



International Center for
Quantum-field Measurement Systems for
Studies of the Universe and Particles
WPI research center at KEK



Two-sided story of sterile neutrinos: Production under the X-ray limit

Muping Chen (QUP, KEK)

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Based on [M. Chen, G. B. Gelmini, P. Lu, and V. Takhistov, Phys.Lett.B 852
(2024) 138609, arXiv: 2309.12258 & JCAP 07 (2024) 059, arXiv:
2312.12136]





Introduction



Sterile neutrino

Sterile neutrinos with no weak interactions (right-handed) can be added to the Standard Model (SM).

- Motivation: massive neutrinos, minimal extension.

- Mixing:

$$|\nu_a\rangle = \cos\theta|\nu_1\rangle + \sin\theta|\nu_2\rangle;$$

$$|\nu_s\rangle = -\sin\theta|\nu_1\rangle + \cos\theta|\nu_2\rangle;$$

- Sterile neutrinos can be created via mixing dependent mechanisms such as active-sterile oscillation.
- Detecting through X-rays due to $\nu_s \rightarrow \nu\gamma$



Production Mechanism

- Doldson & Widrow [S. Dodelson and L. M. Widrow, Phys.Rev.Lett. 72 \(1994\) 17-20 \[hep-ph/9303287\]](#)

- Maximum production temperature: $T_{\max} \simeq 133 \text{ MeV} \left(\frac{m_s}{\text{keV}} \right)^{1/3}$

- Density fraction

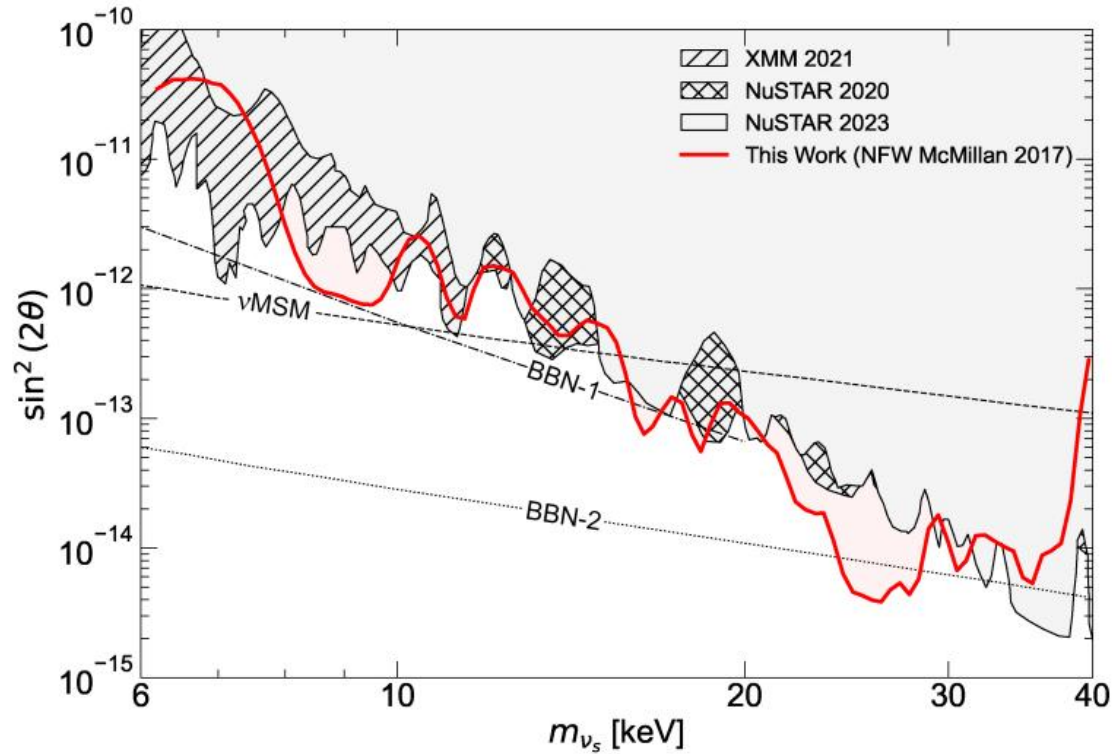
$$f_{s,\text{osc}} = 2.7 \times 10^{-4} \left(\frac{\sin^2 2\theta}{10^{-10}} \right) \left(\frac{m_s}{\text{keV}} \right)^2 \left(\frac{30}{g_*(T_{\max})} \right)^{3/2}$$

$$f_{s,\text{osc}}^{\text{LRT}} \simeq 1 \times 10^{-7} \left(\frac{\sin^2 2\theta}{10^{-10}} \right) \left(\frac{m_s}{\text{keV}} \right) \left(\frac{T_{RH}}{5 \text{ MeV}} \right)^3 \quad \text{if } T_{RH} < T_{\max}$$

- Many other production mechanisms



X-ray Decays



R. A. Krivonos, V. V. Barinov,² Mukhin, and D. S. Gorbunov,
Phys.Rev.Lett. 133 (2024) 26, 261002 [2405.17861]

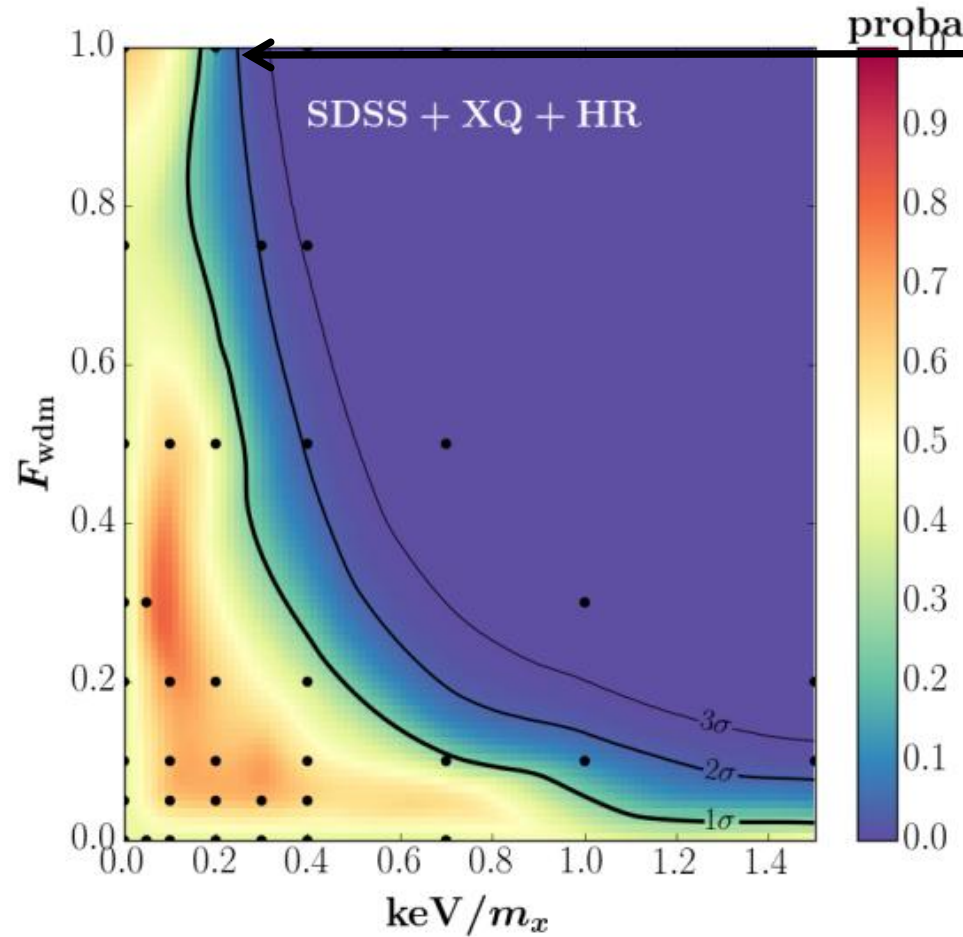
Decay rate:

$$\Gamma = 1.38 \times 10^{-32} \text{s}^{-1} \left(\frac{\sin^2 2\theta}{10^{-10}} \right) \left(\frac{m_\chi}{\text{keV}} \right)^5$$

- Constrain $\sin^2 2\theta$ for $f_s = 1$
- Constrain $f_s \sin^2 2\theta$ for $f_s < 1$



Lyman-alpha bound



DM fermion disfavored

$$m_{\text{therm}} \leq 3\text{keV}$$

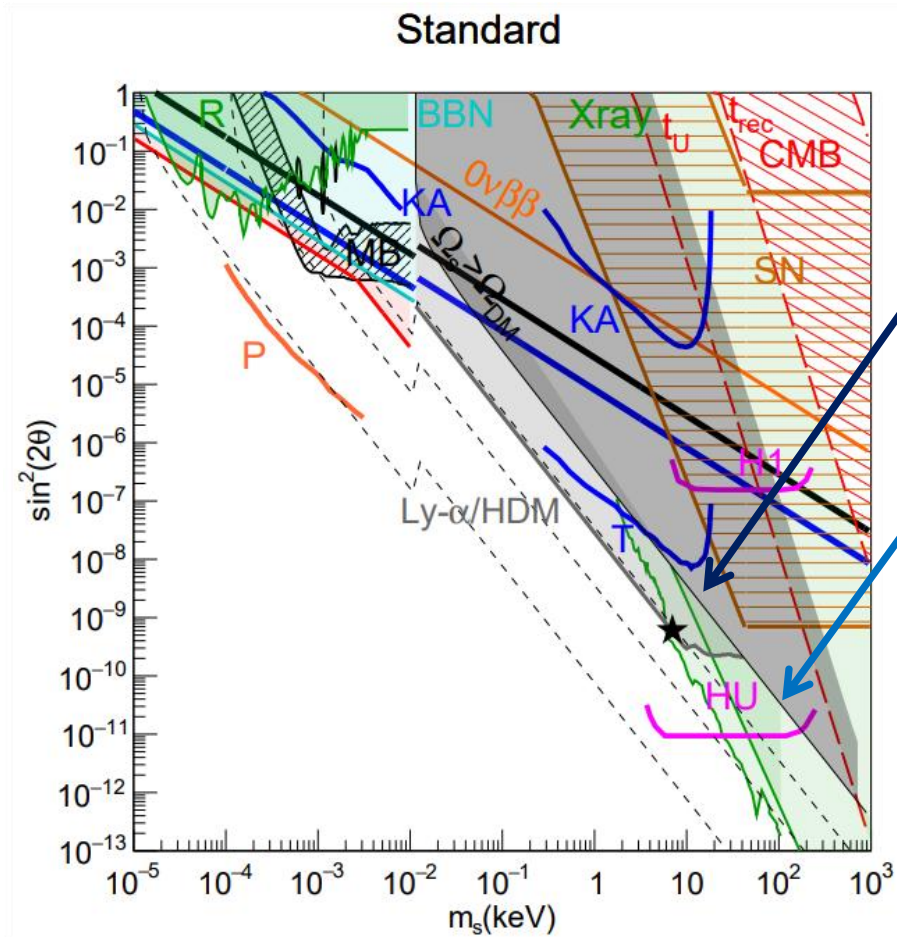
This bound can be extended to sterile neutrinos

$$m_s \simeq 4\text{keV} \left(\frac{\langle \epsilon \rangle}{3.15} \right) \left(\frac{m_{\text{therm}}}{\text{keV}} \right)^{4/3} \left(\frac{10.75}{g_*} \right)^{1/3}$$

J. Baur et al., JCAP 12 (2017) 013 [1706.03118]



Motivation



Too hot to make up the whole of DM

Excluded by X-ray limit

Possible new production mechanism less restricted by neutrino interactions?

G. B. Gelmini, P. Lu, and V. Takhistov, JCAP 12 (2019) 047 [1909.13328]



PBH Neutrinogenesis



PBH Evaporation

PBHs emit Hawking radiation with an approximately blackbody spectrum

- PBHs with mass M_{PBH} and density fraction f_{evap}
 - Radiation domination(RD) if $f_{\text{evap}} < 1$
 - PBH domination(PD) if $f_{\text{evap}} = 1$

- Hawking temperature

$$T_{\text{PBH}} = \frac{M_{\text{Pl}}^2}{8\pi M_{\text{PBH}}} = 1.06 \times 10^5 \text{ GeV} \left(\frac{10^8 \text{ g}}{M_{\text{PBH}}} \right)$$

- Evaporation temperature:

$$T_{\text{evap}} \simeq 43 \text{ MeV} \left(\frac{10^8 \text{ g}}{M_{\text{PBH}}} \right)^{3/2} \left(\frac{10.75}{g_*(T_{\text{evap}})} \right)^{1/4} \left(\frac{g_H}{110} \right)^{1/2}$$



PBH neutrino genesis

Current sterile neutrino density fraction due to PBH:

$$f_s = \frac{\rho_s(T_0)}{\rho_{\text{DM}}} \simeq 2 \times 10^{-6} f_{\text{evap}} \left(\frac{m_s}{\text{keV}} \right) \left(\frac{10^8 \text{g}}{M_{\text{PBH}}} \right)^{1/2} \left(\frac{10.75}{g_*(T_{\text{evap}})} \right)^{1/4}$$
$$\simeq 1 \times 10^{-6} f_{\text{evap}} \left(\frac{m_s}{\text{keV}} \right) \left(\frac{T_{\text{evap}}}{5 \text{ MeV}} \right)^{1/3} \left(\frac{10.75}{g_*(T_{\text{evap}})} \right)^{1/6} .$$

Does not depend on the mixing angle



New parameter space $(m_s, T_{\text{evap}}(M_{\text{PBH}}))$



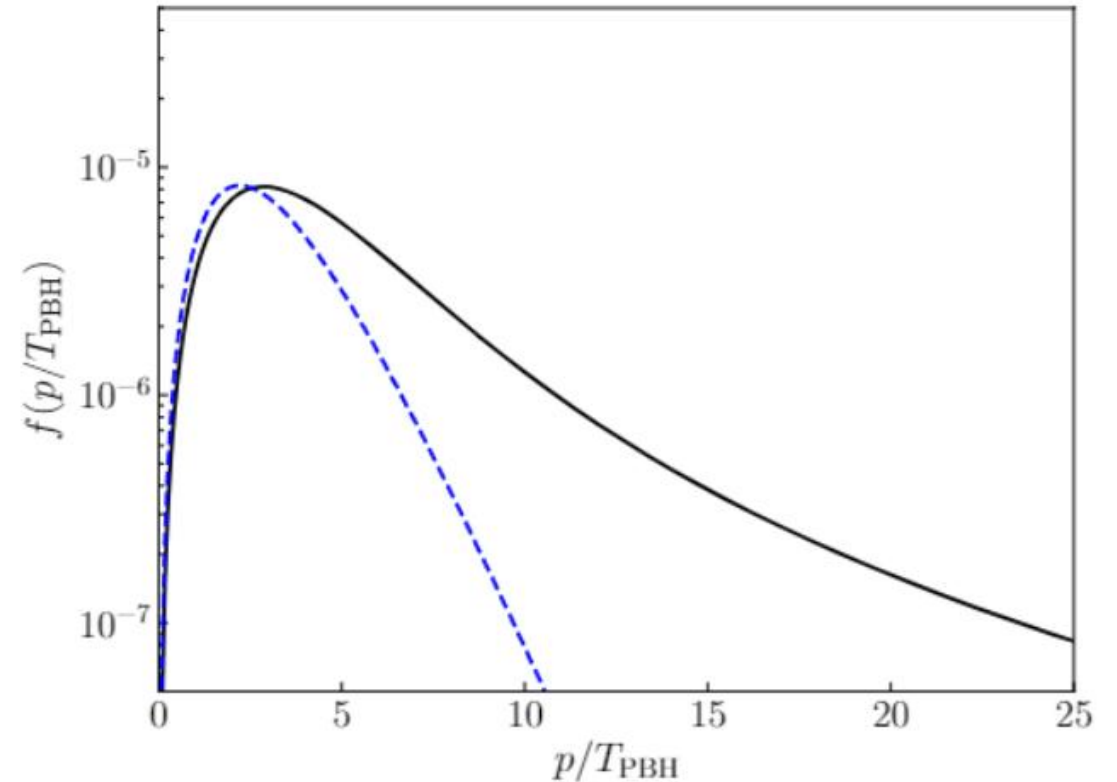
Relic momentum

For spin-1/2 particles, peaks at $E = 4.5T_{\text{PBH}}$

Average scaled momentum at evaporation

$$\begin{aligned} \langle \epsilon \rangle &= \frac{\langle p \rangle}{T_{\text{evap}}} \simeq \frac{6.3 T_{\text{PBH}}}{T_{\text{evap}}} \\ &\simeq 1.5 \times 10^7 \left(\frac{M_{\text{PBH}}}{10^8 \text{g}} \right)^{1/2} \left(\frac{g_*(T_{\text{evap}})}{10.75} \right)^{1/4} \end{aligned}$$

Much larger than $\langle \epsilon \rangle_{\text{FD}} = 3.15$ from a fermi-dirac spectrum.



M. Chen, G. B. Gelmini, P. Lu, and V. Takhistov, JCAP 07 (2024) 059 [2312.12136]



Additional Constraint

- Condition: $f_s > f_{s,\text{osc}}$
- Bound on the mixing:

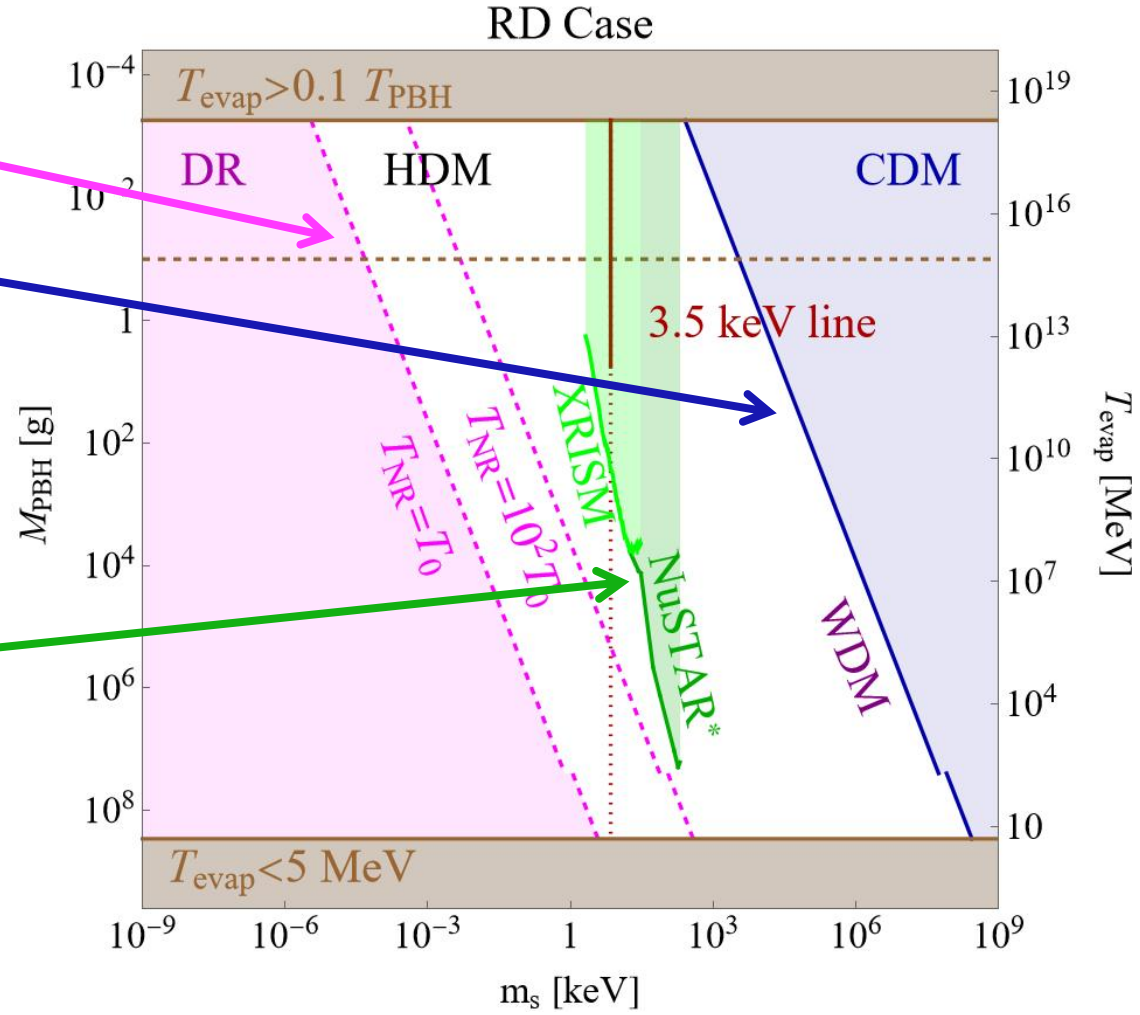
$$\sin^2(2\theta) < 4 \times 10^{-13} f_{\text{evap}} \left(\frac{\text{keV}}{m_s} \right) \left(\frac{10.75}{g_*(T_{\text{evap}})} \right)^{1/6} \left(\frac{g_*(T_{\text{max}})}{30} \right)^{3/2} \left(\frac{T_{\text{evap}}}{5 \text{ MeV}} \right)^{1/3}$$

$$\sin^2(2\theta) < 1 \times 10^{-9} \left(\frac{10.75}{g_*(T_{\text{RH}})} \right)^{1/6} \left(\frac{5 \text{ MeV}}{T_{\text{RH}}} \right)^{8/3} \quad T_{\text{RH}} < T_{\text{max}}$$



Regions and limits-RD

- Nonrelativistic Limit
- Lyman-alpha limit
 - WDM
 - $f_{\text{evap}} \simeq 4 \times 10^{-3} \left(\frac{g_*(T_{\text{evap}})}{10.75} \right)^{1/3}$
 - Above the bound
 - $f_s = 1$
- X-ray signals
 - bounds
 - $f_s > f_{s,osc}$

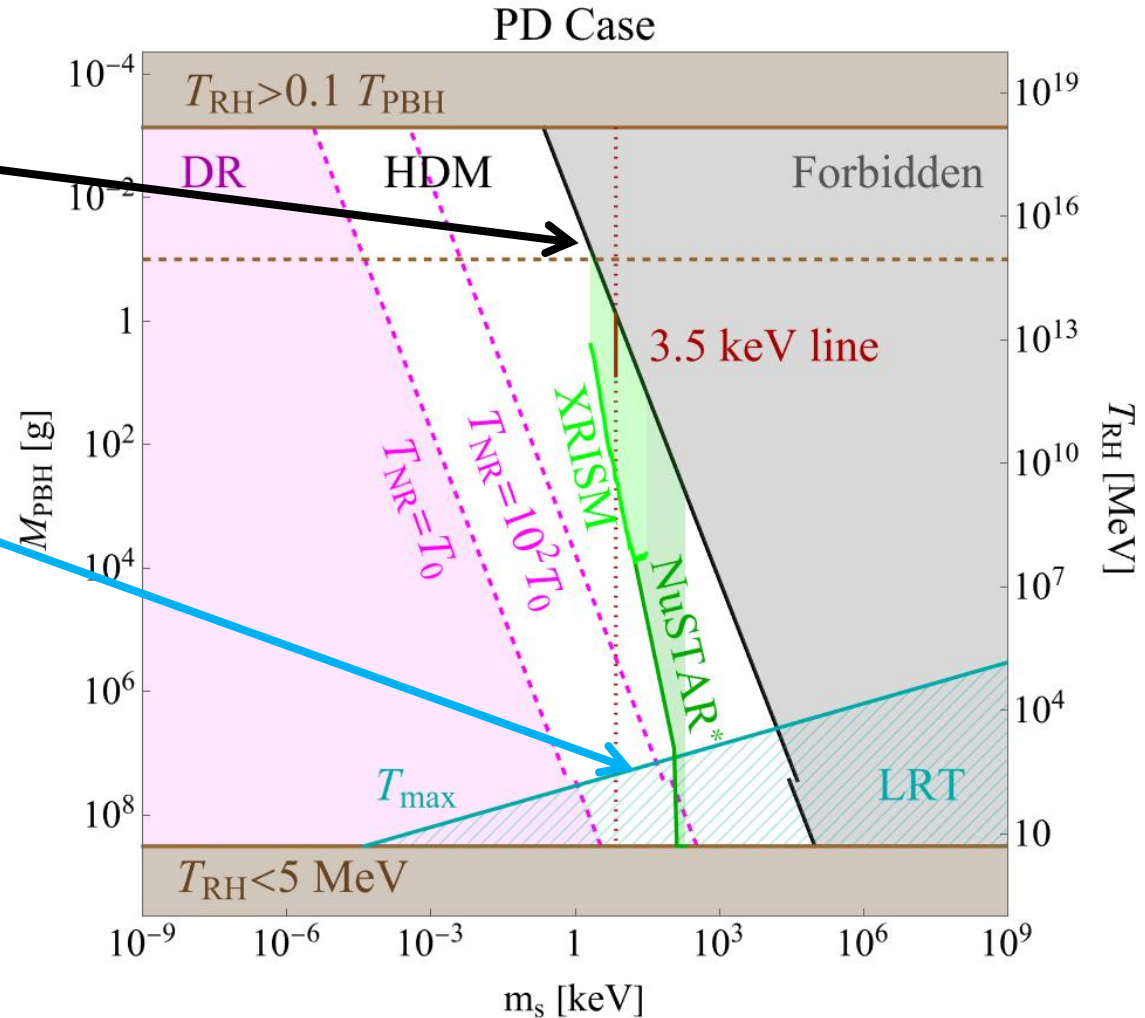


M. Chen, G. B. Gelmini, P. Lu, and V. Takhistov, JCAP 07 (2024) 059 [2312.12136]



Regions and limits-PD

- Forbidden region
 - HDM condition $f_s < 0.1$
- Low Reheating Region



M. Chen, G. B. Gelmini, P. Lu, and V. Takhistov, JCAP 07 (2024) 059 [2312.12136]

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Gravitational Signal

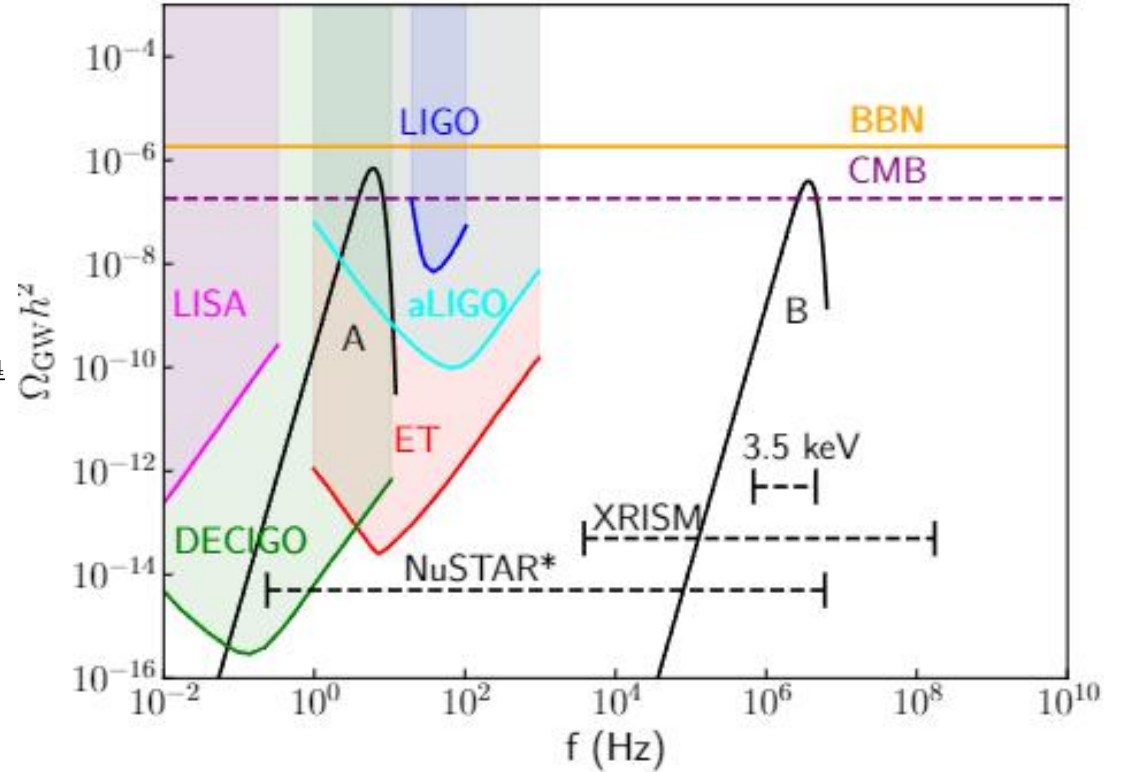
- Induced GWs from PBH evaporation
- UV cut-off frequency:

$$f_{UV} = 1.7 \times 10^3 \text{ Hz} \left(\frac{M_{PBH}}{10^4 \text{g}} \right)^{-5/6}$$

- Present density of GWs at the peak

$$\Omega_{GW,0}^{\text{peak}} h^2 = 1.64 \times 10^{-6} \left(\frac{\gamma}{0.2} \right)^{7/9} \left(\frac{\beta}{10^{-8}} \right)^{16/3} \left(\frac{M_{PBH}}{10^7 \text{g}} \right)^{34/9}$$

- Very optimistic examples:
 - A: $M_{PBH} = 2 \times 10^7 \text{ g}$, $\beta = 6 \times 10^{-9}$
 - B: $M_{PBH} = 1 \text{ g}$, $\beta = 8 \times 10^{-4}$



M. Chen, G. B. Gelmini, P. Lu, and V. Takhistov, JCAP 07 (2024) 059 [2312.12136]

Novel Coincident Signal: X-rays+GWs

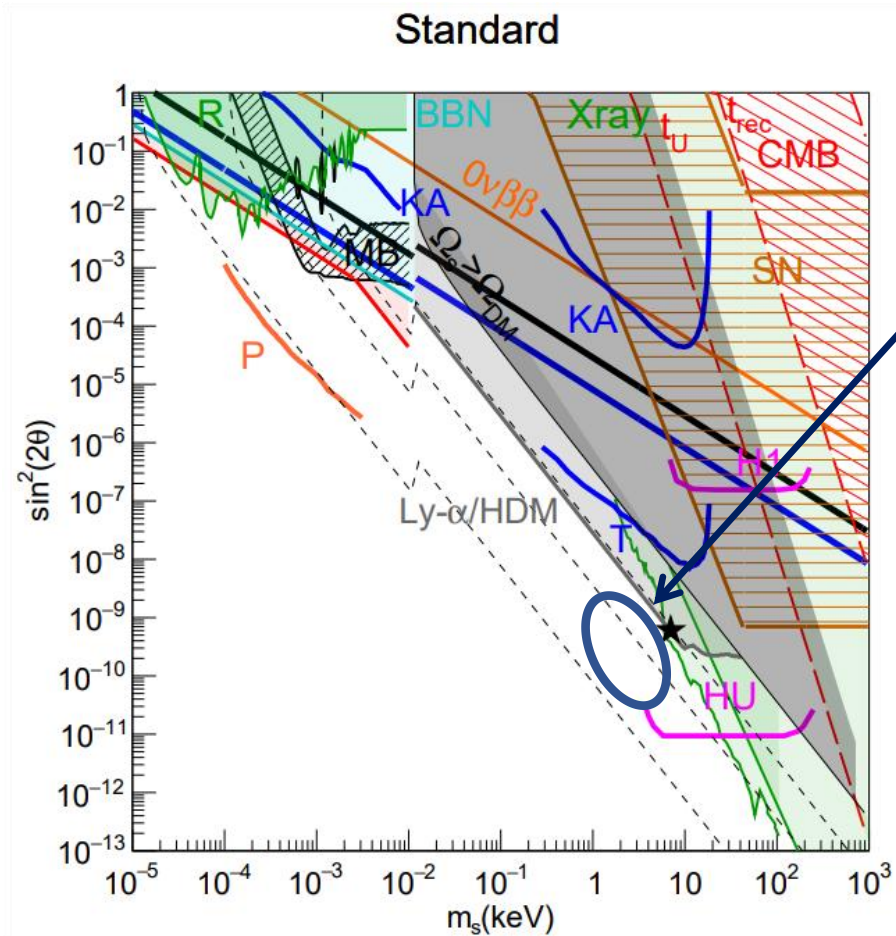


Two-populations

M. Chen, G. B. Gelmini, P. Lu, and V. Takhistov, In preparation



Observation



Low mass area where

- DW density fraction

$$f_{s,DW} \sim 0.1$$

- Other mechanism

- Colder
- Density fraction ~ 0.9

G. B. Gelmini, P. Lu, and V. Takhistov, JCAP 12 (2019) 047 [1909.13328]

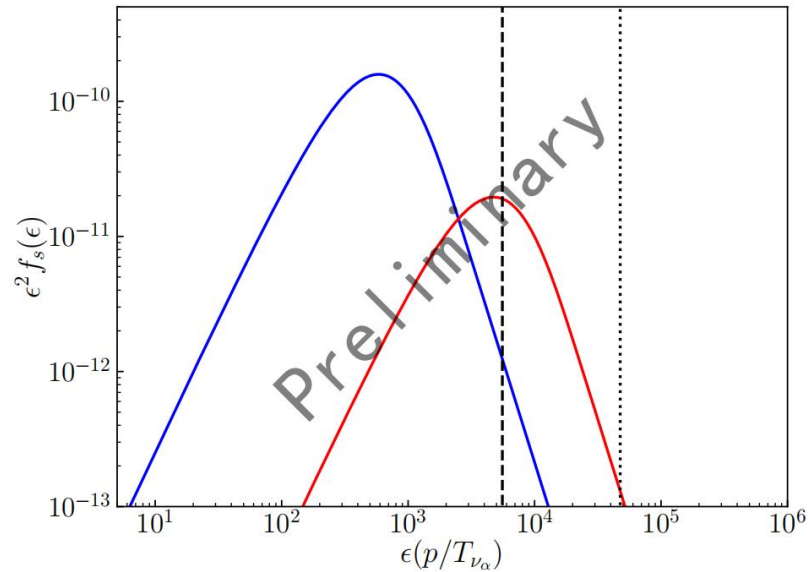


Example

This appears in many models:

PBH neutrino genesis, Dodelson-Widrow, Shi-Fuller...

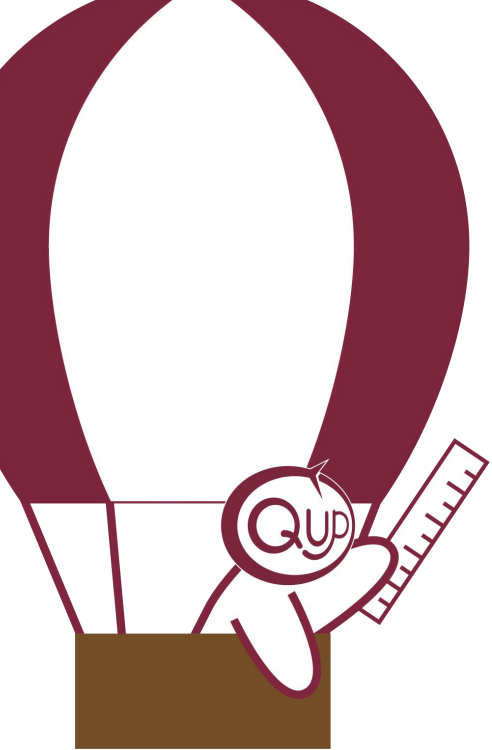
E.g.





Summary

- Sterile neutrinos can be detected by their decay $\nu_s \rightarrow \nu\gamma$. Current X-ray limit ruled out the possibility of producing the whole of DM through DW mechanism.
- Sterile neutrino produced due to PBH evaporation does not depend on the mixing, thus not constrained by the X-ray limit. Sterile neutrinos can be cold/warm in the RD case.
- X-ray signal can be correlated to GW signal to show the possible existence of such mechanism.
- Two sterile neutrinos populations produced with different mechanisms can produce all of the DM under the X-ray limit.



Thank You



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