

FPWS
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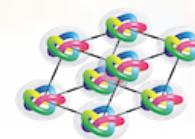
S4対称性と3HDMを用いたフレイバーモデルの構築

広島大学 井澤幸邑

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Graduate School of
Advanced Science and Engineering
Hiroshima University



SKCM²
WPI HIROSHIMA UNIVERSITY

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1. Introduction

4つの力

(強い力, 弱い力, 電磁気力, 重力)

標準模型(SM)

$SU(3) \times SU(2)_L \times U(1)_Y$ gauge対称性

クォークとレプトン (SM粒子)

→ 世代構造

(質量の階層性と世代混合)

特にレプトン→大きな世代混合



SMで理論的な説明が成されていない。
(SMはparameterを充てているのみ)

Standard Model of Elementary Particles and Gravity

three generations of matter (fermions)			interactions / force carriers (bosons)		
	I	II	III		
mass	$\approx 2.2 \text{ MeV}/c^2$	$\approx 1.28 \text{ GeV}/c^2$	$\approx 173.1 \text{ GeV}/c^2$	0	$\approx 124.97 \text{ GeV}/c^2$
charge	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0	0
spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	0
	u up	c charm	t top	g gluon	H higgs
	d down	s strange	b bottom	γ photon	
	e electron	μ muon	τ tau	Z Z boson	
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson	
					G graviton

QUARKS (I, II, III)
 LEPTONS (e, μ, τ, ν_e, ν_μ, ν_τ)
 GAUGE BOSONS VECTOR BOSONS (g, γ, Z, W)
 SCALAR BOSONS (H)
 HYPOTHETICAL TENSOR BOSONS (G)

<https://www.wikiwand.com/>

1. Introduction

AltarelliとFeruglioは非可換離散対称性を世代間に課した(flavor対称性)

G. Altarelli and F. Feruglio, Nucl. Phys. B741 (2006), 215–235.



本研究では世代対称性として S_4 対称性を用いる

加えて、three Higgs doublets model(3HDM)を用いる

新しいflavorモデルを構築し、数値解析を行う

Standard Model of Elementary Particles and Gravity

	three generations of matter (fermions)			interactions / force carriers (bosons)		
	I	II	III			
mass	$\approx 2.2 \text{ MeV}/c^2$	$\approx 1.28 \text{ GeV}/c^2$	$\approx 173.1 \text{ GeV}/c^2$	0	$\approx 124.97 \text{ GeV}/c^2$	0
charge	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	0	0	0
spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	0	2
	u up	c charm	t top	g gluon	H higgs	G graviton
	d down	s strange	b bottom	γ photon		
	e electron	μ muon	τ tau	Z Z boson		
	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	W W boson		

QUARKS (left side of fermion table)
 LEPTONS (left side of fermion table)
 GAUGE BOSONS VECTOR BOSONS (right side of boson table)
 SCALAR BOSONS (right side of boson table)
 HYPOTHETICAL TENSOR BOSONS (right side of boson table)

<https://www.wikiwand.com/>

2. S_4 対称性 S_4 対称性：4次の対称群

$$(x_1, x_2, x_3, x_4) \rightarrow (x_i, x_j, x_k, x_l) \quad 4! = 24 \text{ 要素}$$

表現 $1_1, 1_2, \underline{2}, 3_1, 3_2$  A_4 対称性：4次の交代群表現 $1, 1', 1'', 3_S, 3_A$ 12 要素

掛け算測

$$3_1 \times 3_1 = 1_1 + 2 + 3_1 + 3_2$$

$$3_2 \times 3_2 = 1_1 + 2 + 3_1 + 3_2$$

$$3_1 \times 3_2 = 1_2 + 2 + 3_1 + 3_2$$

$$2 \times 2 = 1_1 + 1_2 + 2$$

$$2 \times 3_1 = 3_1 + 3_2$$

$$2 \times 3_2 = 3_1 + 3_2$$

$$3_1 \times 1_2 = 3_2$$

$$3_2 \times 1_2 = 3_1$$

$$2 \times 1_2 = 2$$

$$\begin{pmatrix} \alpha_1 \\ \alpha_2 \end{pmatrix}_2 \otimes \begin{pmatrix} \beta_1 \\ \beta_2 \end{pmatrix}_2 = (\alpha_1\beta_1 + \alpha_2\beta_2)_{1_1} \oplus (-\alpha_1\beta_2 + \alpha_2\beta_1)_{1_2} \oplus \begin{pmatrix} \alpha_1\beta_2 + \alpha_2\beta_1 \\ \alpha_1\beta_1 - \alpha_2\beta_2 \end{pmatrix}_2$$

$$\begin{pmatrix} \alpha_1 \\ \alpha_2 \\ \alpha_3 \end{pmatrix}_{3_1} \otimes \begin{pmatrix} \beta_1 \\ \beta_2 \\ \beta_3 \end{pmatrix}_{3_1} = (\alpha_1\beta_1 + \alpha_2\beta_2 + \alpha_3\beta_3)_{1_1} \oplus \begin{pmatrix} \frac{1}{\sqrt{2}}(\alpha_2\beta_2 - \alpha_3\beta_3) \\ \frac{1}{\sqrt{6}}(-2\alpha_1\beta_1 + \alpha_2\beta_2 + \alpha_3\beta_3) \end{pmatrix}_2$$

$$\oplus \begin{pmatrix} \alpha_3\beta_2 + \alpha_2\beta_3 \\ \alpha_1\beta_3 + \alpha_3\beta_1 \\ \alpha_2\beta_1 + \alpha_1\beta_2 \end{pmatrix}_{3_1} \oplus \begin{pmatrix} \alpha_3\beta_2 - \alpha_2\beta_3 \\ \alpha_1\beta_3 - \alpha_3\beta_1 \\ \alpha_2\beta_1 - \alpha_1\beta_2 \end{pmatrix}_{3_2}$$

3. 3HDM

SM Higgs doubletを3つに拡張 (12個の 実scalar fields)

$$\phi_1 = \begin{pmatrix} \phi_1^+ \\ \phi_1^0 \end{pmatrix}, \phi_2 = \begin{pmatrix} \phi_2^+ \\ \phi_2^0 \end{pmatrix}, \phi_3 = \begin{pmatrix} \phi_3^+ \\ \phi_3^0 \end{pmatrix}$$

一般的な3HDMのHiggs Potential

$$V = - \sum_{i,j=1}^3 m_{ij}^2 (\phi_i^\dagger \phi_j) + \frac{1}{2} \sum_{i,j,k,l=1}^3 \lambda_{ijkl} (\phi_i^\dagger \phi_j) (\phi_k^\dagger \phi_l)$$

Potentialの最小条件

$$\left(\frac{\partial V}{\partial \phi_1} \right)_{\phi_1=\langle \phi_1 \rangle, \phi_2=\langle \phi_2 \rangle, \phi_3=\langle \phi_3 \rangle} = 0$$

$$\left(\frac{\partial V}{\partial \phi_2} \right)_{\phi_1=\langle \phi_1 \rangle, \phi_2=\langle \phi_2 \rangle, \phi_3=\langle \phi_3 \rangle} = 0$$

$$\left(\frac{\partial V}{\partial \phi_3} \right)_{\phi_1=\langle \phi_1 \rangle, \phi_2=\langle \phi_2 \rangle, \phi_3=\langle \phi_3 \rangle} = 0$$



自発的対称性の破れ

3つの自由度がWとZ bosonsに食われる

→ ϕ は9 (=12-3)つの 実scalar fieldsで表わされる

$$\phi_i = \begin{pmatrix} \rho_i^+ \\ \frac{1}{\sqrt{2}}(v_i + \rho_i + i\chi_i) \end{pmatrix}, i = 1, 2, 3$$

mass eigenstates



- (i) Three CP-even scalar fields
- (ii) Two CP-odd scalar fields
- (iii) Four charged scalar fields

4. Flavor模型

	$\bar{l} = (\bar{l}_e, \bar{l}_\mu, \bar{l}_\tau)$	$l_R = (e_R, \mu_R)$	τ_R	ν_{eR}	$\nu_R = (\nu_{\mu R}, \nu_{\tau R})$	$\phi = (\phi_1, \phi_2, \phi_3)$	X	Θ
$SU(2)_L$	2	1	1	1	1	2	1	1
S_4	3	2	1	1	2	3	2	1
$U(1)_{FN}$	0	+1	0	0	0	0	-1	-1

Lagrangian : $-L_Y = L_l + L_D + L_M + h.c.$

(1) 荷電レプトンの質量項: $L_l = \frac{y_{e\mu}}{\Lambda} \bar{l} \phi l_R \Theta + y_\tau \bar{l} \phi \tau_R + \frac{y_l}{\Lambda} \bar{l} \phi l_R X$

(2) Dirac neutrinoの質量項: $L_D = y_{De} \bar{l} \tilde{\phi} \nu_{eR} + y_{D\mu\tau} \bar{l} \tilde{\phi} \nu_R$

(3) 右巻きMajorana neutrinoの質量項: $L_M = \frac{1}{2} M_{eR} \bar{\nu}_{eR}^c \nu_{eR} + \frac{1}{2} M_{\mu\tau R} \bar{\nu}_R^c \nu_R$



荷電レプトンと左巻きMajorana Neutrinoの質量行列を求める

質量行列の計算

(1) 荷電レプトンの質量項

$$L_l = \frac{y_{e\mu}}{\Lambda} \bar{l}\phi l_R \Theta + y_\tau \bar{l}\phi \tau_R + \frac{y_l}{\Lambda} \bar{l}\phi l_R X$$

$$\frac{y_l}{\Lambda} \begin{pmatrix} \bar{l}_e \\ \bar{l}_\mu \\ \bar{l}_\tau \end{pmatrix}_3 \otimes \begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \end{pmatrix}_3 \otimes \begin{pmatrix} e_R \\ \mu_R \end{pmatrix}_2 \Theta = \frac{y_{e\mu}}{\Lambda} \begin{pmatrix} \frac{1}{\sqrt{2}}(\bar{l}_\mu \phi_2 - \bar{l}_\tau \phi_3) \\ \frac{1}{\sqrt{6}}(-2\bar{l}_e \phi_1 + \bar{l}_\mu \phi_2 + \bar{l}_\tau \phi_3) \end{pmatrix}_2 \otimes \begin{pmatrix} e_R \\ \mu_R \end{pmatrix}_2 \Theta$$

$$= \frac{y_{e\mu}}{\Lambda} \frac{1}{\sqrt{2}} (\bar{l}_\mu \phi_2 - \bar{l}_\tau \phi_3) e_R \Theta + \frac{y_{e\mu}}{\Lambda} \frac{1}{\sqrt{6}} (-2\bar{l}_e \phi_1 + \bar{l}_\mu \phi_2 + \bar{l}_\tau \phi_3) \mu_R \Theta$$

$$\Downarrow \quad \langle \Theta \rangle = \Theta_0, \langle \phi \rangle = (v_1, v_2, v_3)$$

$$= \frac{y_{e\mu} \Theta_0}{\Lambda} \frac{1}{\sqrt{2}} (\bar{\mu}_L v_2 - \bar{\tau}_L v_3) e_R + \frac{y_{e\mu} \Theta_0}{\Lambda} \frac{1}{\sqrt{6}} (-2\bar{e}_L v_1 + \bar{\mu}_L v_2 + \bar{\tau}_L v_3) \mu_R$$

$$\begin{pmatrix} \alpha_1 \\ \alpha_2 \end{pmatrix}_2 \otimes \begin{pmatrix} \beta_1 \\ \beta_2 \end{pmatrix}_2 = (\alpha_1 \beta_1 + \alpha_2 \beta_2)_1$$

質量行列の計算

(1) 荷電レプトンの質量項

$$L_l = \frac{y_{e\mu}}{\Lambda} \bar{l}\phi l_R \Theta + \boxed{y_\tau \bar{l}\phi \tau_R} + \frac{y_l}{\Lambda} \bar{l}\phi l_R X$$

$$y_\tau \underbrace{\begin{pmatrix} \bar{l}_e \\ \bar{l}_\mu \\ \bar{l}_\tau \end{pmatrix}_3}_{1} \otimes \begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \end{pmatrix}_3 \otimes \tau_R = y_\tau (\bar{l}_e \phi_1 + \bar{l}_\mu \phi_2 + \bar{l}_\tau \phi_3) \tau_R$$

↓ $\langle \phi \rangle = (v_1, v_2, v_3)$

$$= y_\tau (\bar{e}_L v_1 \tau_R + \bar{\mu}_L v_2 \tau_R + \bar{\tau}_L v_3 \tau_R)$$

1項目と2項目からの
荷電レプトンの質量行列

$$M_{l1} = \begin{pmatrix} 0 & -\frac{2y_{e\mu}\Theta_0}{\sqrt{6}\Lambda} v_1 & y_\tau v_1 \\ \frac{y_{e\mu}\Theta_0}{\sqrt{2}\Lambda} v_2 & \frac{y_{e\mu}\Theta_0}{\sqrt{6}\Lambda} v_2 & y_\tau v_2 \\ -\frac{y_{e\mu}\Theta_0}{\sqrt{2}\Lambda} v_3 & \frac{y_{e\mu}\Theta_0}{\sqrt{6}\Lambda} v_3 & y_\tau v_3 \end{pmatrix}_{LR}$$

質量行列の計算

(1) 荷電レプトンの質量項

$$L_l = \frac{y_{e\mu}}{\Lambda} \bar{l}\phi l_R \Theta + y_\tau \bar{l}\phi \tau_R + \frac{y_l}{\Lambda} \bar{l}\phi l_R X$$

$$\begin{aligned} & \frac{y_l}{\Lambda} \begin{pmatrix} \bar{l}_e \\ \bar{l}_\mu \\ \bar{l}_\tau \end{pmatrix}_3 \otimes \begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \end{pmatrix}_3 \otimes \begin{pmatrix} e_R \\ \mu_R \end{pmatrix}_2 \otimes \begin{pmatrix} X_1 \\ X_2 \end{pmatrix}_2 \quad \text{--- } 1,2 \\ & = \frac{y_{l1}}{\Lambda} (\bar{l}_e \phi_1 + \bar{l}_\mu \phi_2 + \bar{l}_\tau \phi_3)_1 \otimes (e_R X_1 + \mu_R X_2)_1 + \frac{y_{l2}}{\Lambda} \begin{pmatrix} \frac{1}{\sqrt{2}}(\bar{l}_\mu \phi_2 - \bar{l}_\tau \phi_3) \\ \frac{1}{\sqrt{6}}(-2\bar{l}_e \phi_1 + \bar{l}_\mu \phi_2 + \bar{l}_\tau \phi_3) \end{pmatrix}_2 \otimes \begin{pmatrix} e_R X_2 + \mu_R X_1 \\ e_R X_1 - \mu_R X_2 \end{pmatrix}_2 \\ & = \frac{y_{l1}}{\Lambda} (\bar{l}_e \phi_1 e_R X_1 + \bar{l}_e \phi_1 \mu_R X_2 + \bar{l}_\mu \phi_2 e_R X_1 + \bar{l}_\mu \phi_2 \mu_R X_2 + \bar{l}_\tau \phi_3 e_R X_1 + \bar{l}_\tau \phi_3 \mu_R X_2) \\ & \quad + \frac{y_{l2}}{\Lambda} \left[\frac{1}{\sqrt{2}} (\bar{l}_\mu \phi_2 e_R X_2 + \bar{l}_\mu \phi_2 \mu_R X_1 - \bar{l}_\tau \phi_3 e_R X_2 - \bar{l}_\tau \phi_3 \mu_R X_1) \right. \\ & \quad \left. + \frac{1}{\sqrt{6}} (-2\bar{l}_e \phi_1 e_R X_1 + 2\bar{l}_e \phi_1 \mu_R X_2 + \bar{l}_\mu \phi_2 e_R X_1 - \bar{l}_\mu \phi_2 \mu_R X_2 + \bar{l}_\tau \phi_3 e_R X_1 - \bar{l}_\tau \phi_3 \mu_R X_2) \right] \end{aligned}$$

$$\begin{pmatrix} \alpha_1 \\ \alpha_2 \\ \alpha_3 \end{pmatrix}_3 \otimes \begin{pmatrix} \beta_1 \\ \beta_2 \\ \beta_3 \end{pmatrix}_3 = \begin{pmatrix} \frac{1}{\sqrt{2}}(\alpha_2 \beta_2 - \alpha_3 \beta_3) \\ \frac{1}{\sqrt{6}}(-2\alpha_1 \beta_1 + \alpha_2 \beta_2 + \alpha_3 \beta_3) \end{pmatrix}_2$$

$$\begin{pmatrix} \alpha_1 \\ \alpha_2 \end{pmatrix}_2 \otimes \begin{pmatrix} \beta_1 \\ \beta_2 \end{pmatrix}_2 = (\alpha_1 \beta_1 + \alpha_2 \beta_2)_1 \oplus (\alpha_1 \beta_2 + \alpha_2 \beta_1)_2 \oplus (\alpha_1 \beta_1 - \alpha_2 \beta_2)_2$$

質量行列の計算

(1) 荷電レプトンの質量項

$$L_l = \frac{y_{e\mu}}{\Lambda} \bar{l} \phi l_R \Theta + y_\tau \bar{l} \phi \tau_R + \frac{y_l}{\Lambda} \bar{l} \phi l_R X$$

$$\langle \phi \rangle = (v_1, v_2, v_3), \quad \langle X \rangle = (X_1, 0)$$

3つ目の項からの荷電レプトンの質量行列

$$M_{l2} = \begin{pmatrix} \left(\frac{y_{l1}}{\Lambda} - \frac{2y_{l2}}{\sqrt{6}\Lambda} \right) v_1 X_1 & 0 & 0 \\ \left(\frac{y_{l1}}{\Lambda} + \frac{y_{l2}}{\sqrt{6}\Lambda} \right) v_2 X_1 & \frac{y_{l2}}{\sqrt{2}\Lambda} v_2 X_1 & 0 \\ \left(\frac{y_{l1}}{\Lambda} - \frac{y_{l2}}{\sqrt{6}\Lambda} \right) v_3 X_1 & -\frac{y_{l2}}{\sqrt{2}\Lambda} v_3 X_1 & 0 \end{pmatrix}_{LR}$$

荷電レプトンの質量行列

$$M_l = M_{l1} + M_{l2} = \begin{pmatrix} 0 & -\frac{2y_{e\mu}\Theta_0}{\sqrt{6}\Lambda} v_1 & y_\tau v_1 \\ \frac{y_{e\mu}\Theta_0}{\sqrt{2}\Lambda} v_2 & \frac{y_{e\mu}\Theta_0}{\sqrt{6}\Lambda} v_2 & y_\tau v_2 \\ -\frac{y_{e\mu}\Theta_0}{\sqrt{2}\Lambda} v_3 & \frac{y_{e\mu}\Theta_0}{\sqrt{6}\Lambda} v_3 & y_\tau v_3 \end{pmatrix}_{LR} + \begin{pmatrix} \left(\frac{y_{l1}}{\Lambda} - \frac{2y_{l2}}{\sqrt{6}\Lambda} \right) v_1 X_1 & 0 & 0 \\ \left(\frac{y_{l1}}{\Lambda} + \frac{y_{l2}}{\sqrt{6}\Lambda} \right) v_2 X_1 & \frac{y_{l2}}{\sqrt{2}\Lambda} v_2 X_1 & 0 \\ \left(\frac{y_{l1}}{\Lambda} - \frac{y_{l2}}{\sqrt{6}\Lambda} \right) v_3 X_1 & -\frac{y_{l2}}{\sqrt{2}\Lambda} v_3 X_1 & 0 \end{pmatrix}_{LR}$$

質量行列の計算

(2) Dirac neutrinoの質量項

$$L_D = y_{De} \bar{l} \tilde{\phi} \nu_{eR} + y_{D\mu\tau} \bar{l} \tilde{\phi} \nu_R \quad \text{Dirac neutrinoの質量行列}$$

$$M_D = \begin{pmatrix} y_{De} v_1 & 0 & -2/\sqrt{6} y_{D\mu\tau} v_1 \\ y_{De} v_2 & 1/\sqrt{2} y_{D\mu\tau} v_2 & 1/\sqrt{6} y_{D\mu\tau} v_2 \\ y_{De} v_3 & -1/\sqrt{2} y_{D\mu\tau} v_3 & 1/\sqrt{6} y_{D\mu\tau} v_3 \end{pmatrix}_{LR}$$

(3) 右巻きのMajorana neutrinoの質量項

$$L_M = \frac{1}{2} M_{eR} \bar{\nu}_{eR}^c \nu_{eR} + \frac{1}{2} M_{\mu\tau R} \bar{\nu}_R^c \nu_R$$

右巻きのMajorana neutrino
の質量行列

$$M_R = \begin{pmatrix} M_{eR} & 0 & 0 \\ 0 & M_{\mu\tau R} & 0 \\ 0 & 0 & M_{\mu\tau R} \end{pmatrix}$$

type-I seesaw機構を用いて、左巻きのMajorana neutrinoの質量項

$$m_\nu = -M_D M_R^{-1} M_D^T$$

$$m_\nu = \begin{pmatrix} -\frac{e^{2i\phi} y_{De} y_{De}^2 v_1^2}{M_{eR}} - \frac{2e^{2i\phi} y_{\mu\tau} y_{D\mu\tau}^2 v_1^2}{3M_{\mu\tau R}} & -\frac{e^{2i\phi} y_{De} y_{De}^2 v_1 v_2}{M_{eR}} + \frac{e^{2i\phi} y_{\mu\tau} y_{D\mu\tau}^2 v_1 v_2}{3M_{\mu\tau R}} & -\frac{e^{2i\phi} y_{De} y_{De}^2 v_3 v_1}{M_{eR}} + \frac{e^{2i\phi} y_{\mu\tau} y_{D\mu\tau}^2 v_3 v_1}{3M_{\mu\tau R}} \\ -\frac{e^{2i\phi} y_{De} y_{De}^2 v_1 v_2}{M_{eR}} + \frac{e^{2i\phi} y_{\mu\tau} y_{D\mu\tau}^2 v_1 v_2}{3M_{\mu\tau R}} & -\frac{e^{2i\phi} y_{De} y_{De}^2 v_2^2}{M_{eR}} - \frac{2e^{2i\phi} y_{\mu\tau} y_{D\mu\tau}^2 v_2^2}{3M_{\mu\tau R}} & -\frac{e^{2i\phi} y_{De} y_{De}^2 v_2 v_3}{M_{eR}} + \frac{e^{2i\phi} y_{\mu\tau} y_{D\mu\tau}^2 v_2 v_3}{3M_{\mu\tau R}} \\ -\frac{e^{2i\phi} y_{De} y_{De}^2 v_3 v_1}{M_{eR}} + \frac{e^{2i\phi} y_{\mu\tau} y_{D\mu\tau}^2 v_3 v_1}{3M_{\mu\tau R}} & -\frac{e^{2i\phi} y_{De} y_{De}^2 v_2 v_3}{M_{eR}} + \frac{e^{2i\phi} y_{\mu\tau} y_{D\mu\tau}^2 v_2 v_3}{3M_{\mu\tau R}} & -\frac{e^{2i\phi} y_{De} y_{De}^2 v_3^2}{M_{eR}} - \frac{2e^{2i\phi} y_{\mu\tau} y_{D\mu\tau}^2 v_3^2}{3M_{\mu\tau R}} \end{pmatrix}$$

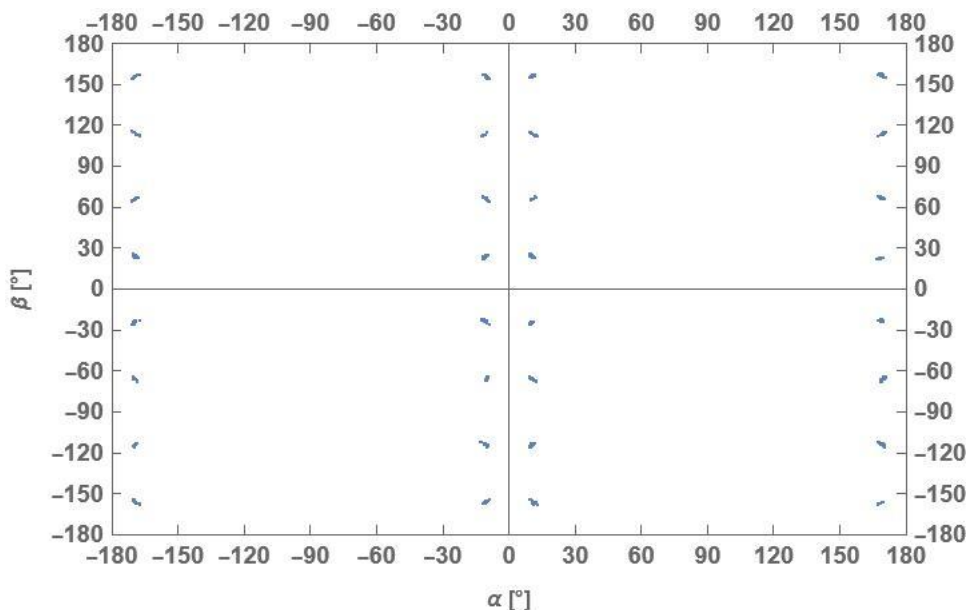
5. 数値解析

Satisfy $m_e, m_\mu, m_\tau, \Delta m_{21}^2, \Delta m_{31}^2, \sin^2 \theta_{12}, \sin^2 \theta_{23}, \sin^2 \theta_{13}$

Parameter $\alpha, \beta, y_{e\mu}, y_\tau, y_{l1}, y_{l2}, X_1, m_1, y_{De}, y_{D\mu\tau}, \phi_{yDe}$

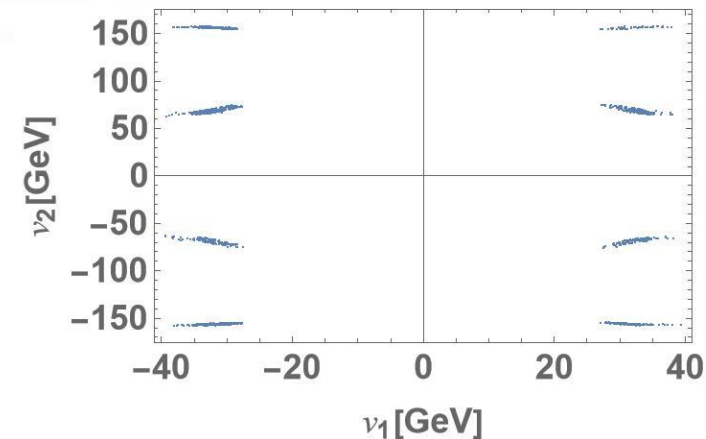
Prediction $\sin^2 \theta_{12}, \sin^2 \theta_{23}, \sin^2 \theta_{13}, \delta_{CP}, m_{light}, m_1 + m_2 + m_3,$
 m_{ee}, η_1, η_2

α と β の関係図

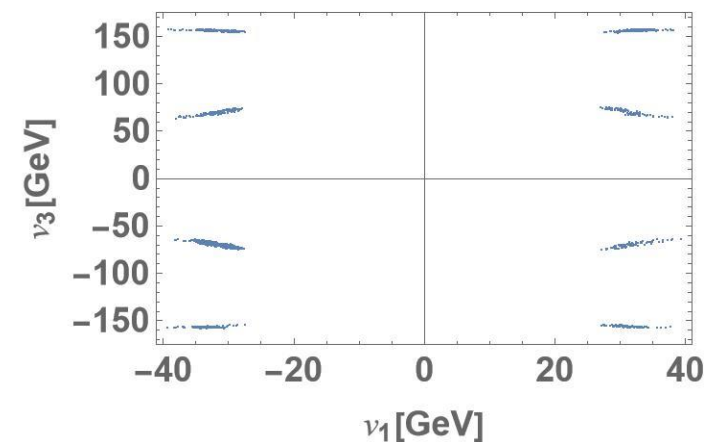


$$\begin{pmatrix} v_1 \\ v_2 \\ v_3 \end{pmatrix} = \begin{pmatrix} v \sin \alpha \\ v \cos \alpha \sin \beta \\ v \cos \alpha \cos \beta \end{pmatrix}$$

v_1 と v_2 の関係図

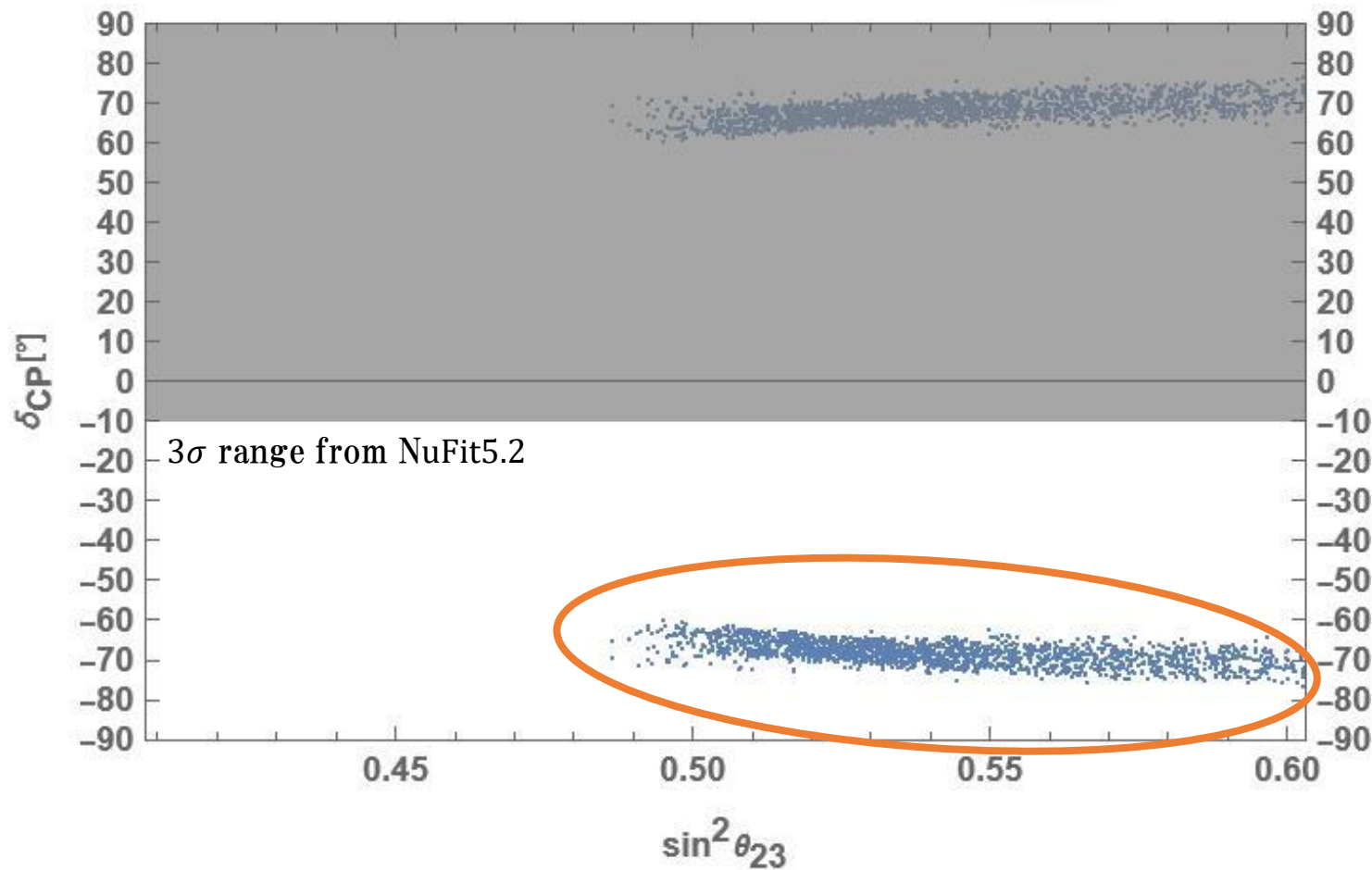


v_1 と v_3 の関係図



数値解析

δ_{CP} と $\sin^2\theta_{23}$ の予測



NuFIT 5.2
 $0.408 \leq \sin^2\theta_{23} \leq 0.603$

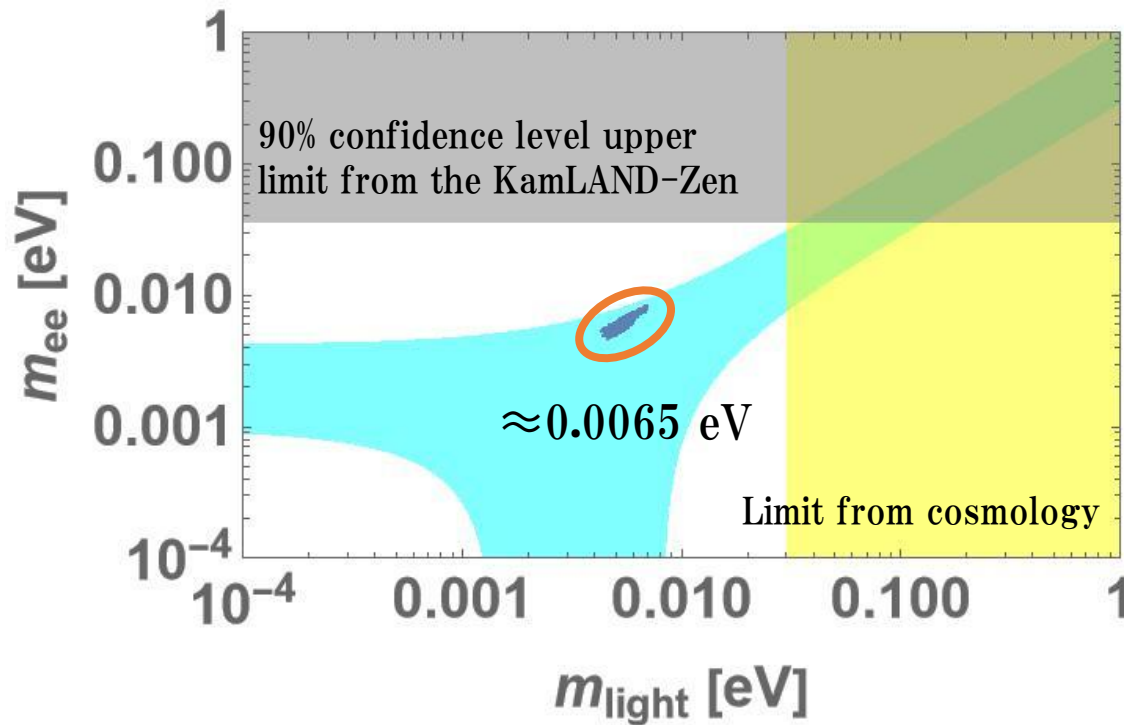
$0.488 \leq \sin^2\theta_{23} \leq 0.603$

$\delta_{CP} \approx -67.7^\circ$

δ_{CP} に対して強い予言

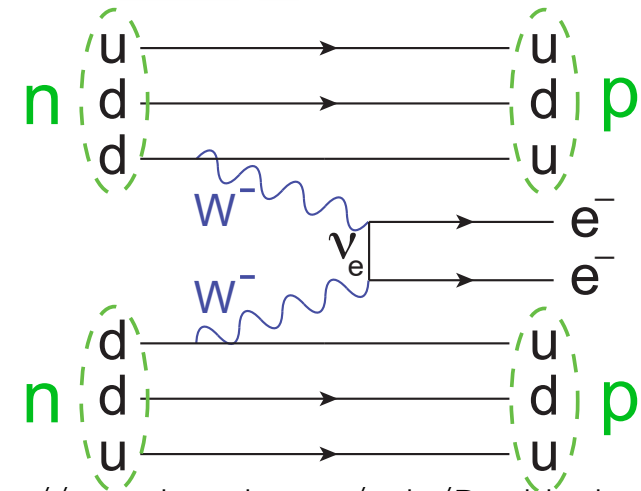
数値解析

Neutrinoの有効質量 m_{ee} と最も軽いneutrinoの質量 m_{light} の予測



比較的制限に近い場所に値が得られた

$0\nu\beta\beta$ decay



https://en.wikipedia.org/wiki/Double_beta_decay

Decay rate

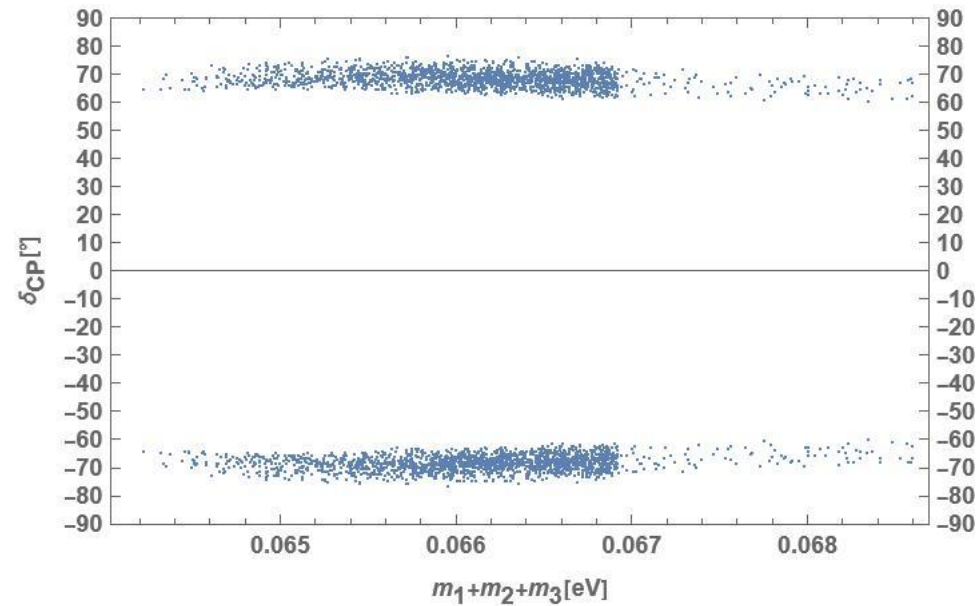
$$\Gamma \propto m_{ee}^2$$

Neutrinoの有効質量

$$m_{ee} = |m_1 U_{e1}^2 + m_2 U_{e2}^2 + m_3 U_{e3}^2|$$

数値解析

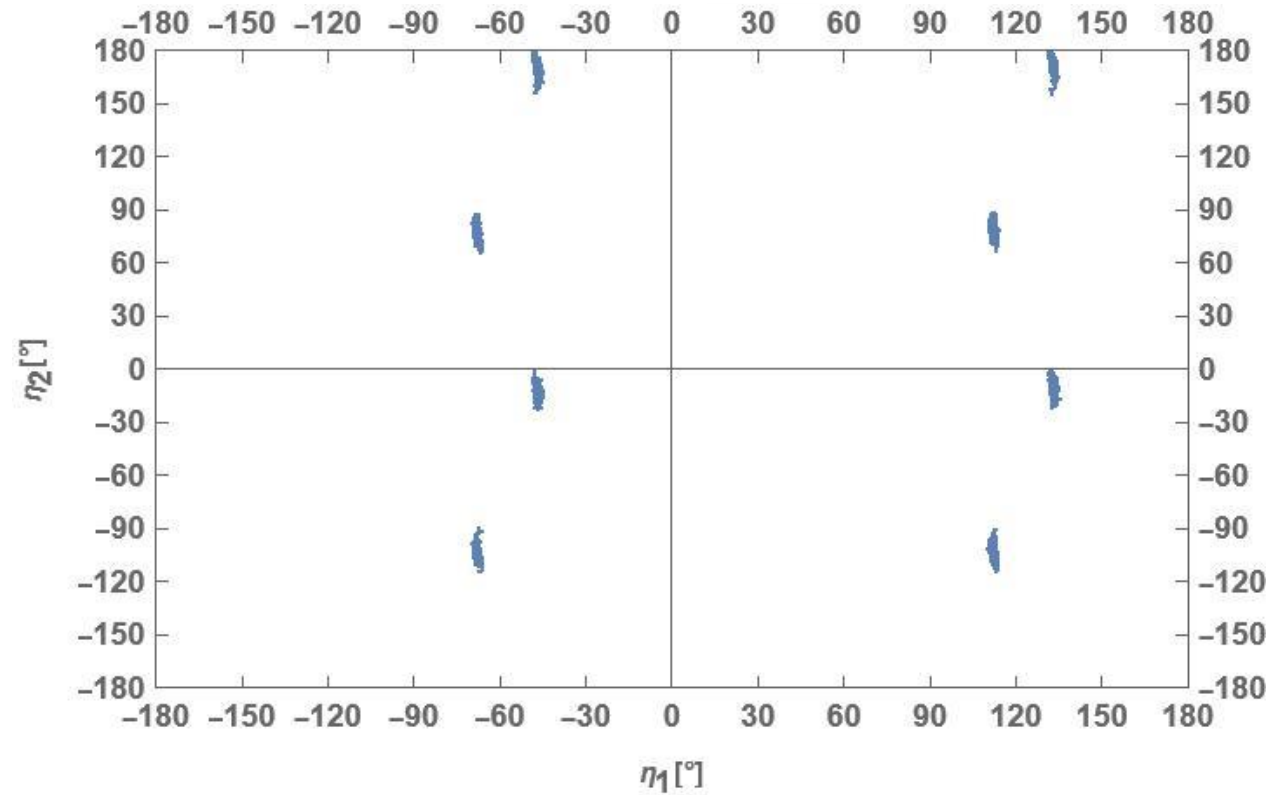
neutrinoの質量和 $m_1 + m_2 + m_3$ と δ_{CP} の予言



実験の制限(宇宙論)
(arXiv:1807.06209)
 $m_1 + m_2 + m_3 \leq 0.12\text{eV}$

数値解析

Majorana phases η_1, η_2 の予測



$$\eta_1 = \arg \left[\frac{U_{e1} U_{e3}^*}{\cos\theta_{12} \cos\theta_{13} \sin\theta_{13} e^{i\delta_{CP}}} \right], \quad \eta_2 = \arg \left[\frac{U_{e2} U_{e3}^*}{\sin\theta_{12} \cos\theta_{13} \sin\theta_{13} e^{i\delta_{CP}}} \right]$$

6. Potential解析

Higgs VEV $\langle \phi \rangle = (v_1, v_2, v_3)$

3HDM + S_4 対称性

ϕ を S_4 triplet と考える $\phi = (\phi_1, \phi_2, \phi_3)$

Higgs potential の計算 $V = -\mu^2 \phi^\dagger \phi + \lambda (\phi^\dagger \phi)^2 + c \phi^\dagger \phi X^\dagger X + k \phi^\dagger \phi \Theta^\dagger \Theta + (g \phi^\dagger \phi X^\dagger \Theta + h.c.)$

$$\phi^\dagger \phi = \begin{pmatrix} \phi_1^\dagger \\ \phi_2^\dagger \\ \phi_3^\dagger \end{pmatrix}_3 \otimes \begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \end{pmatrix}_3 = (\phi_1^\dagger \phi_1 + \phi_2^\dagger \phi_2 + \phi_3^\dagger \phi_3)_1 = |\phi_1|^2 + |\phi_2|^2 + |\phi_3|^2$$

Potential解析

Higgs potentialの計算 $V = -\mu^2 \phi^\dagger \phi + \lambda (\phi^\dagger \phi)^2 + c \phi^\dagger \phi X^\dagger X + k \phi^\dagger \phi \Theta^\dagger \Theta + (g \phi^\dagger \phi \Theta^\dagger X + h.c.)$

$$\begin{aligned}
 (\phi^\dagger \phi)^2 &= \underbrace{\begin{pmatrix} \phi_1^\dagger \\ \phi_2^\dagger \\ \phi_3^\dagger \end{pmatrix}_3}_{1,2,3,3'} \otimes \underbrace{\begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \end{pmatrix}_3}_{1,2,3,3'} = (|\phi_1|^2 + |\phi_2|^2 + |\phi_3|^2)^2 \oplus \underbrace{\begin{pmatrix} \frac{1}{\sqrt{2}}(\phi_2^\dagger \phi_2 - \phi_3^\dagger \phi_3) \\ \frac{1}{\sqrt{6}}(-2\phi_1^\dagger \phi_1 + \phi_2^\dagger \phi_2 + \phi_3^\dagger \phi_3) \end{pmatrix}_2}_{1,2,3,3'} \otimes \underbrace{\begin{pmatrix} \frac{1}{\sqrt{2}}(\phi_2^\dagger \phi_2 - \phi_3^\dagger \phi_3) \\ \frac{1}{\sqrt{6}}(-2\phi_1^\dagger \phi_1 + \phi_2^\dagger \phi_2 + \phi_3^\dagger \phi_3) \end{pmatrix}_2}_{1,2,3,3'} \\
 &\quad \oplus \underbrace{\begin{pmatrix} \phi_3^\dagger \phi_2 + \phi_2^\dagger \phi_3 \\ \phi_1^\dagger \phi_3 + \phi_3^\dagger \phi_1 \\ \phi_2^\dagger \phi_1 + \phi_1^\dagger \phi_2 \end{pmatrix}_3}_{1,2,3,3'} \otimes \underbrace{\begin{pmatrix} \phi_3^\dagger \phi_2 + \phi_2^\dagger \phi_3 \\ \phi_1^\dagger \phi_3 + \phi_3^\dagger \phi_1 \\ \phi_2^\dagger \phi_1 + \phi_1^\dagger \phi_2 \end{pmatrix}_3}_{1,2,3,3'} \oplus \underbrace{\begin{pmatrix} \phi_3^\dagger \phi_2 - \phi_2^\dagger \phi_3 \\ \phi_1^\dagger \phi_3 - \phi_3^\dagger \phi_1 \\ \phi_2^\dagger \phi_1 - \phi_1^\dagger \phi_2 \end{pmatrix}_{3'}}_{1,2,3,3'} \otimes \underbrace{\begin{pmatrix} \phi_3^\dagger \phi_2 - \phi_2^\dagger \phi_3 \\ \phi_1^\dagger \phi_3 - \phi_3^\dagger \phi_1 \\ \phi_2^\dagger \phi_1 - \phi_1^\dagger \phi_2 \end{pmatrix}_{3'}}_{1,2,3,3'} \\
 &= (|\phi_1|^2 + |\phi_2|^2 + |\phi_3|^2)^2 \\
 &\quad \oplus \frac{2}{3} (|\phi_1|^4 + |\phi_2|^4 + |\phi_3|^4 - |\phi_1|^2 |\phi_2|^2 - |\phi_2|^2 |\phi_3|^2 - |\phi_3|^2 |\phi_1|^2) \\
 &\quad \oplus \left[(\phi_1^\dagger \phi_2)^2 + (\phi_2^\dagger \phi_3)^2 + (\phi_3^\dagger \phi_1)^2 + |\phi_1|^2 |\phi_2|^2 + |\phi_2|^2 |\phi_3|^2 + |\phi_3|^2 |\phi_1|^2 + h.c. \right] \\
 &\quad \oplus \left[(\phi_1^\dagger \phi_2)^2 + (\phi_2^\dagger \phi_3)^2 + (\phi_3^\dagger \phi_1)^2 - |\phi_1|^2 |\phi_2|^2 - |\phi_2|^2 |\phi_3|^2 - |\phi_3|^2 |\phi_1|^2 + h.c. \right]
 \end{aligned}$$

Potential解析

Higgs potentialの計算

$$V = -\mu^2 \phi^\dagger \phi + \lambda (\phi^\dagger \phi)^2 + c \phi^\dagger \phi X^\dagger X + k \phi^\dagger \phi \Theta^\dagger \Theta + (g \phi^\dagger \phi \Theta^\dagger X + h.c.)$$

$$\phi^\dagger \phi X^\dagger X = \underbrace{\begin{pmatrix} \phi_1^\dagger \\ \phi_2^\dagger \\ \phi_3^\dagger \end{pmatrix}_3 \otimes \begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \end{pmatrix}_3}_{1,2} \otimes \underbrace{\begin{pmatrix} X_1^\dagger \\ X_2^\dagger \end{pmatrix}_2 \otimes \begin{pmatrix} X_1 \\ X_2 \end{pmatrix}_2}_{1,2} = (\phi_1^\dagger \phi_1 + \phi_2^\dagger \phi_2 + \phi_3^\dagger \phi_3)_1 \otimes (X_1^\dagger X_1 + X_2^\dagger X_2)_1 + \begin{pmatrix} \frac{1}{\sqrt{2}}(\phi_2^\dagger \phi_2 - \phi_3^\dagger \phi_3) \\ \frac{1}{\sqrt{6}}(-2\phi_1^\dagger \phi_1 + \phi_2^\dagger \phi_2 + \phi_3^\dagger \phi_3) \end{pmatrix}_2 \otimes \begin{pmatrix} X_1^\dagger X_2 + X_2^\dagger X_1 \\ X_1^\dagger X_1 - X_2^\dagger X_2 \end{pmatrix}_2$$

$$\langle X \rangle = (X_1, 0)$$

$$= |X_1|^2 (|\phi_1|^2 + |\phi_2|^2 + |\phi_3|^2) + \frac{|X_1|^2}{\sqrt{6}} (-2|\phi_1|^2 + |\phi_2|^2 + |\phi_3|^2)$$

$$\phi^\dagger \phi \Theta^\dagger \Theta = \begin{pmatrix} \phi_1^\dagger \\ \phi_2^\dagger \\ \phi_3^\dagger \end{pmatrix}_3 \otimes \begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \end{pmatrix}_3 \Theta^\dagger \Theta = |\Theta_0|^2 (|\phi_1|^2 + |\phi_2|^2 + |\phi_3|^2)$$

$$\langle \Theta \rangle = \Theta_0$$

$$\phi^\dagger \phi \Theta^\dagger X = \underbrace{\begin{pmatrix} \phi_1^\dagger \\ \phi_2^\dagger \\ \phi_3^\dagger \end{pmatrix}_3 \otimes \begin{pmatrix} \phi_1 \\ \phi_2 \\ \phi_3 \end{pmatrix}_3}_{2} \otimes \begin{pmatrix} X_1 \\ X_2 \end{pmatrix}_2 \Theta^\dagger = \begin{pmatrix} \frac{1}{\sqrt{2}}(\phi_2^\dagger \phi_2 - \phi_3^\dagger \phi_3) \\ \frac{1}{\sqrt{6}}(-2\phi_1^\dagger \phi_1 + \phi_2^\dagger \phi_2 + \phi_3^\dagger \phi_3) \end{pmatrix}_2 \otimes \begin{pmatrix} X_1 \\ X_2 \end{pmatrix}_2 \Theta^\dagger = \frac{\Theta_0^* X_1}{\sqrt{2}} (|\phi_2|^2 - |\phi_3|^2)$$

$$\langle X \rangle = (X_1, 0), \langle \Theta \rangle = \Theta_0$$

真空構造

Potential

$$\begin{aligned}
 V &= -\mu^2 \phi^\dagger \phi + \lambda (\phi^\dagger \phi)^2 + c \phi^\dagger \phi X^\dagger X + k \phi^\dagger \phi \Theta^\dagger \Theta + (g \phi^\dagger \phi \Theta^\dagger X + h.c.) \\
 &= \frac{\Lambda_1}{2} (|\phi_1|^4 + |\phi_2|^4 + |\phi_3|^4) + \Lambda_2 (|\phi_1|^2 |\phi_2|^2 + |\phi_2|^2 |\phi_3|^2 + |\phi_3|^2 |\phi_1|^2) \\
 &\quad + \frac{\Lambda_3}{2} \left[(\phi_1^\dagger \phi_2)^2 + (\phi_2^\dagger \phi_3)^2 + (\phi_3^\dagger \phi_1)^2 + h.c. \right] \\
 &\quad + \left(-\mu^2 + c_1 X_1^2 + \frac{2c_2 X_1^2}{\sqrt{6}} \right) |\phi_1|^2 + \left(-\mu^2 + c_1 X_1^2 + \frac{c_2 X_1^2}{\sqrt{6}} + \frac{g X_1 \Theta_0}{\sqrt{2}} \right) |\phi_2|^2 + \left(-\mu^2 + c_1 X_1^2 + \frac{c_2 X_1^2}{\sqrt{6}} - \frac{g X_1 \Theta_0}{\sqrt{2}} \right) |\phi_3|^2
 \end{aligned}$$

Potentialの最小条件

$$\left(\frac{\partial V}{\partial \phi_i} \right)_{\phi_1=v_1, \phi_2=v_2, \phi_3=v_3} = 0, \quad i = 1, 2, 3$$

vacuum expectation values

$$\begin{aligned}
 v_1 &= \pm \sqrt{\frac{\mu'^2}{\Lambda_1 + 2\Lambda'} - \frac{2\Lambda_1 c_2'}{(\Lambda_1 + 2\Lambda')(\Lambda_1 - \Lambda')}} \\
 v_2 &= \pm \sqrt{\frac{\mu'^2}{\Lambda_1 + 2\Lambda'} - \frac{\Lambda_1 - 2\Lambda'}{(\Lambda_1 + 2\Lambda')(\Lambda_1 - \Lambda')} c'^2_2 - \frac{g'}{\Lambda_1 - \Lambda'}} \\
 v_3 &= \pm \sqrt{\frac{\mu'^2}{\Lambda_1 + 2\Lambda'} - \frac{\Lambda_1 - 2\Lambda'}{(\Lambda_1 + 2\Lambda')(\Lambda_1 - \Lambda')} c'^2_2 + \frac{g'}{\Lambda_1 - \Lambda'}}
 \end{aligned}$$

VEVを
 v, α, β で書く

$$\langle \phi \rangle = \begin{pmatrix} v \sin \alpha \\ v \cos \alpha \sin \beta \\ v \cos \alpha \cos \beta \end{pmatrix}$$

v : Higgs VEV
 α, β : free parameter

7. Summary

flavor symmetryとして S_4 対称性を用いた
Scalar field ϕ を S_4 tripletとした



3HDM と S_4 対称性を用いたflavor模型を構築した
荷電レプトンとneutrinosの質量行列を計算した

荷電レプトンの質量行列

$$M_l = M_{l1} + M_{l2}$$

$$= \begin{pmatrix} 0 & -\frac{2y_{e\mu}\theta_0}{\sqrt{6}\Lambda}v_1 & y_\tau v_1 \\ \frac{y_{e\mu}\theta_0}{\sqrt{2}\Lambda}v_2 & \frac{y_{e\mu}\theta_0}{\sqrt{6}\Lambda}v_2 & y_\tau v_2 \\ -\frac{y_{e\mu}\theta_0}{\sqrt{2}\Lambda}v_3 & \frac{y_{e\mu}\theta_0}{\sqrt{6}\Lambda}v_3 & y_\tau v_3 \end{pmatrix}_{LR} + \begin{pmatrix} \left(\frac{y_{l1}}{\Lambda} - \frac{2y_{l2}}{\sqrt{6}\Lambda}\right)v_1 X_1 & 0 & 0 \\ \left(\frac{y_{l1}}{\Lambda} + \frac{y_{l2}}{\sqrt{6}\Lambda}\right)v_2 X_1 & \frac{y_{l2}}{\sqrt{2}\Lambda}v_2 X_1 & 0 \\ \left(\frac{y_{l1}}{\Lambda} - \frac{y_{l2}}{\sqrt{6}\Lambda}\right)v_3 X_1 & -\frac{y_{l2}}{\sqrt{2}\Lambda}v_3 X_1 & 0 \end{pmatrix}_{LR}$$



数値解析を行い、 δ_{CP} , neutrinoの有効質量 m_{ee} , Majorana phases η_1, η_2 .

δ_{CP} と m_{ee} ($m_{ee} \approx 0.0065$ [eV])に対して強い予言が得られた

→比較的実験で確かめやすい場所に結果が得られた

最後にHiggs potentialの解析を行い、 VEVを $\langle \phi \rangle = (v_1, v_2, v_3)$ の形で得た

neutrinoの質量行列

$$m_\nu = -M_D M_R^{-1} M_D^\dagger$$

$$M_D = \begin{pmatrix} y_{De}v_1 & 0 & -2/\sqrt{6}y_{D\mu\tau}v_1 \\ y_{De}v_2 & 1/\sqrt{2}y_{D\mu\tau}v_2 & 1/\sqrt{6}y_{D\mu\tau}v_2 \\ y_{De}v_3 & -1/\sqrt{2}y_{D\mu\tau}v_3 & 1/\sqrt{6}y_{D\mu\tau}v_3 \end{pmatrix}_{LR} \quad M_R = \begin{pmatrix} M_{eR} & 0 & 0 \\ 0 & M_{\mu\tau R} & 0 \\ 0 & 0 & M_{\mu\tau R} \end{pmatrix}$$