

# Control of equilibrium and non-equilibrium Casimir forces

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# Outline

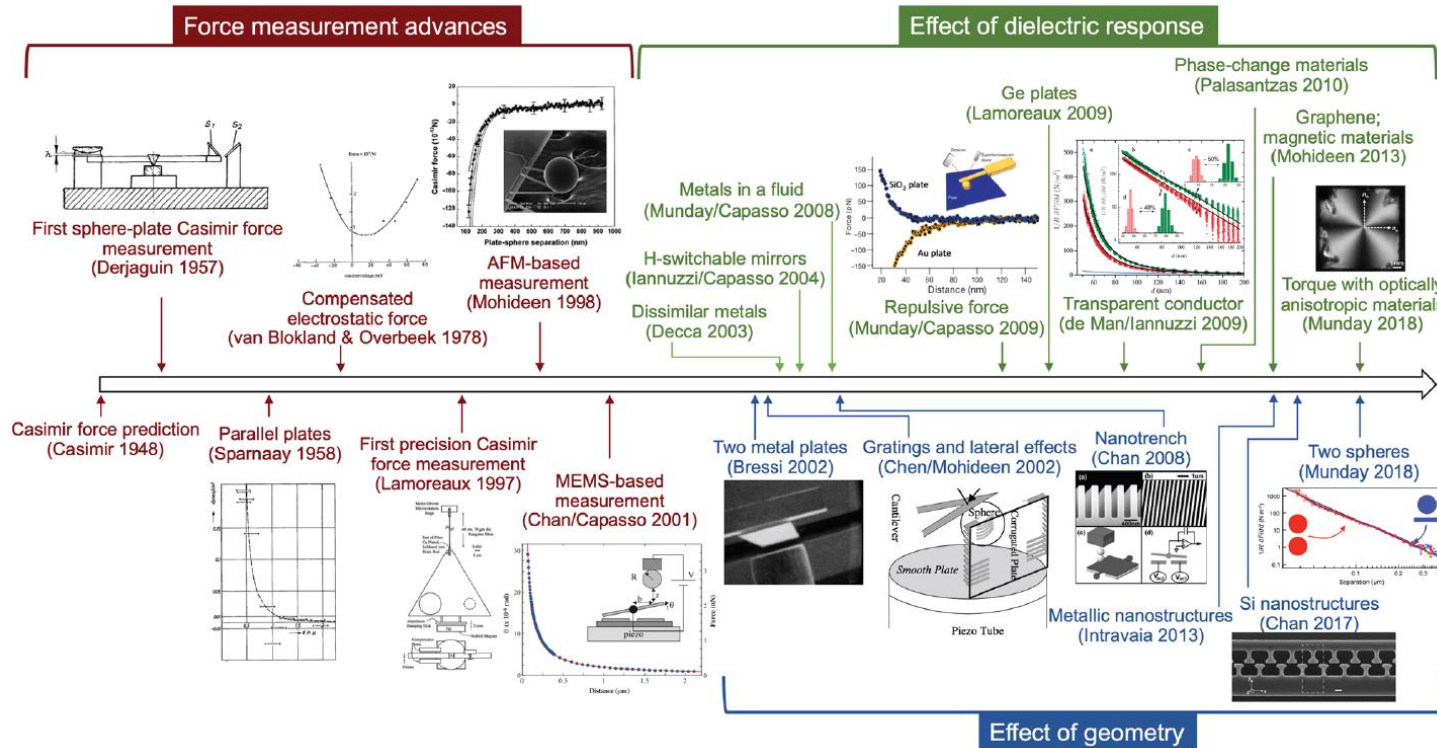
Brief overview of Casimir forces

Control theory for Casimir forces

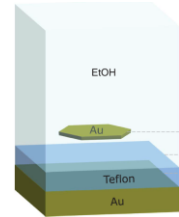
Symmetry of Casimir forces in wavevector space

Conclusions

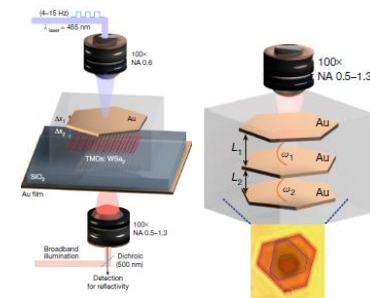
# History of Casimir force research



## Past few years



Levitation in liquid  
[Science 364, 984 (2019)]



Active control, many-body  
in liquid  
[Nature, 597, 214 (2021)]

T. Gong, M. R. Corrado, A. R. Mahbub, C. Shelden, and J. N. Munday, *Nanophotonics* 10, 523 (2021).  
J. N. Munday, KEK IPNS-IMSS-QUP joint workshop Feb. 8-10, 2022.

## Constraining New Forces in the Casimir Regime Using the Isoelectric Technique

R. S. Decca,<sup>1,\*</sup> D. López,<sup>2</sup> H. B. Chan,<sup>3</sup> E. Fischbach,<sup>4</sup> D. E. Krause,<sup>5,†</sup> and C. R. Jamell<sup>1</sup>

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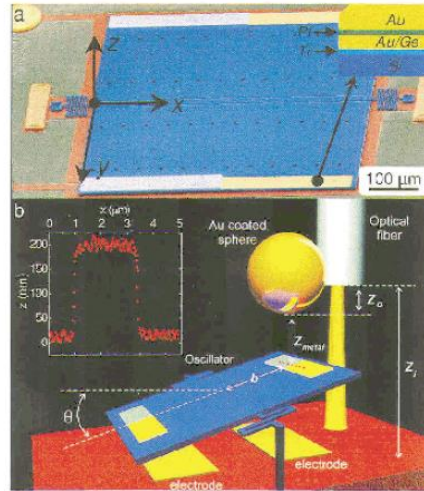


FIG. 1 (color). (a) Scanning electron microscope image of the MTO with the composite sample deposited on it. The coordinate system used in the Letter is indicated. Inset: schematic of the sample deposited on the MTO. The thickness of the different layers are (in order of deposition):  $d_{Ti} = 1$  nm,  $d_{Ge} = 200$  nm,  $d_{Pt} = 1$  nm, and  $d_{Au}^p = 150$  nm. The thickness of the layers deposited on the sphere (not shown) are:  $d_{Cr} = 1$  nm and  $d_{Au}^s = 200$  nm. (b) Experimental setup. The red dotted line indicates where AFM line cuts were taken. Inset: AFM profile of the sample interface.

## Hypothetical force difference

$$\Delta F_h = -4\pi^2 G \alpha \lambda^3 e^{-z/\lambda} R K_s K_p,$$

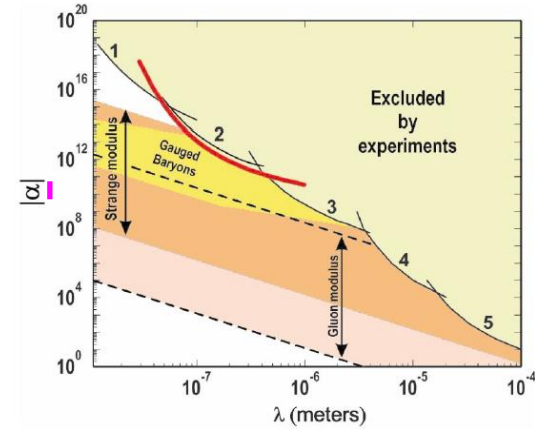


FIG. 4 (color). Values in the  $\{\lambda, \alpha\}$  space excluded by experiments. The red curve represents limits obtained in this work. Curves 1 to 5 were obtained by Mohideen's group [7], our group [1], Lamoreaux [6], Kapitulnik's group [8], and Price's group [5], respectively. Also shown are theoretical predictions [21].

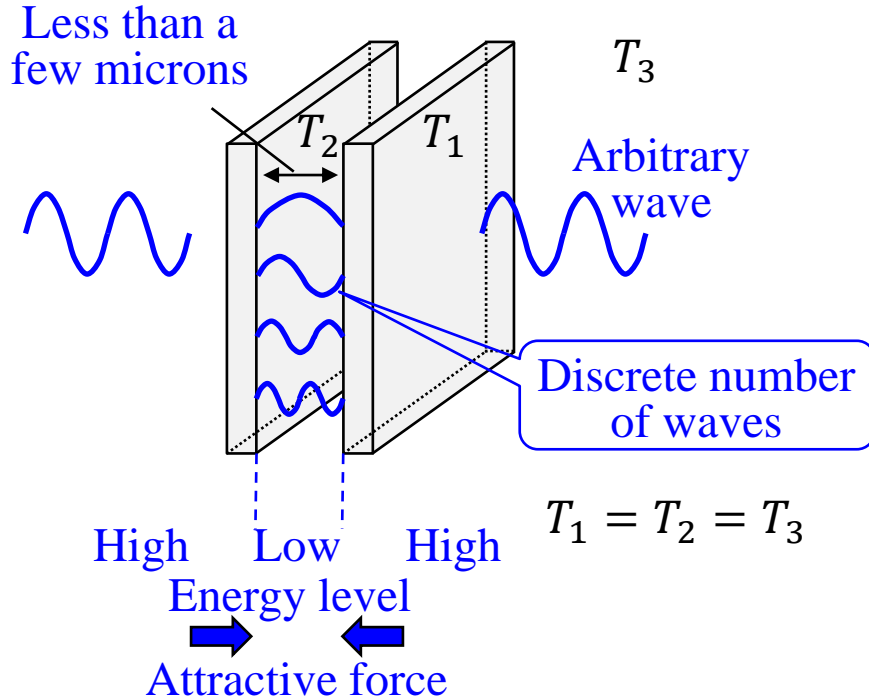
$G$ : Gravitational constant

$K_s, K_p$ : Terms relating to densities of the sphere and the plate

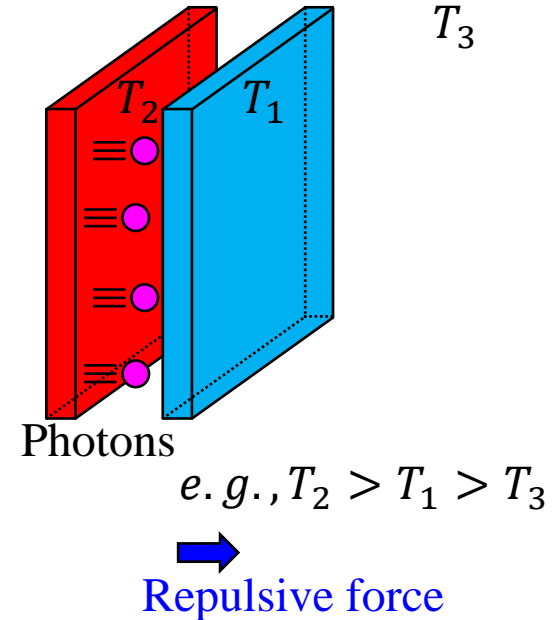
$R$ : Radius of the sphere

# Equilibrium and non-equilibrium Casimir forces

Equilibrium Casimir force  
(Attractive in vacuum)



Non-equilibrium Casimir force  
(Can be repulsive in vacuum)



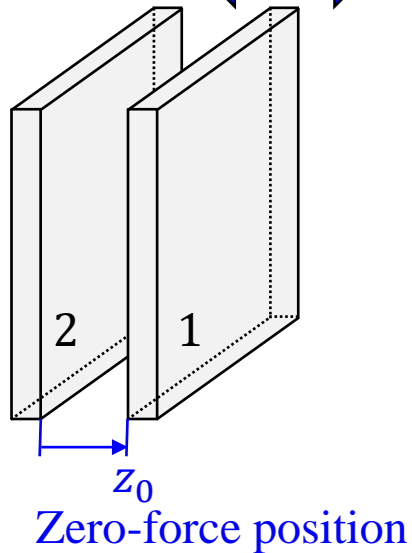
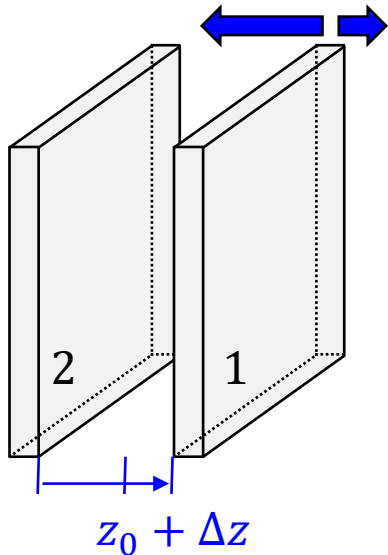
Note: The plates consist of reciprocal materials.

# Stable and unstable Casimir force systems

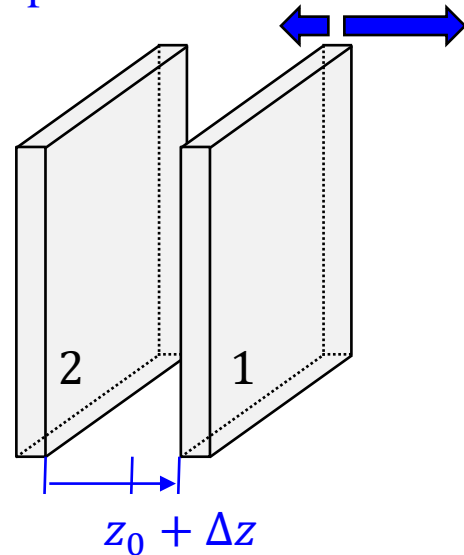
When attractive and repulsive force components acting on body 1 are balanced, zero-force position can exist.

Attractive force ← → Repulsive force

Stable  
Attractive force is enhanced.

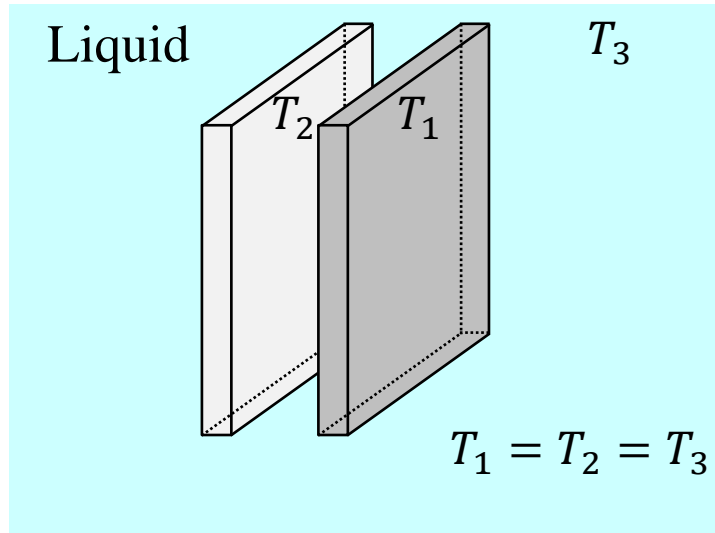


Unstable  
Repulsive force is enhanced.

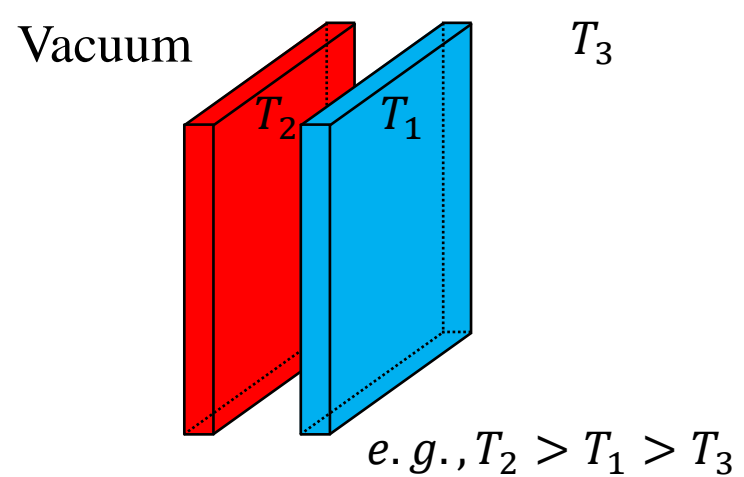


# Examples of stable and unstable systems

Equilibrium Casimir force in liquid  
(Can be stable)



Non-equilibrium Casimir force in vacuum  
(Unstable)

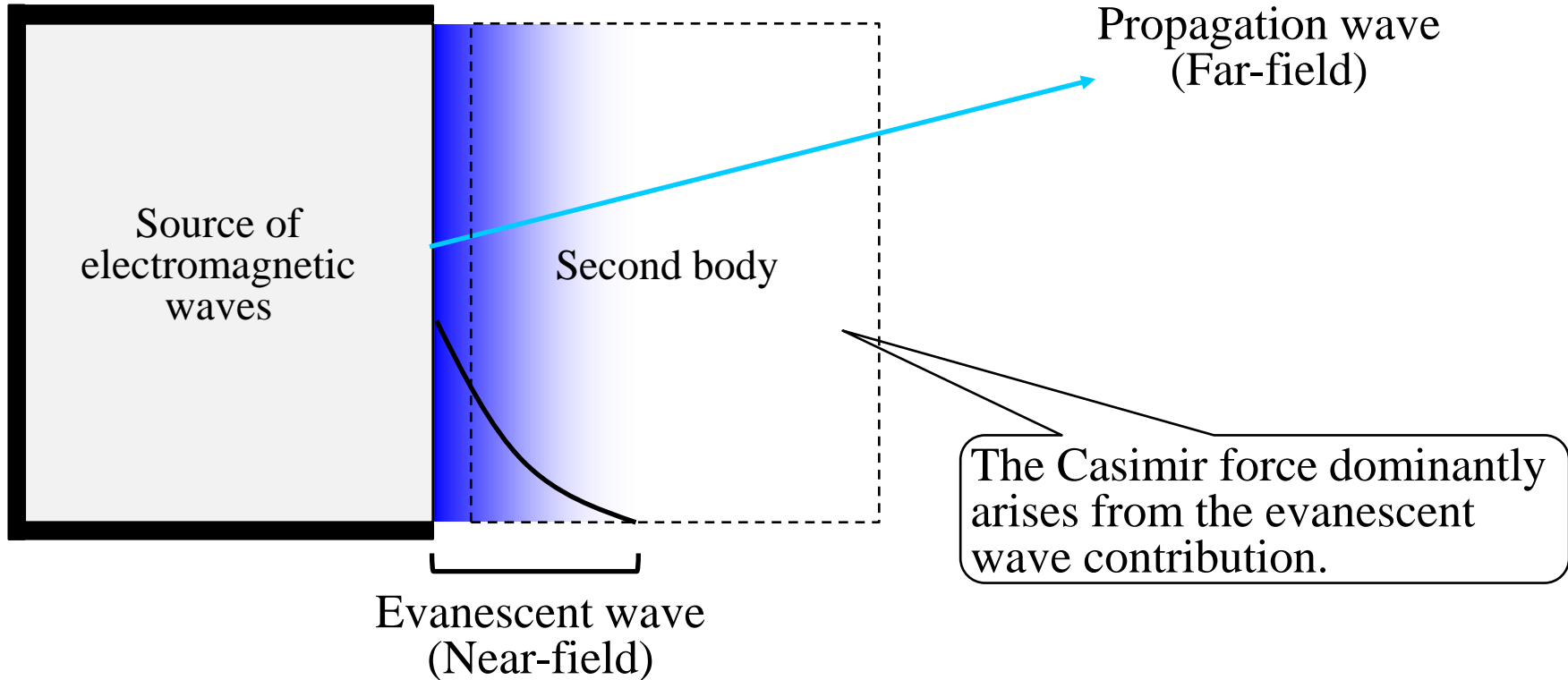


[See *Science* 364, 984 (2019)]

Note: The plates consist of reciprocal materials.

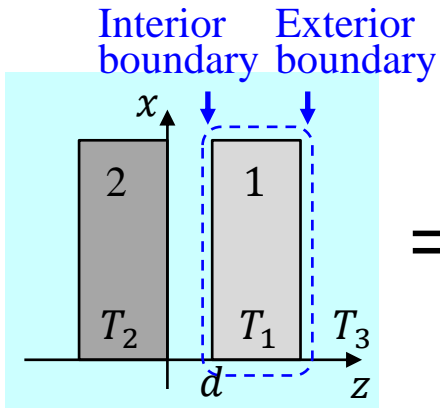
# Propagation and evanescent waves

The propagation wave goes away from the electromagnetic source.  
The evanescent wave stays around the electromagnetic source.

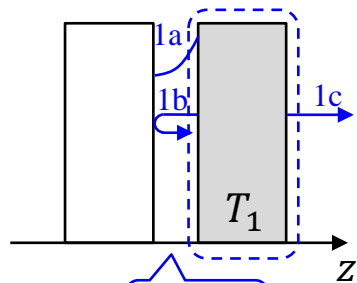




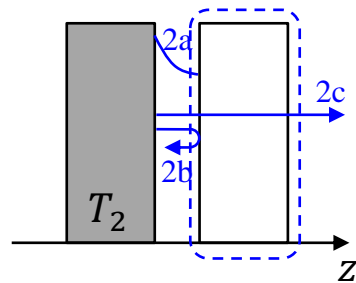
# Casimir force calculation



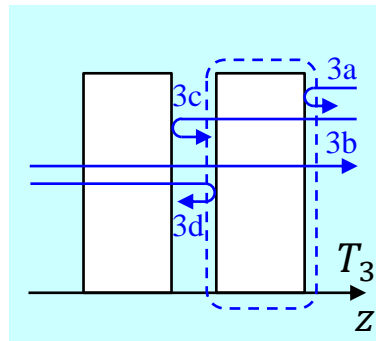
=



+



+



Maxwell stress tensor

$$T_{ij} = \epsilon_0 E_i E_j + \mu_0 H_i H_j - \delta_{ij} \frac{1}{2} (\epsilon_0 E^2 + \mu_0 H^2), \quad (1)$$

Casimir force acting on body 1 in the two-body system (isotropic materials)

$$F_Z = \int_S \langle T_{ZZ} \rangle dz, \quad (2)$$

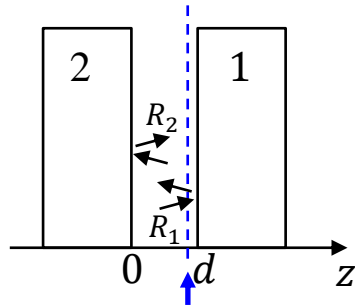
$$\begin{aligned}
 = & - \sum_{p,s} \frac{\hbar}{\pi^2} \int_0^\infty d\omega \left\{ \left[ n(\omega, T_1) + \frac{1}{2} \right] \left[ \int_{k_0}^\infty k_{\parallel} dk_{\parallel} \kappa_0 \frac{\text{Im}(\tilde{r}_{01}) \text{Re}(\tilde{r}_{02}) e^{-2\kappa_0 d}}{|1 - \tilde{r}_{01} \tilde{r}_{02} e^{-2\kappa_0 d}|^2} + \int_0^{k_0} k_{\parallel} dk_{\parallel} k_{z0} \frac{1}{4} \left( - \frac{(1 - |\tilde{r}_{01}|^2 - |\tilde{t}_{01}|^2)(1 + |\tilde{r}_{02}|^2)}{|1 - \tilde{r}_{01} \tilde{r}_{02} e^{i2k_{z0} d}|^2} + 1 - \left| \tilde{r}_{31} + \frac{\tilde{r}_{02} \tilde{t}_{01}^2 e^{i2k_{z0} d}}{1 - \tilde{r}_{01} \tilde{r}_{02} e^{i2k_{z0} d}} \right|^2 - \frac{|\tilde{t}_{01}|^2 (1 - |\tilde{r}_{02}|^2)}{|1 - \tilde{r}_{01} \tilde{r}_{02} e^{i2k_{z0} d}|^2} \right) \right] \right. \\
 & + \left[ n(\omega, T_2) + \frac{1}{2} \right] \left[ \int_{k_0}^\infty k_{\parallel} dk_{\parallel} \kappa_0 \frac{\text{Im}(\tilde{r}_{02}) \text{Re}(\tilde{r}_{01}) e^{-2\kappa_0 d}}{|1 - \tilde{r}_{01} \tilde{r}_{02} e^{-2\kappa_0 d}|^2} + \int_0^{k_0} k_{\parallel} dk_{\parallel} k_{z0} \frac{1}{4} \left( - \frac{(1 - |\tilde{r}_{02}|^2 - |\tilde{t}_{02}|^2)(1 + |\tilde{r}_{01}|^2)}{|1 - \tilde{r}_{01} \tilde{r}_{02} e^{i2k_{z0} d}|^2} + \frac{|\tilde{t}_{01}|^2 (1 - |\tilde{r}_{02}|^2 - |\tilde{t}_{02}|^2)}{|1 - \tilde{r}_{01} \tilde{r}_{02} e^{i2k_{z0} d}|^2} \right) \right] \\
 & \left. + \left[ n(\omega, T_3) + \frac{1}{2} \right] \int_0^{k_0} k_{\parallel} dk_{\parallel} k_{z0} \frac{1}{4} \left( 1 + \left| \tilde{r}_{31} + \frac{\tilde{r}_{02} \tilde{t}_{01}^2 e^{i2k_{z0} d}}{1 - \tilde{r}_{01} \tilde{r}_{02} e^{i2k_{z0} d}} \right|^2 + \frac{|\tilde{t}_{01}|^2 |\tilde{t}_{02}|^2}{|1 - \tilde{r}_{01} \tilde{r}_{02} e^{i2k_{z0} d}|^2} - \frac{|\tilde{t}_{01}|^2 (1 + |\tilde{r}_{02}|^2)}{|1 - \tilde{r}_{01} \tilde{r}_{02} e^{i2k_{z0} d}|^2} - \frac{|\tilde{t}_{02}|^2 (1 + |\tilde{r}_{01}|^2)}{|1 - \tilde{r}_{01} \tilde{r}_{02} e^{i2k_{z0} d}|^2} \right) \right\} \quad (3)
 \end{aligned}$$

$\tilde{r}_{0j}$ : reflection coefficient

$\tilde{t}_{0j}$ : reflection coefficient

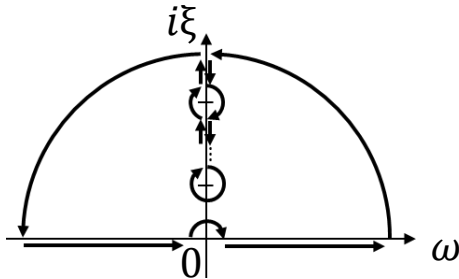
# Casimir force calculation in equilibrium

Calculations of equilibrium Casimir forces can be simplified since photon exchange is balanced at the exterior boundary. (The interior boundary is only considered.)



Interior  
boundary

$$T \equiv T_1 = T_2 = T_3$$



Casimir force formula, integration along the real frequency axis

$$P_t = - \sum_{j=p,s} \int_0^{\infty} \frac{k_{\parallel} dk_{\parallel}}{2\pi} \int_0^{\infty} \frac{d\omega}{2\pi} 4\hbar \left[ n(\omega, T) + \frac{1}{2} \right] \text{Re}[k_z Z(\omega, \beta)], \quad (11a)$$

$$Z(\omega, \beta) = \frac{R_1(\omega, \beta)R_2(\omega, \beta)e^{i2k_z d}}{1 - R_1(\omega, \beta)R_2(\omega, \beta)e^{i2k_z d}}, \quad (11b)$$

- Understanding the mechanism
- Long calculation time

Wick rotation approach, integration along the imaginary frequency axis (E.M. Lifshitz, Sov. Phys. 1956)

$$P_t = \sum_{j=p,s} \int_0^{\infty} \frac{k_{\parallel} dk_{\parallel}}{2\pi} 2k_B T \sum_{n=0}^{\infty} q_{0,n} Z(i\xi_n, \beta), \quad (12a)$$

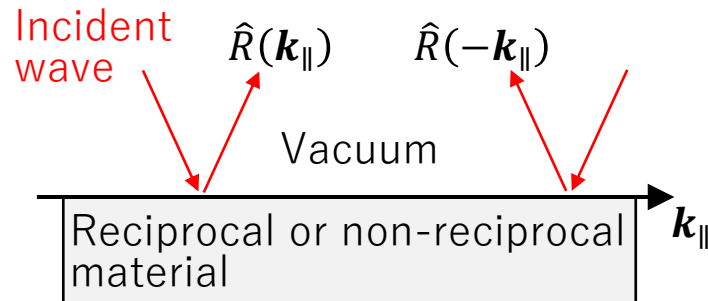
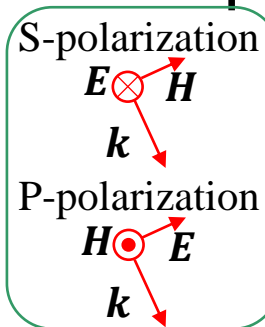
$$Z(i\xi_n, \beta) = \frac{R_1(i\xi_n, \beta)R_2(i\xi_n, \beta)e^{-2q_{0,n}d}}{1 - R_1(i\xi_n, \beta)R_2(i\xi_n, \beta)e^{-2q_{0,n}d}}, \quad (12b),$$

- Significantly reduced calculation time
- Little observation of the mechanism

# Reciprocal and non-reciprocal materials

Reflection matrix

$$\hat{R}(\mathbf{k}_{\parallel}) = \begin{bmatrix} R^{s \rightarrow s}(\mathbf{k}_{\parallel}) & R^{p \rightarrow s}(\mathbf{k}_{\parallel}) \\ R^{s \rightarrow p}(\mathbf{k}_{\parallel}) & R^{p \rightarrow p}(\mathbf{k}_{\parallel}) \end{bmatrix}, \quad (21)$$



Reciprocal materials

$$\hat{R}(-\mathbf{k}_{\parallel}) = \hat{\sigma}_z \hat{R}^T(\mathbf{k}_{\parallel}) \hat{\sigma}_z, \quad (22)$$

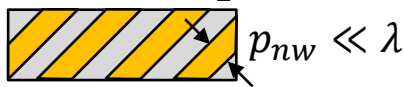
$$\left( \hat{\sigma}_z = \begin{bmatrix} 1 & \\ & -1 \end{bmatrix} \right)$$

Bulk  
(Isotropic)



$$\hat{\epsilon}_i = \begin{bmatrix} \epsilon_p & & \\ & \epsilon_p & \\ & & \epsilon_p \end{bmatrix}, \quad (23)$$

Inclined nanowires  
(Anisotropic)

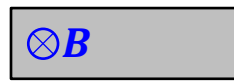


$$\hat{\epsilon}^S = \begin{bmatrix} \epsilon_d & & \epsilon_f \\ & \epsilon_p & \\ \epsilon_f & & \epsilon_d \end{bmatrix}, \quad (24)$$

Non-reciprocal materials

Eq. (22) can be violated.

InSb



$$\hat{\epsilon}^A = \begin{bmatrix} \epsilon_d & & i\epsilon_f \\ & \epsilon_p & \\ -i\epsilon_f & & \epsilon_d \end{bmatrix}, \quad (25)$$

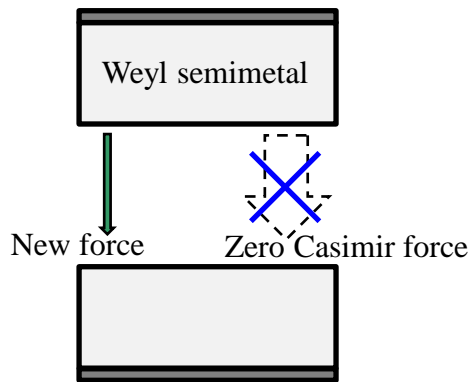
Magnetic Weyl  
semimetals



# Our Casimir force research

## Particle physics

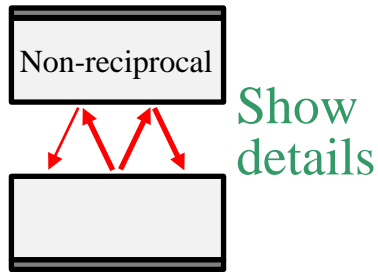
New force search via zero Casimir force[1]



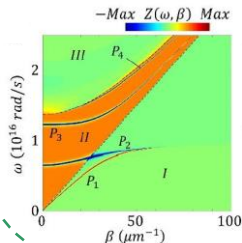
[1] Y. Ema, M. Hazumi, H. Iizuka, K. Mukaida, and K. Nakayama, Phys. Rev. D 108, 016009 (2023).

## Fundamental understanding

[2] Symmetry argument in Casimir forces

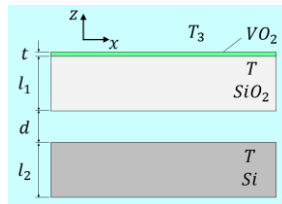


[3] Casimir force is insensitive to material loss



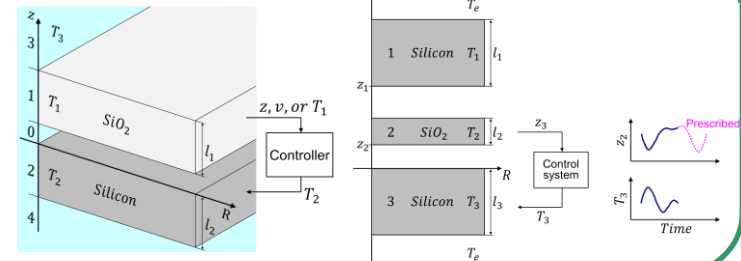
[2] H. Iizuka and S. Fan, Phys. Rev. B 108, 075429 (2023).  
 [3] H. Iizuka and S. Fan, J. Optical Society of America B 36, 2981 (2019).  
 [4] H. Iizuka and S. Fan, J. Optical Society America B 38, 151–158 (2021).

[4] Exterior control of non-equilibrium Casimir force

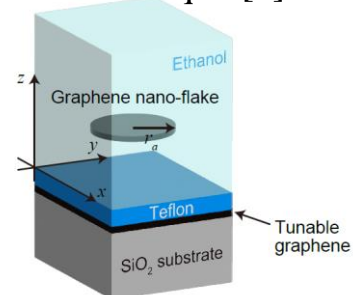


## Toward industry

Dynamic control of Casimir forces in vacuum  
 Control theory[5] Trajectory tracking[6]



Dynamic control of Casimir forces in liquid[7]



[5] H. Iizuka and S. Fan, Applied Physics Letters 118, 144001 (2021).  
 [6] H. Iizuka and S. Fan, J. Quantitative Spectroscopy Radiative Transfer 289, 108281 (2022).  
 [7] H. Toyama, T. Ikeda, and H. Iizuka, Phys. Rev. B 108, 245402 (2023).

# Outline

Brief overview of Casimir forces

**Control theory for Casimir forces**

Symmetry of Casimir forces in wavevector space

Conclusions

# Casimir force to Industries

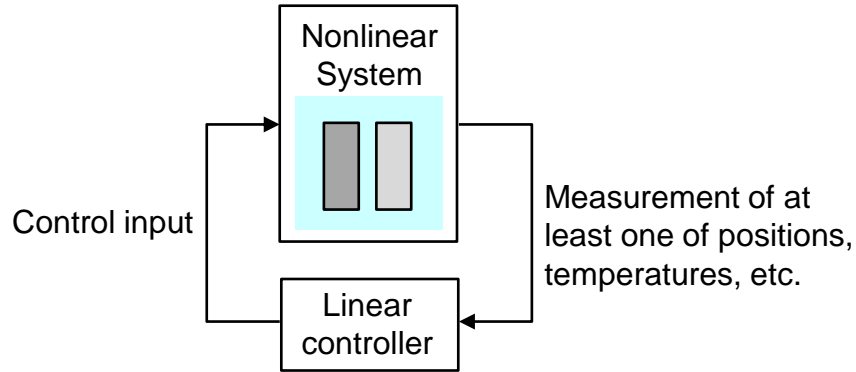
Control theory has been widely used in industries.



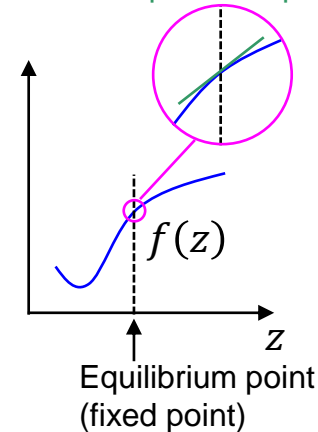
We introduce control theory to manipulate Casimir forces for the first time.

We show that a linear controller enables a nonlinear system interacting through non-equilibrium Casimir forces to be stable around the equilibrium point with a limited number of sensors.

└ zero force, zero net heat transfer



A nonlinear function can be well approximated by the linear function around the equilibrium point.



# Linear control system (1/2)

Dynamics of a linear system is given by

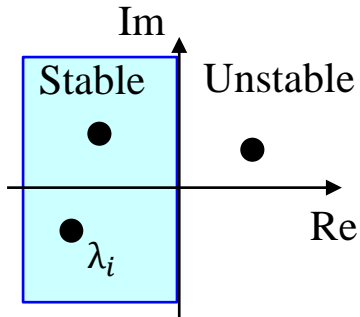
$$\frac{dx}{dt} = Ax, \quad (51)$$

$n \times n$  matrix  $\downarrow$  State vector  $x = \begin{bmatrix} x_1 \\ \vdots \\ x_n \end{bmatrix}$

System  
 $\dot{x} = Ax$

The system is stable if all eigenvalues meet

$$\text{Re}[\lambda_i] < 0, \quad (52)$$



Eigenvalues (unstable)

The controller changes the dynamics.

$$\frac{dx}{dt} = Ax + Bu, \quad (53)$$

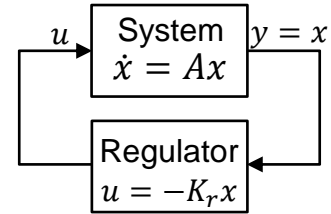
$\downarrow$  Control input

$$u = -K_r x, \quad (54)$$

$\downarrow$  Gain matrix of the regulator

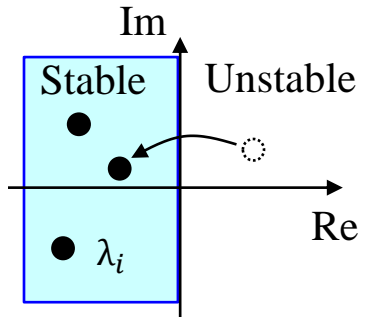
Eqs. (53),(54) give

$$\frac{dx}{dt} = \underline{(A - BK_r)}x, \quad (55)$$



If the controllability matrix is full ranked, the system is “controllable”.  
(All  $\lambda_i$  can have negative real parts.)

$$Q = [A, AB, \dots, AB^{n-1}], \quad (56)$$



Eigenvalues (stable)

# Linear control system (2/2)

The estimator (Kalman filter) enables the system to be stable with limited measurement results.

$$\frac{dx}{dt} = Ax + Bu, \quad (61)$$

$$y = Cx, \quad (62)$$

└ Observation vector      └ Gain matrix of the estimator

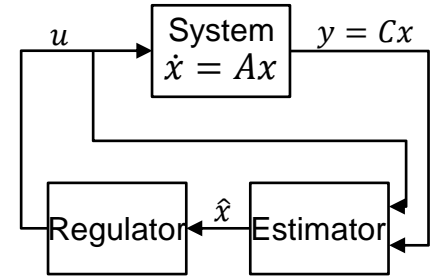
$$\frac{d\hat{x}}{dt} = (A - K_f C)\hat{x} + [B \quad K_f] \begin{bmatrix} u \\ y \end{bmatrix}, \quad (63)$$

└ Estimated state vector

$$u = -K_r \hat{x}, \quad (64)$$

If the observability matrix is full ranked, the system is “observable”.  
(The system is stable with limited measurement results.)

$$N = [C, CA, \dots, CA^{n-1}], \quad (65)$$

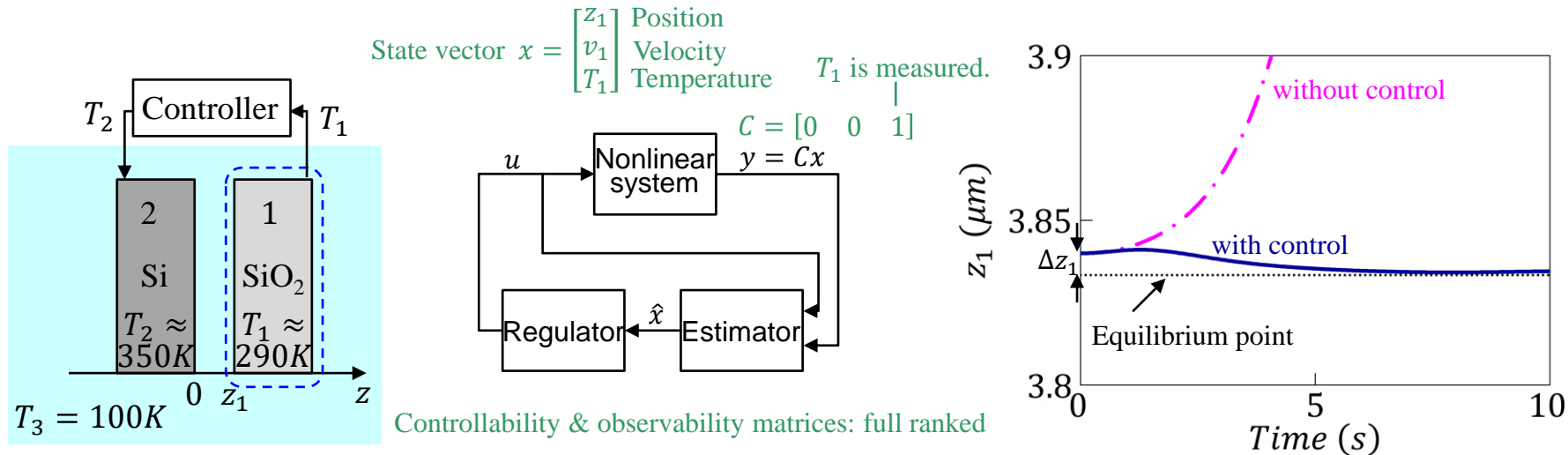




# Control of non-equilibrium Casimir force

Consider a system consisting of a silicon plate and a SiO<sub>2</sub> plate in a low-temperature environment. The silicon plate is fixed and the temperature can be adjusted while the SiO<sub>2</sub> plate is moved and the temperature is determined through the thermal emission exchange.

We theoretically prove that the SiO<sub>2</sub> plate can stay at a distance away from the silicon plate indefinitely by monitoring the temperature of the SiO<sub>2</sub> plate and adjusting the temperature of the silicon plate.



# Trajectory tracking through the control of Casimir force

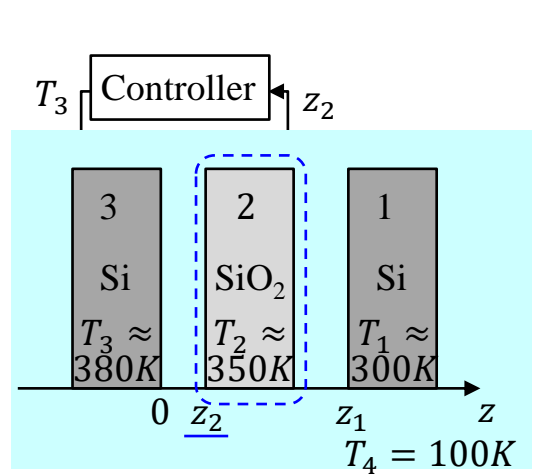
An autonomous vehicle can follow a path as we set by control theory, which is called trajectory tracking.



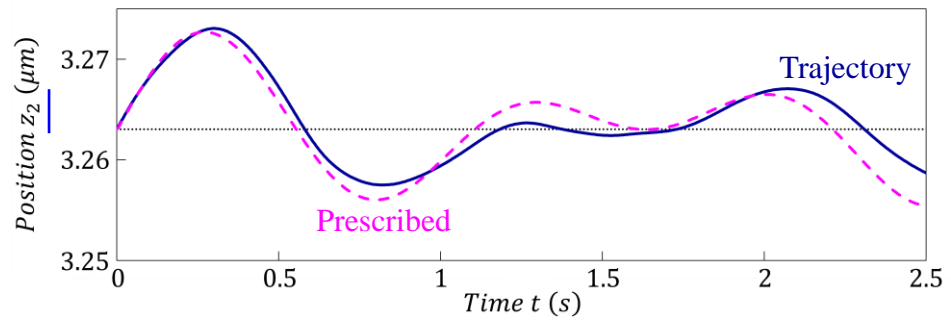
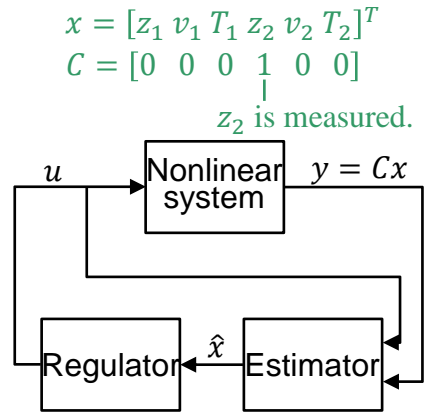
↓ Quantum world

Trajectory tracking is applied to a three-body system interacting through non-equilibrium Casimir forces.

The position of the middle SiO<sub>2</sub> plate (blue line) follows a prescribed trajectory (pink dashed line) by adjusting the temperature of the left silicon plate (single parameter) while the right silicon plate is kept around the equilibrium position.



Controllability & observability matrices: full ranked



# Outline

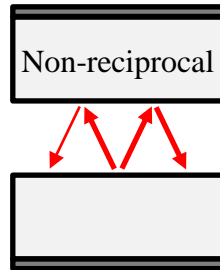
Brief overview of Casimir forces

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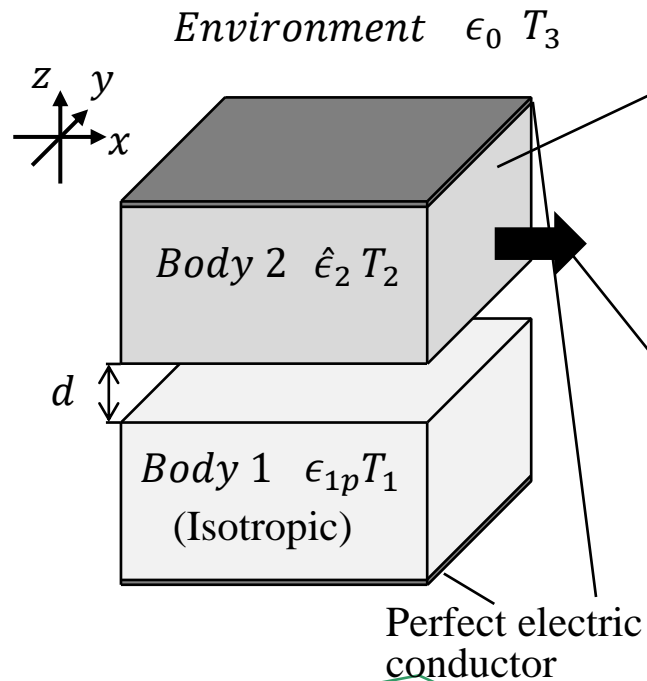
Conclusions

[2] Symmetry argument  
in Casimir forces



[2] H. Iizuka and S. Fan, Phys. Rev. B 108, 075429 (2023).

# Casimir lateral force



Anisotropic reciprocal material



$$\epsilon_2^S = \begin{bmatrix} \epsilon_d & & \epsilon_f \\ & \epsilon_p & \\ \epsilon_f & & \epsilon_d \end{bmatrix}, \quad (24)$$

or

non-reciprocal material



$$\epsilon_2^A = \begin{bmatrix} \epsilon_d & & i\epsilon_f \\ & \epsilon_p & \\ -i\epsilon_f & & \epsilon_d \end{bmatrix}, \quad (25)$$

Casimir lateral force acting on body 2

$$\mathbf{F}_2^{\parallel}(T_1, T_2) = \int_0^{\infty} d\omega \int_{-\infty}^{\infty} d\mathbf{k}_{\parallel} \underline{\underline{\mathbf{F}_2^{\parallel}(\omega, \mathbf{k}_{\parallel}, T_1, T_2)}}, \quad (41)$$

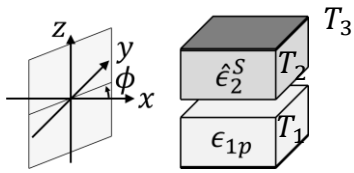
$$\underline{\underline{\mathbf{F}_2^{\parallel}(\omega, \mathbf{k}_{\parallel}, T_1, T_2)}} = -\mathbf{F}_{1 \rightarrow 2}^{\parallel}(\omega, \mathbf{k}_{\parallel}, T_1) - \mathbf{F}_{2 \rightarrow 1}^{\parallel}(\omega, \mathbf{k}_{\parallel}, T_2), \quad (42)$$

$$\mathbf{F}_{1 \rightarrow 2}^{\parallel}(\omega, \mathbf{k}_{\parallel}, T_1) = \left[ n(\omega, T_1) + \frac{1}{2} \right] \frac{\hbar \mathbf{k}_{\parallel}}{8\pi^3} \underline{\underline{\tilde{\mathbf{F}}_{1 \rightarrow 2}^{\parallel}(\omega, \mathbf{k}_{\parallel})}}, \quad (43)$$

$$\tilde{\mathbf{F}}_{l \rightarrow m}^{\parallel}(\omega, \mathbf{k}_{\parallel}) = \begin{cases} \text{Tr} \left[ (-1)^l (\hat{I} - \hat{R}_m^{\dagger} \hat{R}_m) \hat{D}_{lm} (\hat{I} - \hat{R}_l \hat{R}_l^{\dagger}) \hat{D}_{lm}^{\dagger} \right], & (k_{\parallel} < k_0) \\ \text{Tr} \left[ (-1)^l (\hat{R}_m^{\dagger} - \hat{R}_m) \hat{D}_{lm} (\hat{R}_l - \hat{R}_l^{\dagger}) \hat{D}_{lm}^{\dagger} e^{-2\kappa_{z0} d} \right], & (k_{\parallel} > k_0) \end{cases}, \quad (44)$$

No photon exchange between the two-body system and the environment.

$T_1 = T_2 = T_3 = T$  (Equilibrium)  $\tilde{\mathbf{F}}_{l \rightarrow m}^{\parallel}(\omega, \mathbf{k}_{\parallel}) =$  Exchange function  
 $T_1 \neq T_2$  (Non-equilibrium)



# Casimir lateral force, reciprocal

Casimir lateral force is symmetric for reciprocal systems in equilibrium and non-equilibrium.

$$\tilde{F}_{1 \rightarrow 2}^{\parallel} + \tilde{F}_{2 \rightarrow 1}^{\parallel} \quad a\tilde{F}_{1 \rightarrow 2}^{\parallel} + b\tilde{F}_{2 \rightarrow 1}^{\parallel}$$

$$\underline{F_2^{\parallel}(\omega, \mathbf{k}_{\parallel}, T_1, T_2) = -F_2^{\parallel}(\omega, -\mathbf{k}_{\parallel}, T_1, T_2), (45)}$$

Symmetric for  $\mathbf{k}_{\parallel}$

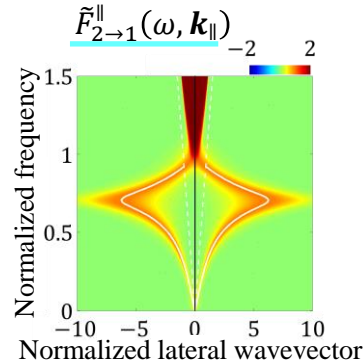
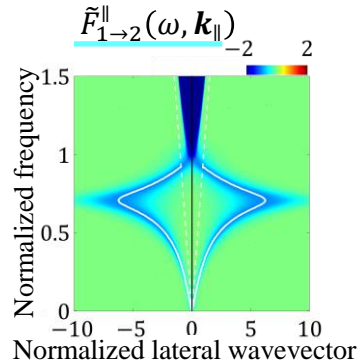
Casimir lateral force  
in  $(\omega, \mathbf{k}_{\parallel})$  space

$$\underline{F_2^{\parallel}(\omega, \mathbf{k}_{\parallel}, T_1, T_2, T_3)}$$

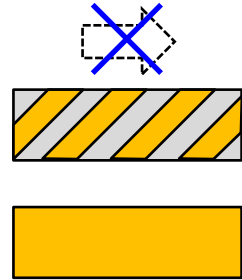
Exchange function

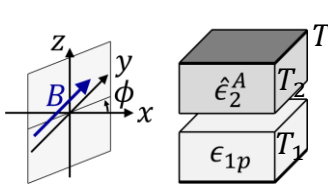
$$\underline{\tilde{F}_{l \rightarrow m}^{\parallel}(\omega, \mathbf{k}_{\parallel})}$$

White lines:  
Dispersion curves



Lateral force does not occur in  
equilibrium and nonequilibrium  
for reciprocal systems.





# Casimir lateral force, non-reciprocal

Symmetry of the Casimir lateral force is broken for non-reciprocal systems in equilibrium and non-equilibrium.

$$\tilde{F}_{1 \rightarrow 2}^{\parallel} + \tilde{F}_{2 \rightarrow 1}^{\parallel} \quad a\tilde{F}_{1 \rightarrow 2}^{\parallel} + b\tilde{F}_{2 \rightarrow 1}^{\parallel}$$

Cancelled out

Asymmetric for  $k_{\parallel}$

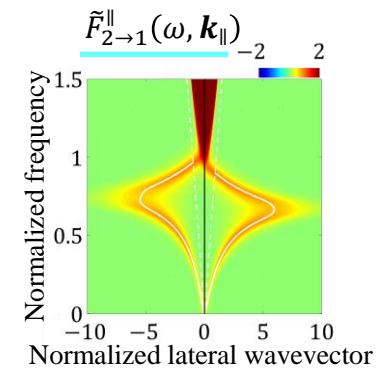
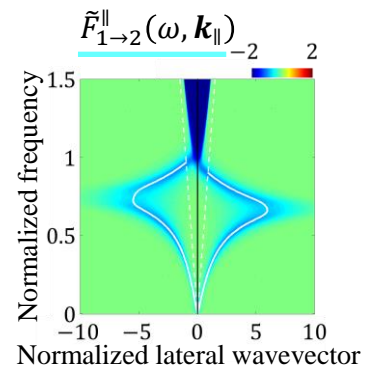
Casimir lateral force in  $(\omega, k_{\parallel})$  space

$$F_2^{\parallel}(\omega, \mathbf{k}_{\parallel}, T_1, T_2, T_3)$$

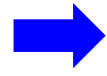
Exchange function

$$\tilde{F}_{l \rightarrow m}^{\parallel}(\omega, \mathbf{k}_{\parallel})$$

White lines:  
Dispersion curves



Lateral force can occur in non-equilibrium for non-reciprocal systems.



InSb/Weyl

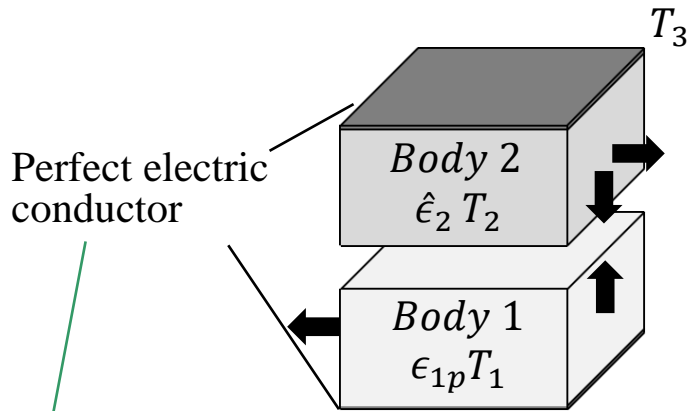


# Newton's third law

Newton's third law holds for every frequency and wavevector, as long as no exchange of photons occurs between the two-body system and the environment.

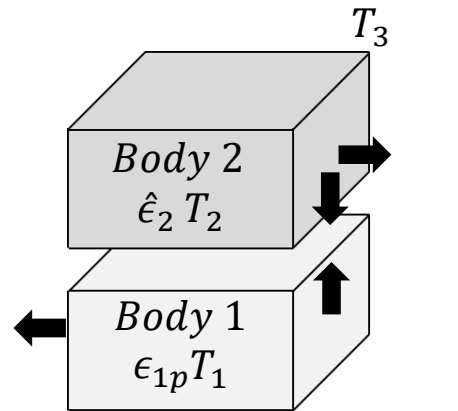
Casimir pressure:  $F_1^Z(\omega, \mathbf{k}_\parallel, T_1, T_2, T_3) = -F_2^Z(\omega, \mathbf{k}_\parallel, T_1, T_2, T_3)$

Casimir lateral force:  $F_1^\parallel(\omega, \mathbf{k}_\parallel, T_1, T_2) = -F_2^\parallel(\omega, \mathbf{k}_\parallel, T_1, T_2)$



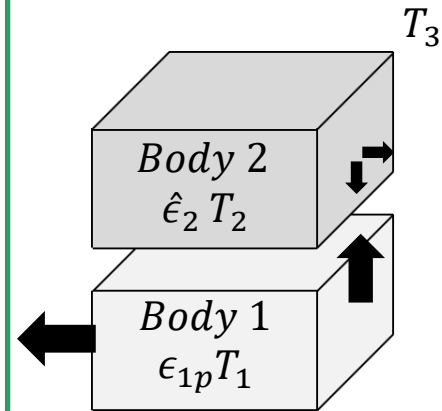
Equilibrium ( $T_1 = T_2 = T_3$ ) and non-equilibrium ( $T_1 \neq T_2$ )

No exchange of photons



Equilibrium ( $T_1 = T_2 = T_3$ )

Newton's third law does not hold.



Non-equilibrium  
( $T_1 \neq T_3$  or  $T_2 \neq T_3$ )

Exchange of photons

The above is true for both reciprocal  $\hat{\epsilon}_2 = \hat{\epsilon}_2^S$  and non-reciprocal  $\hat{\epsilon}_2 = \hat{\epsilon}_2^A$  materials.

# Conclusions

A brief overview of Casimir forces was presented.

Parallel-plate systems for non-equilibrium Casimir forces are controllable and observable by introducing a linear controller.

Symmetry of Casimir forces in wavevector space was discussed. This understanding is helpful for investigating Casimir forces using Weyl semimetals.

## Acknowledgement

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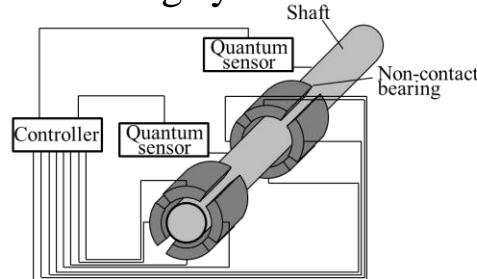
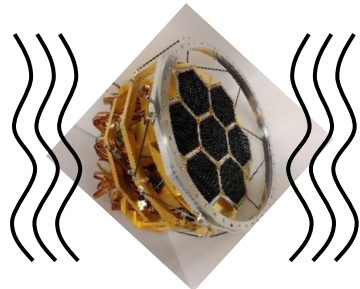
# Research scope of QUP quantum sensor cluster <sup>25/26</sup>

Active vibration isolation for CMB

New physics search

Circuit diagnosis

Non-contact shaft-bearing system



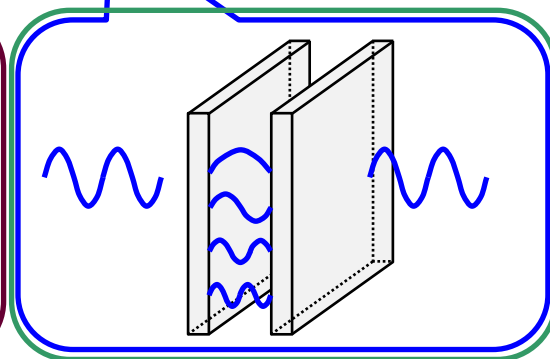
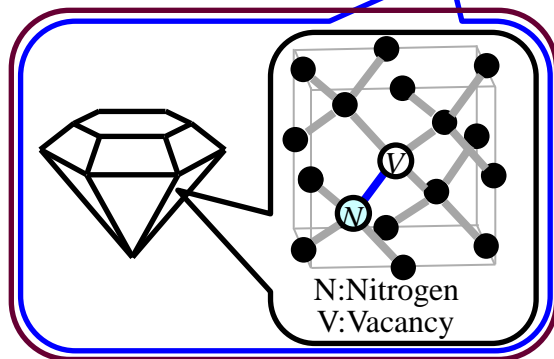
Particle physics/Universe

Industry/Social implementation

Casimir forces

Quantum sensors

Internship program



Today's talk

# QUP Internship program

<https://www2.kek.jp/qup/en/jobs/qupip.html>

Jobs

## QUPIP

Announcement

2023.10.18

### QUP Internship Program (QUPIP)

The International Center for Quantum-field Measurement Systems for Studies of the Universe and Particles (WPI-QUP) invites applications for young researchers (post-doctoral fellows and graduate students) to stay and work with QUP researchers. QUP is developing a new measurement system using existing quantum fields to explore unknown quantum fields under "Bringing New Eyes to Humanity."

**Hideo Iizuka** (Research Location: KEK Tsukuba, Japan)

Our aim is to search for new quantum fields enabled by quantum effects in a single lab space. We have started setting up a measurement system using nitrogen-vacancy centers in diamond for light dark matter search.

Example study subjects for the young researchers

- Setup of the ODMR (optically detected magnetic resonance) measurement system.
- Evaluation of optical properties of diamond samples.

## Quantum sensors

