

A New Mechanism for Leptogenesis in three Higgs doublet model

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Based on K. Mukaida, HW, M. Yamada, JCAP 09 (2024) 063 [arXiv:2405.14332].

Introduction

From Cosmic Microwave Background and Big Bang Nucleosynthesis

$$\frac{q_B}{s} := \frac{n_B - n_{\bar{B}}}{s} = 8.6 \times 10^{-11}$$

n_B : baryon number density
 $n_{\bar{B}}$: anti-baryon number density
 s : entropy density

We need this baryon asymmetry.



Is there possibility that baryon asymmetry existed in the beginning?

Because of inflation and reheating, pre-existing baryon asymmetry is diluted.

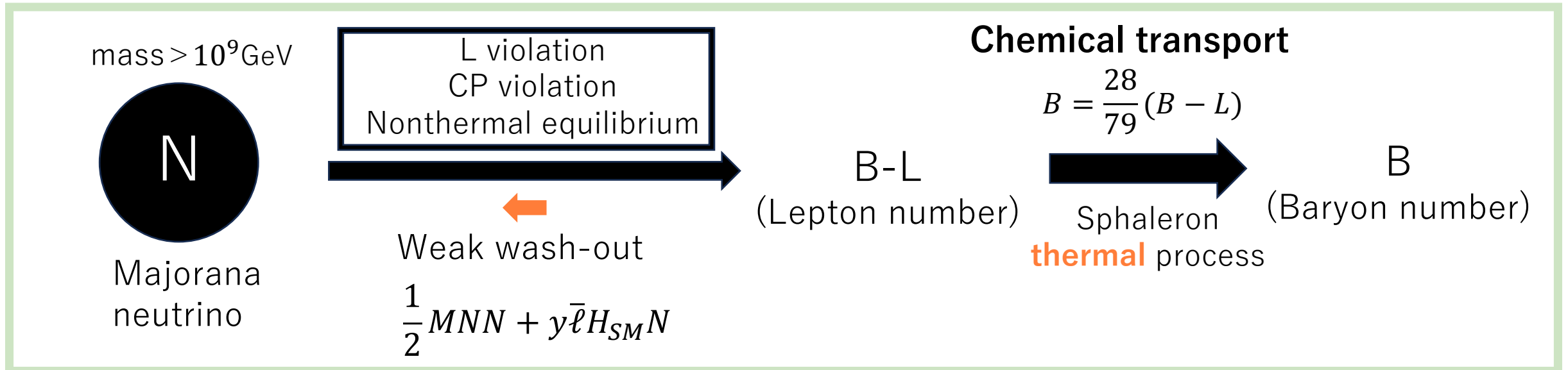


The problem of baryon asymmetry is inevitable
Beyond Standard Model is necessary.

Leptogenesis

Vanilla leptogenesis (Thermal leptogenesis)

M. Fukugita, T. Yanagida Phys.Lett.B 174 (1986) 45-47

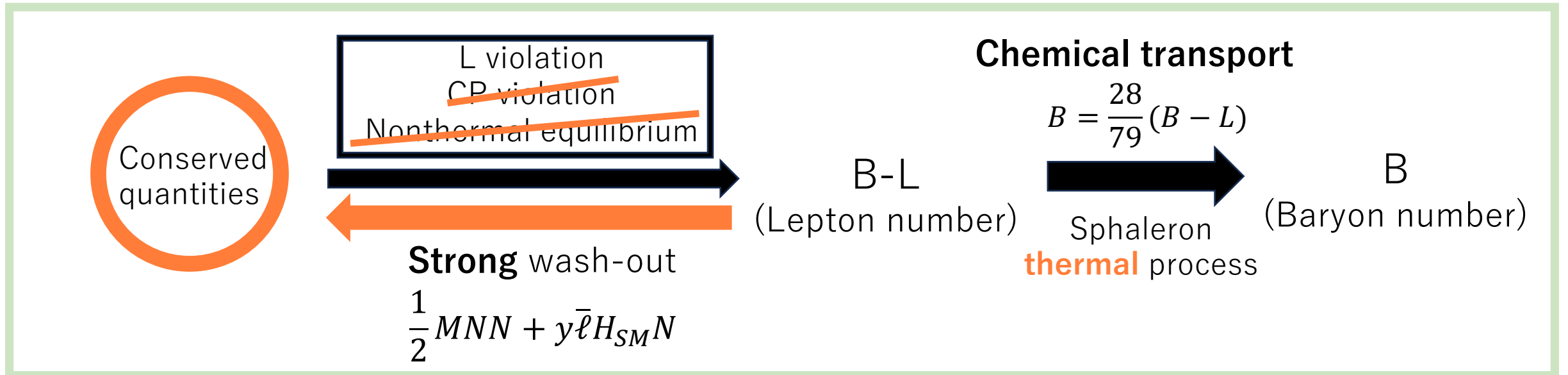


Conditions to produce the enough baryon asymmetry

- Weak wash-out and nonthermal equilibrium
- CP violation in the neutrino sector
- Large Majorana mass
- L violation when non-equilibrium decay

Our work

Wash-in leptogenesis + Three Higgs doublet model



~~Conditions to produce the enough baryon asymmetry~~

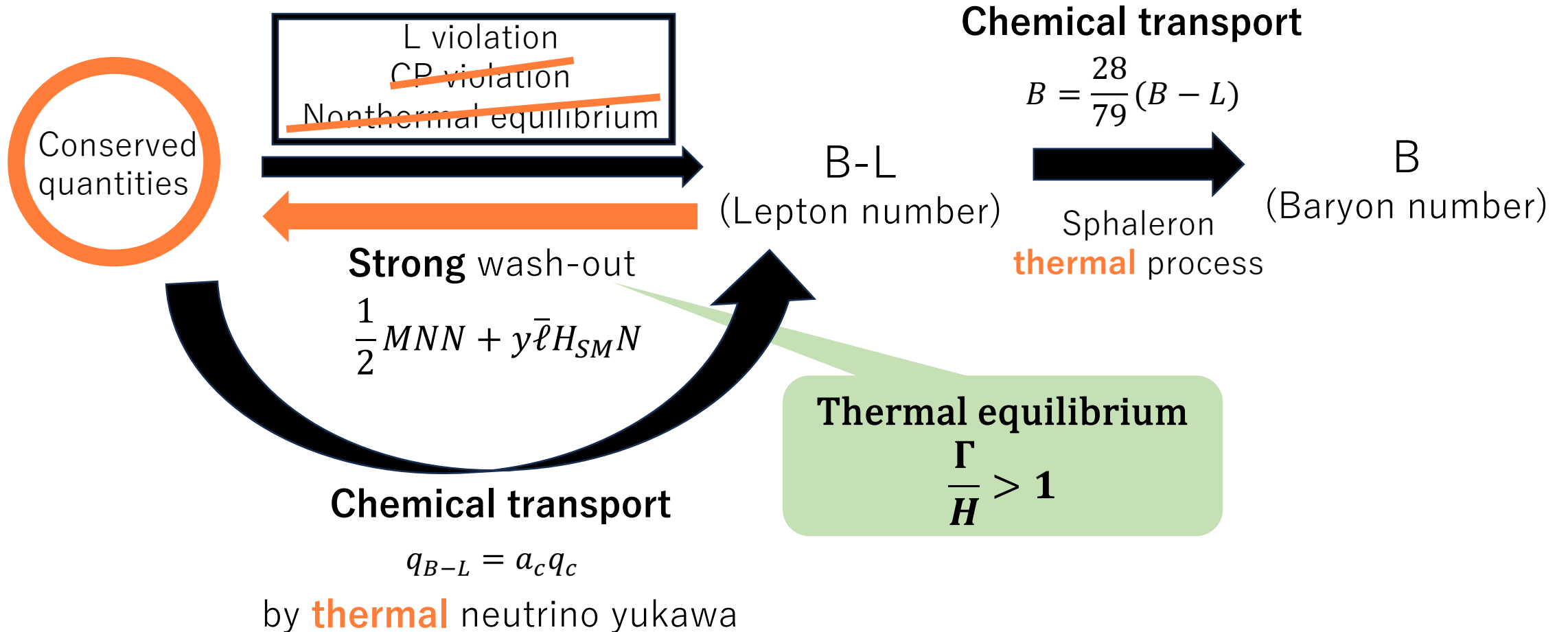
- ~~• Weak wash-out and nonthermal equilibrium~~
- ~~• CP violation in the neutrino sector~~
- ~~• Large Majorana mass~~
- ~~• L violation when non-equilibrium decay~~

These conditions are not necessary.

Wash-in Leptogenesis

V. Domcke *et al.*, Phys. Rev. Lett. 126 (2021) 201802.

Wash-in leptogenesis



Decoupling and Conserved Quantities

Interaction rate Γ
and Hubble parameter H

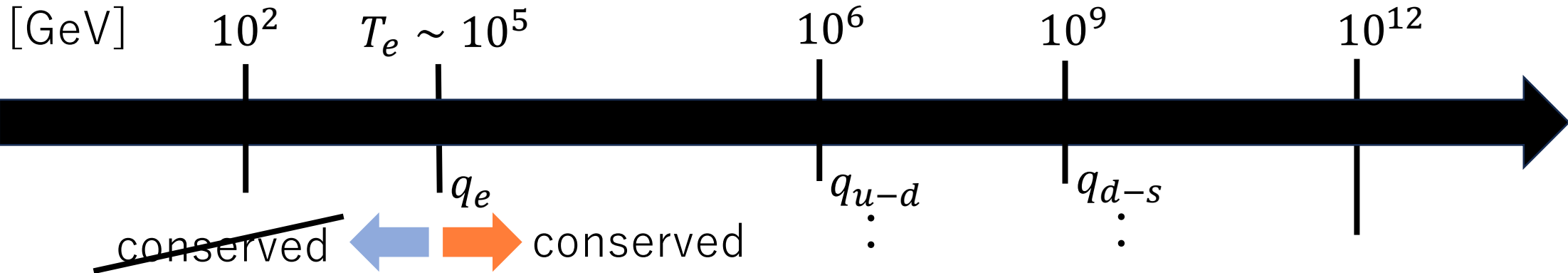
$\frac{\Gamma}{H} < 1$ **→** Conserved quantities

decoupling

Γ and H when all particles are relativistic

$\Gamma \propto g^2 T$ $H \propto T^2$ g : coupling constant

$\frac{\Gamma}{H} \propto \frac{g^2}{T}$ **→** weaker coupling } easier decoupling
 higher temperature }



The right-handed electron asymmetry q_e is conserved until lowest temperature.

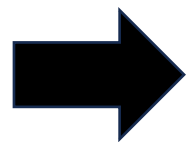
q_e can be a source of $B - L$ $q_{B-L} = -\frac{3}{10} q_e$ **→** The rest problem is how to produce q_e

UV model to realize Wash-in Leptogenesis

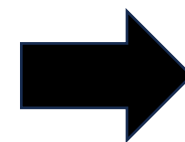
- GUT baryogenesis [V. Domcke *et al.*, Phys. Rev. Lett. 126 (2021) 201802.]
- Axion Inflation [V. Domcke, *et al.*, 2210.06412]
- Evaporation of PBH [K. Schmitz, *et al.*, 2311.01089]

etc

These models can violate charges other than the required conserved quantities.



These violations are **not** necessary.



3HDM

Three-Higgs doublet model

V. Keus, S. F. King, and S. Moretti, *JHEP* **01** (2014) 052, arXiv:1310.8253 [hep-ph].

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \mathcal{L}_{\text{RH}\nu} + \mathcal{L}_{\text{HH}}$$

$$\mathcal{L}_{\text{RH}\nu} = \bar{N}_i i \not{\partial} N_i - \frac{1}{2} M_i \bar{N}_i N_i - \left(\bar{\ell}_f \tilde{H}_{\text{SM}} \lambda_{fi}^{\text{SM}} N_i + \text{H.c.} \right)$$

The thermal equilibrium process which produces q_{B-L}

$$\mathcal{L}_{\text{HH}} = |DH_\alpha|^2 - M_{H_\alpha}^2 |H_\alpha|^2 - \left(\bar{\ell}_f H_\alpha Y_{ff'}^\alpha e_{f'} + \bar{\ell}_f \tilde{H}_\alpha \lambda_{fi}^\alpha N_i + \text{H.c.} \right)$$

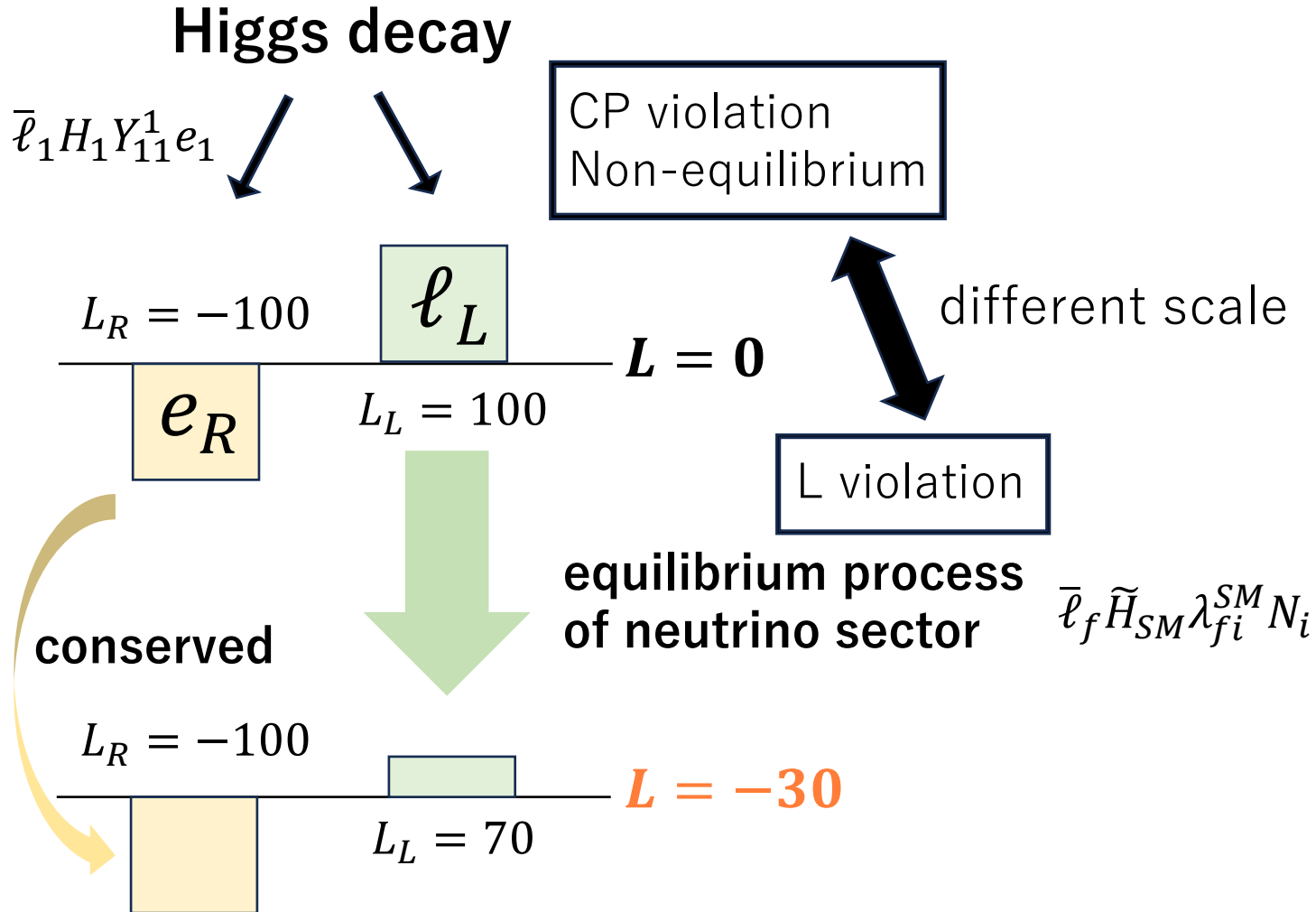
decay of Higgs to produce q_e

- violate CP
- not violate L

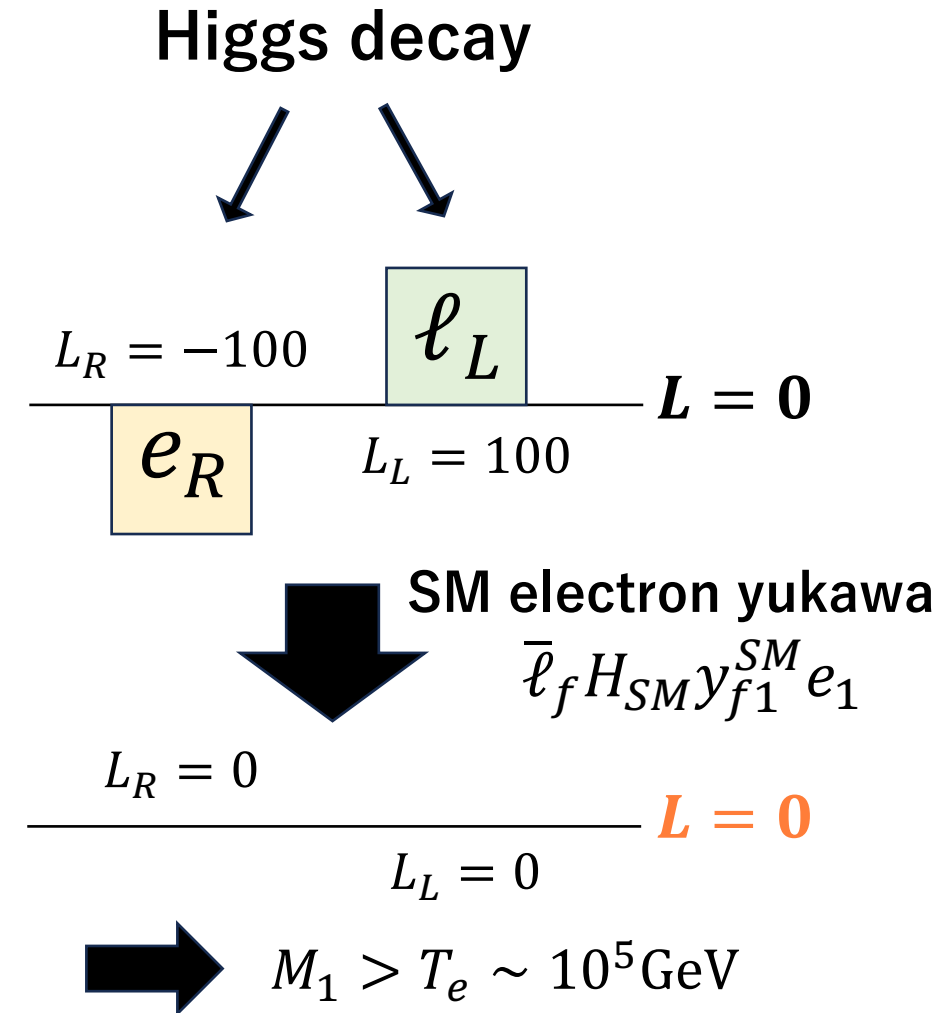
N_i	: Majorana neutrinos
M_i	: Mass of Majorana
ℓ_f	: lepton doublet
H_{SM}	: SM Higgs doublet
$H_\alpha (\alpha = 1, 2)$: new Higgs doublet
M_{H_α}	: Mass of H_α
$\lambda^{\text{SM}}, Y, \lambda$: coupling constant

The flow of this model to produce B

If q_e is conserved,



If q_e is not conserved,



result

Mass of each particle in this scenario

Majorana Mass

$$M_1 \sim 10^5 \text{ GeV}$$

Higgs Mass

$$M_{H_1} \sim 10^{12} \text{ GeV}$$

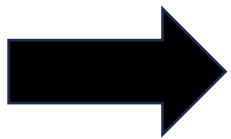
Non-thermal equilibrium condition of Higgs decay

Final baryon asymmetry

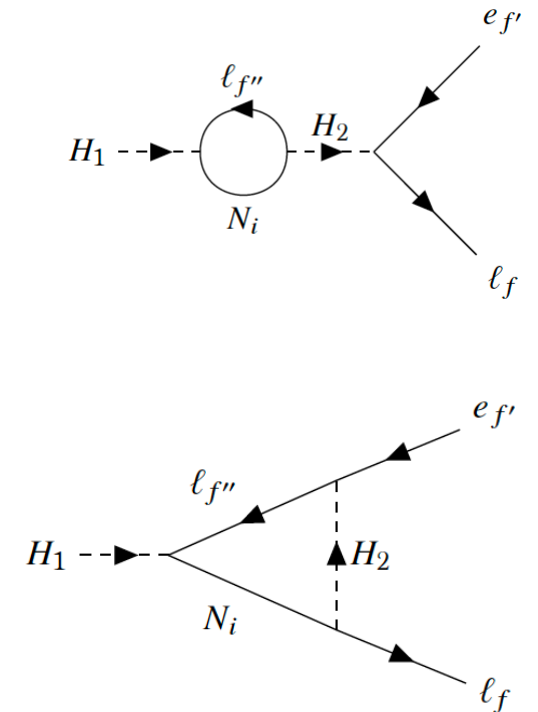
observation : $\frac{q_B}{s} \sim 8.6 \times 10^{-11}$

$$\frac{q_B}{s} \sim 8 \times 10^{-11} \delta_{CP} \left(\frac{Y}{0.01} \right)^2$$

δ_{CP} : CP phase
 Y : yukawa coupling



We probe baryon asymmetry in 3HDM



Summary and Future Work

Summary

- In vanilla leptogenesis, some conditions such as “weak wash-out” are needed to produce enough amount of baryon asymmetry. However, we probe it **without** such conditions by using wash-in leptogenesis and three Higgs doublet model.
- In our scenario, Majorana mass and Higgs mass are 10^5GeV and 10^{12}GeV to produce enough baryon asymmetry.

According to the recent research [K. Schmitz, *et al.*, (2501.07634)],
 M_1 can be as small as 7TeV

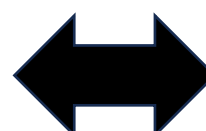
Future Work

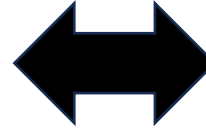
- How about the cases of other conserved quantities?(Higgs can be coupled to quarks.)
- Quantitative analysis by using Boltzmann equation

Back up

The condition to determine M_{H_1}

Suppose weak wash-out to produce enough q_e

 $1 \geq \frac{\Gamma}{H} \Big|_{T=M_{H_1}}$

 $M_{H_1} \geq 10^{12} \text{GeV} \left(\frac{Y}{0.01} \right)^2$

Decay rate

$$\Gamma \sim \frac{1}{16\pi} \text{tr}(Y^{1\dagger} Y^1) M_{H_1}$$

Hubble parameter

$$H = \sqrt{\frac{\pi^2}{90} g_*} \frac{T^2}{M_{pl}} \quad (\text{RD})$$

Relation between B and B-L

Conserved quantities

$$q_Y = 0 \quad \longleftrightarrow \quad \sum(\mu_{qi} + 2\mu_{ui} - \mu_{di} - \mu_{li} - \mu_{ei}) + 2\mu_H = 0$$

$$\frac{B}{3} - L_f \quad f = e, \mu, \tau \quad \longleftrightarrow \quad \frac{1}{3} \sum(2\mu_{qi} + \mu_{ui} + \mu_{di}) - 2\mu_{lf} - \mu_f$$

Equilibrium process

$$\text{strong sphaleron} \quad \longleftrightarrow \quad \sum(2\mu_{qi} - \mu_{ui} - \mu_{di}) = 0$$

$$\text{Weak sphaleron} \quad \longleftrightarrow \quad \sum(3\mu_{qi} + \mu_{li}) = 0$$

$$\text{Up yukawa} \quad \longleftrightarrow \quad \mu_{qi} + \mu_H - \mu_{uj} = 0$$

$$\text{Down yukawa} \quad \longleftrightarrow \quad \mu_{qi} - \mu_H - \mu_{dj} = 0$$

$$\text{Charged lepton yukawa} \quad \longleftrightarrow \quad \mu_{li} - \mu_H - \mu_{ej} = 0$$

$$\longrightarrow \quad q_B = \frac{28}{79} q_{B-L}$$

$$q_X = \frac{g}{6} T^2 \begin{cases} \mu_X/T & \text{fermions} \\ 2\mu_X/T & \text{bosons} \end{cases}$$

Relation between B-L and q_e

$10^5 \text{ GeV} < T < 10^6 \text{ GeV}$

Conserved quantity

$$q_Y = 0$$



$$\sum(\mu_{qi} + 2\mu_{ui} - \mu_{di} - \mu_{li} - \mu_{ei}) + 2\mu_H = 0$$

$$q_{e1}$$



$$\mu_{e1}$$

Equilibrium process

strong sphaleron



$$\sum(2\mu_{qi} - \mu_{ui} - \mu_{di}) = 0$$

Weak sphaleron



$$\sum(3\mu_{qi} + \mu_{li}) = 0$$

Up yukawa



$$\mu_{qi} + \mu_H - \mu_{uj} = 0$$

Down yukawa



$$\mu_{qi} - \mu_H - \mu_{dj} = 0$$

Charged yukawa

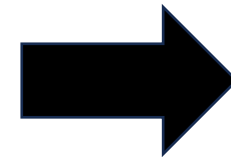


$$\mu_{li} - \mu_H - \mu_{ej} = 0$$

Neutrino yukawa



$$\mu_{li} - \mu_H = 0 \quad \mu_N = 0$$



$$q_{B-L} = -\frac{3}{10} q_{e1}$$

The amount of right-handed electron asymmetry

$$\frac{q_B}{s} \stackrel{=}{=} \frac{28}{79} \frac{q_{B-L}}{s} \stackrel{=}{=} -\frac{42}{395} \frac{q_e}{s}$$

Sphaleron equilibrium

$$q_B = \frac{28}{79} q_{B-L}$$

Neutrino yukawa equilibrium

$$q_{B-L} = -\frac{3}{10} q_e$$

$$\frac{q_B}{s} \sim 8.6 \times 10^{-11}$$



$$\frac{q_e}{s} \simeq -8.1 \times 10^{-10}$$

This is the value to explain today's baryon asymmetry ! !

The production of q_e in 3HDM

$$\begin{aligned} \frac{q_e}{s} &:= \frac{n_e}{s} - \frac{n_{\bar{e}}}{s} = \frac{\Gamma_{\bar{H}_1 \rightarrow e\bar{\ell}} n_{\bar{H}_1}}{\Gamma_{H_1}} - \frac{\Gamma_{H_1 \rightarrow e\bar{\ell}} n_{H_1}}{\Gamma_{H_1}} \\ &= -\frac{\Gamma_{H_1 \rightarrow e\bar{\ell}} - \Gamma_{\bar{H}_1 \rightarrow e\bar{\ell}} n_{H_1}}{\Gamma_{H_1}} \\ &\simeq -\epsilon_{\bar{e}\ell} \times \frac{45\zeta(3)}{\pi^4} \frac{1}{g_*} \end{aligned}$$

Assumption : Thermal production of H_1
 $(T_{rh} > M_{H_1})$

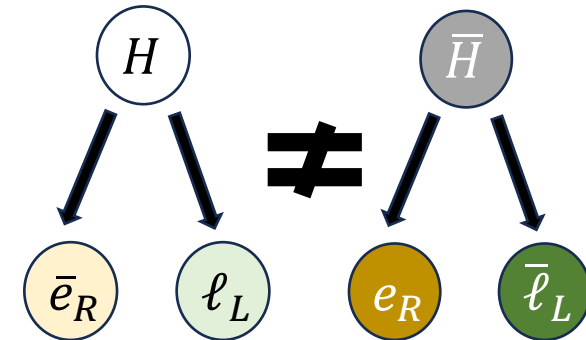
$$\frac{n_{H_1}}{s} = \frac{n_{\bar{H}_1}}{s} = \frac{45\zeta(3)}{\pi^4} \frac{1}{g_*}$$

g_* : effective degrees of freedom

$$\Gamma_{H_1} := \Gamma_{H_1 \rightarrow e\bar{\ell}} + \Gamma_{H_1 \rightarrow \ell N}$$

$\epsilon_{\bar{e}\ell}$: asymmetry parameter (CP violation)

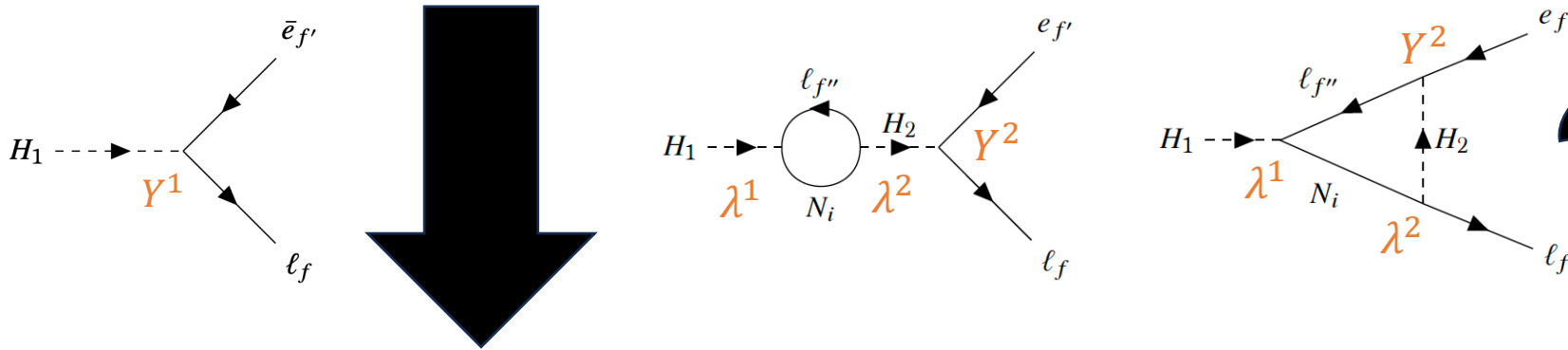
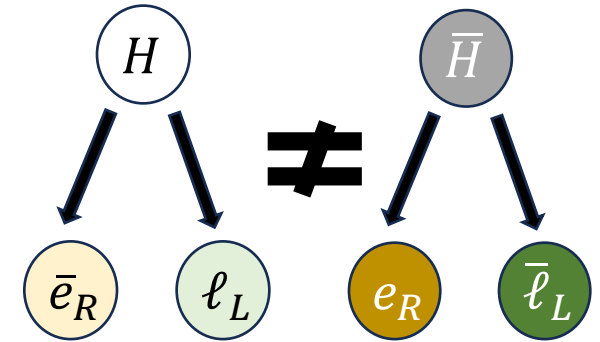
$$\epsilon_{\bar{e}\ell} := \frac{\Gamma_{H_1 \rightarrow e\bar{\ell}} - \Gamma_{\bar{H}_1 \rightarrow e\bar{\ell}}}{\Gamma_{H_1 \rightarrow e\bar{\ell}} + \Gamma_{\bar{H}_1 \rightarrow e\bar{\ell}}}$$



Asymmetry parameter

$\epsilon_{\bar{e}\ell}$: asymmetry parameter (CP violation)

$$\epsilon_{\bar{e}\ell} := \frac{\Gamma_{H_1 \rightarrow \bar{e}\ell} - \Gamma_{\bar{H}_1 \rightarrow e\bar{\ell}}}{\Gamma_{H_1 \rightarrow \bar{e}\ell} + \Gamma_{\bar{H}_1 \rightarrow e\bar{\ell}}}$$



We need to calculate one-loop to get ϵ

Asymmetry parameter comes from interference

between tree and one-loop

$$\epsilon_{\bar{e}\ell} \propto \frac{\text{Im} \left(Y_{11}^1 Y_{11}^{2\dagger} \lambda_{11}^1 \lambda_{11}^{2\dagger} \right)}{|Y^1|^2}$$

$\text{Im}(\dots)$ \longleftrightarrow CP violation

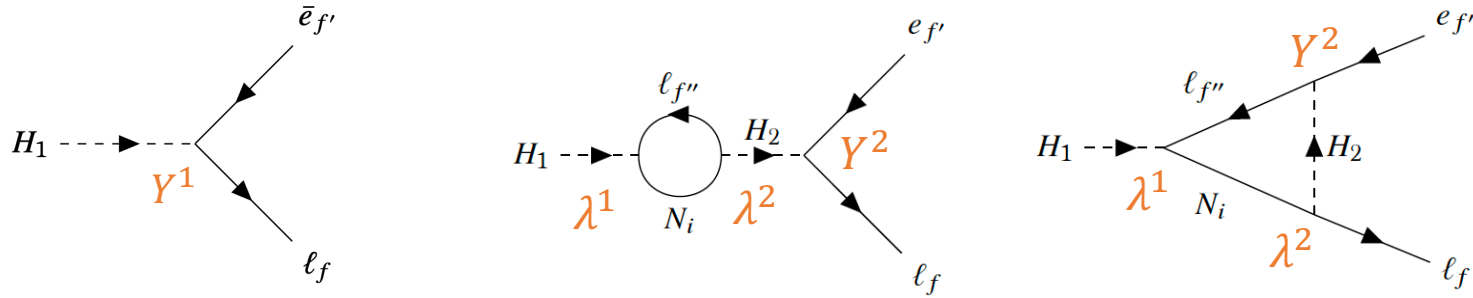
$$\Gamma_{H_1 \rightarrow \bar{e}\ell}^{(0)} \sim |Y^1|^2 M_{H_1}$$



CP transformation

$$\Gamma_{\bar{H}_1 \rightarrow e\bar{\ell}}^{(0)} \sim |Y^{1*}|^2 M_{H_1}$$

Asymmetry parameter



$$\epsilon_{\bar{e}\ell} = \frac{1}{8\pi} \frac{\text{Im}(Y_{11}^1 Y_{11}^{2\dagger} \lambda_{11}^1 \lambda_{11}^{2\dagger})}{\text{tr}(Y^{1\dagger} Y^1)} \left(\underbrace{\frac{1}{x^2 - 1}}_{\text{self}} + \underbrace{1 - x^2 \log\left(1 + \frac{1}{x^2}\right)}_{\text{vertex}} \right)$$

$$x = M_{H_2}/M_{H_1}$$