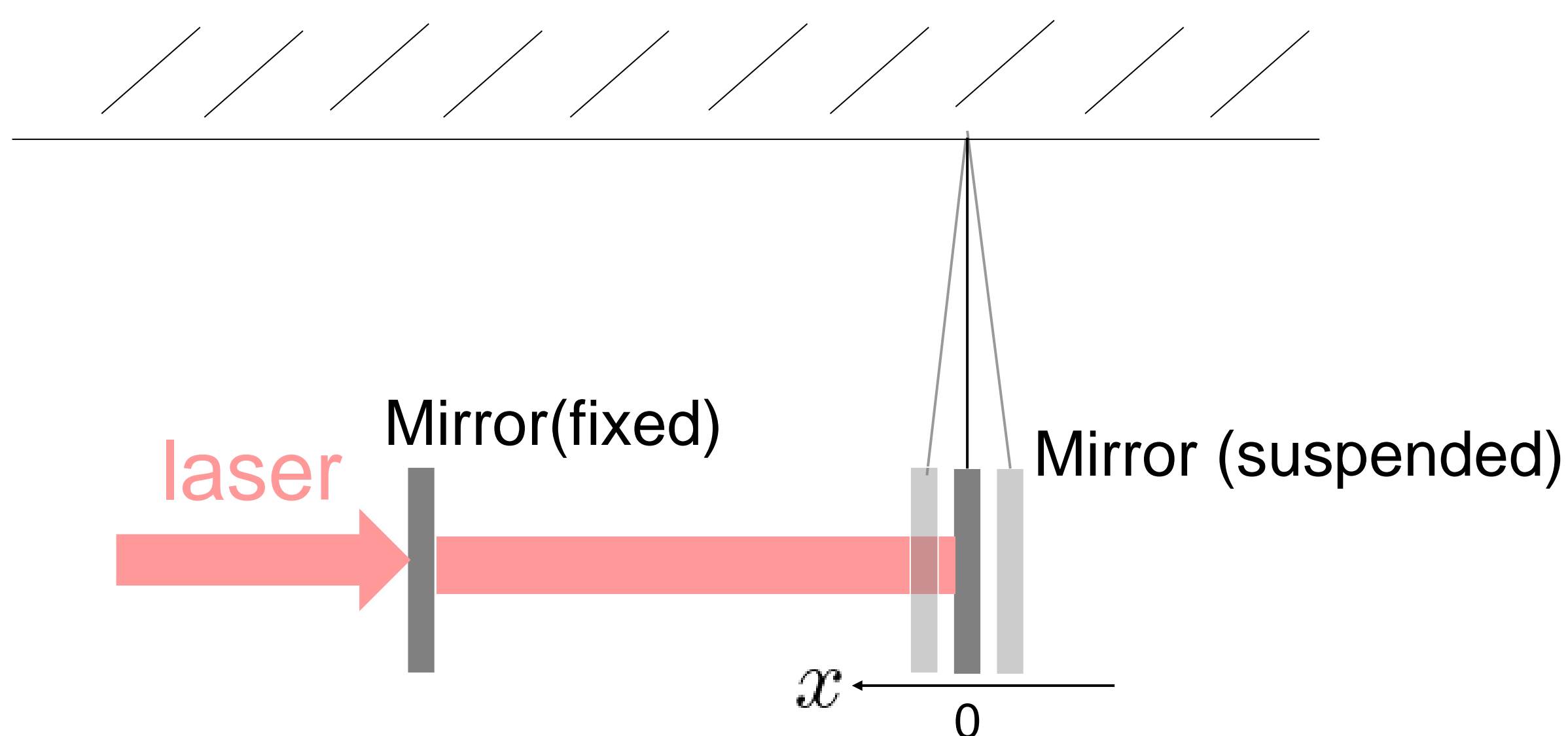


Possibility of macroscopic spatial superposition via double-well optomechanical potential

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1. Optomechanics by moving mirror



- Optical cavity with a suspended mirror → **Optomechanical system**

Hamiltonian in rotating frame:

$$\frac{H}{\hbar} = \delta \hat{a}^\dagger \hat{a} + \frac{\Omega}{2} (\hat{p}^2 + \hat{x}^2) + g \hat{x} \hat{a}^\dagger \hat{a} + iF(\hat{a}^\dagger - \hat{a})$$

Cavity mode
Suspended mirror mode
Optomech. coupling
Laser driving

- Linearized dynamics... well studied

$$g \hat{x} \hat{a}^\dagger \hat{a} = g \hat{x} (\alpha^* + \hat{a}^\dagger) (\alpha + \hat{a}) \quad \hat{x} = \frac{\hat{b} + \hat{b}^\dagger}{\sqrt{2}}$$

$$\approx \begin{cases} g(\alpha \hat{a}^\dagger \hat{b} + \alpha^* \hat{a} \hat{b}^\dagger) & \text{For red-detuned: } \delta = \Omega \\ g(\alpha \hat{a}^\dagger \hat{b}^\dagger + \alpha^* \hat{a} \hat{b}) & \text{For blue-detuned: } \delta = -\Omega \end{cases}$$

Linearized optomech. is used e.g. for sideband cooling of massive mirror, quantum transducer with optomech. crystal, creation of entangled photon-phonon pair, (anti)optical spring etc.

- Can we make use of **nonlinear dynamics**, in particular to realize a **quantum superposition of a massive mirror**? (typically, 1mg or heavier!)

2. Optomechanical potential force and friction for mirror

- Stationary solution for the Heisenberg eqs. with relaxations:

$$\langle \hat{a} \rangle = \frac{F}{i(\delta + g x_s) + \kappa/2} \equiv \sqrt{n} e^{i\phi} \quad \langle \hat{x} \rangle = -\frac{gn}{\Omega^2} \equiv x_s$$

- By solving the eq. for the fluctuation of the cavity mode, and by substituting it to the eq. for the mirror mode, we can read off the **effective potential force** and **friction** affecting the mirror.

$$\frac{d}{dt} \hat{p} = -\Omega \hat{x} - \frac{\Gamma}{2} \hat{p} - gn \left(\frac{(\delta + g x_s)^2 + \kappa^2/4}{1 + 2g(\delta + g x_s) \hat{x} + g^2 \hat{x}^2} - 1 \right) - gn \Omega \operatorname{Im} \left(\frac{1}{(\delta + g x_s - i\kappa/2 + g \hat{x})^2} \right) \hat{p} + (\text{higher order terms in } \Omega/|\delta - i\kappa/2|)$$

Although it is an adiabatic expansion, each term is of all-order in g !

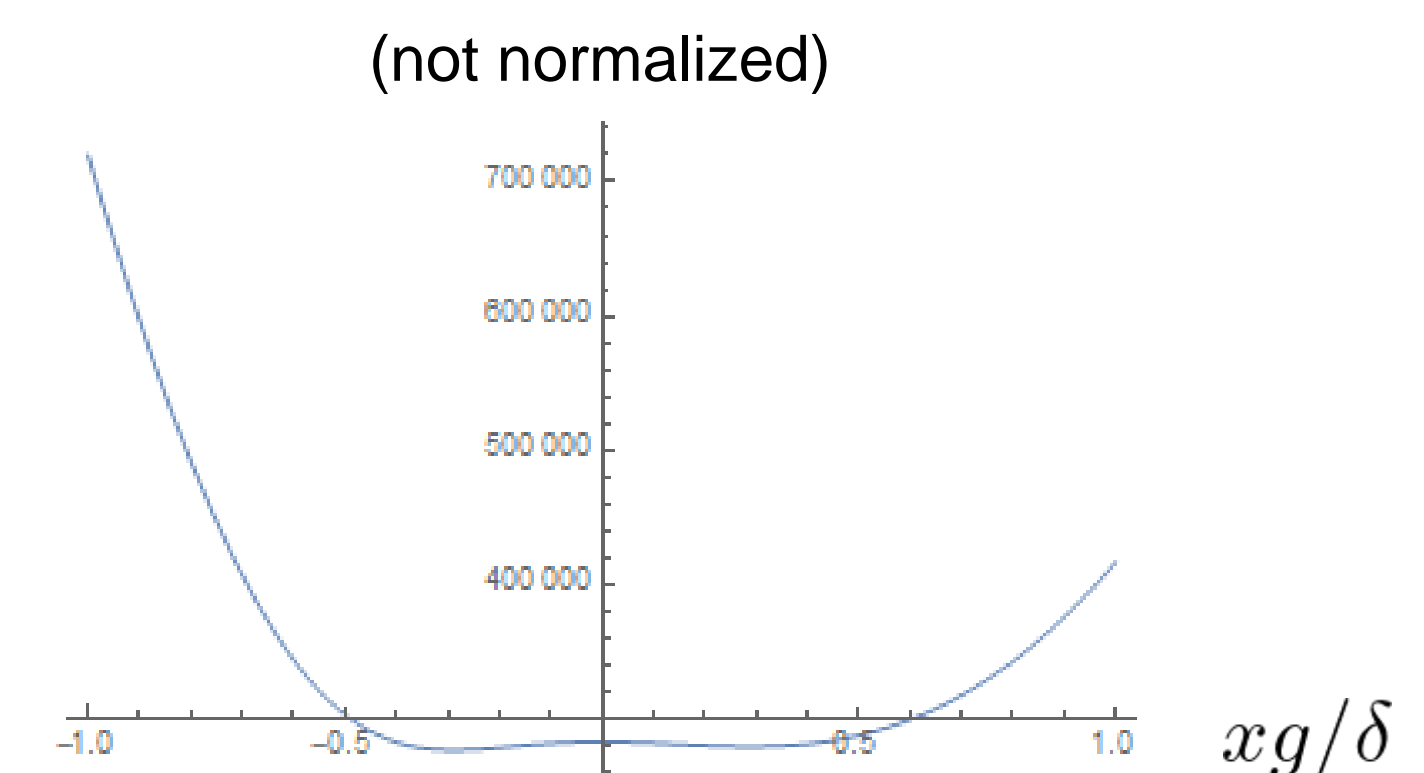
Nonlinear extensions of **(anti)optical spring** & **(anti)friction**.

3. Effective potential

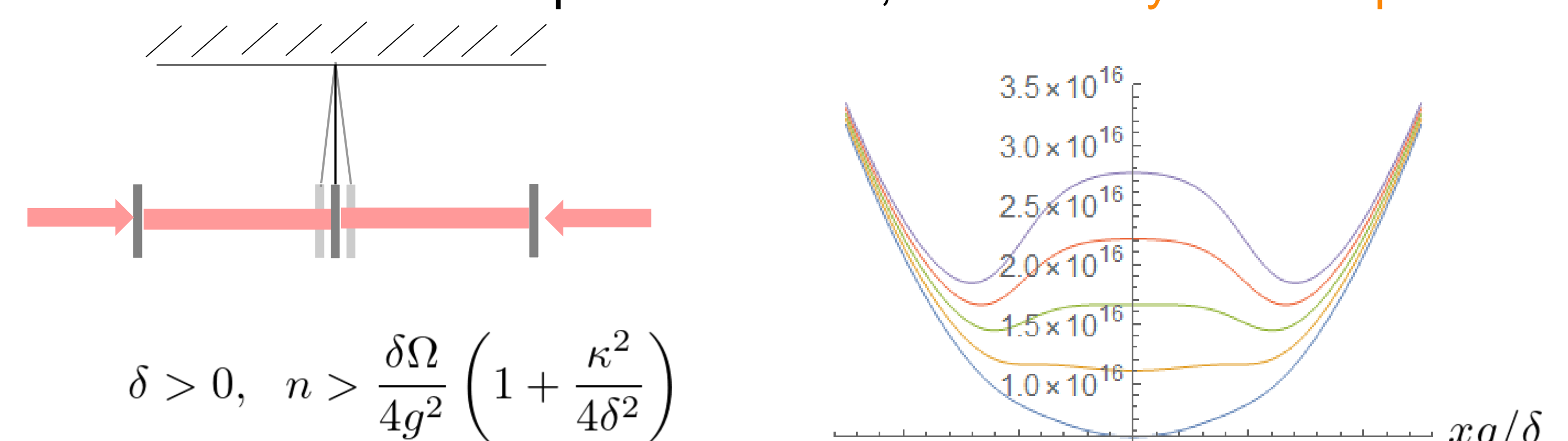
- The origin is always stable/unstable since x is the fluctuation from stationary pt. $x_s \rightarrow$ Easy to evaluate an existence of multiple minimal pt.s.
- The following conditions for laser detuning, loss rate and cavity photon number realize a double-well like potential

$$\delta > 0 \quad \kappa < \frac{2}{\sqrt{3}} \delta,$$

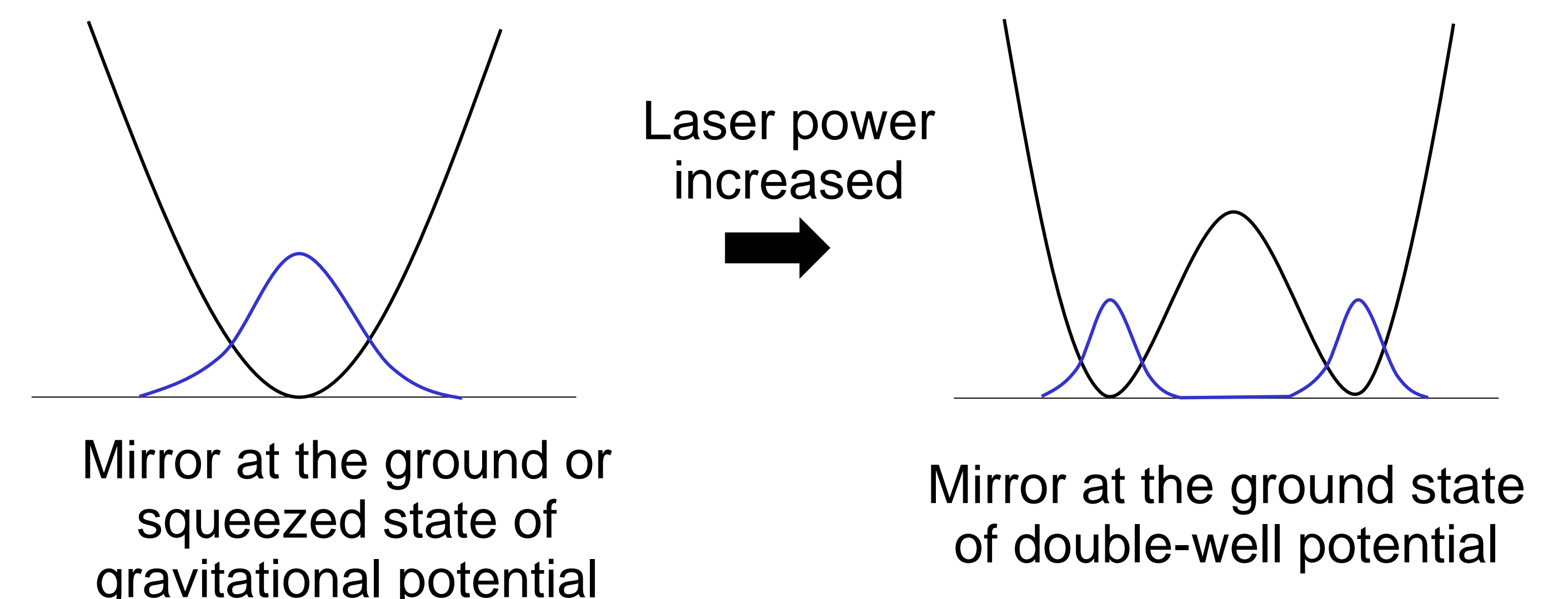
$$\frac{\delta \Omega}{3g^2} \left(2 - \sqrt{1 - \frac{3^2}{4\delta}} \right) < n < \frac{\delta \Omega}{3g^2} \left(2 - \sqrt{1 + \frac{3^2}{4\delta}} \right)$$



- If we input **double laser beams** with the same detuning and power on both sides of the suspended mirror, we have **symmetric potential**



- Consider the following adiabatic tuning:



Spatial superposition!

Such a superposition of a massive object is expected to be useful in detection of the dark matter or quantum graviton by decoherence

4. Parameters for experiments

Suspended mirror freq.:	$\Omega/2\pi = O(1)\text{Hz}$ (1mg mirror, 10cm suspension)
Optomech. Coupling:	$g/2\pi = O(1)\text{Hz}$ (normalized for quantum fluctuation length of mirror)
Laser detuning:	$\delta/2\pi \gtrsim O(10)\text{kHz}$
Laser loss rate:	$\kappa/2\pi \gtrsim O(1)\text{MHz}$

For the symmetric potential case, the requirement on the photon num. for the double-well is $n \gtrsim O(10^8)$.
→ O(1) μW power is sufficient(!?) for O(100THz) laser.

5. Outlook

- Feasibility of the adiabatic change in the potential and the wavefunction (in a case $\delta \simeq 10g$, we confirmed it is **easily possible**)
- Detection scheme of the dark matter and graviton by decoherence of the mirror (cf. poster by Jinyang Li)