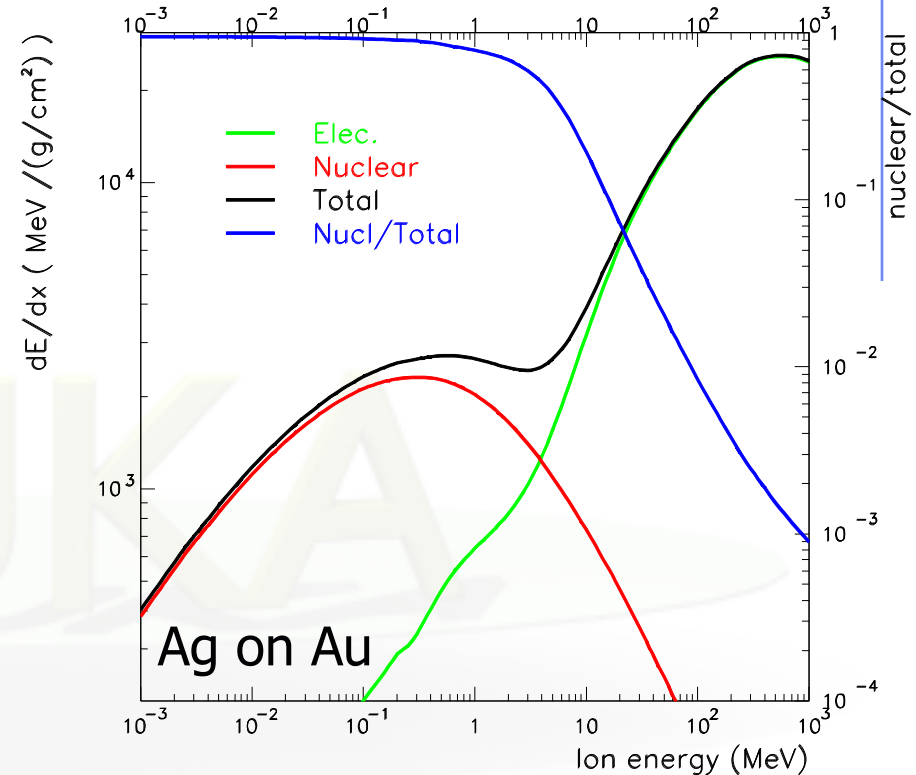
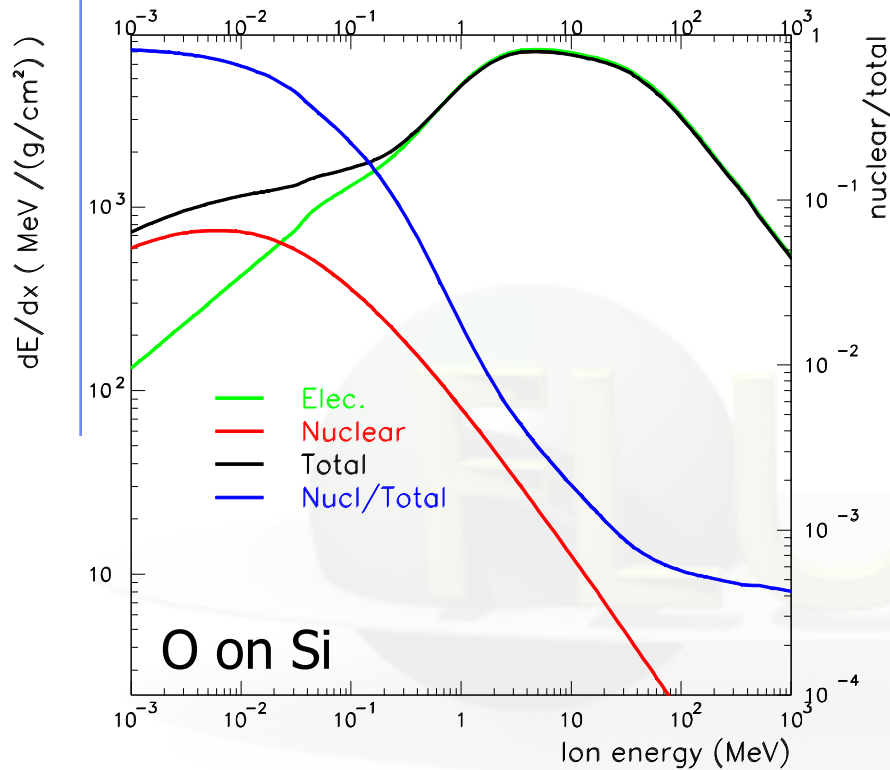


# κ Stoller vs Nordlund





# Nuclear Stopping Power



The total ( $S$ ), nuclear ( $S_n$ ) and electronic ( $S_e$ ) stopping power.

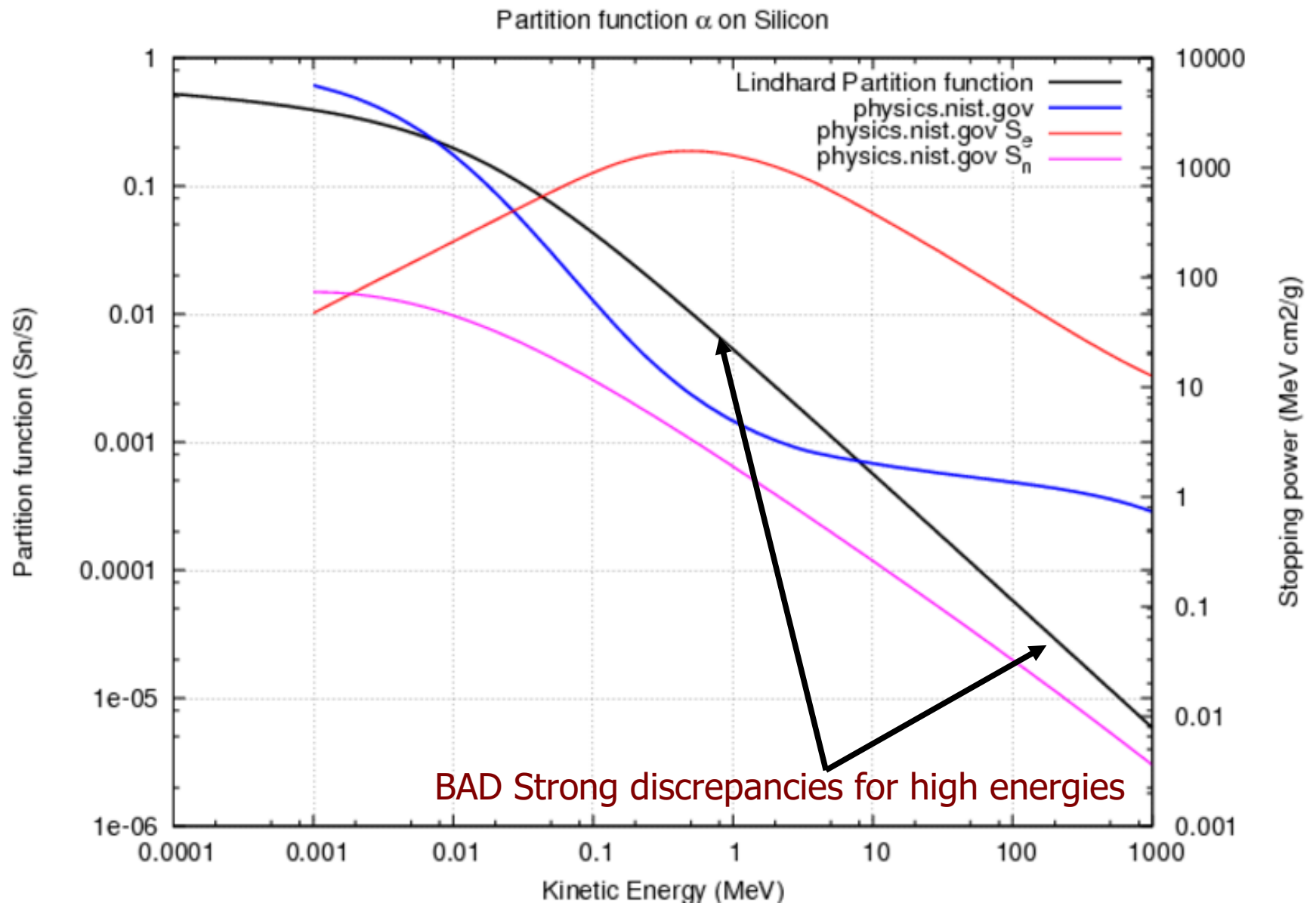
The abscissa is the ion total kinetic energy.

The **partition function**  $S_n/(S_n+S_e)$  is also plotted.

**$S_n/S$  is going down with energy (and up with charge)**  
**→ NIEL/DPA are dominated by low energy (heavy) recoils!!**

$$N_F = \kappa \frac{\xi(T)T}{2E_{th}}$$

# Lindhard partition function $\xi$ [2/2]



# Restricted Nuclear Stopping Power

- Lindhard approximation uses the **unrestricted NIEL**. Including all the energy losses also those below the threshold  $E_{th}$
- FLUKA is using a more accurate way by employing the **restricted nuclear losses**

$$S(E, E_{th}) = N \int_{E_{th}}^{\gamma E} T \left( \frac{d\sigma}{dT} \right) dT$$

where:

$S(E, E_{th})$  is the restricted energy loss  
 $N$  atomic density  
 $T$  energy transfer during ion-solid interaction  
 $d\sigma/dT$  differential scattering cross section

$$\gamma = \frac{2E(2m + E)}{M + \frac{m^2}{M} + 2(m + E)}$$

maximum fraction of energy transfer during collision

# FLUKA Implementation

## Charged particles and heavy ions

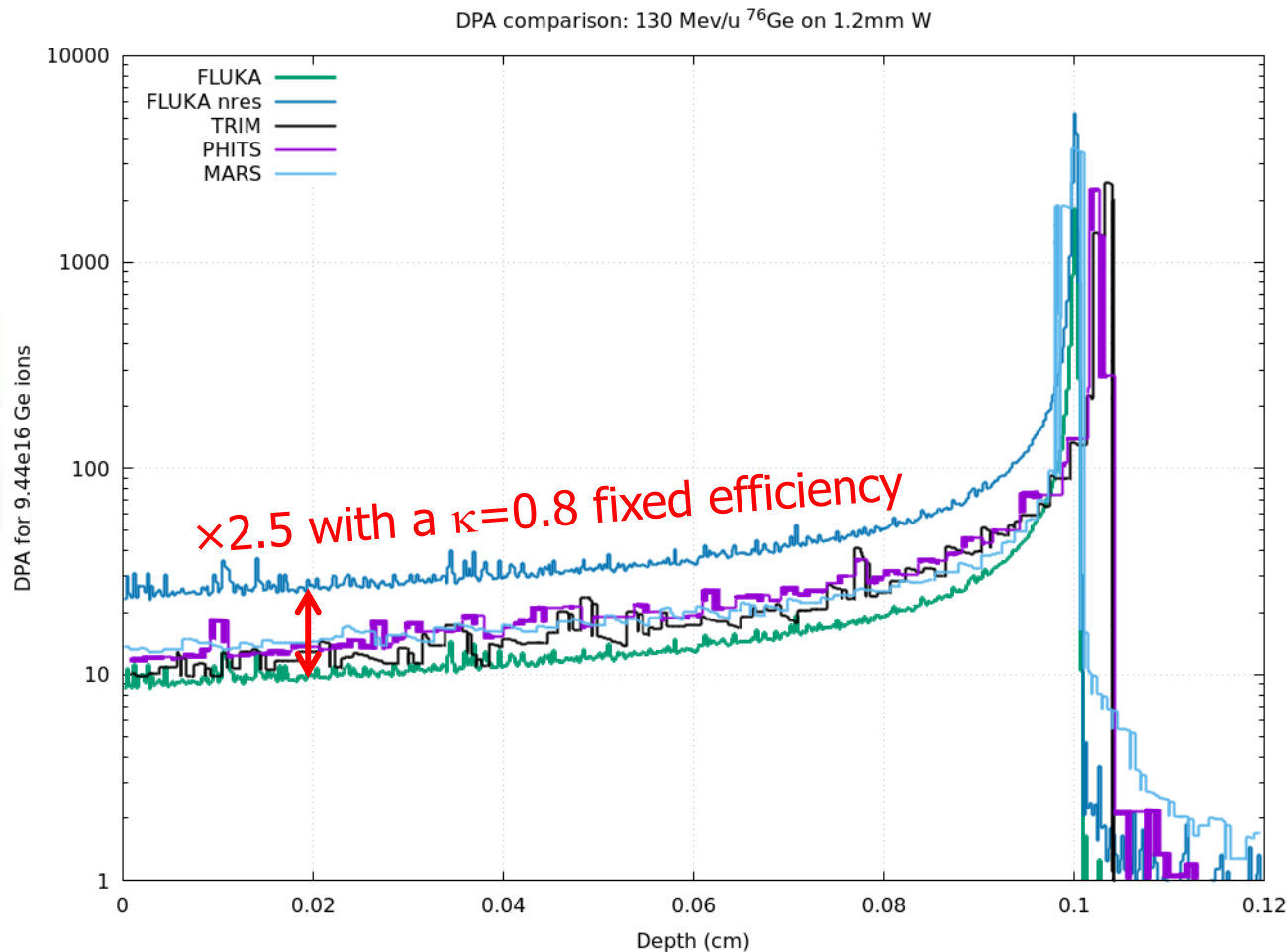
- **During transport**  
Calculate the restricted non ionizing energy loss
- **Below threshold**  
Calculate the integrated nuclear stopping power with the Lindhard partition function
- **At (elastic and inelastic) interactions**  
Calculate the recoil, to be transported or treated as below threshold

## Neutrons:

- **High energy  $E_n > 20$  MeV**
  - Calculate the recoils after interaction  
Treat recoil as a “normal” charged particle/ion
- **Low energy  $E_n \leq 20$  MeV (group-wise)**
  - Calculate the NIEL from NJOY
- **Low energy  $E_n \leq 20$  MeV (point-wise)**
  - Calculate the recoil if possible  
Treat recoil as a “normal” charged particle/ion

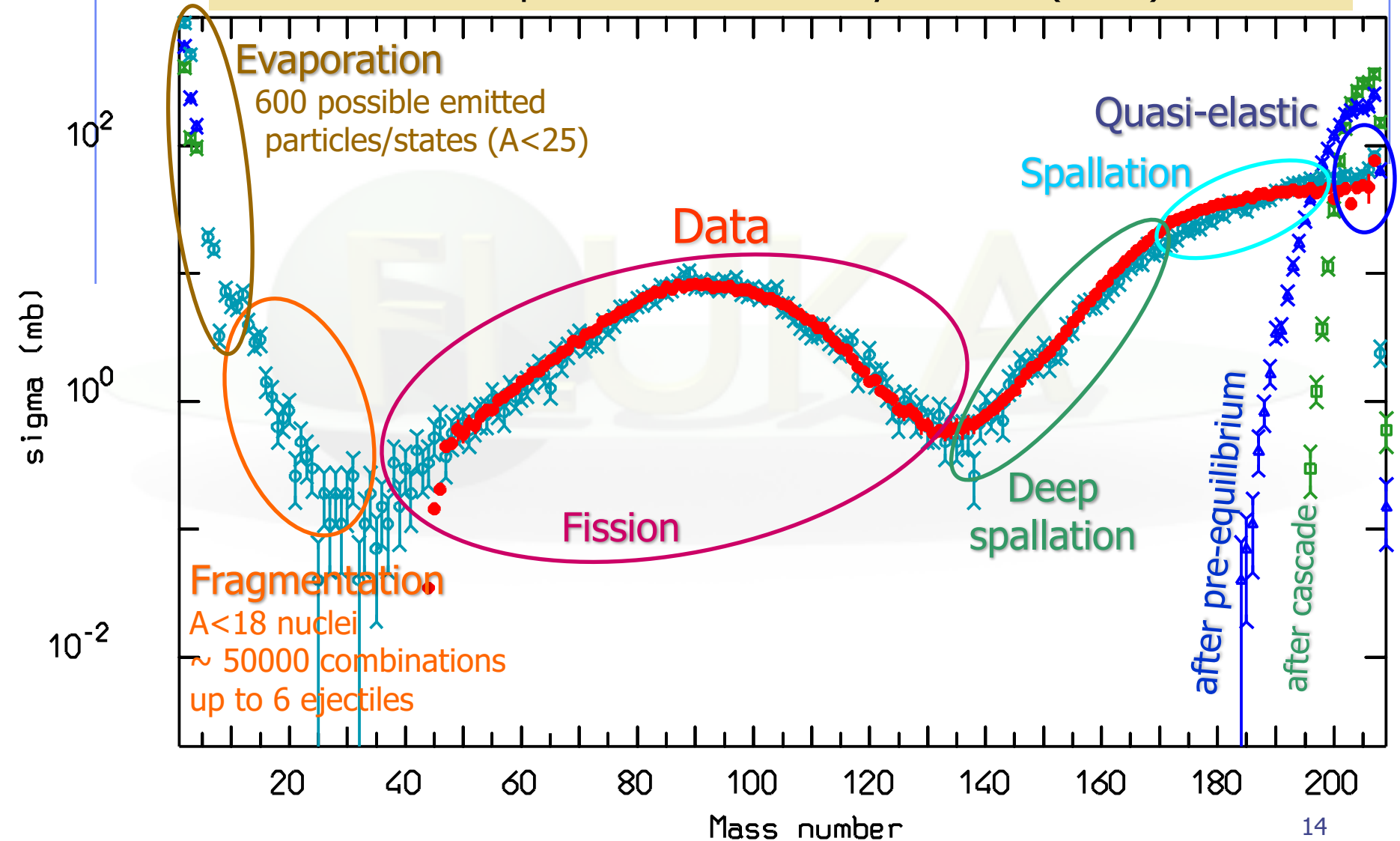
# $^{76}\text{Ge}$ ion pencil beam of 130 MeV/A on W

- $^{76}\text{Ge}$  ion pencil beam of 130 MeV/A uniform in W target a disc of  $R=0.3568$  mm, 1.2 mm thickness

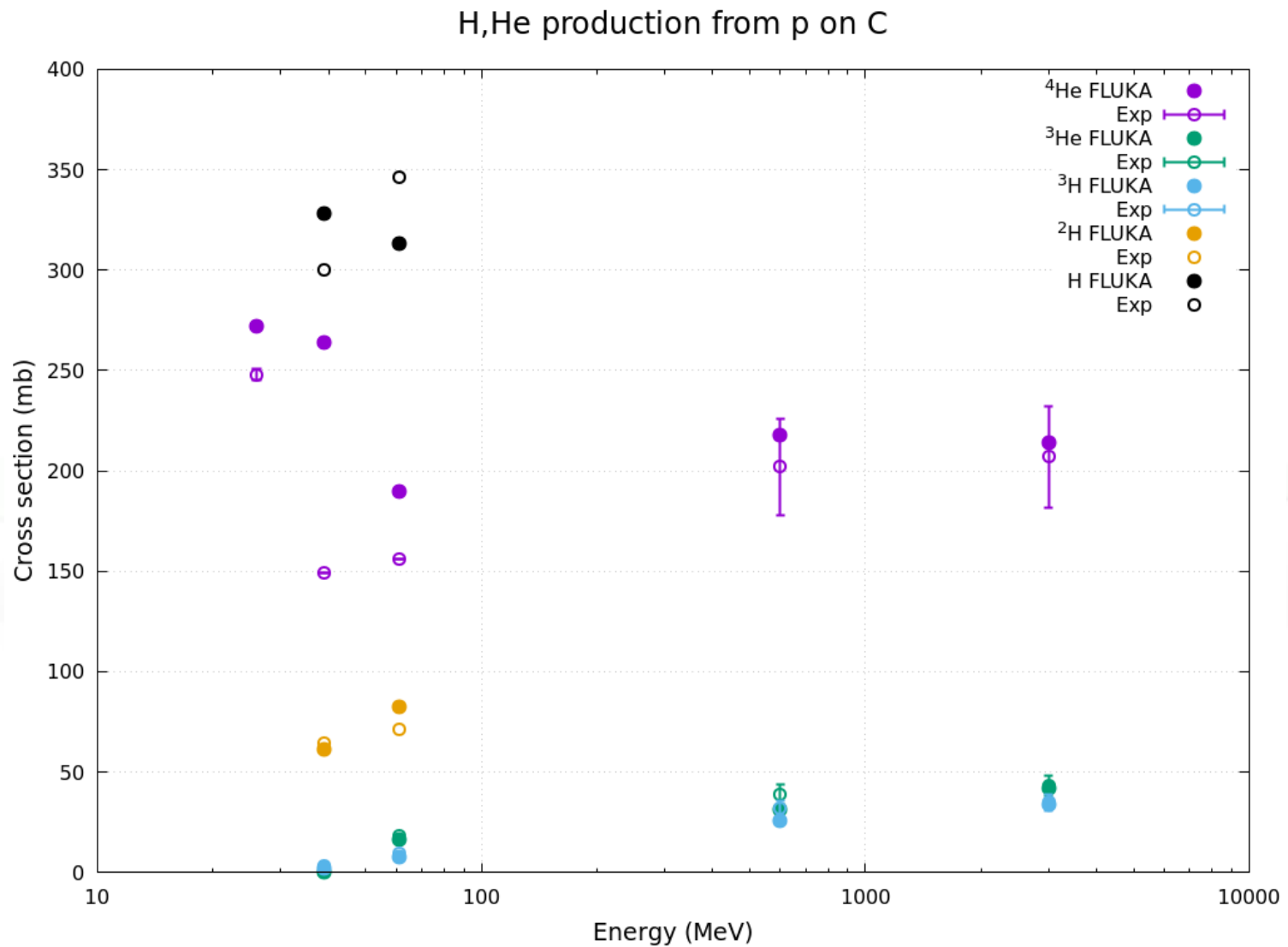


# Example of fission/evaporation

1 A GeV  $^{208}\text{Pb} + \text{p}$  reactions Nucl. Phys. A 686 (2001) 481-524



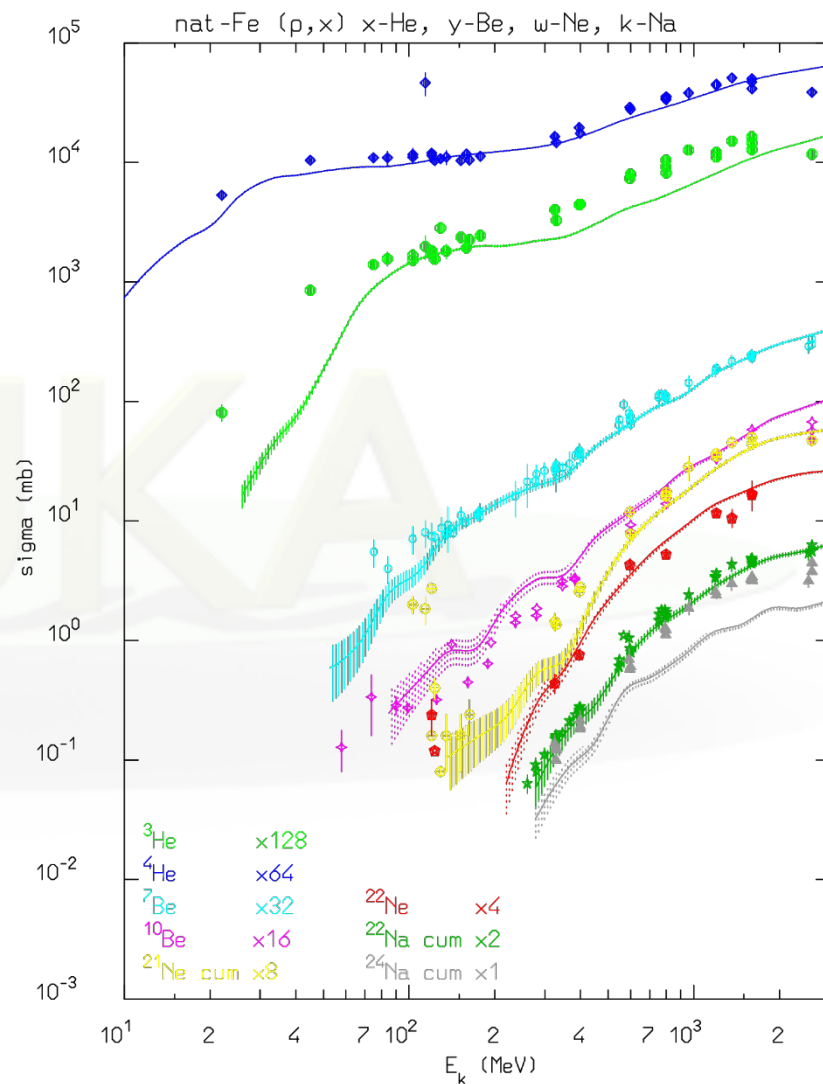
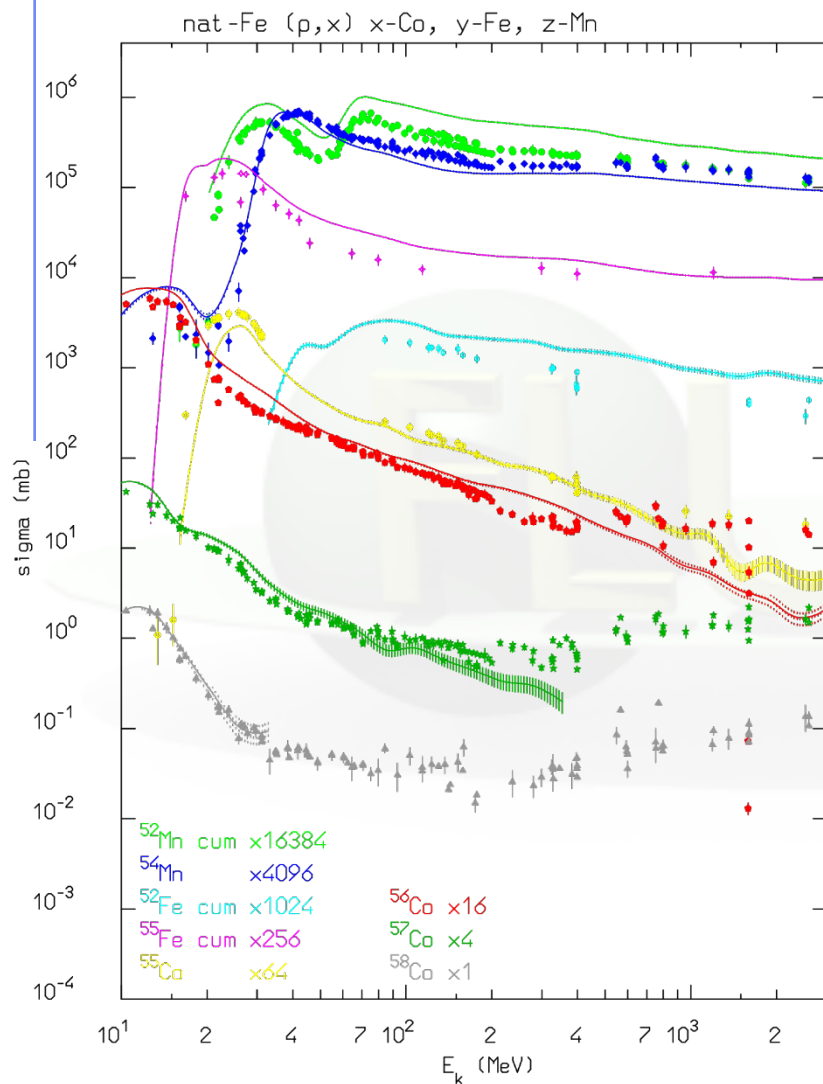
# Particle production C(p,x)



Data: JNST36 313 1999, PRC7 2179 1973



# Isotope production for $^{nat}\text{Fe}(p,x)$ :



Data: Michel et al. 1996 and 2002

# Beam Dump Facility (BDF)

## Beam:

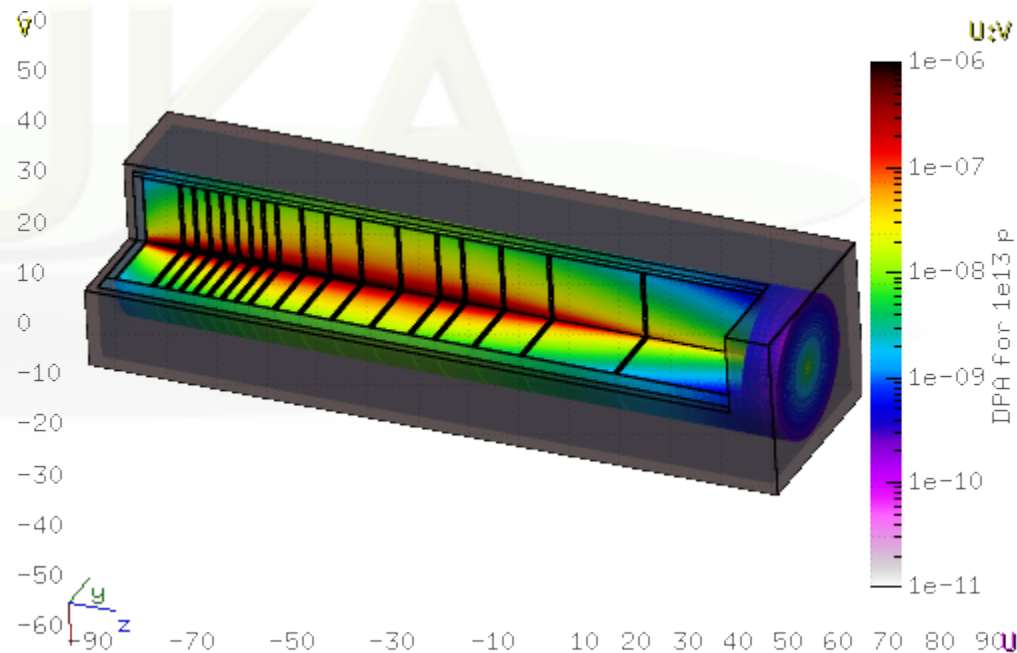
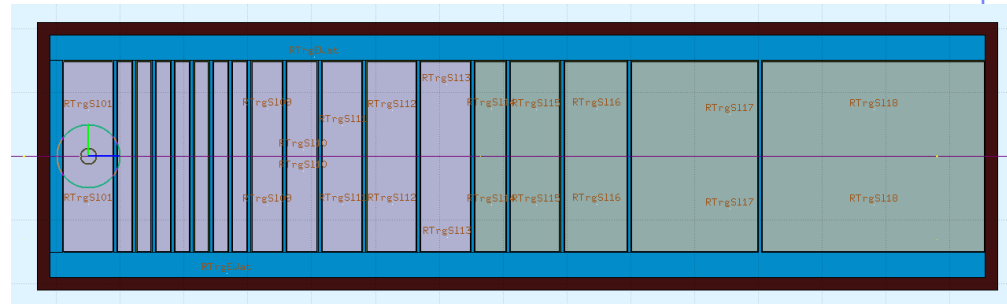
- Protons: 400 GeV/c
- Sweep pattern:
  - radius 3cm
  - $1\sigma$  0.6cm

## Geometry:

- 1.4m long cylinder discs of
- TZM enclosed in Ta
- W enclosed in Ta

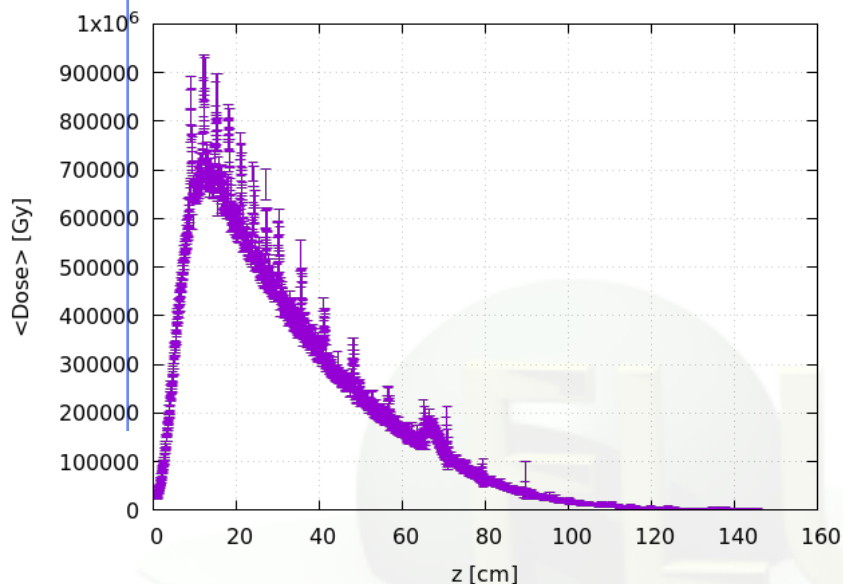
## Materials:

- Tungsten  $E_d=90$  eV
- SS 316N  $E_d=40$  eV
- Tantalum  $E_d=53$  eV
- TZM(Mo,Zr,Ti...)  $E_d=60$  eV

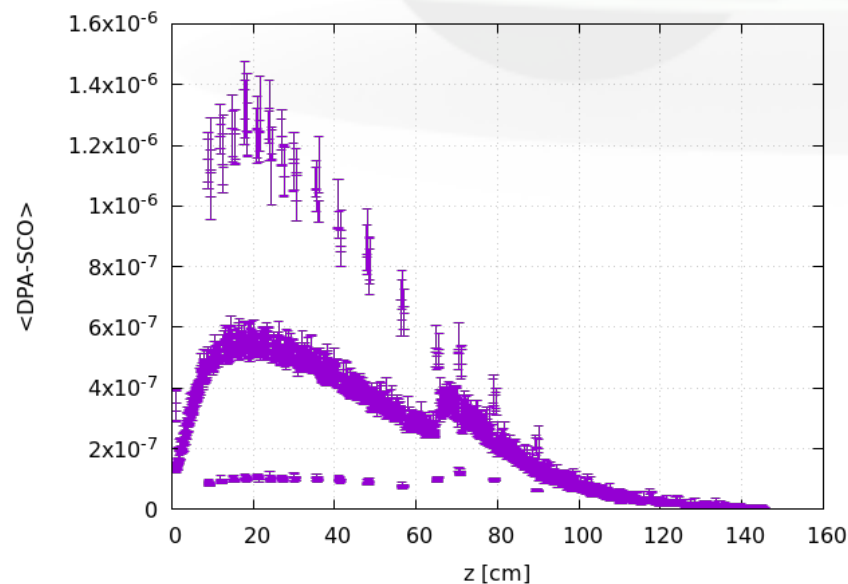


# BDF H/He[appm] vs DPA

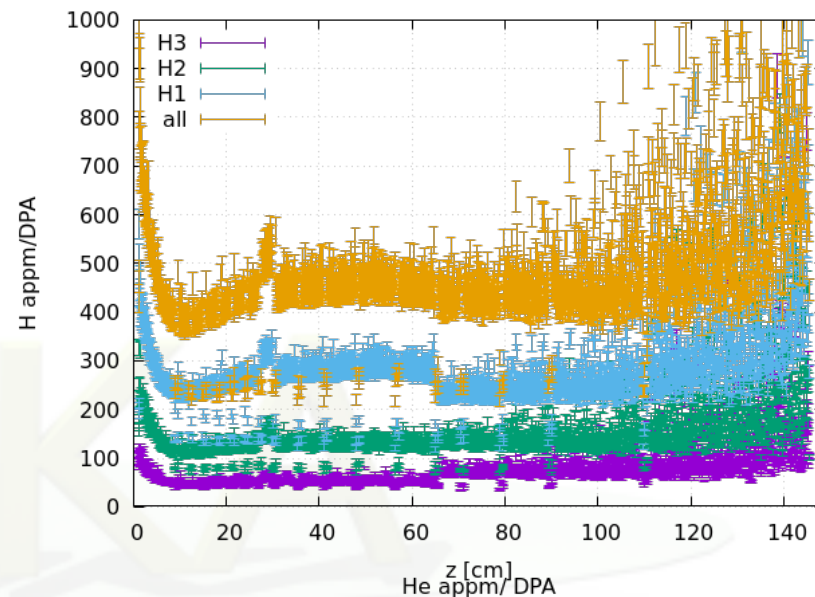
Dose (energy deposited per unit mass),  $R < 0.25\text{cm}$ , for  $3\text{e}13$  POT



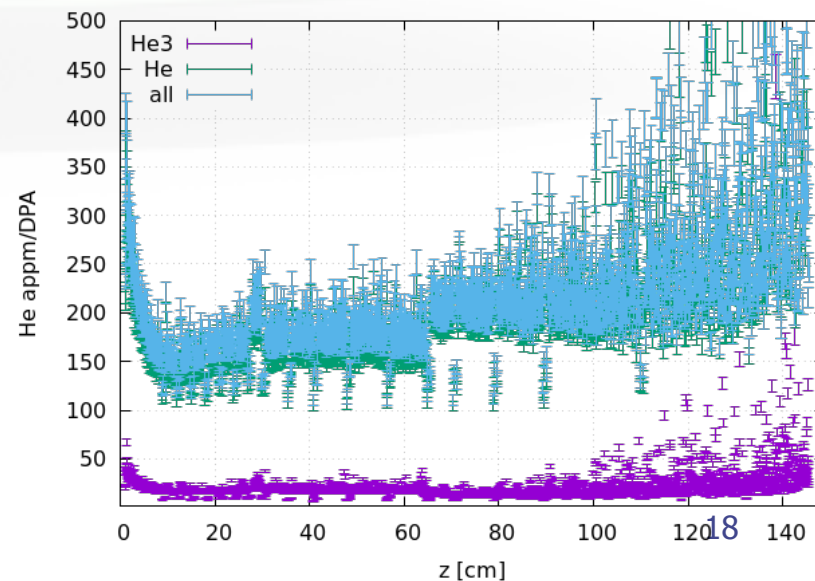
Displacements per atoms,  $R < 0.25\text{cm}$ , for  $3\text{e}13$  POT



H appm/ DPA

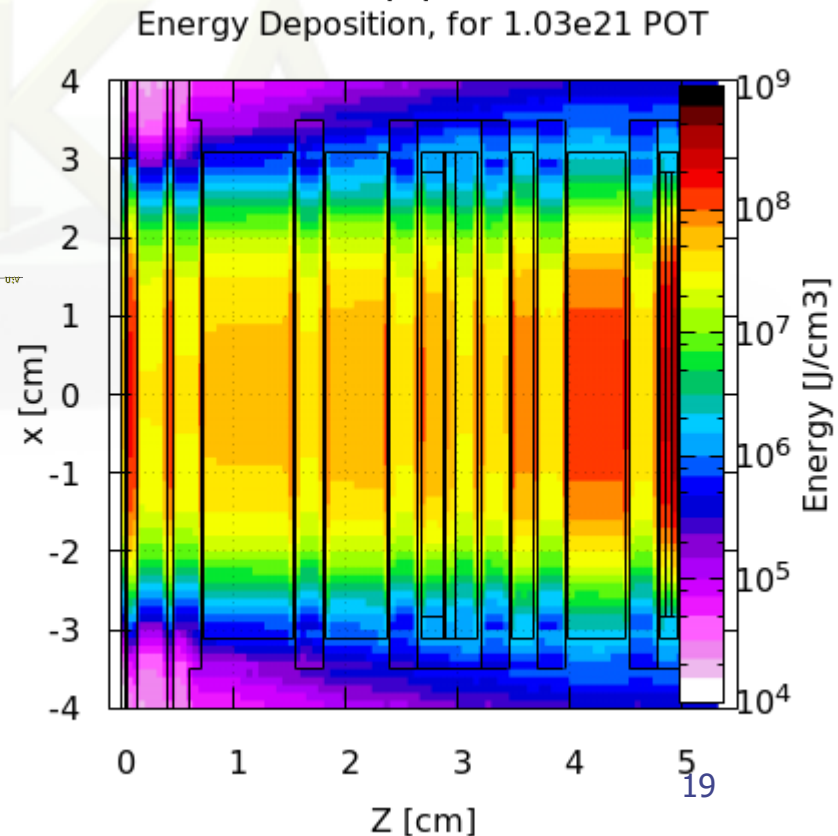
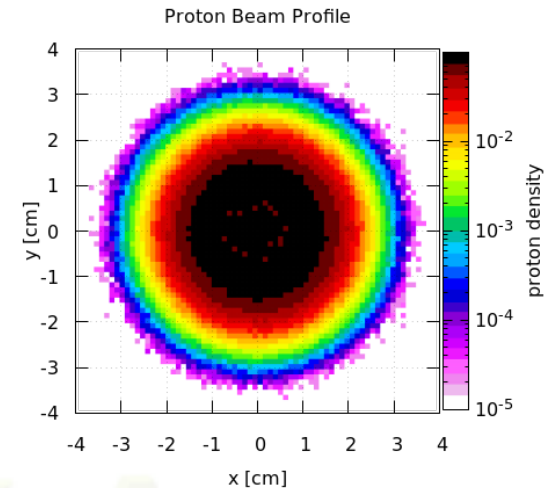
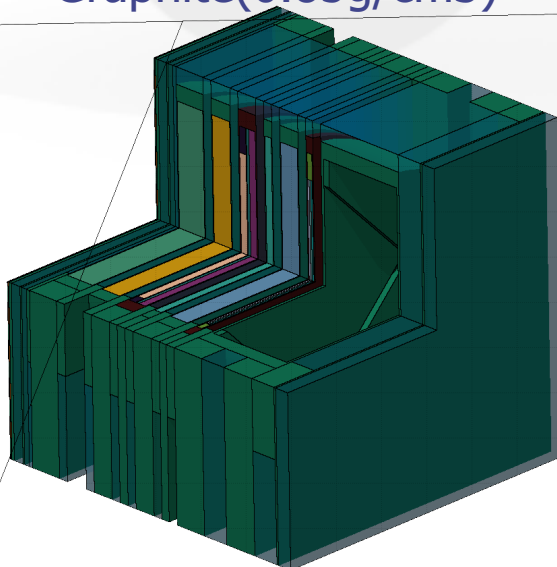


He appm/ DPA

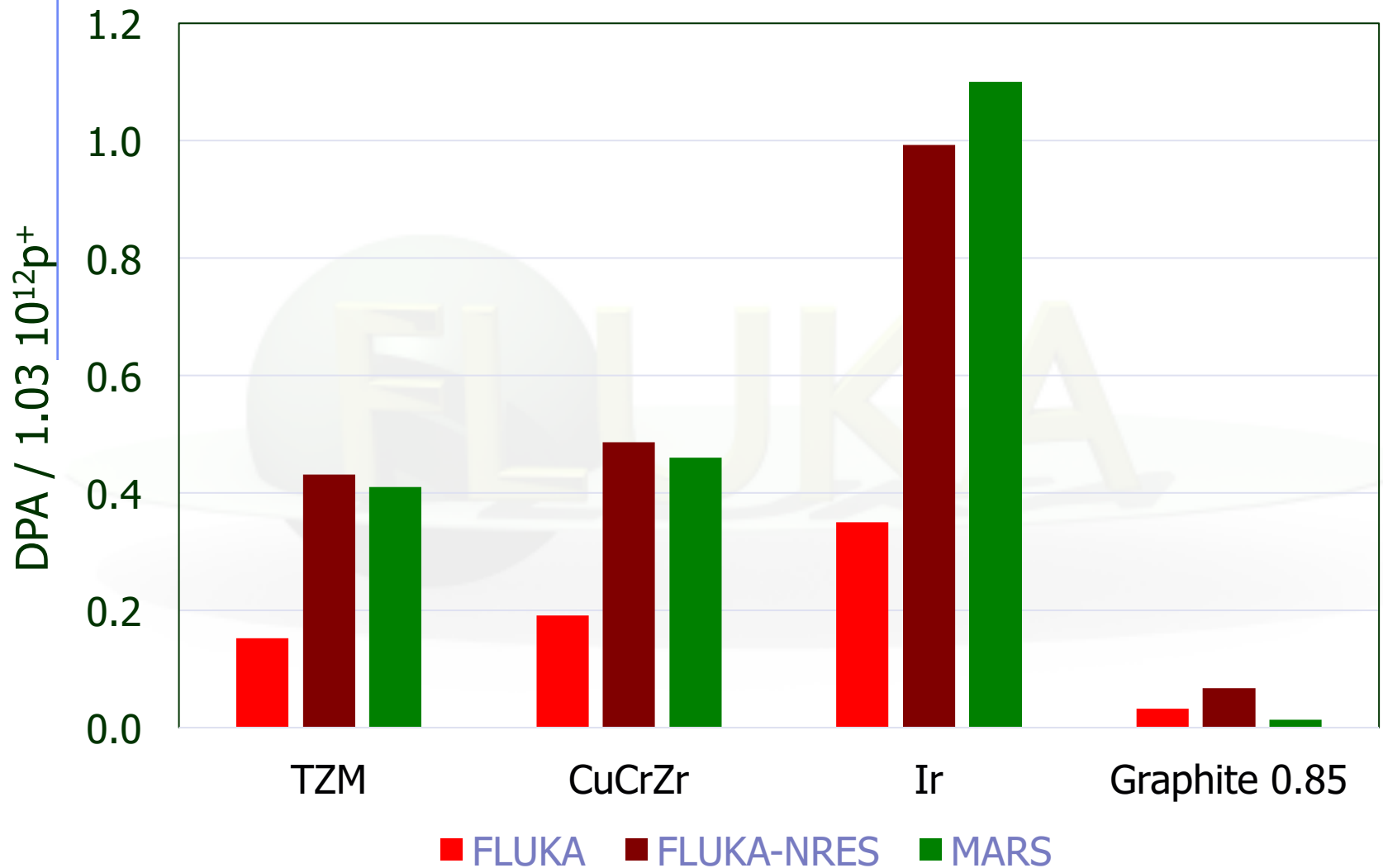


# BLIP capsule

- Beam
  - Proton E=181 MeV
  - $\sigma_{x,y} = 5.1$  mm
- Geometry: Layers of
  - Window SS304L 0.3mm
  - TZM 0.5mm
  - CuCrZr 0.5mm
  - Ir 0.5mm
  - Graphite(0.85g/cm<sup>3</sup>) 0.1mm

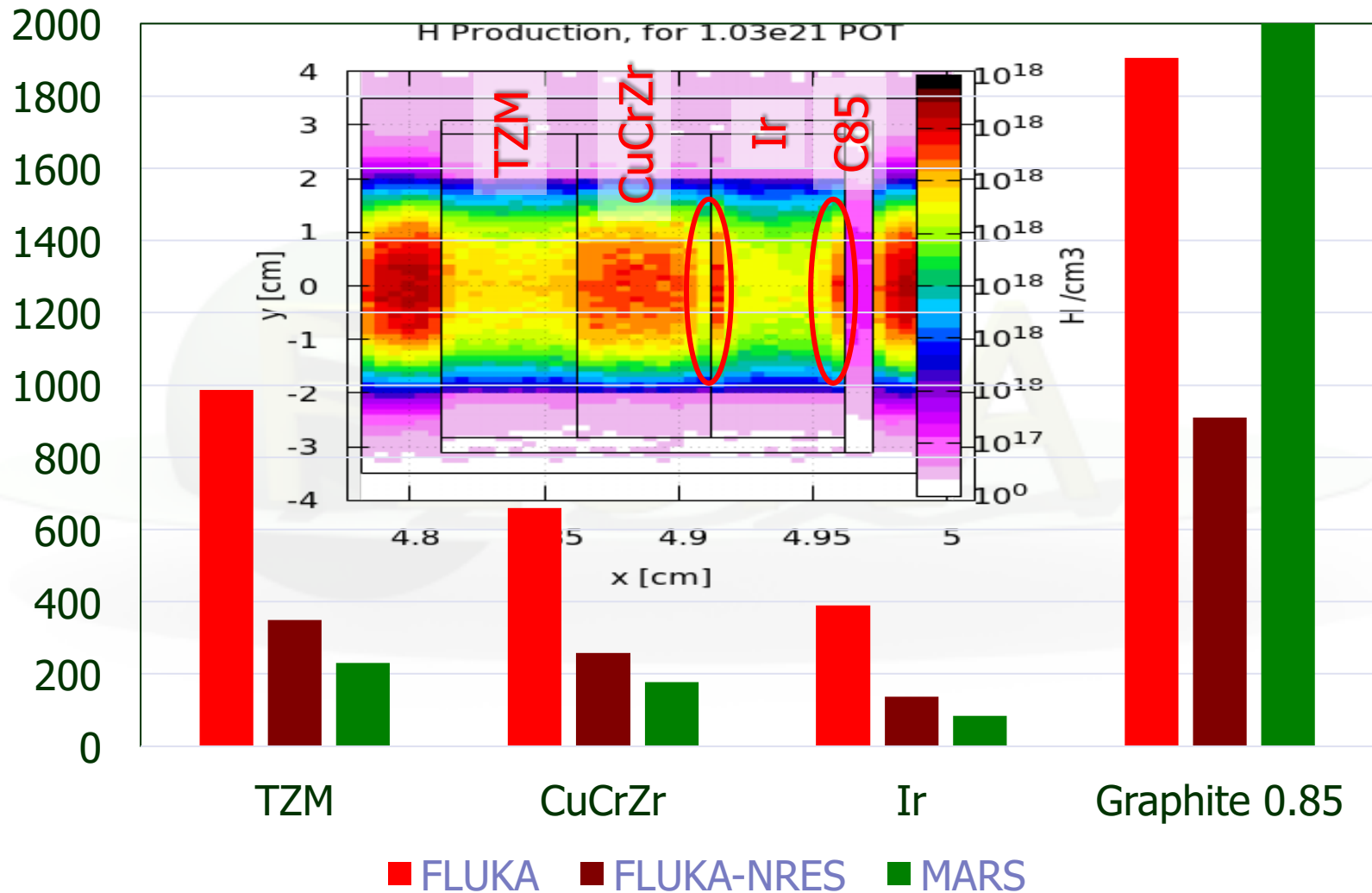


# DPA High-Z BLIP [FLUKA vs MARS]



Note: NRT model of MARS

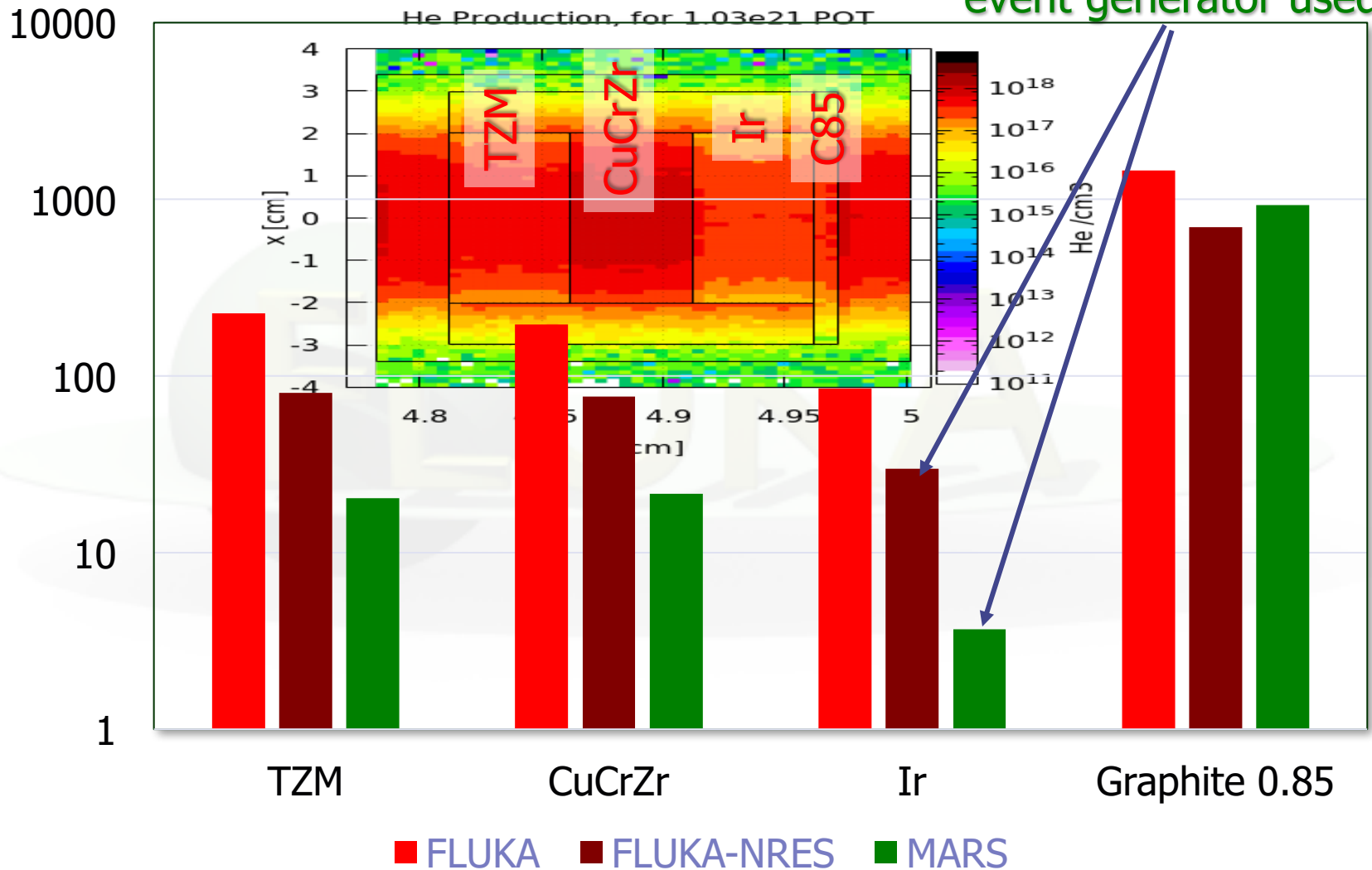
# H appm/DPA High-Z BLIP



# He appm/DPA High-Z BLIP

Probably due to "Old" MARS event generator used

Warning:  
Log-scale



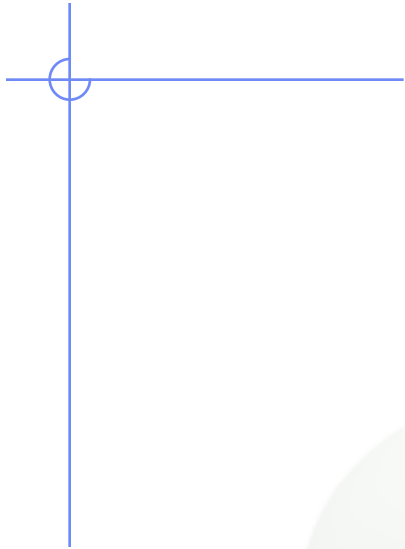


# Summary

- FLUKA dpa model uses a restricted NIEL computed during initialization and run time.
- Not based on Lindhard but reworked all formulas
- The **only free parameter** for the user is the **damage threshold**
- Uniform treatment from the transport threshold up to the highest energies
- Use of Stoller displacement efficiency instead of a fixed 0.8 as NRT suggest
- H/He production cross sections “in agreement” with available data

## Possible Future improvements:

- Implementation of the Nordlund arc-dpa
- More accurate recoil momentum cross section for **pair production** and **Bremsstrahlung**
- **Point wise** treatment of **low energy neutrons** will provide correct recoil information
- Multiple damage thresholds for compounds



$$N_F = \kappa \frac{\xi(T)T}{2E_{th}}$$

# $E_{th}$ Damage Threshold Energy

- $E_{th}$  is the value of the threshold displacement energy **averaged over all crystallographic directions** or a minimum energy to produce a defect

Element	$E_{th}$ (eV)	Element	$E_{th}$ (eV)
Lithium	10	Co	40
C in SiC	20	Ni	40
Graphite	30..35	Cu	40
Al	27	Nb	40
Si	25	Mo	60
Mn	40	W	90
Fe	40	Pb	25

Typical values used in NJOY99 code

- The only variable requested for FLUKA

**MAT-PROP**    *WHAT(1)*                    =  $E_{th}$  (eV)  
                       *WHAT(4,5,6)*                = Material range  
                       *SDUM*                            = **DPA-ENER**

$$N_F = \kappa \frac{\xi(T)T}{2E_{th}}$$

# $\kappa$ displacement efficiency

- $\kappa=0.8$  value deviates from the **hard sphere model** (K&P), and compensates for the forward scattering in the displacement cascade
- The displacement efficiency  $\kappa$  can be considered as independent of  $T$  only in the range of  $T \leq 1-2 \text{ keV}$ . At higher energies, the development of collision cascades results in **defect migration** and **recombination of Frenkel pairs** due to overlapping of different branches of a cascade which translates into decay of  $\kappa(T)$ .
- From molecular dynamics (MD\*) simulations of the primary cascade the number of surviving displacements,  $N_{MD}$ , normalized to the number of those from NRT model,  $N_{NRT}$ , decreases down to the values about 0.2–0.3 at  $T \approx 20-100 \text{ keV}$ . The efficiency in question only slightly depends on atomic number  $Z$  and the temperature.

$$N_{MD}/N_{NRT} = 0.3-1.3$$

$$N_{MD} / N_{NRT} = 0.3 - 1.3 \left( -\frac{9.57}{X} + \frac{17.1}{X^{4/3}} - \frac{8.81}{X^{5/3}} \right)$$

where  $X \equiv 20 T$  (in keV).

•Roger E. Stoller, *J. Nucl. Mat.*, 276 (2000) 22

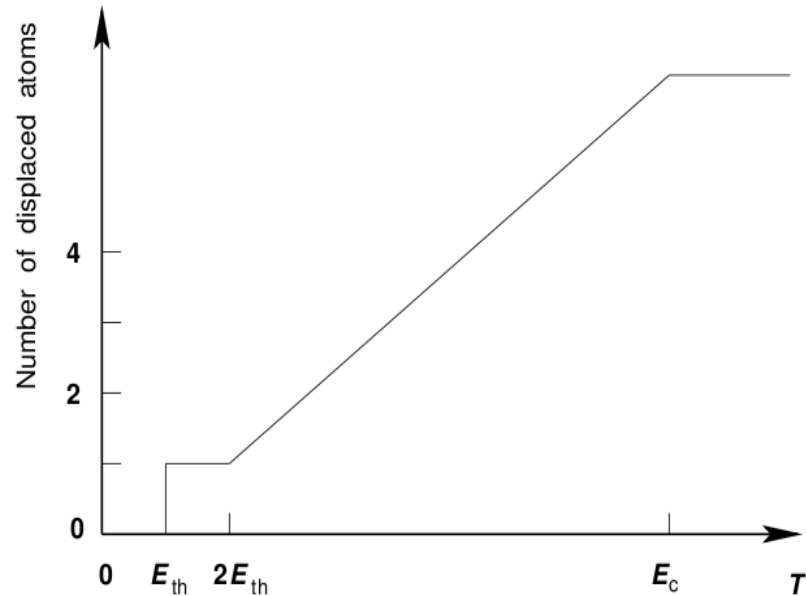
•D.J. Bacon, F. Gao and Yu.N. Osetsky, *J. Comp.-Aided Mat. Design*, 6 (1999) 225.

$$N_F = \kappa \frac{\xi(T)T}{2E_{th}}$$

# Factor of 2 (Kinchin & Pease)

- The cascade is created by a sequence of two-body elastic collisions between atoms
- In the collision process, the energy transferred to the lattice is zero
- For all energies  $T < E_c$  electronic stopping is ignored and only atomic collisions take place. No additional displacement occur above the cut-off energy  $E_c$
- The energy transfer cross section is given by the **hard-sphere** model.

$$\begin{aligned} v(T) &= 0 && \text{for } 0 < T < E_{th} \text{ (phonons)} \\ v(T) &= 1 && \text{for } E_{th} < T < 2E_{th} \\ v(T) &= T/2E_{th} && \text{for } 2E_{th} < T < E_c \\ v(T) &= E_c/2E_{th} && \text{for } T > E_c \end{aligned}$$



Schematic relation between the number of displaced atoms in the cascade and the kinetic energy  $T$  of the primary knock-on atom

Energy is equally shared between two atoms after the first collision  
Compensates for the energy lost to sub threshold reactions

$$N_F = \kappa \frac{\xi(T)T}{2E_{th}}$$

# Lindhard partition function $\xi$ [1/2]

- The partition function gives the fraction of **stopping power**  $S$  that goes to NIEL
- Approximations used: Electrons do not produce recoil nuclei with appreciable energy, lattice binding energy is neglected, etc...

$$(S_n + S_e)E'_n(E) = \int E_n(T) \frac{d\sigma_n}{dT} dT$$

where

$$S_{n,e}(E) = \int T_{n,e} d\sigma_{n,e}$$

- approximated to

$$\xi(T) = \frac{1}{1 + F_L \cdot (3.4008 \cdot \varepsilon(T)^{1/6} + 0.40244 \cdot \varepsilon(T)^{3/4} + \varepsilon(T))}$$

$$F_L = 30.724 \cdot Z_1 \cdot Z_2 \sqrt{Z_1^{2/3} + Z_2^{2/3}}$$

$$\varepsilon(T) = \frac{T}{0.0793 \frac{Z_1^{2/3} \cdot \sqrt{Z_2}}{(Z_1^{2/3} + Z_2^{2/3})^{3/4}} \cdot \frac{(A_1 + A_2)^{3/2}}{A_1^{3/2} \sqrt{A_2}}}$$

Z,A	charge and mass
1	projectile
2	medium
T	recoil energy (eV)

Nice feature: It can handle any projectile  $Z_1, A_1$  whichever charged particle

# Nuclear Stopping power

- Nuclear stopping power (unrestricted)

$$\frac{1}{\rho} S_n(E, E_{th}) = -2\pi N \int_0^{b_{\max}} b \frac{db}{d\theta} W(\theta, E) d\theta$$

- Energy transferred to recoil atom

$$W(\theta, T) = \gamma T \sin^2(\theta / 2)$$

- Deflection angle, by integrating over all impact parameters  $b$

$$\theta = \pi - 2 \int \frac{b dr}{r^2 \sqrt{1 - \frac{V(r)}{E_{cms}} - \frac{b^2}{r^2}}}$$

- Universal potential

$$V(r) = \frac{Z_1 Z_2 e^2}{r} F_s \left( \frac{r}{r_s} \right)$$

where:

$$\begin{aligned} F_s(x) &= \sum a_i \exp(-c_i x) \\ r_s &= 0.88534 r_B / (Z_1^{0.23} + Z_2^{0.23}) \\ r_s &= 0.88534 r_B Z_1^{-1/3} \end{aligned}$$

screening function  
screening length  
in case of particle



# Ziegler approximation

- Reduced kinetic energy  $\varepsilon$  ( $T$  in keV)

$$\varepsilon = \frac{32.536 T}{(Z_1^{0.23} + Z_2^{0.23}) \left( 1 + \frac{M_1}{M_2} \right) Z_1 Z_2}$$

- Reduced stopping power

$$\text{if } \varepsilon < 30 \quad \hat{S}_n(\varepsilon) = \frac{0.5 \ln(1 + 1.1383\varepsilon)}{\varepsilon + 0.01321 \varepsilon^{0.21226} + 0.19593\sqrt{\varepsilon}}$$

$$\text{if } \varepsilon \geq 30 \quad \hat{S}_n(\varepsilon) = \frac{\ln(\varepsilon)}{2\varepsilon}$$

## Important features of Reduced Stopping Power

- Independent from the **projectile** and **target** combination
- Accurate within **1%** for  $\varepsilon < 1$  and to within **5%** or better for  $\varepsilon > 3$
- Stopping power (MeV/g/cm<sup>2</sup>)

$$\frac{1}{\rho} S_n(T) = \frac{5105.3 Z_1 Z_2 \hat{S}_n(\varepsilon)}{(Z_1^{0.23} + Z_2^{0.23}) \left( 1 + \frac{M_2}{M_1} \right) A}$$

# Restricted Stopping Power

- The restricted nuclear stopping power is calculated the same way only integrating from 0 impact parameter up to a maximum  $b_{max}$  which corresponds to a transfer of energy equal to the

$$E_{th} = W_{min}(\theta_{min}, T)$$

$$\frac{1}{\rho} S_n(E, E_{th}) = -2\pi N \int_0^{b_{max}} b \frac{d\theta}{db} W(\theta, E) db$$

- To find  $b_{max}$  we have to approximately solve the previous  $\theta$  integral using an iterative approach for

$$\theta_{min} = 2 \arcsin \left( \sqrt{\frac{E_{th}}{\gamma T}} \right)$$

This can be done either by integrating numerically for  $\theta$  or using the magic scattering formula from Biersack-Hagmark that gives a fitting to  $\sin^2(\theta/2)$

# Implementation: Charged Particles

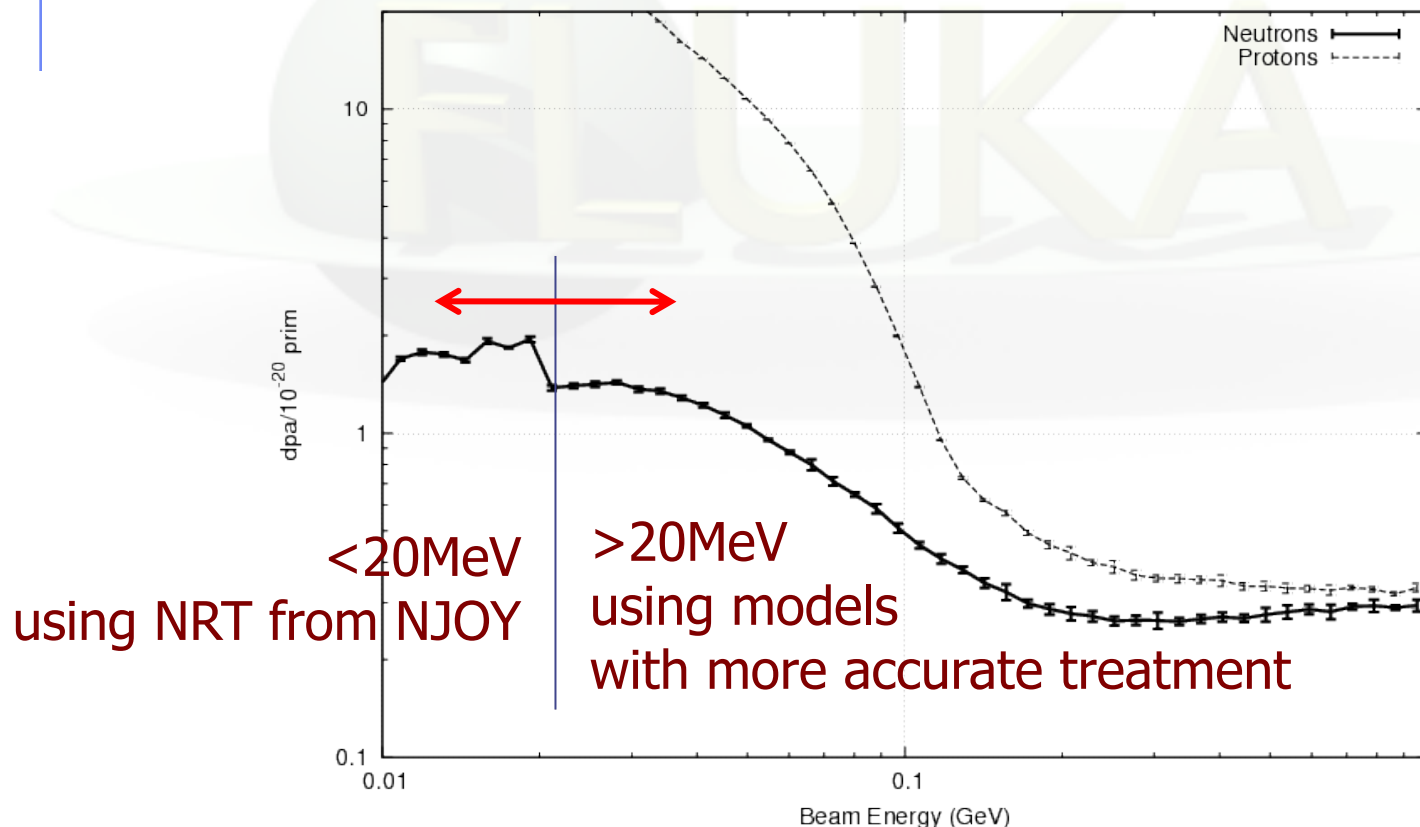
- During the transport of all charged particles and heavy ions the dpa estimation is based on the restricted nuclear stopping power while for NIEL on the unrestricted one.
- For every charged particle above the transport threshold and for every Monte Carlo step, the number of defects is calculated based on a modified multiple integral
- Taking into account also the **second level of sub-cascades initiated by the projectile**

$$N(E) = \int_{E_{th}}^{\gamma E} \left[ \underset{\substack{\text{restricted} \\ \text{partition function}}}{\xi_r(T, E_{th})} \left( \frac{d\sigma}{dT} \right)_E \int_{E_{th}}^{\gamma T} \left( \underset{\substack{\text{Lindhard} \\ \text{partition function}}}{\kappa(T') \xi(T')} T' \left( \frac{d\sigma}{dT'} \right)_{T'} dT' \right) dT$$

- Below the transport threshold (1 keV) it employs the Lindhard approximation

# Group Wise Neutron Artifacts

- Due to the group treatment of low-energy neutrons, there is no direct way to calculate properly the recoils.
- Therefore the evaluation is based on the KERMA factors calculated by NJOY, which in turn is based on the Unrestricted Nuclear losses from using the NRT model.



# Implementation: others

For Bremsstrahlung and pair production the recoil is sampled randomly from an approximation of the recoil momentum cross section

## Bremsstrahlung

$$\frac{d\sigma}{dp_{\perp}} = \frac{32a(Za)^2}{kp_{\perp}^3} \left[ 1 - \frac{k}{E} + \frac{1}{2} \left( \frac{k}{E} \right)^2 \right] \ln \left( \frac{p_{\perp}}{m_e} \right)$$

## Pair production

$$\frac{d\sigma}{dp} = \frac{0.183 \cdot 10^{-2} Z^2}{p^3} (\ln(p) + 0.5)$$

both can be written in the same approximate way as

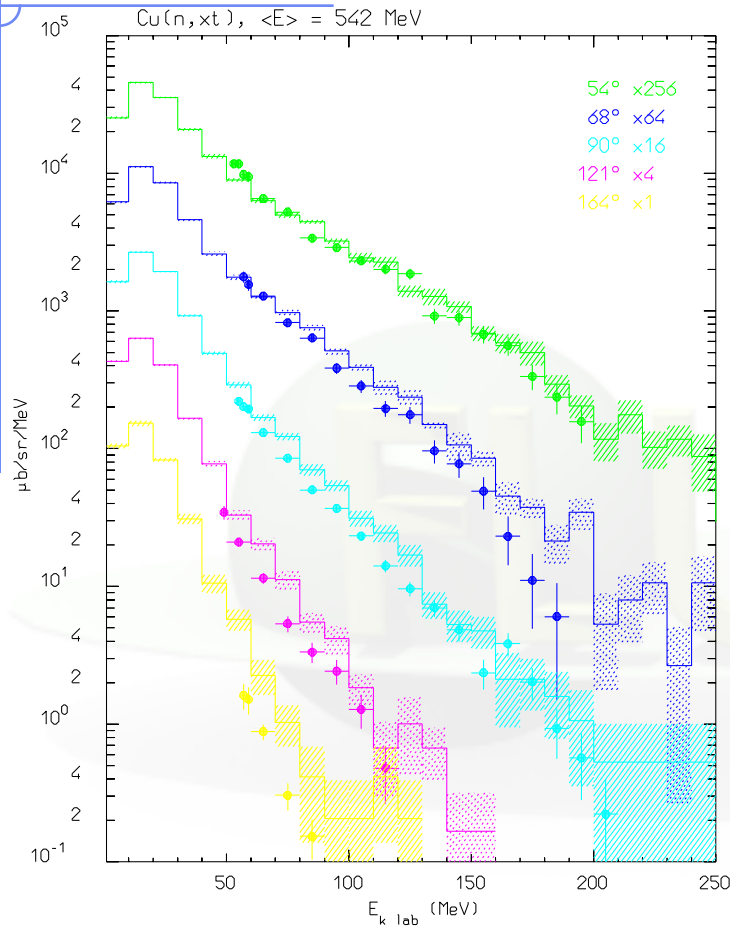
$$\frac{d\sigma}{dp} \propto \frac{\ln(p/c)}{p^3}$$

where the recoil momentum is sampled randomly by rejection from a similar function

# Coalescence:

- d, t,  $^3\text{He}$ , and alpha's generated during the (G)INC and preequilibrium stage
- All possible combinations of (unbound) nucleons and/or light fragments checked at each stage of system evolution
- FOM evaluation based on phase space "closeness" used to decide whether a light fragment is formed rather than not
  - ❑ FOM evaluated in the CMS of the candidate fragment at the time of minimum distance
  - ❑ Naively a momentum or position FOM should be used, but not both due to quantum non commutation
  - ❑ ... however the best results are obtained with a Wigner transform FOM (assuming gaussian wave packets) which should be the correct way of considering together positions and momenta
- Binding energy redistributed between the emitted fragment and residual excitation (exact conservation of 4-momenta)

# Coalescence



High energy light fragments are emitted through the coalescence mechanism: “put together” emitted nucleons that are near in phase space.

Example : double differential t production from 542 MeV neutrons on Copper

**Warning:** coalescence is OFF by default  
Can be important, ex for . residual nuclei.  
To activate it:

PHYSICS 1.

COALESCE

If coalescence is on, switch on Heavy ion transport and interactions (see later)