

HADRONIC τ DATA AND LATTICE QCD+QED SIMULATIONS FOR THE MUON ($g - 2$)

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work in collab. with T. Izubuchi, C. Lehner, A. Meyer, X. Tuo
for the RBC/UKQCD collaborations



Seventh plenary workshop of the muon $g-2$ theory initiative
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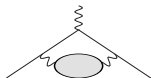
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$(g - 2)_\mu$
Lattice



Hadronic Vacuum Polarization (HVP) contribution to a_μ

Lattice

Time-momentum representation

[Bernecker, Meyer, '11]

$$G^\gamma(t) = \frac{1}{3} \sum_k \int d\vec{x} \langle j_k^\gamma(x) j_k^\gamma(0) \rangle \quad \rightarrow \quad a_\mu = 4\alpha^2 \sum_t w_t G^\gamma(t)$$

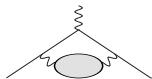
Windows in Euclidean time

[RBC/UKQCD '18]

$$a_\mu^W = 4\alpha^2 \sum_t w_t G^\gamma(t) [\Theta(t, t_0, \Delta) - \Theta(t, t_1, \Delta)]$$

$t_0 = 0.4 \text{ fm} \quad t_1 = 1.0 \text{ fm} \quad \Delta = 0.15 \text{ fm}$

$(g - 2)_\mu$
Dispersive



Hadronic Vacuum Polarization (HVP) contribution to a_μ

Dispersive

$$a_\mu = \frac{\alpha}{\pi} \int \frac{ds}{s} K(s, m_\mu) \frac{\text{Im}\Pi(s)}{\pi}$$

[Brodsky, de Rafael '68]

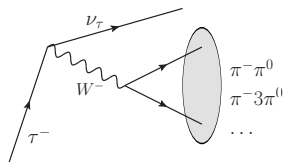
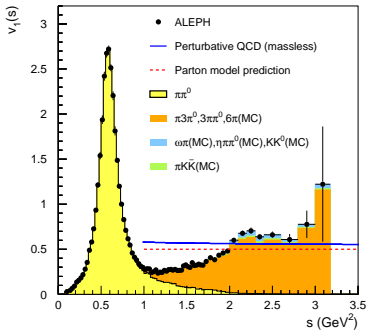
$$\text{Im} \left[\text{Diagram} \right] = \sum_X \left| \left[\text{Diagram} \right] \right|^2 \quad \frac{4\pi^2 \alpha}{s} \frac{\text{Im}\Pi(s)}{\pi} = \sigma_{e^+e^- \rightarrow \gamma^* \rightarrow \text{had}}$$

The diagram on the left is a photon loop with a shaded oval. The diagram on the right is a photon line with a shaded oval and three external lines labeled X.

Windows can be estimated dispersively as well and compared

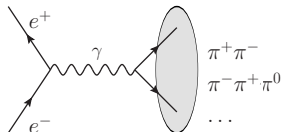
MOTIVATIONS

τ decays



$V - A$ current

Final states $I = 1$ charged



EM current

Final states $I = 0, 1$ neutral

τ data can improve $a_\mu[\pi\pi]$

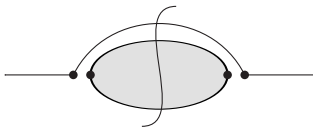
→ 72% of total Hadronic LO

→ competitive precision on a_μ^W

HADRONIC τ DECAYS

Fermi theory

$$\mathcal{M}_f(P, q, p_1 \cdots p_{n_f}) = \frac{G_F V_{ud}}{\sqrt{2}} \bar{u}_\nu(-q) \gamma_\mu^L u_\tau(P) \langle \text{out}, p_1 \cdots p_{n_f} | \mathcal{J}_\mu^-(0) | 0 \rangle$$



$$\begin{aligned} d\Gamma &= \frac{1}{4m} d\Phi_q \sum_f d\Phi_f \sum_{\text{spin}} |\mathcal{M}_f|^2 \\ &= \frac{1}{4m} d\Phi_q \frac{G_F^2 |V_{ud}|^2}{2} \mathcal{L}_{\mu\nu}(P, q) \rho_{\mu\nu}^w(p) \end{aligned}$$

Charged spectral density isospin limit = $\rho^{w,0}$ $\left[d\Phi_q = \frac{d^3q}{(2\pi)^3 2\omega_q} \right]$

$$\begin{aligned} \frac{d\Gamma(s)}{ds} &= G_F^2 |V_{ud}|^2 \frac{m^3}{16\pi^2} \left(1 + \frac{2s}{m^2}\right) \left(1 - \frac{s}{m^2}\right)^2 \rho^{w,0}(s) \\ &= G_F^2 |V_{ud}|^2 \frac{m^3}{16\pi^2} \kappa(s) \rho^{w,0}(s) \end{aligned}$$

ELECTRONIC RATE

Eliminating G_F

from experiment we get $\frac{1}{\Gamma} \frac{d\Gamma}{ds} \rightarrow \frac{\Gamma}{\Gamma_e} \frac{1}{\Gamma} \frac{d\Gamma}{ds} = \frac{1}{\Gamma_e} \frac{d\Gamma}{ds}$

$$\Gamma_e = \Gamma(\tau \rightarrow e\bar{\nu}\nu) = \frac{\mathcal{B}_e \Gamma}{\mathcal{B}} = \frac{G_F^2 m_\tau^5}{192\pi^3}$$

$$\text{conventionally } \rho^{w,0}(s) = \frac{m_\tau^2}{12\pi^2 |V_{ud}|^2 \kappa(s)} \times \frac{\mathcal{B}}{\mathcal{B}_e} \times \frac{1}{\Gamma} \frac{d\Gamma}{ds}$$

$O(\alpha)$ correction fo Γ_e finite in Fermi theory

[Kinoshita, Sirlin '59]

$$\Gamma_e = \frac{G_F^2 m_\tau^5}{192\pi^3} \left[1 + \frac{\alpha}{2\pi} \left(\frac{25}{4} - \pi^2 \right) \right] \left[1 + O(m_W^2/m_\tau^2) + O(m_e^2/m_\tau^2) \right]$$

→ 0.4% correction

W REGULARIZATION

Short-distance effects

[Sirlin '82][Marciano, Sirlin '88][Braaten, Li '90]

Effective Hamiltonian $H_W \propto G_F O_{\mu\nu}$

G_F low-energy constant; 4-fermion operator $O_{\mu\nu}$

At $O(\alpha)$ new divergences in EFT \rightarrow need regulator, Z factors



$$\frac{1}{k^2} = \frac{1}{k^2 - m_W^2} - \frac{m_W^2}{k^2(k^2 - m_W^2)}$$

[Sirlin '78]

1. universal UV divergences re-absorbed in G_F
2. process-specific corrections in S_{EW} , like a Z factor

Effective Hamiltonian at $O(\alpha)$: $H_W \propto G_F S_{EW}^{1/2} O_{\mu\nu}$

matching required as noted by [Carrasco et al '15][Di Carlo et al '19]

ISOSPIN BREAKING

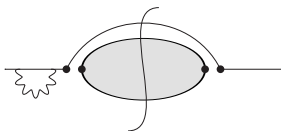
Initial state

Wave-function renormalization

$$Z_\tau = 1 + \frac{\alpha}{2\pi} \left[\log \frac{m_\tau}{\mu} + 2 \log \frac{m_\gamma}{m_\tau} + \dots \right]$$

$$\frac{d\Gamma}{ds} \simeq 2 \times \frac{1}{2} [Z_\tau - 1] |\mathcal{M}|^2$$

$$\delta Z_\tau \equiv \frac{\alpha}{2\pi} \log(m_W/m_\tau) \quad [\text{Sirlin '82}]$$



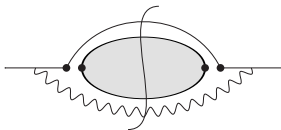
τ Bremsstrahlung

[Cirigliano et al '00, '01][MB et al, in prep]

$$\frac{d\Gamma}{ds} \frac{\alpha}{\pi} [G_{\log}(s, m_\gamma) + \dots]$$

$$G_{\log}(s, m_\gamma) = \log \frac{m_\gamma}{m_\tau} + \dots$$

$$\delta\kappa(s) \equiv G_{\log}(s, m_\tau) + \dots$$



$$\frac{d\Gamma}{ds} \simeq G_F^2 |V_{ud}|^2 \frac{m^3}{16\pi^2} \kappa(s) \rho^{w,0}(s) [\delta Z_\tau + \delta\kappa(s)]$$

ISOSPIN BREAKING

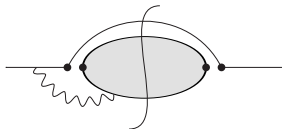
Initial-final state

Virtual photon loop

$$\delta Z_{\kappa\rho} \propto \frac{\alpha}{\pi} \log(m_W/m_\tau) \quad [\text{Sirlin '82}]$$

[Cirigliano et al '01]

Finite parts EFT and 2π



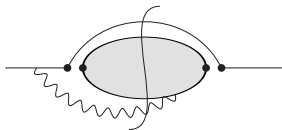
$\tau - \pi$ bremsstrahlung interference

From EFT and 2π [Cirigliano et al' 00, '01]

Structure-independent captured by EFT

Structure-dependent meson dominance

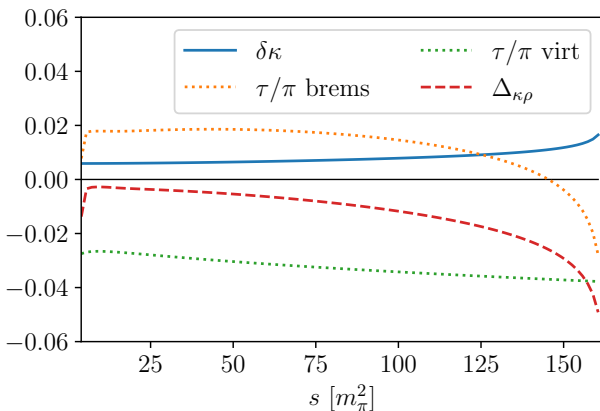
[Flores-Talpa et al. '06, '07]



$$\frac{d\Gamma}{ds} \text{ += } G_F^2 |V_{ud}|^2 \frac{m^3}{16\pi^2} \kappa(s) \rho^{w,0}(s) [\delta Z_{\kappa\rho} + \Delta_{\kappa\rho}(s)]$$

LONG-DISTANCE CORRECTIONS

Let's take a look



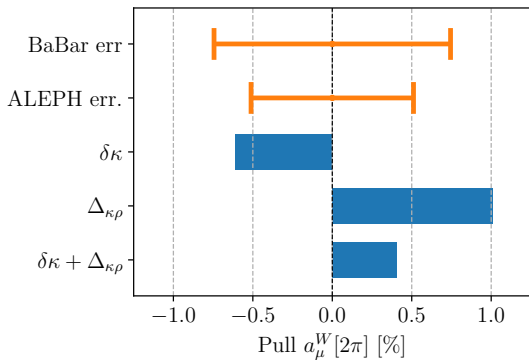
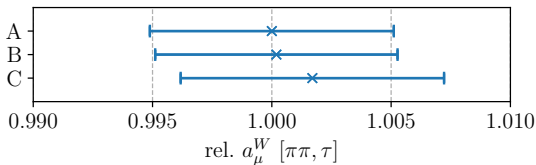
$\delta\kappa$ is channel and m_γ independent [MB et al, in prep]

$\Delta_{\kappa\rho} \rightarrow 2\pi$, point-like, m_γ independent [Cirigliano et al '01, '02]

TOWARDS a_μ^W

2 π channel

ALEPH'13 data
3 analysis groups
relative unblinding



$\Delta\kappa\rho$ bulk from m_ρ region

$S_{EW} \approx 2\%$ largest effect

ISOSPIN BREAKING

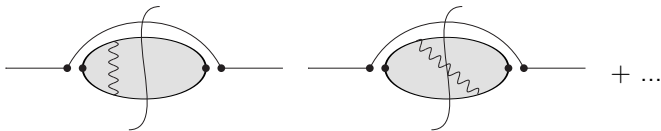
Strategy

1. take experimental $d\Gamma/ds$ (e.g. Aleph13, Belle08, BelleI ?)
2. $\delta\kappa$ initial state corrections: analytic, under control
3. $\Delta_{\kappa\rho}$ initial-finite mixed rad. corr:
 - analytically known for intermediate two-pion channel
 - effective field theory ($R\chi T$) [Cirigliano et al '01, '02]
 - meson dominance models [Flores-Talpa et al. '06, '07]
 - new results from ($R\chi T$) using pheno input [Roig et al '23]
4. define $\delta\Gamma_{EM} \equiv \delta\kappa(s) + \Delta_{\kappa\rho}(s)$ and calculate:

$$\frac{m_\tau^2}{12\pi^2 G_F^2 |V_{ud}|^2 \kappa(s)} \frac{1}{S_{EW}} \frac{1}{1 + \frac{\alpha}{\pi} \delta\Gamma_{EM}(s)} \left[\frac{\mathcal{B}_e}{\mathcal{B}} \frac{1}{\Gamma} \frac{d\Gamma}{ds} \right]_{\text{exp}} = \rho^{w,0}(s) + \delta\rho(s)$$

ISOSPIN BREAKING

Final state



[Braaten, Li '90]

$$-Q_u Q_d \times \begin{array}{c} \text{---} \nearrow \\ \text{---} \searrow \\ \text{---} \nearrow \\ \text{---} \searrow \end{array} + (Q_u^2 + Q_d^2) \times \frac{1}{2} \begin{array}{c} \text{---} \nearrow \\ \text{---} \searrow \\ \text{---} \nearrow \\ \text{---} \searrow \end{array}$$

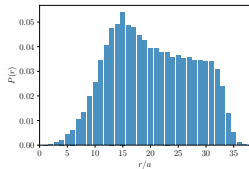
Naive scaling $\frac{\alpha}{\pi} \log(aM_w) \simeq 0.8 - 1\%$ so delicate matching required

FIRST RESULTS

Connected strong-isospin breaking

Ideas from stochastic locality [Lüscher '17][RBC/UKQCD '23][MB, Cé et al '23]

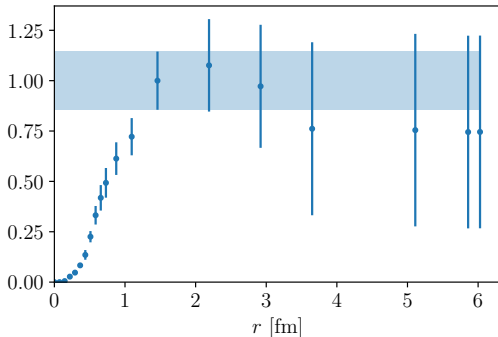
$O(10^3)$ point sources
→ $O(10^6)$ pairs



r = spatial separation vector
and mass operators

t^4 interm. window

[preliminary 96I]



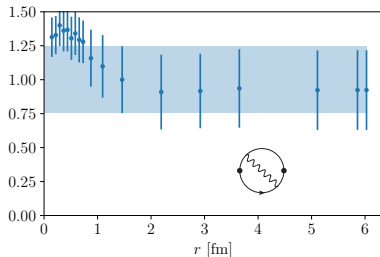
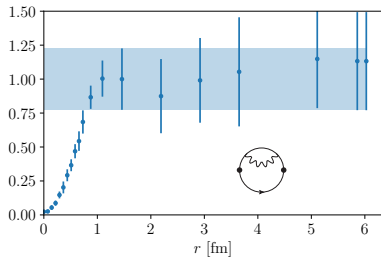
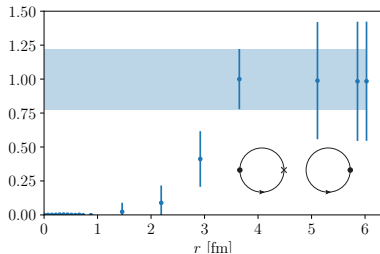
ongoing study 48I, 64I, 96I
3 lattice spacings, phys. mass

FIRST RESULTS

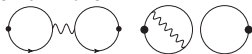
Leading isospin-breaking

[preliminary 96!]

t^4 intermediate window



and more ...




data for all qed/sib sea
diagrams on disk

WHAT'S NEXT

Lattice is fully inclusive...

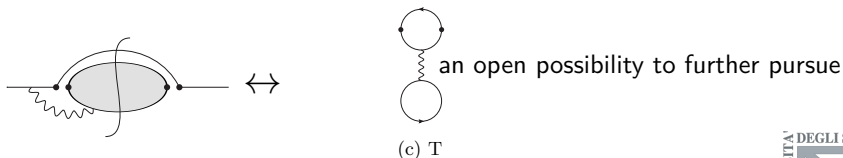
Lattice calculation fully inclusive in energy (cut at m_τ) and channels

 \rightarrow isospin-breaking from both 2π and 3π

[Colangelo et al '22][Hoferichter et al '23]

IB correction of $a^W[3\pi] \approx -1 \cdot 10^{-10}$, $a^W[2\pi] \approx +1 \cdot 10^{-10}$

Possibilities for τ -data + LQCD: (A) fully inclusive vs (B) 2π exclusive

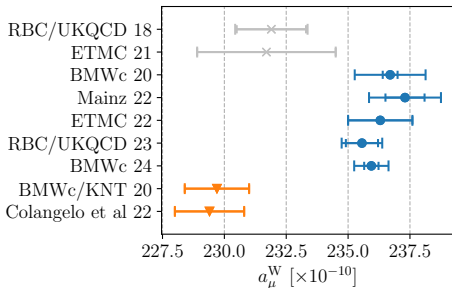


Remaining systematic effects $[m_\tau, \infty)$ in $\rho^{w,0}$ from e^+e^- or lattice*
suppressed for low and intermediate windows

CONCLUSIONS

...and outlooks

hadronic τ -decays can shed light on tension lattice vs e^+e^-



τ data **competitive** on intermediate window

blinded analysis of Aleph

initial+mixed rad.cors. analytic

final radiative from LQCD+QED

Remaining work (in progress) to finalize full formalism [MB et al, in prep]

W-regularization and short-distance corrections

non-factorizable effects: beyond EFT?

Thanks for your attention

DEFINITIONS

Hadronic currents

$$\begin{aligned}\mathcal{J}_\mu^\gamma &= Q_u \bar{u} \gamma_\mu u + Q_d \bar{d} \gamma_\mu d \\ \mathcal{J}_\mu^- &= \bar{u} \gamma_\mu d, \quad \mathcal{J}_\mu^1 = \frac{Q_u - Q_d}{\sqrt{2}} \bar{u} \gamma_\mu d\end{aligned}$$

Hadronic phase-space factor, i labels hadrons

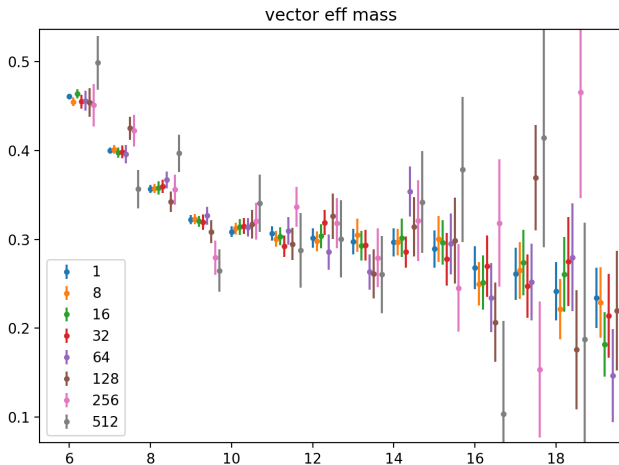
$$d\Phi_f(p) \equiv (2\pi)^4 \delta^4(p - \sum_i p_i) S_f \prod_i \frac{d^3 p_i}{(2\pi)^3 2\omega_i}$$

Charged spectral densities

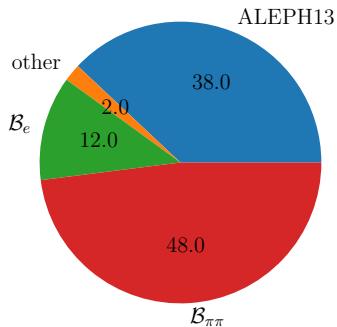
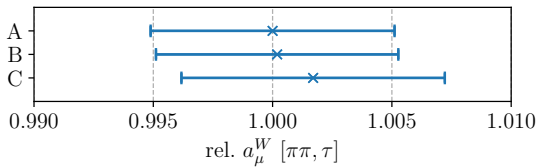
$$\begin{aligned}\rho_{\mu\nu}^w(p) &= \frac{1}{2\pi} \int d^4 x e^{ipx} \langle 0 | \mathcal{J}_\mu^+(x) \mathcal{J}_\nu^-(0) | 0 \rangle \\ &= \frac{1}{2\pi} \sum_f \int d\Phi_f \langle 0 | \mathcal{J}_\mu^+(0) | p_1 \cdots, \text{out} \rangle \langle p_1 \cdots, \text{out} | \mathcal{J}_\nu^-(0) | 0 \rangle \\ &= (p^2 g_{\mu\nu} - p_\mu p_\nu) \rho^w(s) \quad [s = p^2]\end{aligned}$$

SPARSE PROPAGATORS

Save on disk sparse props \rightarrow efficient, more point sources [RBC/UKQCD '18]
side effects? observable dependent, we tested vector correlator



ERROR BREAKDOWN



PHASE SPACE

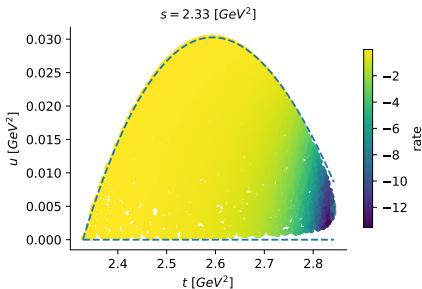
A numerical n -particle phase-space integrator

Grid/GPT backend, support for several parallelization schemes

partial support for 1-loop Passarino-Veltman functions

no support for MCMC yet (needed for ≥ 6 particles)

currently private, soon public github.com/mbruno46



Used to cross-check analytic formulae

Example: Dalitz plot τ Bremsstrahlung

→ wrong boundary: finite m_γ effects