

Correlation functions from homotopy algebras and the applications

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based on arXiv 2305.11634 **K.K.**, Y. Okawa

arXiv 2305.13103 **K.K.**

in preparation **K.K.**, Y. Okawa

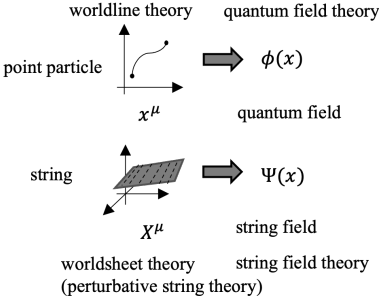
Motivations

We want to know the non-perturbative definition of string theory!

Candidates

- ▶ matrix models
- ▶ string field theory (SFT)

etc....



String Field Theory and Homotopy algebra

Homotopy algebras are effective tools to construct SFT.

- ▶ (quantum) A_∞ algebra \rightarrow open SFT
- ▶ (quantum) L_∞ algebra \rightarrow closed SFT
- ▶ open-closed homotopy algebra \rightarrow open-closed SFT

etc....

Advantage

- ▶ Systematic
- ▶ Algebraic: We do not need path integral.
- ▶ Universal

etc....

Construct the formulas for QFT and extend them to SFT!!

Homotopy Algebra and Field Theory

There are many things we can do using homotopy algebras!

- ▶ relating covariant and light-cone string field theories
[Erler and Matsunaga, arXiv:2012.09521]
- ▶ **calculating scattering amplitudes**
[Kajiura, math/0306332] etc.
- ▶ **integrating out fields**
[Sen, arXiv:1609.00459],
[Erbin, Maccaferri, Schnabl and Vošmera, arXiv:2006.16270],
[Koyama, Okawa and Suzuki, arXiv:2006.16710] etc.
- ▶ **correlation functions**
scalar field [Okawa, arXiv:2203.05366],
free scalar and Dirac field [K.K. and Okawa, arXiv:2305.11634],
general scalar and Dirac field [K.K., arXiv:2305.13103] etc.

etc....

String Field Theory and Homotopy algebra

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Homotopy algebra and the universal property

Formula for on-shell amplitudes

$$\pi_1 \mathbf{P} \mathbf{Q} \mathbf{P} + \pi_1 \mathbf{P} \mathbf{m} \frac{1}{\mathbf{I} + \hbar \mathbf{m} + i\hbar \mathbf{h} \mathbf{U}} \mathbf{P},$$

\mathbf{I} : identity on $T\mathcal{H}$

\mathbf{P} : on-shell projection

$\mathbf{h}, \mathbf{m}, \mathbf{U}$: determined by the theory

Homotopy algebra

Vector space

$$\mathcal{H} = \bigoplus_{i \in \mathbb{Z}} \mathcal{H}_i,$$

$$T\mathcal{H} = \bigoplus_{i=0}^{\infty} \mathcal{H}^{\otimes n},$$

Action

$$S = -\frac{1}{2}\omega(\Phi, \pi_1 \mathbf{Q}\Phi) - \sum_{n=0}^{\infty} \frac{1}{n+1} \omega(\Phi, \pi_1 \mathbf{m}_n(\Phi \otimes \dots \otimes \Phi)),$$

where π_n is the projection from $T\mathcal{H}$ onto $\mathcal{H}^{\otimes n}$.

Gauge transformation

$$\delta_\lambda \Phi = \pi_1 \mathbf{Q}\lambda + \sum_{k=1}^{\infty} \sum_{l=0}^{k-1} \pi_1 \mathbf{m}_k(\underbrace{\Phi \otimes \dots \otimes \Phi}_l \otimes \lambda \otimes \underbrace{\Phi \otimes \dots \otimes \Phi}_{k-l-1}),$$

Gauge invariance (quantum A_∞ relations)

$$\left(\mathbf{Q} + \sum_{n=0}^{\infty} \mathbf{m}_n + i\hbar \mathbf{U} \right)^2 = 0$$

Homotopy algebra and Correlation Function

In the theories without gauge symmetry, it is sufficient to consider two spaces:

$$\mathcal{H} = \mathcal{H}_1 \oplus \mathcal{H}_2 .$$

We take Φ to be

$$\Phi = \int d^d x \varphi(x) c(x) + \int d^d x (\bar{\theta}_\alpha(x) \Psi_\alpha(x) + \bar{\Psi}_\alpha(x) \theta_\alpha(x)) ,$$

where $\varphi(x)$ is the scalar field and $\Psi(x)$ is the Dirac field.
Then, correlation functions are given by

$$\langle \Phi^{\otimes n} \rangle = \pi_n \mathbf{f} \mathbf{1} ,$$

$$\begin{aligned} \mathbf{f} &= \frac{1}{\mathbf{I} + \mathbf{h} \mathbf{m} + i\hbar \mathbf{h} \mathbf{U}} \\ &= \mathbf{I} + \sum_{n=1}^{\infty} (-1)^n (\mathbf{h} \mathbf{m} + i\hbar \mathbf{h} \mathbf{U})^n . \end{aligned}$$

Homotopy algebra and Correlation Function

For simplicity, we consider only the scalar part.

Then,

$$\begin{aligned}\langle \Phi^{\otimes n} \rangle &= \langle \underbrace{\Phi \otimes \Phi \otimes \dots \otimes \Phi}_n \rangle \\ &= \int d^d x_1 d^d x_2 \dots d^d x_n \langle \varphi(x_1) \varphi(x_2) \dots \varphi(x_n) \rangle \\ &\quad c(x_1) \otimes c(x_2) \otimes \dots \otimes c(x_n),\end{aligned}$$

and correlation functions are obtained as the coefficients of the bases of $\mathcal{H}^{\otimes n}$.

The information of Schwinger-Dyson equations are contained by

$$f^{-1} f \mathbf{1} = \mathbf{1}.$$

This fact confirms that our formula represents the correlation functions in the ordinary quantum field theory.

Nonperturbative extension

In zero-dimensional case, we can represent the homotopy operators by infinite dimensional matrices.

In the quartic Euclidean theory, it is given by

$$\mathbf{hU} = \frac{1}{m^2} \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 & \dots \\ 0 & 0 & 0 & 0 & 0 & 0 & \dots \\ 1 & 0 & 0 & 0 & 0 & 0 & \dots \\ 0 & 2 & 0 & 0 & 0 & 0 & \dots \\ 0 & 0 & 3 & 0 & 0 & 0 & \dots \\ 0 & 0 & 0 & 4 & 0 & 0 & \dots \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \ddots \end{pmatrix},$$
$$\mathbf{hm}_3 = \frac{\lambda}{m^2} \begin{pmatrix} 0 & 0 & 0 & 0 & 0 & 0 & \dots \\ 0 & 0 & 0 & 1 & 0 & 0 & \dots \\ 0 & 0 & 0 & 0 & 1 & 0 & \dots \\ 0 & 0 & 0 & 0 & 0 & 1 & \dots \\ 0 & 0 & 0 & 0 & 0 & 0 & \dots \\ 0 & 0 & 0 & 0 & 0 & 0 & \dots \\ \vdots & \vdots & \vdots & \vdots & \vdots & \vdots & \ddots \end{pmatrix}.$$

Nonperturbative extension

Recall that the key operator \mathbf{f} is defined perturbatively:

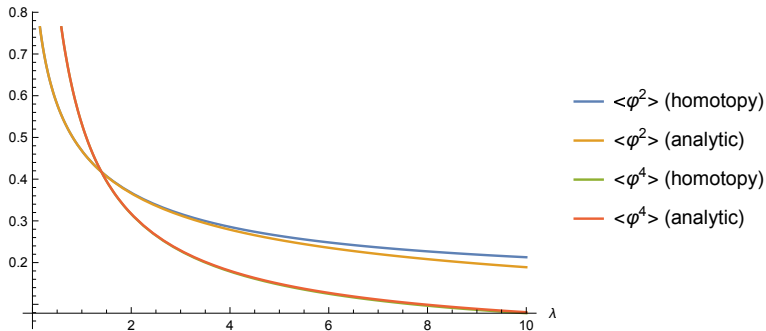
$$\mathbf{f} = \frac{1}{\mathbf{I} + \mathbf{h} \mathbf{m} + i\hbar \mathbf{h} \mathbf{U}} = \mathbf{I} + \sum_{n=1}^{\infty} (-1)^n (\mathbf{h} \mathbf{m} + i\hbar \mathbf{h} \mathbf{U})^n .$$

In this case, we can calculate exact inverse of $\mathbf{I} + \mathbf{h} \mathbf{m} + i\hbar \mathbf{h} \mathbf{U}$ by truncating infinite matrices to N by N matrices.

Nonperturbative extension

the quartic Euclidean theory

$$S = \frac{1}{2}\varphi^2 + \frac{1}{4}\lambda\varphi^4$$



Nonperturbative extension

Even when $\lambda = 0.2$, the nonperturbative effects dominate.

$$\langle \varphi^2 \rangle = 0.724059$$

and the estimated value from our formula for $N = 100$

$$f_{20} = 0.724059.$$

the number of loops	perturbative values for $\langle \varphi^2 \rangle$
1	0.4
2	1.36
3	-1.016
4	6.8176
5	-25.2714
6	131.548

Summary and Outlook

We proposed the formula for correlation functions involving the scalar fields and the Dirac field in homotopy algebras:

$$\langle \Phi^{\otimes n} \rangle = \pi_n \mathbf{f} \mathbf{1}.$$

Moreover, nonperturbative extension of the formula was discussed.

The future works are as follows:

- ▶ general description for nonperturbative correlation functions in homotopy algebras
- ▶ symmetries: K.K. and J.Yoshinaka *in preparation*
- ▶ extensions of perturbative formula to general gauge theory e.g.) Yang-Mills theory: *working in progress* with J.Yoshinaka
- ▶ extensions to string field theory