

A Semi-classical Spacetime Region with Maximum Entropy

RIKEN iTHEMS

Yuki Yokokura

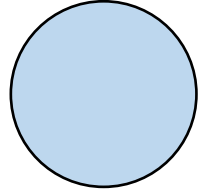
[\[arXiv: 2309.00602\]](#)

Motivation

Q1: What is the maximum entropy that can be given to a finite region?

Q2: What is the structure of such a spacetime $g_{\mu\nu}^*$?

⇒ Important for finding the fundamental d.o.f. in quantum gravity

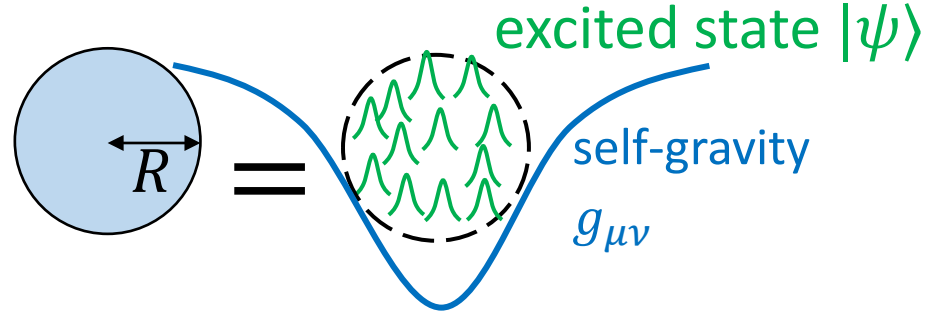


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- Consider these in a semi-classical level.



4D spherical static spacetime region = a collection of excited quanta in $(|\psi\rangle, g_{\mu\nu})$

semi-classical Einstein eq

$$G_{\mu\nu} = 8\pi G \langle \psi | T_{\mu\nu} | \psi \rangle$$

“weight of information”

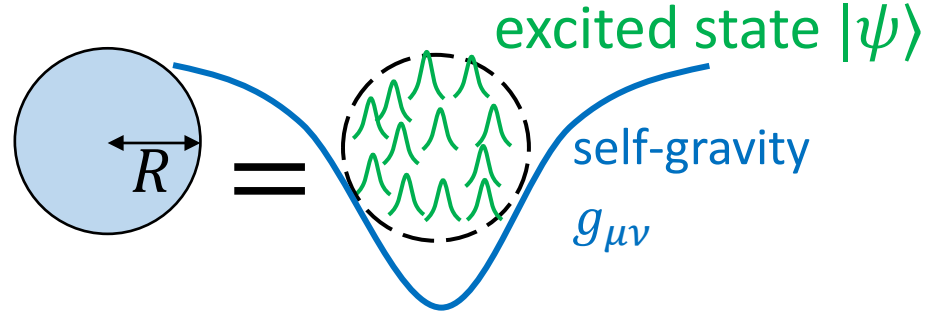
←self-consistent eq
 in a mean field approximation of QG
 gravity = classical $g_{\mu\nu}$
 matter = quantum $\hat{\phi}_i(x)$

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in a mean field approximation of QG
gravity = classical $g_{\mu\nu}$
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- We will estimate the entropy including self-gravity

$$S[g_{\mu\nu}] = \int_{\Sigma} d\Sigma_{\mu} s^{\mu}$$

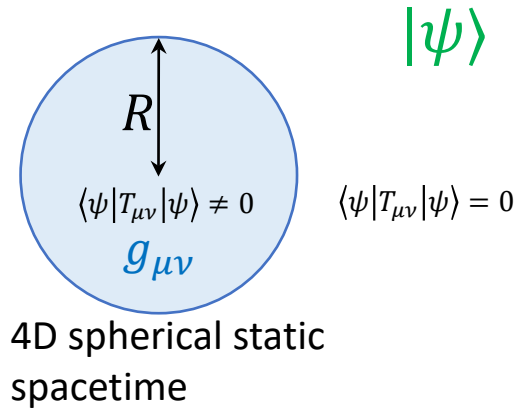
- many local d.o.f. with local interactions
- sufficiently excited states $\{|\psi\rangle\}$
- $\nabla_{\mu} s^{\mu} = 0$

and find $g_{\mu\nu}^*$ s.t.

$$S[g_{\mu\nu}] \leq S_{max} \equiv S[g_{\mu\nu}^*]$$

Setup: 1-bit quantum

- Consider $(g_{\mu\nu}, |\psi\rangle)$ satisfying $G_{\mu\nu} = 8\pi G \langle \psi | T_{\mu\nu} | \psi \rangle$ self-consistently.



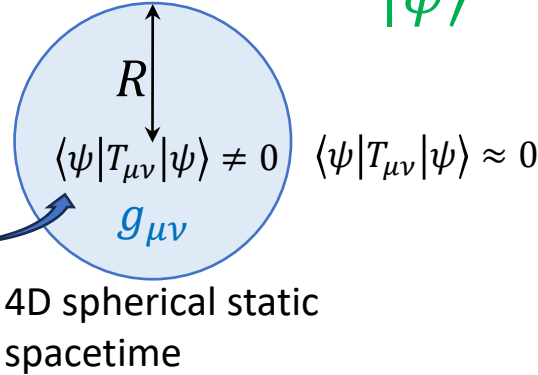
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metric

$$ds^2 = \begin{cases} -\left(1 - \frac{a_0}{r}\right) dt^2 + \frac{1}{1 - \frac{a_0}{r}} dr^2 + r^2 d\Omega^2 & \text{for } R \leq r \\ -\left(1 - \frac{a(r)}{r}\right) e^{A(r)} dt^2 + \frac{1}{1 - \frac{a(r)}{r}} dr^2 + r^2 d\Omega^2 & \text{for } r \leq R \end{cases}$$

$\frac{a(r)}{2G} \equiv m(r)$: Misner-Sharp energy at r
 $r \geq R \rightarrow \frac{a_0}{2G} \equiv M_0$: total ADM energy



4D spherical static spacetime

$|\psi\rangle$

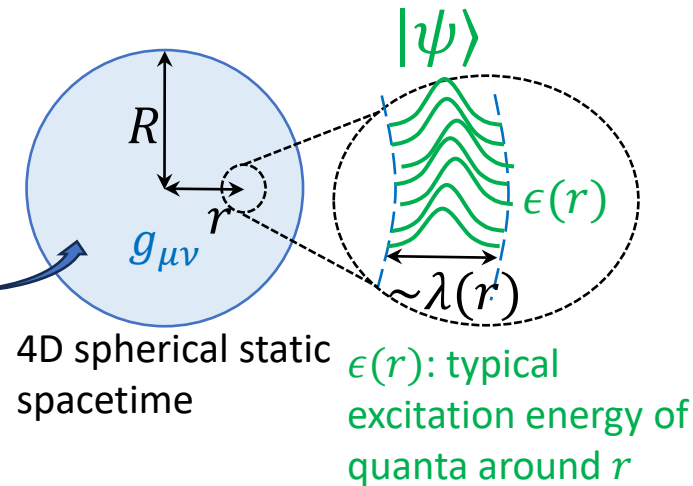
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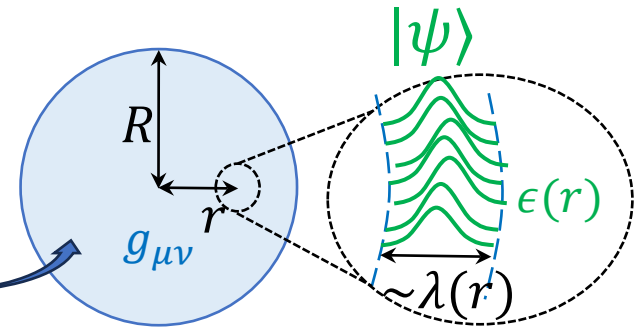
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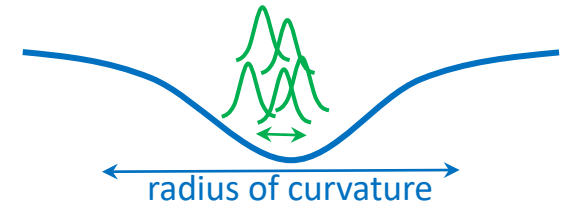
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$\epsilon(r)$: typical excitation energy of quanta around r

- For a sufficiently excited state $|\psi\rangle$,

$$\lambda(r) \sim \frac{\hbar}{\epsilon(r)} \lesssim \mathcal{R}(r)^{-\frac{1}{2}}$$

curvature



\Rightarrow WKB-like particle without feeling gravity

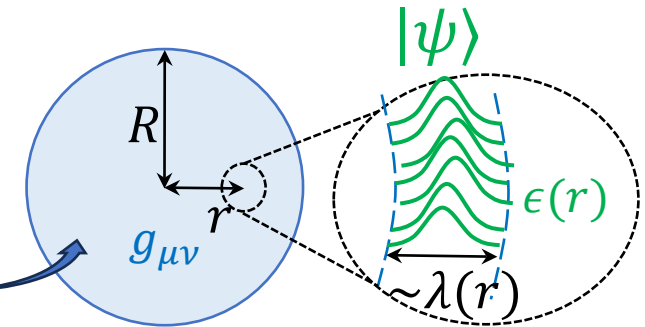
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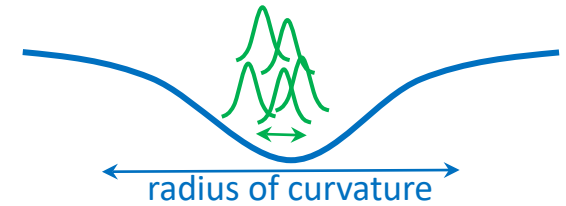
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radius of curvature

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- Many such quanta behave typically like a local thermal equilibrium:

$$\epsilon(r) \sim T_{loc}(r)$$

Thermodynamical Typicality
e.g. [Goldstein, et al. 2006]

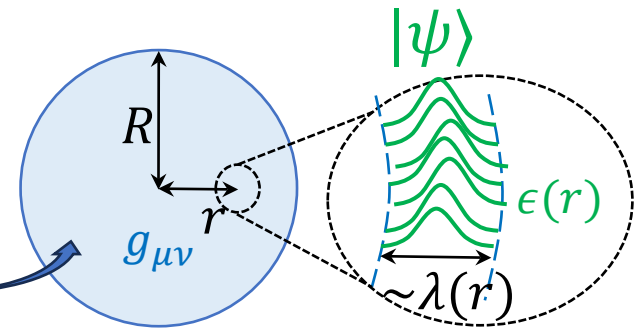
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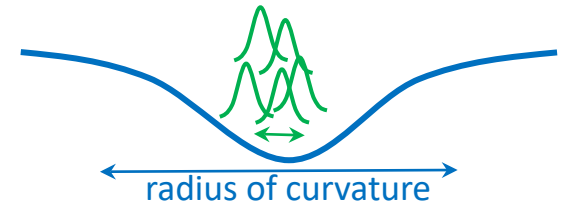


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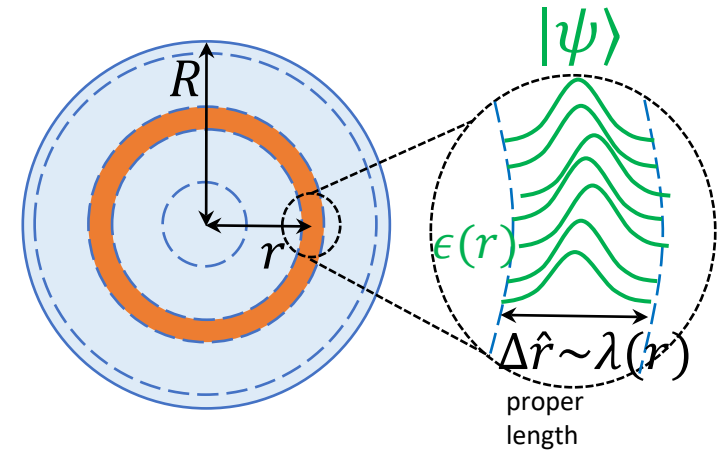
⇒ Such a quantum has 1-bit of entropy.

$$T = \frac{dE}{dS} = \text{energy/1bit}$$

“1-bit quantum”

WKB-like estimation of $S[g_{\mu\nu}]$

- Decompose the system into the smallest spherical subsystems.



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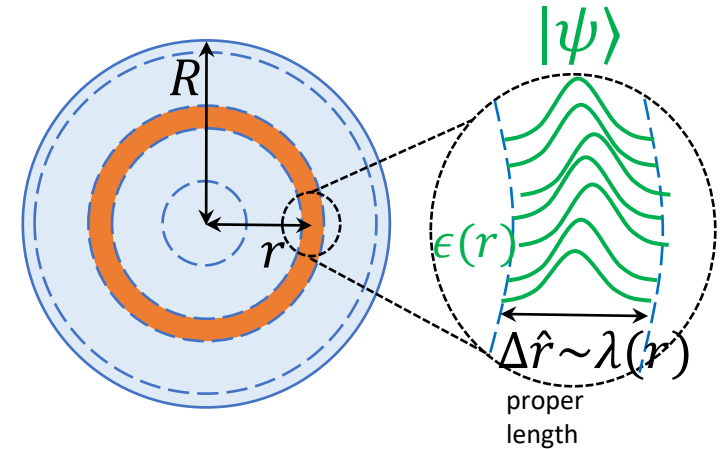
- Decompose the system into the smallest spherical subsystems.

- Define a quasi-local function

$$N(r) \equiv \frac{4\pi r^2 \langle T^{\hat{t}\hat{t}}(r) \rangle \Delta\hat{r}}{\epsilon(r)} \quad \text{for } \Delta\hat{r} \sim \lambda(r)$$

\sim # of 1-bit quanta in the subsystem for $N(r) \gg 1$

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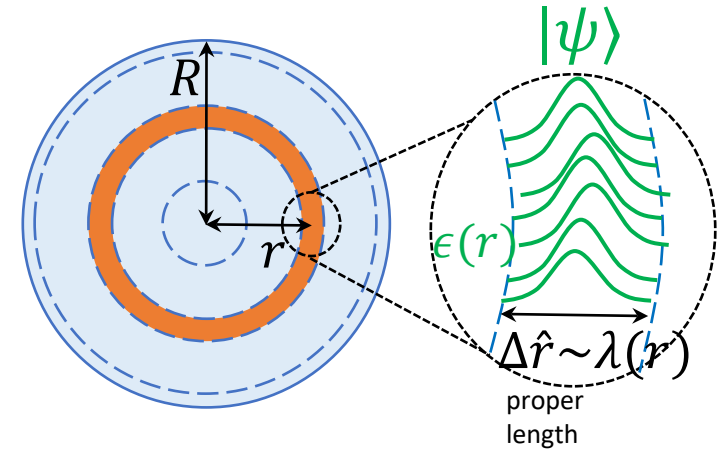
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$$s(r) = \frac{N(r)}{\Delta\hat{r}} \quad \text{for } \Delta\hat{r} \sim \lambda(r)$$



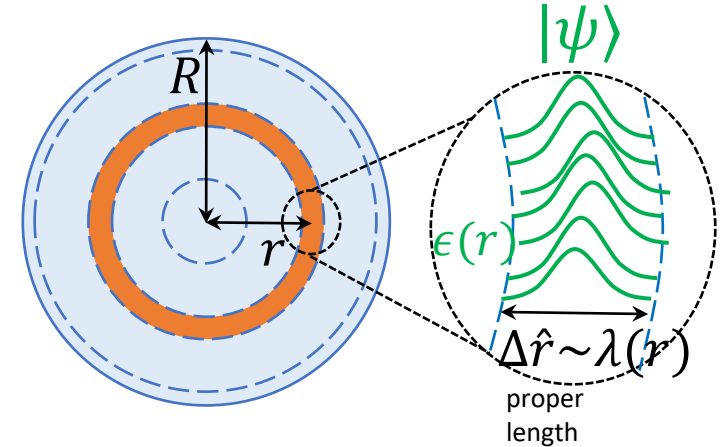
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$$S[g_{\mu\nu}] \sim \frac{1}{\hbar} \int_0^R dr N(r) \epsilon(r) \left[1 - \frac{2}{r} \int_0^r dr' N(r') \frac{\epsilon(r')^2}{m_p^2} \right]^{-\frac{1}{2}}$$

Hamiltonian constraint $\mathcal{H} = 0$:

$$\partial_r a(r) = 8\pi G r^2 \langle -T_t^t(r) \rangle$$

self-gravity

\Rightarrow For a given $N(r)$, ϵ_{max} provides S_{max} .

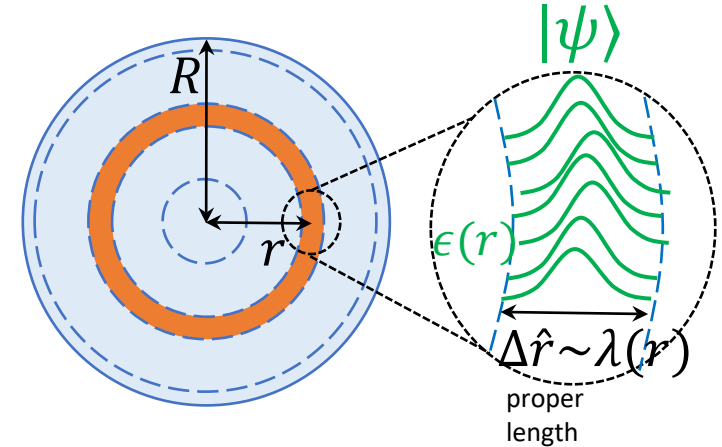
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For the semi-classical description,

$$\epsilon(r) \leq \epsilon_{max} \sim \frac{m_p}{\sqrt{n}} \quad (n: \text{O}(1) \text{ constant } \gg 1)$$

Upper bound

- A static spacetime has a timelike Killing vector globally.
⇒ No trapped surface exists. [\[Mars-Senovilla 2003\]](#)

Ex. $k = \partial_t$ in Schwarzschild metric

$$k^2 = -\left(1 - \frac{a_0}{r}\right)$$

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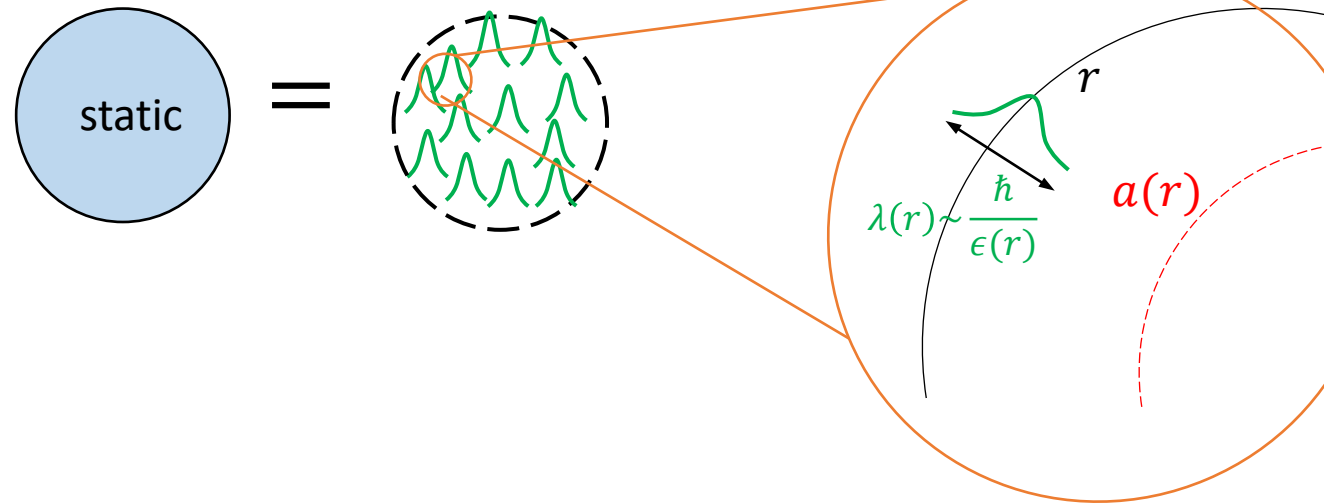
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$$\lambda(r) \lesssim \sqrt{g_{rr}(r)}(r - a(r))$$

$a(r)$: Schwarzschild radius for r



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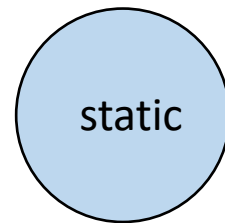
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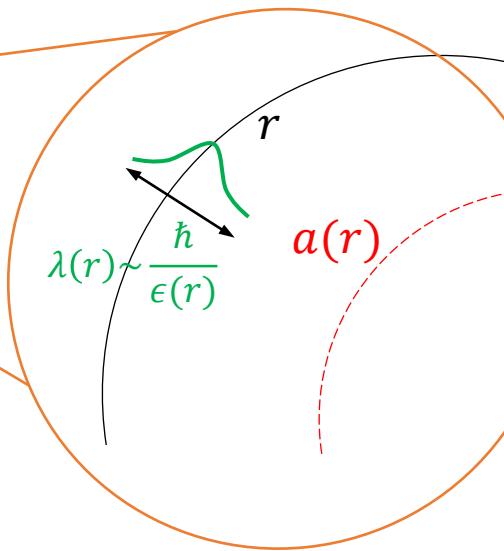
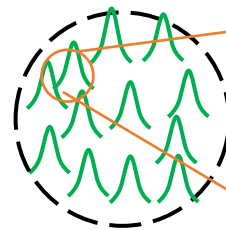
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- Then, we can get the upper bound:

$$S \lesssim \frac{1}{l_p^2} \int_0^R dr r \frac{\partial_r a(r)}{=} 8\pi G r^2 \langle -T_t^t(r) \rangle$$

Entropy-maximized spacetime (1/2)

- To get the saturating configuration $g_{\mu\nu}^*$, we solve

$$\lambda(r) \sim \sqrt{g_{rr}(r)}(r - a(r)) \text{ for } \epsilon(r) = \epsilon_{max} \sim \frac{m_p}{\sqrt{n}}$$

and use two consistencies:

- thermodynamics,
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- Thus, we reach **uniquely** the physical saturating configuration:

$$ds^2 = -\frac{\sigma\eta^2}{2r^2} e^{-\frac{R^2-r^2}{2\sigma\eta}} dt^2 + \frac{r^2}{2\sigma} dr^2 + r^2 d\Omega^2. \quad \boxed{\begin{array}{l} \sigma = O(nl_p^2), \\ 1 \leq \eta < 2 \end{array}}$$

- Can be obtained in various ways and robust.

→ $\lambda(r) \lesssim \sqrt{g_{rr}(r)}(r - a(r))$ can be verified in a dynamical model.

- Non-perturbative solution of $G_{\mu\nu} = 8\pi G \langle \psi | T_{\mu\nu} | \psi \rangle$ for \hbar [Kawai-Yokokura 2020]

- $n = \#$ of d.o.f. in the theory

[Kawai-Matsuo-Yokokura 2013,
Kawai-Yokokura 2014, 2015, 2016, 2020, 2021,
Yokokura 2022, Ho-Kawai-Liyao-Yokokura 2023]

Entropy-maximized spacetime (2/2)

- The entropy-maximized spacetime

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represents a dense configuration without horizon or singularity.

$$\langle T_\theta^\theta \rangle = O(1)$$

- Stabilize the system
- Not fluid

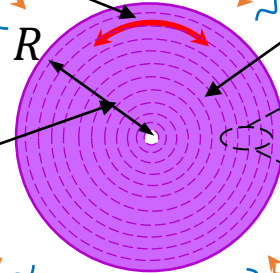
Surface from thermodynamics

$$R = a_0 + \frac{\sigma\eta^2}{2a_0} (> a_0)$$

$(a_0 \equiv 2GM_0)$

strong pressure

$$\langle T_\theta^\theta \rangle_*$$



heat bath of T_H

- $\mathcal{R} = O\left(\frac{1}{nl_p^2}\right) \ll O\left(\frac{1}{l_p^2}\right)$
 - almost flat around $r = 0$
- \Rightarrow No singularity

- Maximum entropy

$$S_{max} =$$

Entropy-maximized spacetime (2/2)

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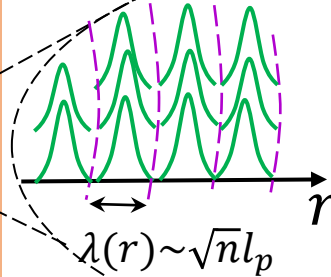
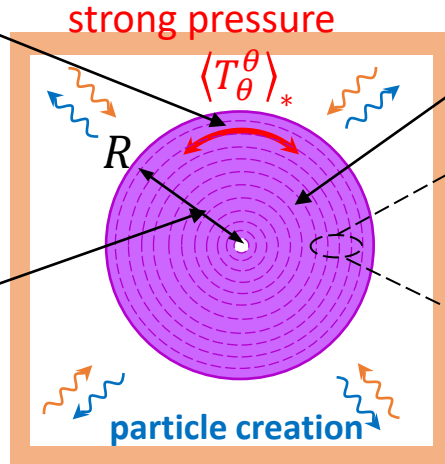
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uniform in r direction
 $\Rightarrow N(r) = \text{const.} \sim n$
 \Rightarrow entropy density

$$s(r) = \frac{2\pi\sqrt{2\sigma}}{l_p^2}$$

- Maximum entropy

$$S_{max} = \int_0^R dr \sqrt{g_{rr}} s(r) = \int_0^R dr \sqrt{\frac{r^2}{2\sigma}} \frac{2\pi\sqrt{2\sigma}}{l_p^2} = \frac{A}{4l_p^2}$$

[Yokokura 2022]

$$A \equiv 4\pi R^2 = 4\pi a_0^2 + O(1)$$

Verification of Bousso bound

- We have obtained

$$S \leq S_{max} = \frac{A}{4l_p^2}$$

Verification of Bousso bound

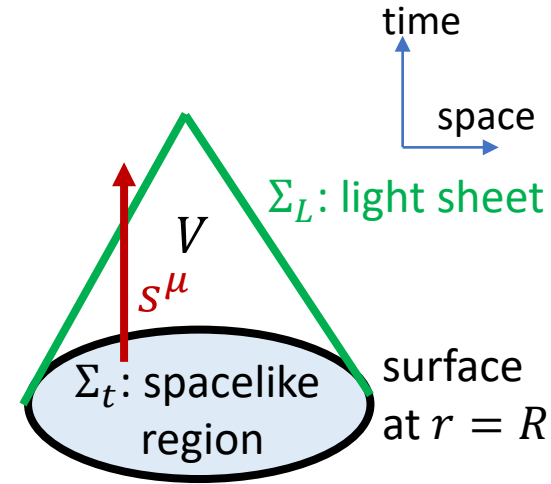
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- This is the **Bousso bound**: [Bousso 1999]

$$S = \int_{\Sigma_t} d\Sigma_\mu S^\mu \stackrel{\text{ours}}{=} \int_{\Sigma_L} d\Sigma_\mu S^\mu \stackrel{\text{Bousso's}}{=} \int_{\Sigma_L} d\Sigma_\mu S^\mu$$

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Verification of Bousso bound

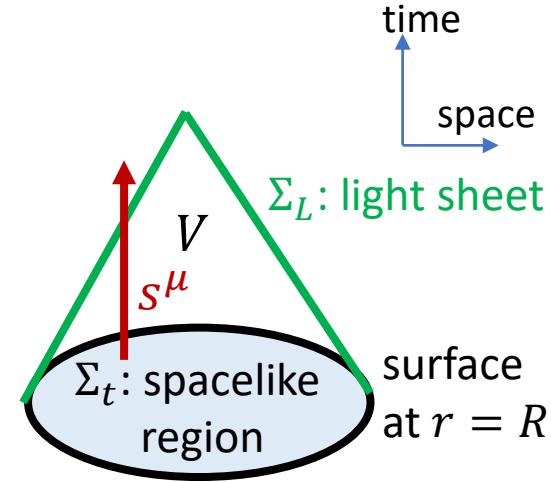
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- $g_{\mu\nu}^*$ saturates **local** sufficient conditions for the Bousso bound:

$$\left\{ \begin{array}{l} | -s_\mu k^\mu | \leq \frac{1}{\hbar} \langle T_{\mu\nu} \rangle k^\mu k^\nu \Delta\zeta, \text{ [Flanagan-Marolf-Wald 2000],} \\ | k^\mu k^\nu \nabla_\mu s_\nu | \leq \frac{2\pi}{\hbar} \langle T_{\mu\nu} \rangle k^\mu k^\nu \text{ [Bousso-Flanagan-Marolf 2003]} \end{array} \right.$$

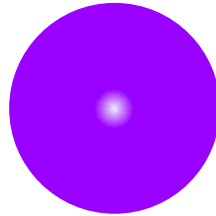
- Thus, we have verified the Bousso bound by constructing explicitly the saturating configuration.

Conclusions

- Considered 4D spherical static spacetime for highly excited states $\{|\psi\rangle\}$.
- Estimated the entropy $S[g_{\mu\nu}]$ including the self-gravity.
- Found the entropy-maximized spacetime uniquely:

$$ds^2 = -\frac{\sigma\eta^2}{2r^2} e^{-\frac{R^2-r^2}{2\sigma\eta}} dt^2 + \frac{r^2}{2\sigma} dr^2 + r^2 d\Omega^2.$$

Non-perturbative solution of
 $G_{\mu\nu} = 8\pi G \langle \psi | T_{\mu\nu} | \psi \rangle$ for \hbar



Dense configuration
 (without horizon or singularity)

- Verified the Bousso bound:

$$S \leq S_{max} = \frac{A}{4l_p^2}, \quad \leftarrow \text{the result of the self-gravity}$$

where the information is stored inside.

- Q1: What is the maximum entropy S_{max} that can be given a finite region?
 $\Rightarrow S_{max} = \frac{A}{4l_p^2}$ (in this class)
- Q2: What is the structure of such a spacetime?
 \Rightarrow Necessarily, the above metric (in this class)

Future directions

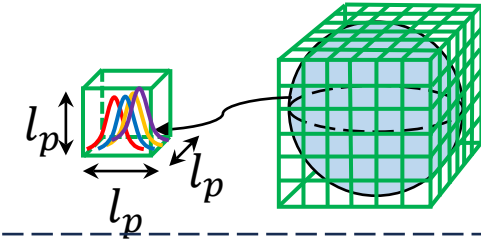
- Origin of Holography?

$$S_{max} = 4\pi \int_0^R dr r^2 \sqrt{g_{rr}} \bar{s}(r) = \frac{A}{4l_p^2} \text{ from } \bar{s}(r) \sim \frac{\sqrt{n}}{l_p r^2} \ll \bar{s}_{naive}(x) \sim \frac{n}{l_p^3}$$

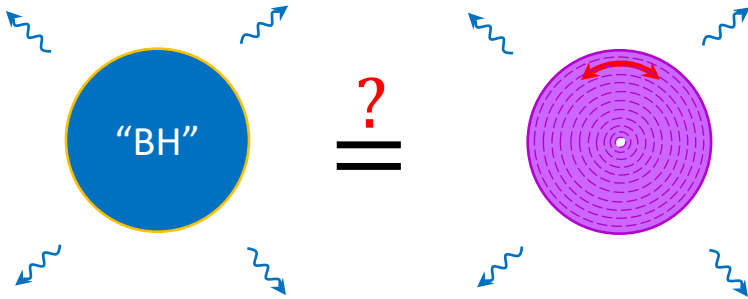
⇒ Self-gravity suppresses excitations of local d.o.f. ? [work in progress]

A naïve estimation of entropy density

Full excitations of n local d.o.f.



- Quantum BH = the dense configuration?



Both have

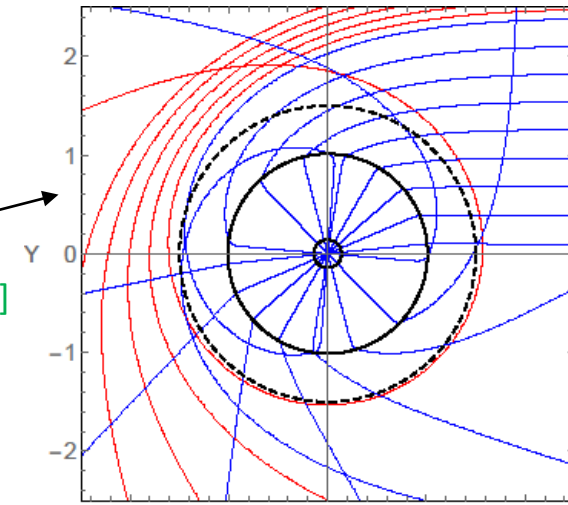
$$T = \frac{\hbar}{4\pi a_0}$$

$$S = \frac{A}{4l_p^2}$$

$$R \approx a_0$$

- Phenomenology as a BH mimicker?

- BH shadow image [work in progress with C.Y. Chen (iTHEMS)]
- Gravitational Waves [work in progress with N. Oshita (Hakubi-YITP)]



Thank you!