

Black String in the Standard Model

Yu Hamada¹, Yuta Hamada^{2,3}, Hayate Kimura³

¹DESY, ²KEK, ³SOKENDAI

KEK-PH2025, Feb.19, 2025

Short summary

Our work is ...

constructing black string solutions in the SM numerically.

Its existence is predicted by the swampland conjecture.

There is a no-go theorem prohibiting the existence of black strings in 4d theory ...

→ by considering **Casimir energy** in energy-momentum tensor, the no-go theorem can be avoided.

The black strings are intrinsically “quantum” object.

* based on arXiv:2501.05678 [hep-th]

Motivation

Motivation

- Many charged objects should be contained in **QG theory** in order to be coupled with the gravity in the consistent way.
- However, we do not know the full QG theory. The possible and reasonable argument is that we impose this conjecture on effective theories.
- This kind of argument predicts **\mathbb{Z}_2 charged string-like objects in SM.**

The horizon supported by Casimir energy

How is a black string horizon supported ?

It is known that **only spherical topology** is allowed as horizon topology in 4D theory at least classically.

Then, how we can construct a black string solution as vacuum solution ... ?

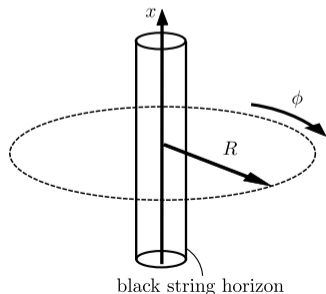
Actually, this theorem is valid under **dominant energy condition**.

The quantum correction for vacuum, **Casimir energy**, can violate the dominant energy condition !

Casimir energy

Casimir energy associated the S_1 surrounding the string is below.

$$V_{\text{Casimir}} = - \sum_{\text{particle}} (-1)^{2s_p n_p} \frac{m_p^4}{2\pi^2} \sum_{n=1}^{\infty} \frac{\cos(2\pi n \theta_p)}{(2\pi R m_p n)^2} K_2(2\pi R m_p n)$$



θ_p : corresponds to the boundary condition of the particle when going around the string.

Fermions have periodic boundary condition ($\theta_p = 0$)
instead of anti-periodic, because this BS is \mathbb{Z}_2 charged.
(corresponding flipping the sign of the fermion fields)

What we have to do

$$Rm_p \ll 1 \Rightarrow V_{\text{Casimir}} \sim (-1)^{2s_p} \cos(2\pi\theta_p) \frac{1}{(Rm_p)^4}.$$

$Rm_p \gg 1 \Rightarrow K_2(Rm_p)$ exponentially dump $\Rightarrow V_{\text{Casimir}}$ also dump.

In our low energy case, only **light particles** are dominant \rightarrow **graviton, photon, neutrino**
assumption : neutrinos are Majorana / Normal hierarchy / lightest mass is 0 ($m_3 = 0.05[\text{eV}]$)

What we have to do is ...

Solving Einstein eq. with Casimir energy-momentum tensor **with appropriate metric ansatz**
and boundary conditions.

Metric ansatz and boundary conditions

Extremal ansatz

The simplest string-like metric ansatz is

$$ds^2 = A^2(z)(-dt^2 + dx^2) + \left(\frac{M_P}{m_3^2}\right)^2 dz^2 + R(z)^2 d\phi^2.$$

z : parametrize the radial direction (proper distance). R is now a function of z .

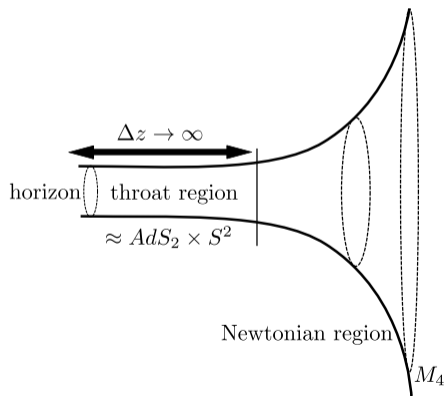
How about the boundary condition ?

Actually, the good candidate for such solution was suggested by Arkani-Hamed et. al. in 2007 (but they didn't construct it specifically).

They asserted that **the extremal charged black string interpolate 4D vacua.**

Reissner-Nordstrom BH interpolating 4D vacua (review)

Recall the extremal Reissner-Nordstrom BH : an exact solution of Einstein-Maxwell theory.



This BH interpolates between $AdS_2 \times S^2$ and M_4 .

What is the $AdS_2 \times S^2$?

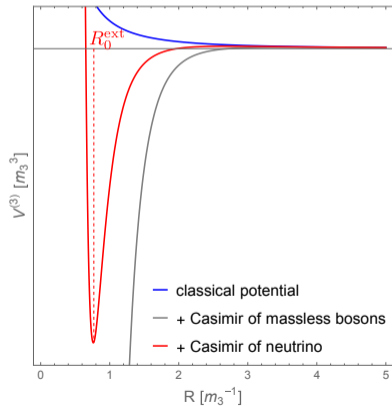
→ The vacuum solution of KK compactified
2-dim effective theory :

$$g_{\mu\nu}^{(4)} \xrightarrow{\text{decomposed}} g_{ij}^{(2)}, \chi \text{ (area of } S^2, \text{ moduli)}$$

Its moduli is stabilized by electric flux.

The black string interpolating 4D vacua (review)

They expected the same thing is true for extremal \mathbb{Z}_2 charged black string.



We consider 1-dim compactified 4d SM

: the spacetime is $N_3 \times S^1$.

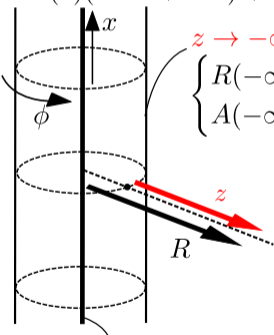
$$S = \int d^4x \sqrt{-g^{(4)}} \left(\frac{1}{2} M_P^2 \mathcal{R}^{(4)} - \Lambda^{(4)} - V_{\text{Casimir}}(R) \dots \right).$$

$$g_{\mu\nu}^{(4)} \xrightarrow{\text{decomposed}} g_{ij}^{(3)}, \quad R \text{ (radius of } S^1, \text{ moduli).}$$

Actually the Casimir energy (periodic neutrinos) in the action works as the radion potential.

In this case, $N_3 = AdS_3$.

Boundary conditions for extremal BS

$$ds^2 = A^2(z)(-dt^2 + dx^2) + \left(\frac{M_P}{m_3^2}\right)^2 dz^2 + R^2(z)d\phi^2$$


$$\begin{cases} z \rightarrow -\infty \\ R(-\infty) = R_0^{\text{ext}} \\ A(-\infty) \sim e^{(M_P/m_3^2)z/l_{\text{AdS}}} \rightarrow 0 \\ \text{: horizon, } AdS_3 \times S^1 \end{cases}$$

$$\begin{cases} z \rightarrow \infty \\ R(\infty) \propto \frac{M_P}{m_3^2} z \rightarrow \infty \\ A(\infty) = \text{const.} \end{cases}$$

$R = 0$: conical singularity : infinite distance, M_4

Expected boundary conditions :

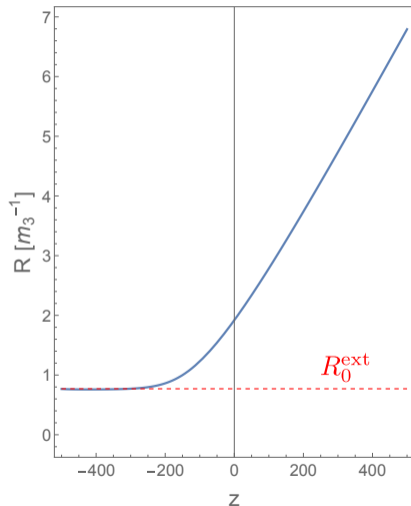
- $AdS_3 \times S^1$ near horizon
- M_4 at infinity.

infinite throat region of extremal
RN-BH

\Leftrightarrow the horizon is put at $z = -\infty$

Numerical Results

Extremal solution : R



$$ds^2 = A^2(z)(-dt^2 + dx^2) + \left(\frac{M_P}{m_3^2}\right)^2 dz^2 + R(z)^2 d\phi^2.$$

- $z \rightarrow -\infty$ (horizon) : $R \rightarrow R_0^{\text{ext}} (S^1)$.
- $z \rightarrow \infty$ (infinity) : $R \sim (M_P/m_3^2)z/\sqrt{a}$ (flat).

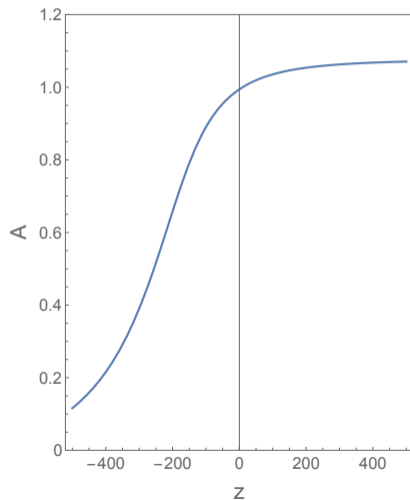
Especially,

$$\frac{1}{\sqrt{a}} = \mathcal{O}(10^{-2}) \times \frac{m_3}{M_P}$$

This means the deficit angle $\theta = 2\pi(1 - 1/\sqrt{a})$ is

$$\theta \sim \mathcal{O}(1).$$

Extremal solution : A



$$ds^2 = A^2(z)(-dt^2 + dx^2) + \left(\frac{M_P}{m_3^2}\right)^2 dz^2 + R(z)^2 d\phi^2.$$

- $z \rightarrow -\infty$ (horizon)
: $A \sim e^{(M_P/m_3^2)z/l_{\text{AdS}}} (AdS_3) \rightarrow 0.$

where

$$l_{\text{AdS}} = 6.2 \times 10^{-3} \times \frac{M_P}{m_3^2}$$

- $z \rightarrow \infty$ (infinity) : $A \rightarrow \text{const}$ (flat).

Conclusion

Conclusion

- The cobordism conjecture predicts a \mathbb{Z}_2 charged string-like object in 4D SM.
- We could actually construct this object as the “quantum” black string in SM.

Future direction

- checking the stability of the string
- cosmological construction ?
- $E_8 \times E_8$ heterotic SUGRA