

# From algorithms to applications: potential roles of quantum computing

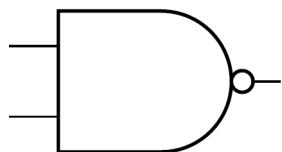
---

Kosuke Mitarai  
Osaka University

# Quantum computing

## Conventional computers

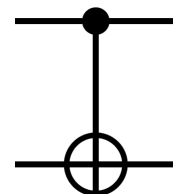
{NAND} is a universal gateset



Input	Output
00	1
01	1
10	1
11	0

## Quantum computers

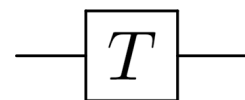
{CNOT, T, H} is a universal gateset



Input	Output
$ 00\rangle$	$ 00\rangle$
$ 01\rangle$	$ 01\rangle$
$ 10\rangle$	$ 11\rangle$
$ 11\rangle$	$ 10\rangle$



$ 0\rangle$	$( 0\rangle +  1\rangle)/\sqrt{2}$
$ 1\rangle$	$( 0\rangle -  1\rangle)/\sqrt{2}$



$ 0\rangle$	$ 0\rangle$
$ 1\rangle$	$\exp(i\pi/4)  1\rangle$

Quantum computers works in a similar manner to classical ones, but with different gates

# Power of quantum computers

- Simulating dynamics of interacting  $n$   $1/2$ -spins [S. Lloyd, Science, **273**, 1073-1078 (1996)]

$$O(2^n) \rightarrow \text{poly } n$$

- Factoring of  $n$  bit integers [P. W. Shor, Proceedings 35th Annual Symposium on Foundations of Computer Science, 124-134 (1994)]

$$O\left(e^{1.9n^{1/3}(\log n)^{2/3}}\right) \rightarrow O(n^2 \log n \log \log n)$$

- Searching among  $N$  possibilities [L. K. Grover, Proceedings, 28th Annual ACM Symposium on the Theory of Computing, 212-219, (1996)]

$$O(N) \rightarrow O(\sqrt{N})$$

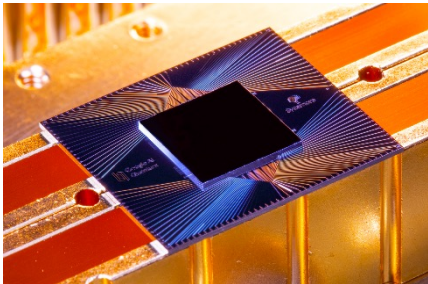
- Inversion of sparse  $N \times N$  matrix (sparseness  $s$ , condition number  $\kappa$ , precision  $1/\epsilon$ ) [A. Harrow et al., PRL, **103**, 150502 (2009)]

$$O(Ns\sqrt{\kappa} \log 1/\epsilon) \rightarrow \tilde{O}(\log N s^2 \kappa^2 / \epsilon)$$

**Many applications, but needs hardware for realizing them**

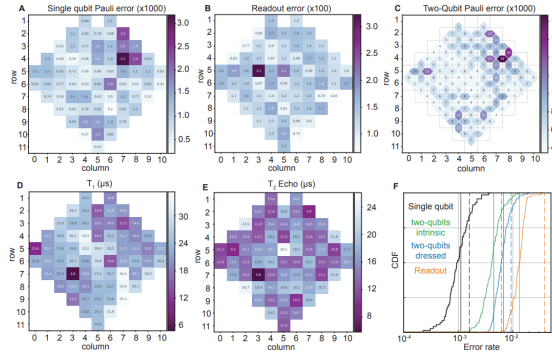
# Current quantum hardware as of 2024

**2019 53 qubit  
(transmon qubits)**



F. Arute et al., Nature 2019

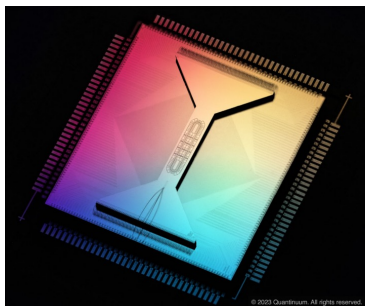
**2023 70 qubit  
(transmon qubits)**



Google Quantum AI, arXiv: 2304.11119

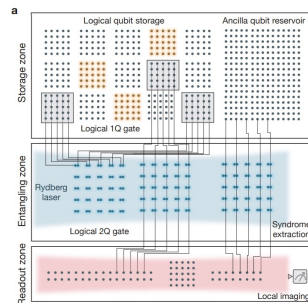
**Quantum supremacy demonstrated:**  
classical supercomputer could not simulate dynamics of a programmable, gate-based, quantum device.

**2023 56 qubit  
(ion trap)**



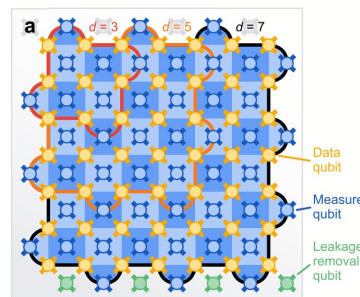
S. A. Moses et al., PRX 2023

**2023 280 qubit  
(neutral atoms)**



D. Bluvstein et al., Nature 2023

**2024 105 qubit  
(transmon qubits)**

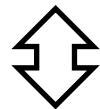


Google Quantum AI, arXiv: 2408.13687

**Quantum error correction demonstrated:**  
error correction seems to be possible in real world, for the first time.

# Our ultimate goal: fault-tolerant quantum computing

Current error rate of qubits  $\sim$  **0.1%** [Arute et. al., Nature (2019)]



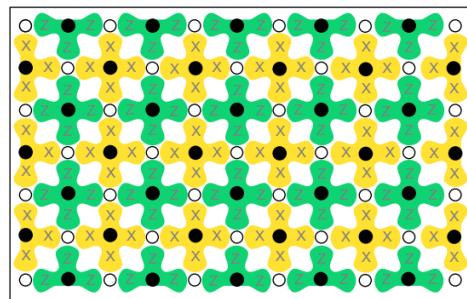
Error rate of classical bits  $\sim$   **$10^{-17}$  %** [Oliveira et al, SC17 (2017)]

\* Converting FIT to error rate from the number of clocks

Error correction is essential for "normal" calculations



Surface code



Make clean 1 qubit with  $\sim$ 1000 qubits

[Phys. Rev. A **86**, 032324 (2012)]

# My talk today

---

- Developing more efficient algorithms and frameworks are very important to harness the power of quantum.
- First part: efficient simulation algorithm for Schwinger model and its applications.
- Second part: a novel quantum machine learning framework

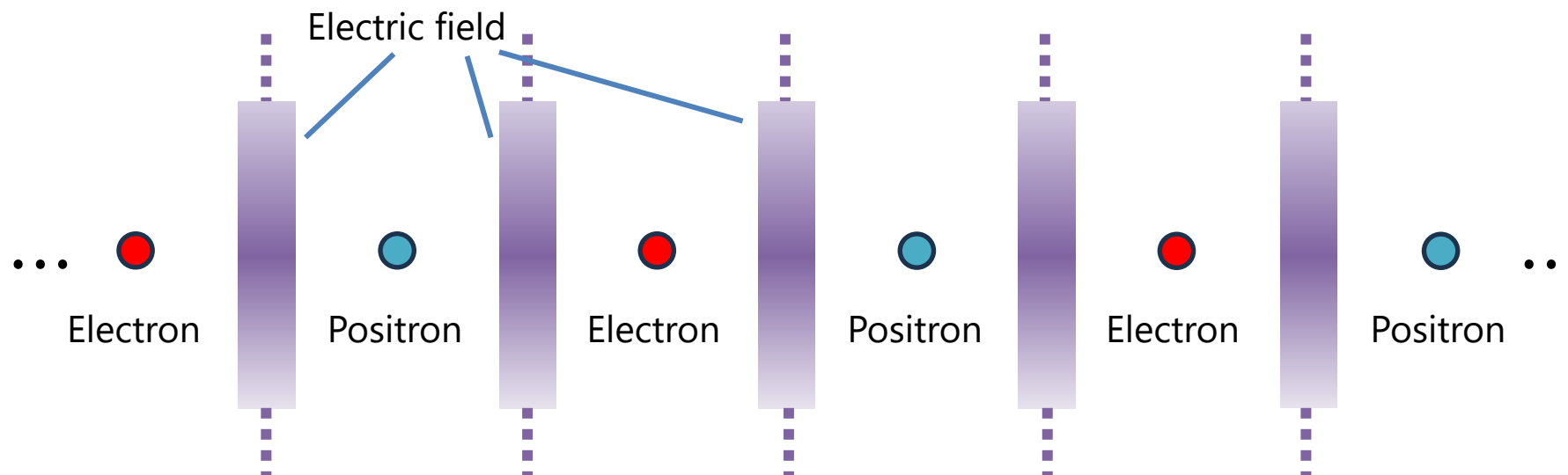
# First part: quantum algorithm for Schwinger model

---

K. Sakamoto, Hayata Morisaki, Junichi Haruna, Etsuko Itou, Keisuke Fujii, Kosuke Mitarai,  
“End-to-end complexity for simulating the Schwinger model on quantum computers”,  
arXiv:2311.17388

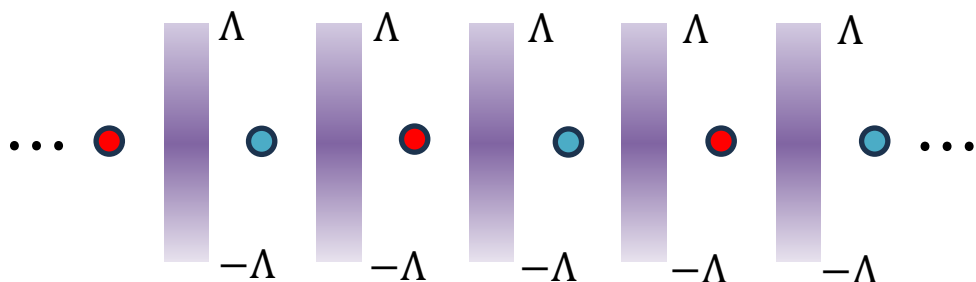
# Schwinger model

One of the simplest yet non-trivial gauge theories



Truncate the electric field at  $\Lambda$

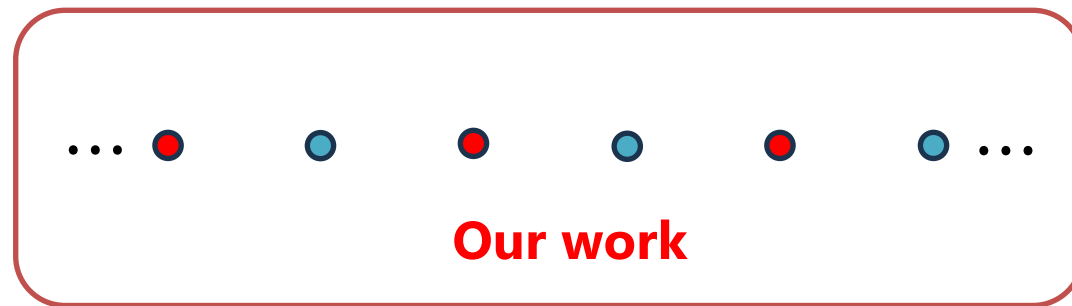
J. Kogut, et al, Phys. Rev. D (1975).



Two types of Hamiltonian formulation

Remove the electric field with Gauss's law

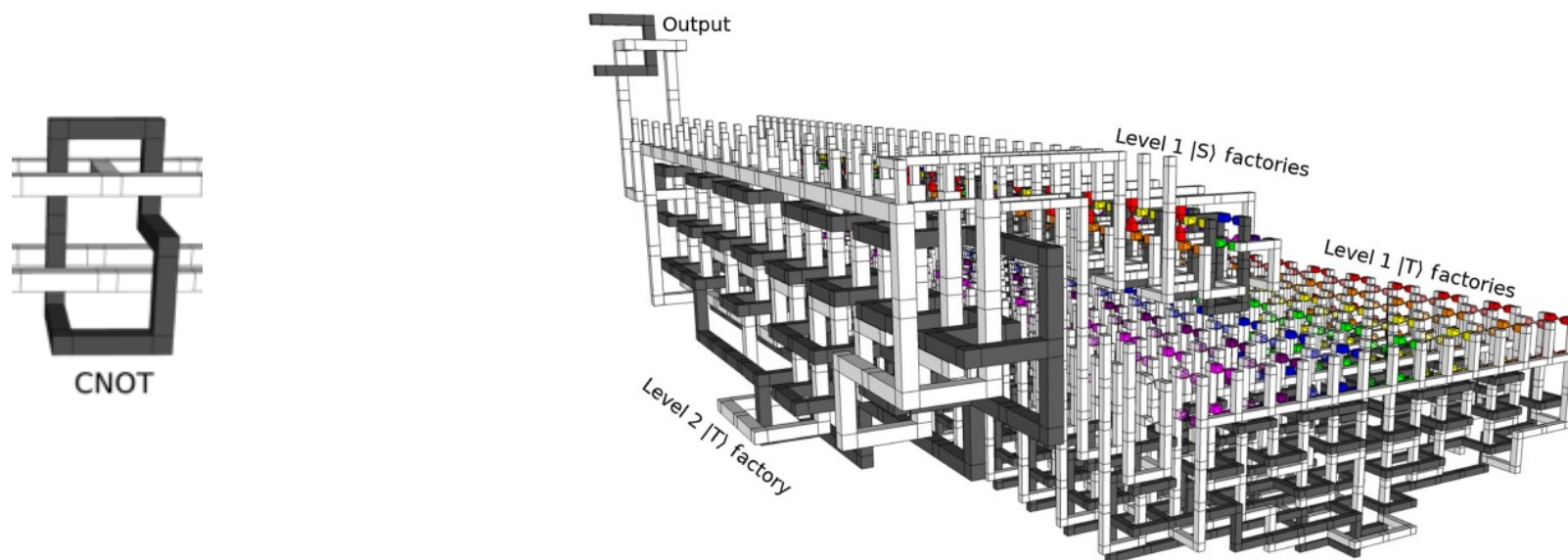
T. Banks, et al, Phys. Rev. D (1976).





# The cost is estimated via number of T gates

- In FTQC setting, T gates are the most costly.
  - FTQC usually allows  $\{H, \text{CNOT}, T\}$  gates, which are universal.
  - H and CNOT gates are very easy, but T gates need large space-time cost.
  - It is because of the structure of error-correction codes defined via commuting Pauli operators.



R. Babbush, et al, Phys. Rev. X, (2018).

- *Note added: recent works (Itogawa et al., arXiv: 2403.03991, Gidney et al., arXiv:2409.17595) might change the situation. Number of T gates, however, still roughly represents how many gates we need.*

# Previous works on Schwinger model for $e^{-iHt}$

- |  |                 |  |
|--|-----------------|--|
|  |                 | System size : $N$                              |
|  |                 | Precision : $\varepsilon$                      |
|  |                 | Evolution time : $t$                           |
| ➤ The Hamiltonian formulation which does not have electric field                             |                 |  |
| E. A. Martinez, et al, Nature 534, 516 (2016).   |                 |  |
| N. H. Nguyen, et al, PRX Quantum 3, 020324 (2022).   | .....           | $O(N^{4.5}t^{1.5}/\varepsilon^{0.5})$          |
| - Based on Trotter formula   |                 |  |
|  | <b>Our work</b> | .....  |
|  |                 | $\tilde{O}(N^4t + \log(1/\varepsilon))$        |
| ➤ The Hamiltonian formulation which has electric field                                       |                 |  |
| A. F. Shaw, et al, Quantum 4, 306 (2020).  | .....           | $\tilde{O}(N^{2.5}t^{1.5}/\varepsilon^{0.5})$  |
| - Based on Trotter formula   |                 |  |
| - Provides rigorous cost analysis  |                 |  |
| Y. Tong, et al, Quantum 6, 816 (2022).   | .....           | $\tilde{O}(Nt \text{ polylog}(1/\varepsilon))$ |
| - The smallest query complexity at present   |                 |  |
| - Probably needs a huge number of qubits   |                 |  |
| ➤ <b>Our work improves in every factor from the previous Trotter-based one.</b>              |                 |  |
| ➤ <b>Compared to ones with electric field, our algorithm needs smaller number of qubits.</b> |                 |  |

# Block-encoding

- Block-encoding of a Hamiltonian  $H$  is defined as,

$$U = |0^b\rangle\langle 0^b| \otimes H + \dots = \begin{pmatrix} H & \cdot \\ \cdot & \cdot \end{pmatrix}$$

We will see how to implement such  $U$  on the next page.

- Here we assume  $U^2 = I$ . This holds for popular block-encoding implementations.
- Let  $R = 2|0^b\rangle\langle 0^b| \otimes I - I \otimes I$ ;  $R$  adds phase -1 when first  $b$  qubits are not  $|0\rangle$ .
- Surprisingly, the following holds:

$$(RU)^n = RU \dots RURU = \begin{pmatrix} T_n(H) & \cdot \\ \cdot & \cdot \end{pmatrix}$$

where  $T_n(H)$  is the Chebyshev polynomial.

- **Advantages:**

- **Block-encoding  $U$  of  $H$  with error  $\epsilon$  only requires  $O(\log(1/\epsilon))$  gates in most cases.** (Trotter expansion needs  $\text{poly}(1/\epsilon)$  gates.)
- **We can get any information about  $H$  with  $T_n(H)$ ; Most functions can be efficiently approximated by linear combination of  $T_n(x)$ .**

# Block-encoding of Pauli-sum Hamiltonians

- Assume Hamiltonian is decomposed as sum of Pauli operators  $P \in \pm\{I, X, Y, Z\}^{\otimes n}$ :

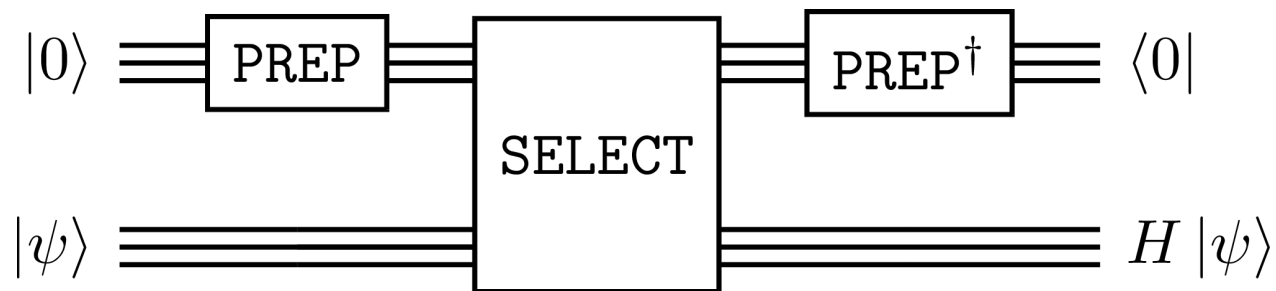
$$H = \sum_{l=0}^{L-1} a_l P_l$$

Furthermore, assume  $a_i > 0$  and it is normalized such that  $\sum_i a_i = 1$ .

- Let the PREPARE operator PREP and SELECT operator SELECT be ones that satisfies:

$$\text{PREP}|0^b\rangle = \sum_{l=0}^{L-1} \sqrt{a_l} |l\rangle, \text{SELECT}(|l\rangle|\psi\rangle) = |l\rangle \otimes (P_l|\psi\rangle)$$

- The following gives a block-encoding:



- $P$  and  $V$  can be implemented  $O(L + \log 1/\epsilon)$  gates (using ancillary qubits). [R. Babbush, et al, Phys. Rev. X, \(2018\)](#)
- This technique is called the **linear combination of unitaries (LCU)**.

# Our idea to efficiently implement block-encoding

- The Schwinger model Hamiltonian after Jordan-Wigner transformation looks like:

$$H_S = J \sum_{n=0}^{N-2} \left( \sum_{i=0}^n \frac{Z_i + (-1)^i}{2} + \frac{\theta_0}{2\pi} \right)^2 + \frac{w}{2} \sum_{n=0}^{N-2} (X_n X_{n+1} + Y_n Y_{n+1}) + \frac{m}{2} \sum_{n=0}^{N-1} (-1)^n Z_n$$

- It has  $O(N^2)$  terms, we naively need  $O(N^2)$  gates to block-encode  $H_S$ .

- **Our strategy to realize it with  $O(N)$  gates:**

- Uniform superposition states  $\frac{1}{\sqrt{N}} \sum_{i=0}^{N-1} |i\rangle$  can be prepared efficiently with  $O(\log N)$  T gates.

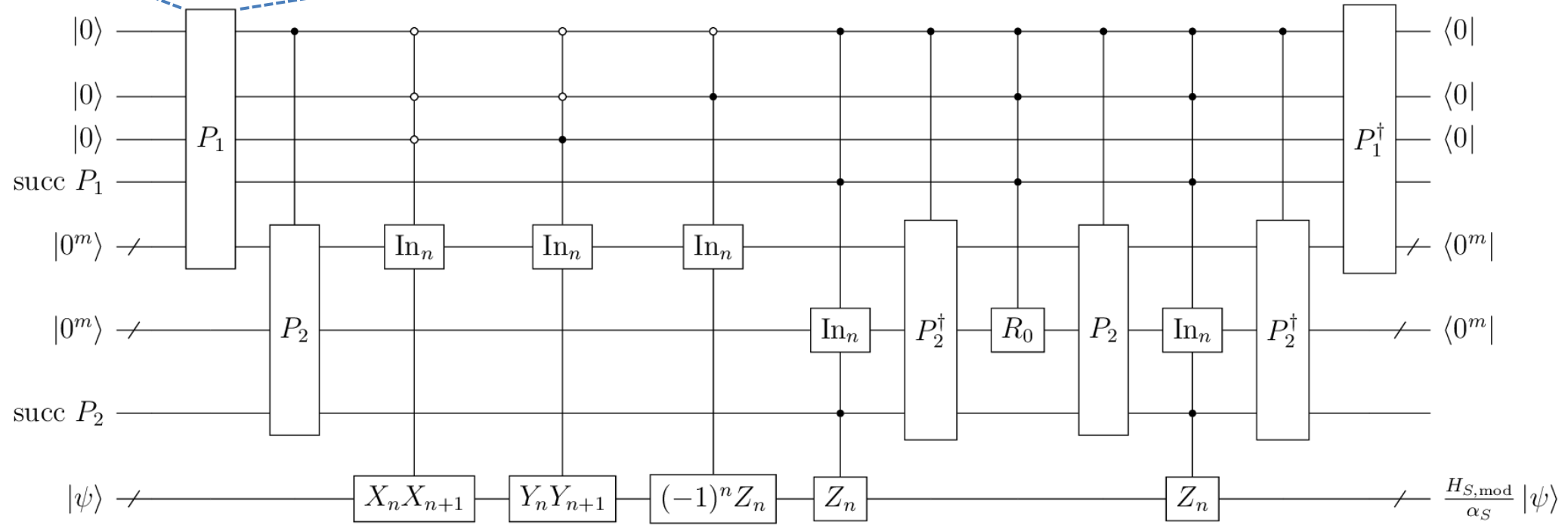
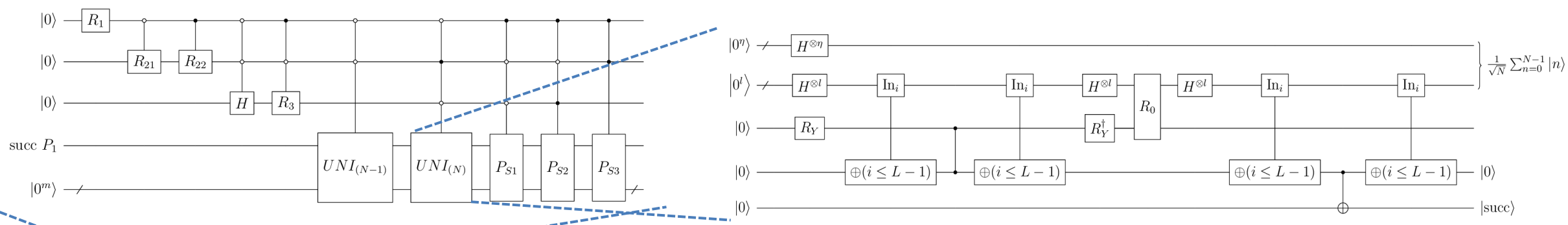
Y. R. Sanders, et al, PRX Quantum (2020)

- We can take a linear combination of the block-encodings via LCU.

- Noting the above, group the terms as follows:

$$H_S = \boxed{\frac{J}{4} \sum_{n=1}^{N-1} \left( \sum_{i=0}^{n-1} Z_i \right)^2} + \boxed{J \frac{\theta}{2\pi} \sum_{n=1}^{N-1} \sum_{i=0}^{n-1} Z_i} + \boxed{J \left( \frac{1}{2} + \frac{\theta}{2\pi} \right) \sum_{n=1}^{N-1} \sum_{i=0}^{n-1} Z_i} + \boxed{\frac{w}{2} \sum_{n=0}^{N-2} X_n X_{n+1}} + \boxed{\frac{w}{2} \sum_{n=0}^{N-2} Y_n Y_{n+1}} + \boxed{\frac{m}{2} \sum_{n=0}^{N-1} (-1)^n Z_n}$$

# Quantum circuit for block-encoding looks like...



# Resource estimates for computing $\langle \text{vac} | e^{-iHt} | \text{vac} \rangle$

- $|\text{vac}\rangle = |1010 \dots\rangle$  is the ground state of  $H_S$  for  $J = \theta_0 = w = 0, m = m_0$ , representing vacuum without any particle.
- $\langle \text{vac} | e^{-iHt} | \text{vac} \rangle$  is *vacuum persistent amplitude*, representing the creation and annihilation of electron-positron pairs
- Based on our block-encoding, how long does it take to compute it with FTQC?

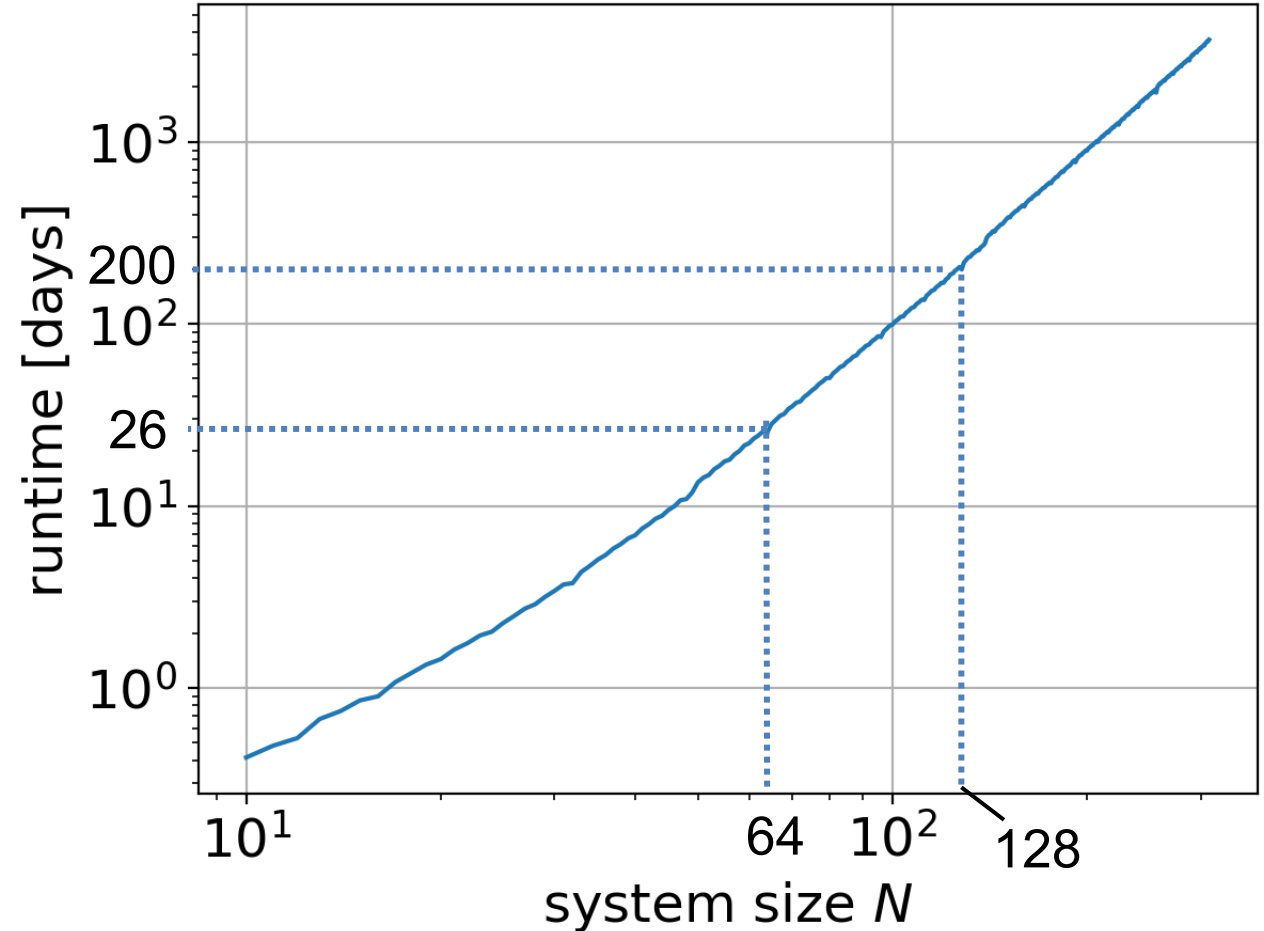
J. Schwinger, Phys. Rev. (1951)

# Resource estimate result - runtime

- Parameters
  - Precision (additive error) :  $\varepsilon = 0.01$
  - Evolution time :  $t = 4$
  - T gate consumption rate : 1MHz
  - Lattice spacing :  $a = 0.2$
  - electron mass :  $m = 0.1$
  - $w = \frac{1}{2a} = 2.5$
  - $J = \frac{g^2 a}{2} = 0.1, (g = 1)$
  - $\theta_0 = \pi$

Examples

System size	Runtime [days]
64	26
128	200



Runtime for calculating the vacuum persistence amplitude.



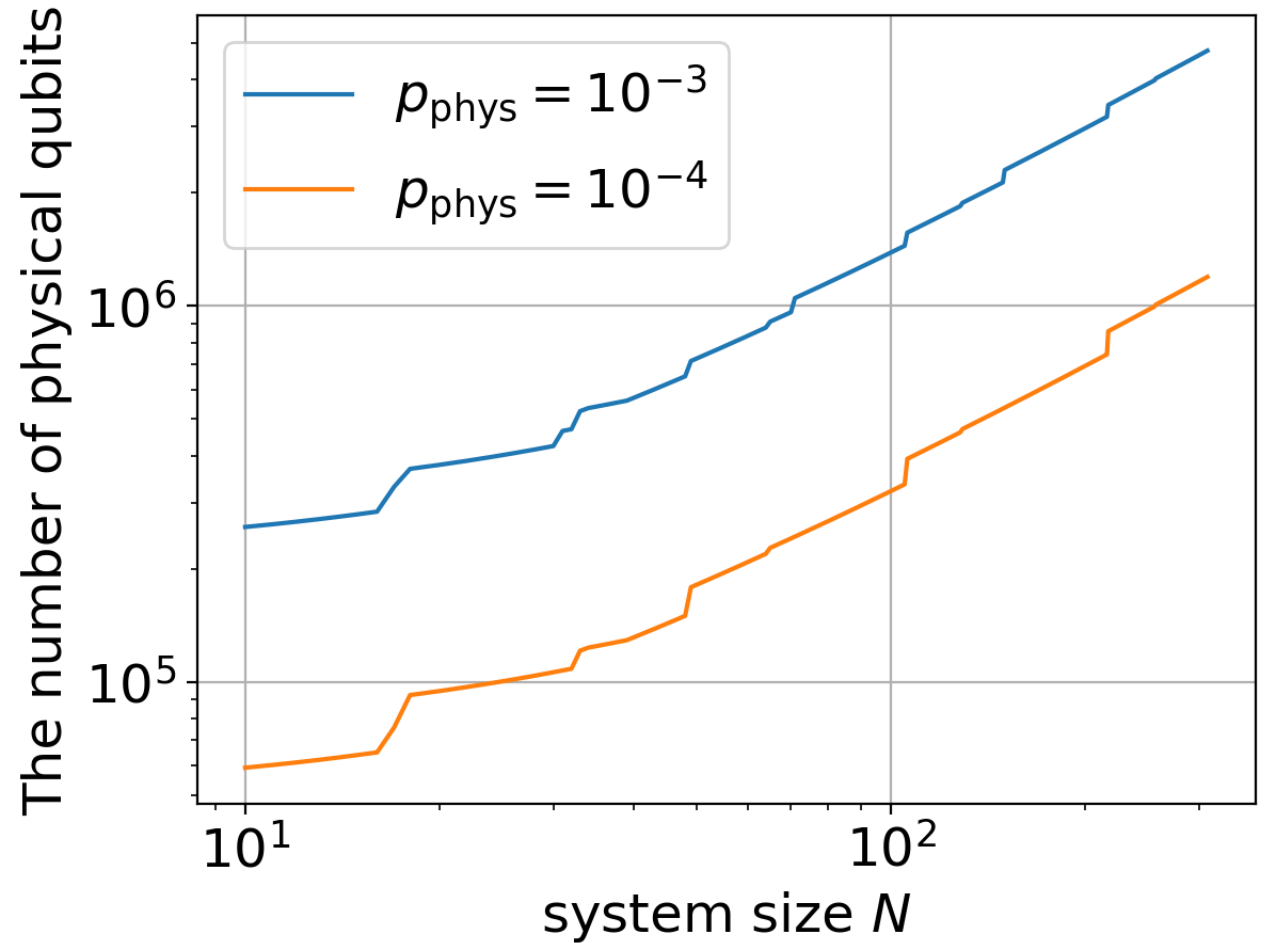
# Resource estimate result: qubit requirements

- Parameters

- Precision (additive error) :  $\varepsilon = 0.01$
- Evolution time :  $t = 4$
- Lattice spacing :  $a = 0.2$
- electron mass :  $m = 0.1$
- $w = \frac{1}{2a} = 2.5$
- $J = \frac{g^2 a}{2} = 0.1, (g = 1)$
- $\theta_0 = \pi$

Examples ( $N = 64$ )

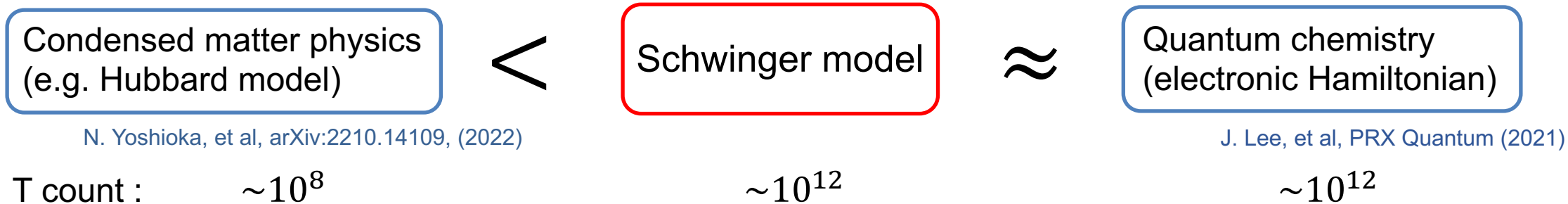
Physical error rate	Physical qubits
$10^{-3}$	$9 \times 10^5$
$10^{-4}$	$2 \times 10^5$



The number of physical qubits for calculating the vacuum persistence amplitude.

# Summary

## Comparing resource to other applications



### Technical contributions:

- An efficient block-encoding of the Schwinger model Hamiltonian
  - Decompose the Hamiltonian into several parts.
  - Use  $O(\log^2 N)$  T gates for  $P$ ,  $O(N)$  T gates for  $V$ , with a normalization factor of  $O(N^3)$ .
- End-to-end complexity for the Schwinger model

### Future challenges:

- More precise resource estimates. Maybe using libraries such as qualtran or Qiskit, which have implementations of reversible arithmetics.

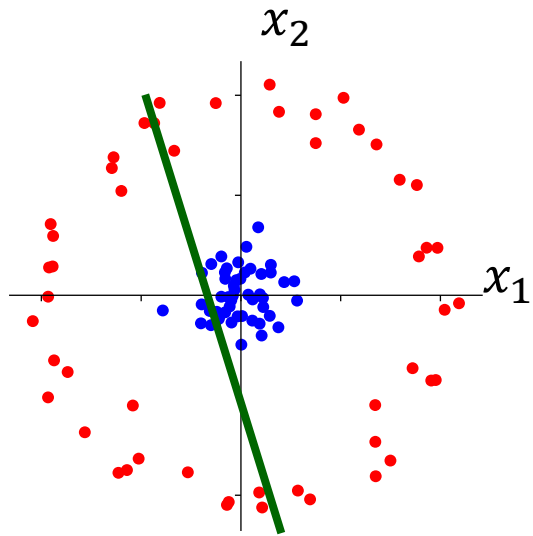
# A new quantum machine learning framework: Explicit quantum surrogate

---

Akimoto Nakayama, Hayata Morisaki, Kosuke Mitarai, Hiroshi Ueda, Keisuke Fujii,  
"Explicit quantum surrogates for quantum kernel models", arXiv:2408.03000

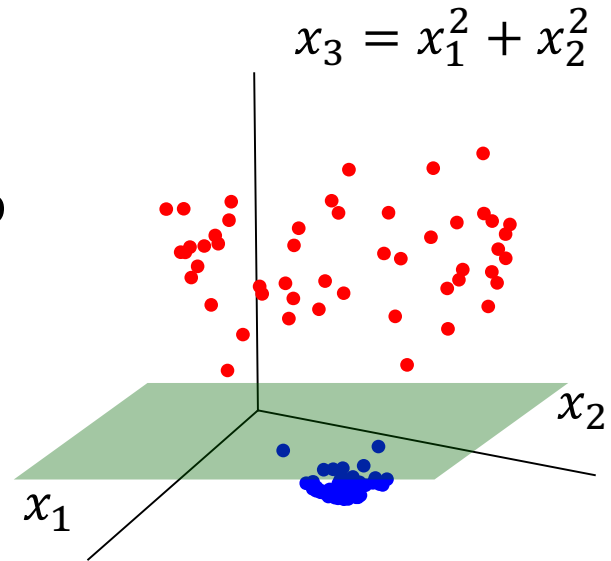
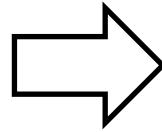
# Feature map

Transform data  $x$  to  $\phi(x)$  to extract "pattern" in the data.



Linear model fails

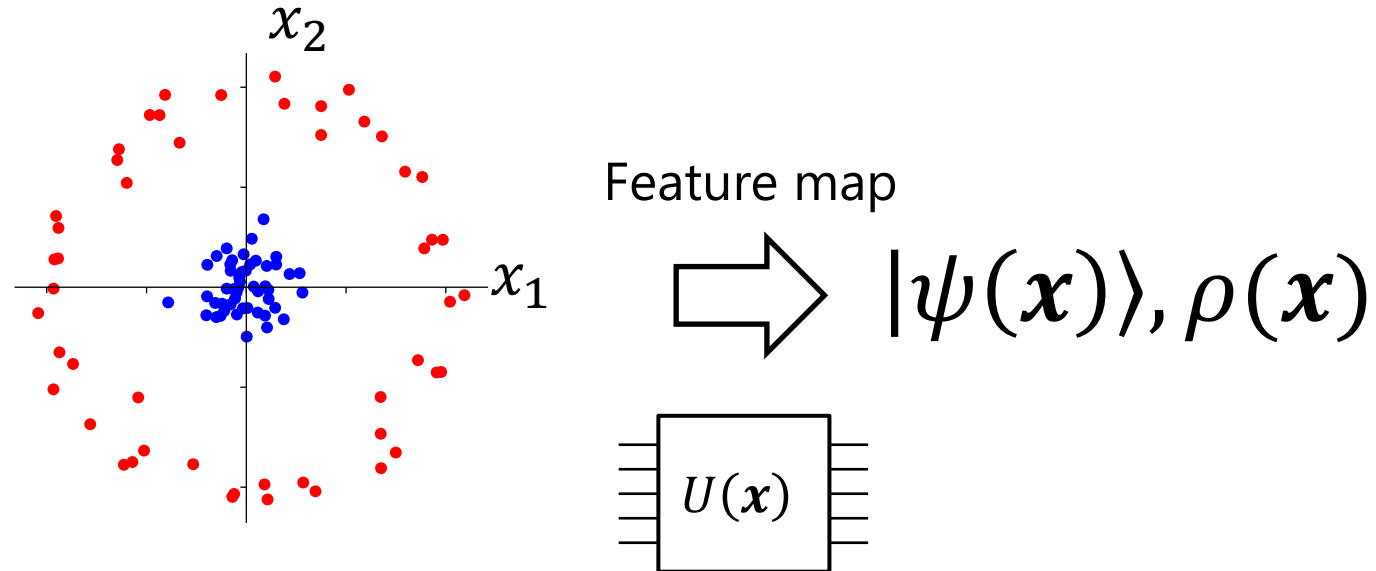
Feature map



Linear model succeeds

# Quantum feature

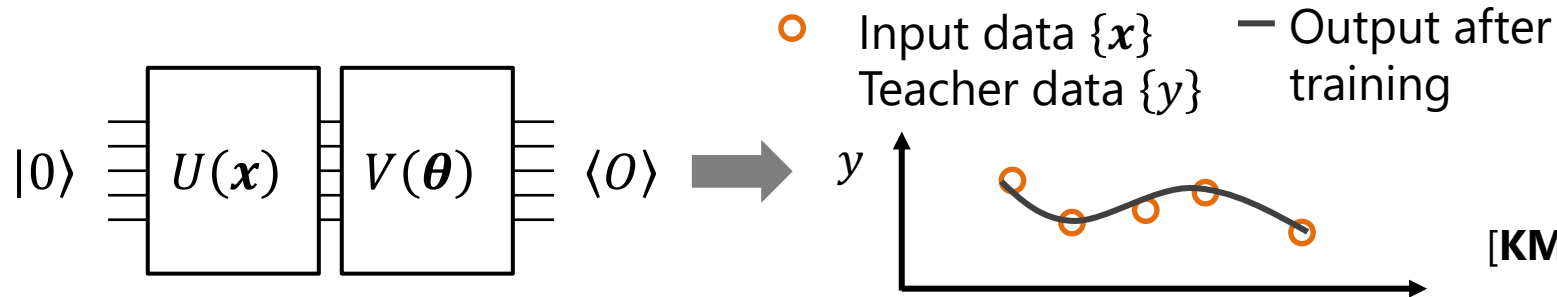
- Map a data  $\mathbf{x}$  to a quantum state  $|\psi(\mathbf{x})\rangle$



- Difficult to classically simulate  $|\psi(\mathbf{x})\rangle$ 
  - **We can construct a model which cannot be treated with classical computers**
- Note!: it doesn't mean there is practical advantage.
- Models based on quantum feature can be categorized into two major class:  
**explicit models and implicit models**

# Explicit quantum models

- Apply parameterized unitary  $V(\theta)$  and use some expectation value as a prediction.



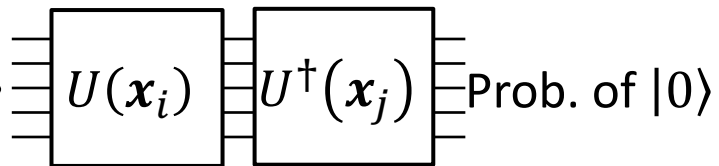
[KM, et al, Phys. Rev. A **98**, 032309 (2018)]

- Training is performed by tuning  $\theta$  via e.g. gradient decent.
- Advantage:
  - Can use optimizers from neural networks, such as Adam.
  - One training iteration needs only  $O(N)$  resource for  $N$  data.
- Disadvantage:
  - Theoretical performance not guaranteed.
  - Difficult optimization, *barren plateau* (gradient vanishes when using random initialization.)

McClean et al., Nat. Comm. **9**, 4612 (2018)

# Implicit quantum models

- Using  $k(x_i, x_j) = |\langle \psi(x_i) | \psi(x_j) \rangle|^2$  as a kernel function, rely on classical kernel techniques for constructing a model. Model in this case is represented by

$$y = \sum_{i=1}^N \alpha_i k(x_i, x_j) = \sum_{i=1}^N \alpha_i$$


The diagram shows a quantum circuit starting with an input state  $|0\rangle$  on the left. This state passes through a unitary operation  $U(x_i)$ . The output of  $U(x_i)$  then passes through its adjoint  $U^\dagger(x_j)$ . The final output is the probability of measuring the state  $|0\rangle$ , labeled "Prob. of  $|0\rangle$ ".

- Training is performed by tuning  $\alpha$  via, e.g., solving linear system of equations.
- Advantage:
  - *Convergence to global optimum guaranteed.*
  - Friendly to experiments, >20 qubits experiments are possible.
- Disadvantage:
  - Needs at least  $O(N^2)$  cost for training,  **$O(N)$  cost for prediction.**

**So we thought, can we train a model *implicitly* and then convert it to *explicit* ones?**

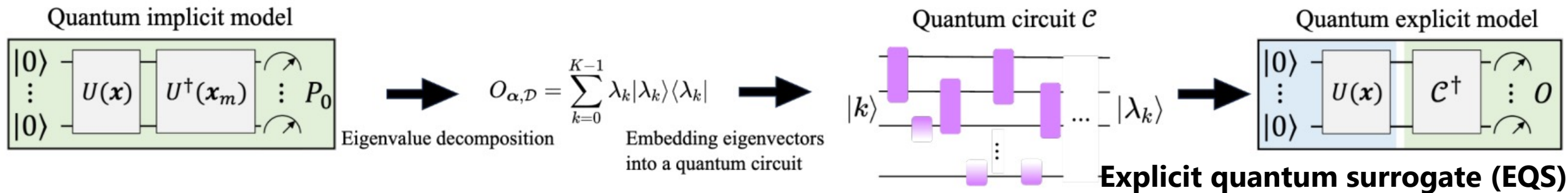
# Implicit to explicit conversion

## Algorithm

1. Train a quantum implicit model with your choice of quantum feature map  $U(\mathbf{x})$
2. Diagonalize  $O = \sum_i \alpha_i |\psi(\mathbf{x}_i)\rangle\langle\psi(\mathbf{x}_i)|$  and identify  $K$  important eigenvectors  $|\lambda_k\rangle$ .
3. Construct a circuit  $C$  that satisfies  $C|k\rangle \approx |\lambda_k\rangle$  using AQCE algorithm [Shirakawa et al., Phys. Rev. Research 6, 043008 (2024)].

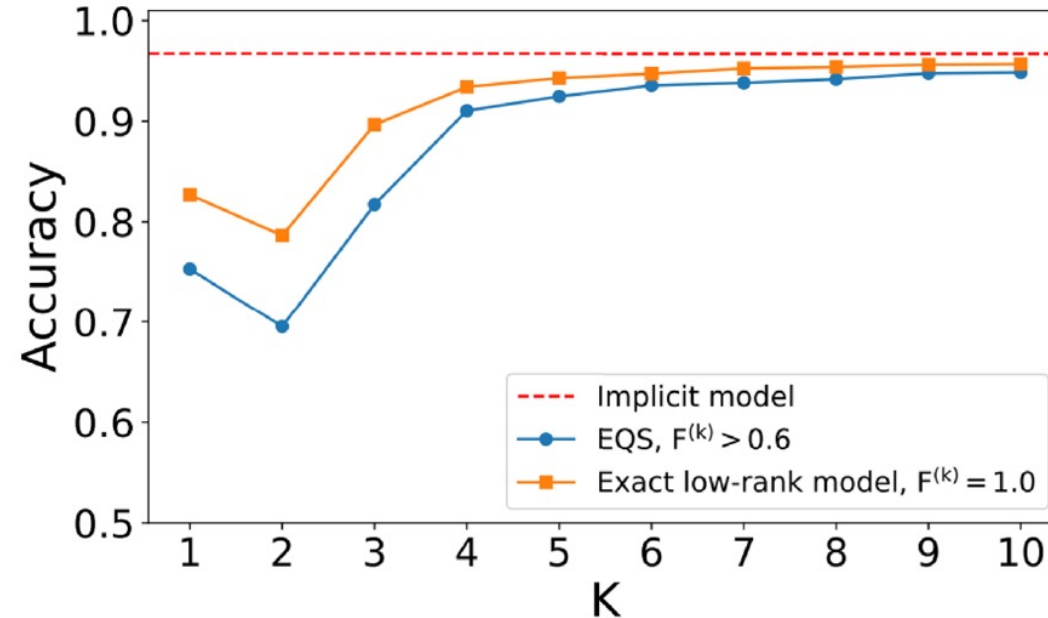
## Notes

- Step 2 can be done efficiently on a classical computer when given  $\langle\psi(\mathbf{x}_i)|\psi(\mathbf{x}_j)\rangle$  from quantum, because  $\text{rank } O \leq N$
- AQCE algorithm brute-force searches possible circuits using ideas from tensor network.



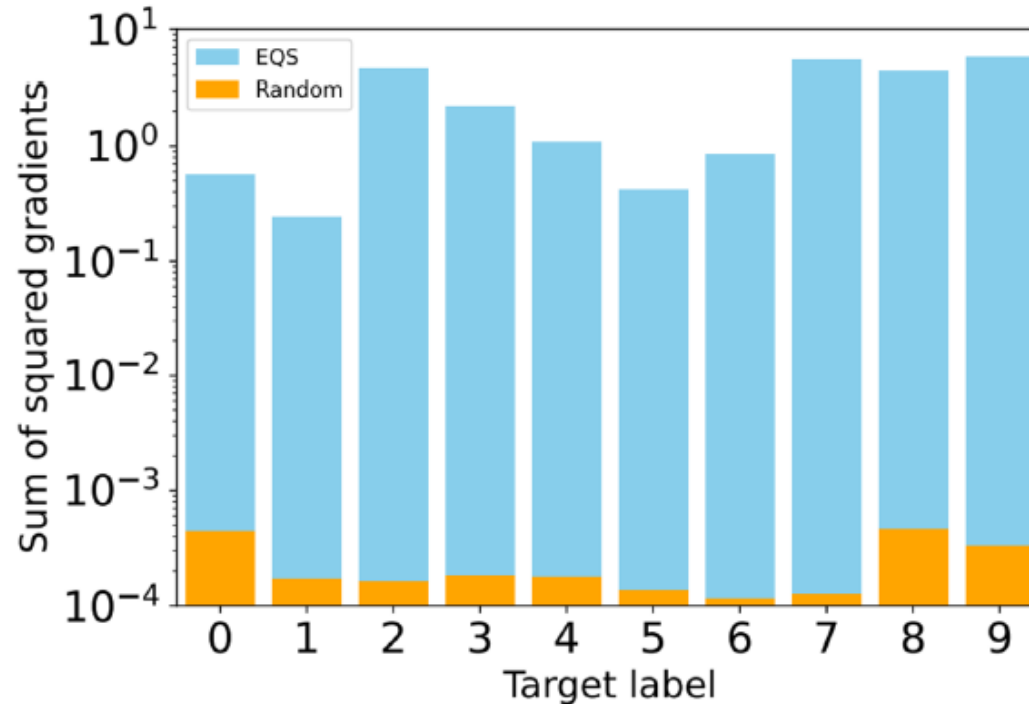


# Numerics: Accuracy for MNISQ-MNIST dataset



- ✓ MNISQ dataset [Placidi et al., arXiv:2306.16627] is a dataset developed by us, which consists of quantum circuits that approximately encodes MNIST handwritten digits.
- ✓ The conversion needs relatively small  $K$  and low fidelity ( $>0.6$ ).
- ✓ **Numerics demonstrate that we can suppress the prediction cost by this approach.**

# Numerics: EQS as an initialization strategy



- ✓ Motivation Random circuit initialization leads to barren plateau, but EQS construction is not random.  
**Can our approach mitigate the barren plateau?**
- ✓ Numerics show that magnitude of gradients are (EQS)  $\gg$  (Random initialization)
- ✓ EQS might open a way to mitigate the barren plateau.

# My talk today

---

- Developing more efficient algorithms and frameworks are very important to harness the power of quantum.
- First part: efficient simulation algorithm for Schwinger model and its applications.
  - We constructed algorithm based on block-encoding framework with detailed resource estimates.
  - We need around 1 million physical qubits and for  $10^{12}$  gates for  $\sim 100$  site Schwinger model.
- Second part: a novel quantum machine learning framework
  - Converting trained implicit models to explicit models has various benefits: Shorter prediction time, potential to mitigate barren plateau, etc.

# Some ads

---

# Qcoder: competitive quantum programming

- An IPA MITOU project (for which I am doing technical adviser)
- Competitive programming using qiskit.

An example problem from the latest contest:

## Problem Statement

You are given an integer  $n$ . Implement the operation of preparing the state  $|\psi\rangle$  from the zero state on a quantum circuit  $qc$  with  $n$  qubits.

The state  $|\psi\rangle$  is defined as

$$|\psi\rangle = \frac{1}{\sqrt{n}} (|10\dots 0\rangle_n + |010\dots 0\rangle_n + \dots + |0\dots 01\rangle_n).$$

## Constraints

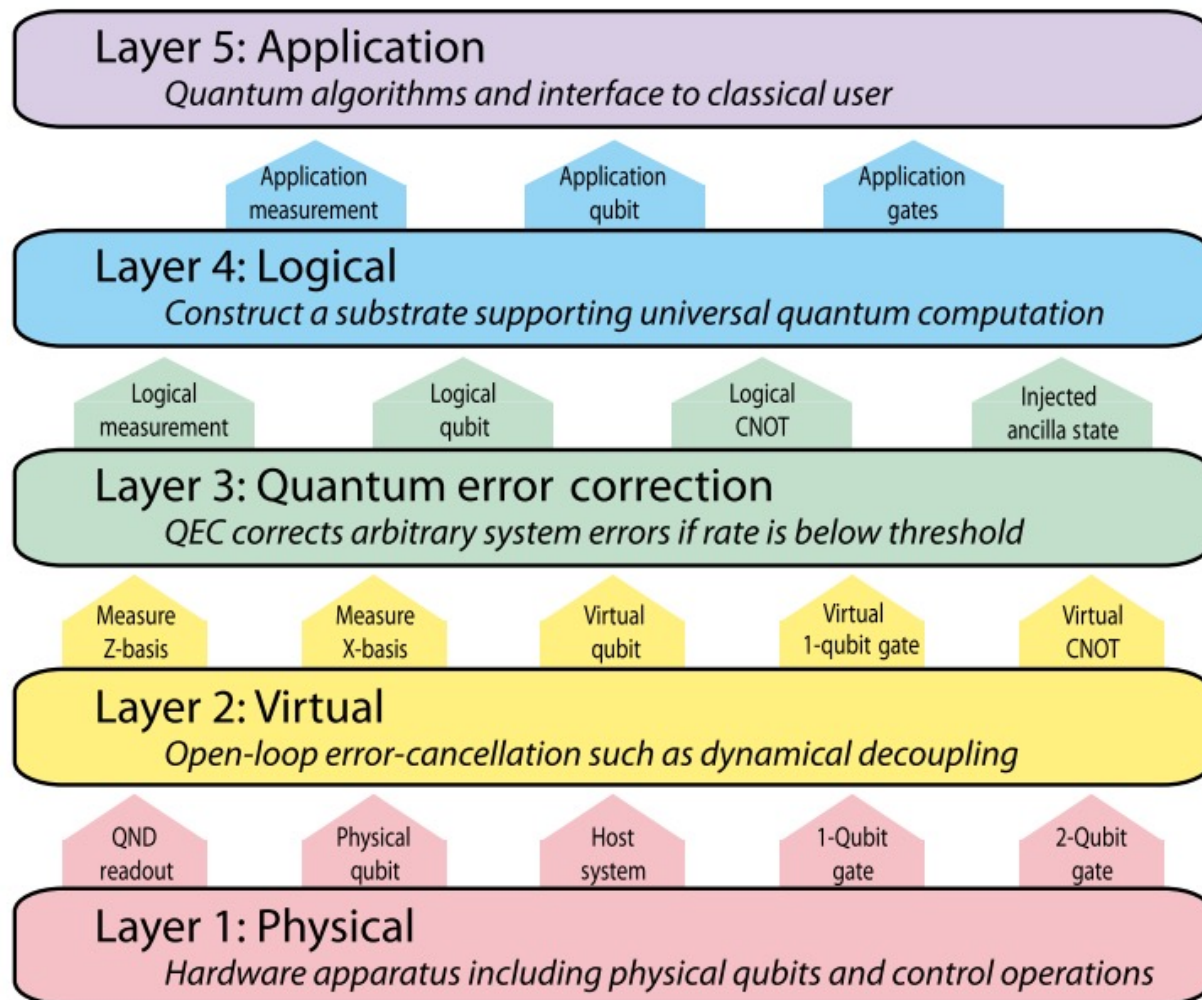
- $2 \leq n \leq 15$
- The circuit depth must not exceed 10.
- Global phase is ignored in judge.

Access **qcoder.jp**



# We are conducting full-stack research

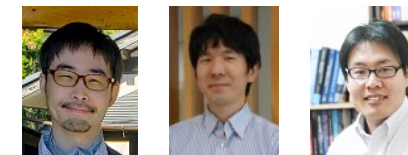
## Layers for practical quantum computing



Phys. Rev. X **2**, 031007 (2016)

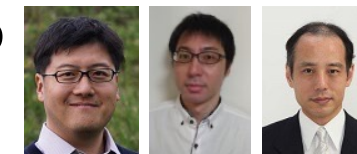
Chemistry: Prof. Mizukami  
Many-body/cond-mat: Prof. Ueda  
Machine learning: Mitarai

Mitarai, Prof. Fujii



Prof. Fujii

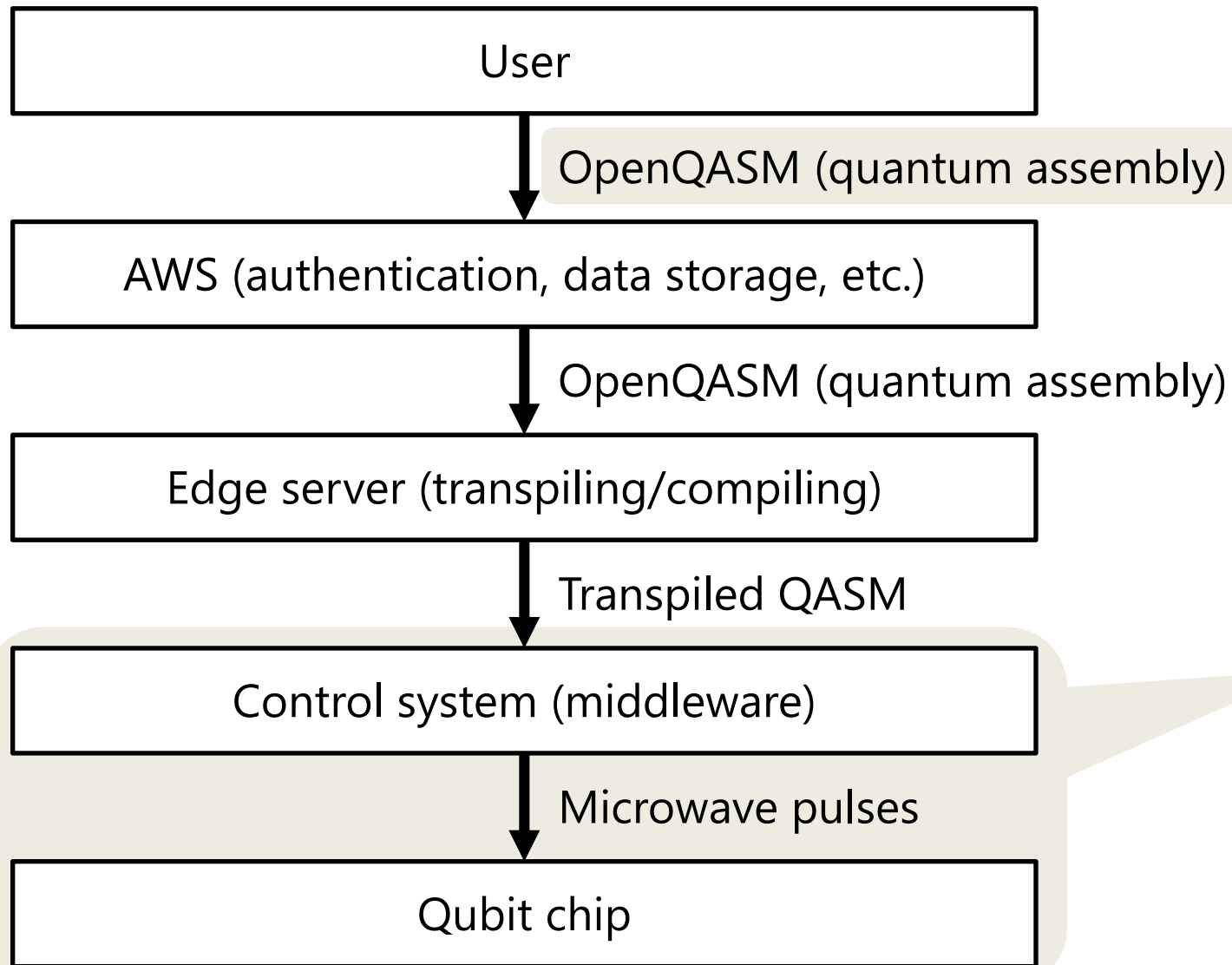
Prof. Negoro, Miyoshi, Ogawa with QuEL,  
quantum middleware startup



Superconducting qubits: Prof. Negoro  
Ion trap: Prof. Toyoda



# Current quantum computer system @ OU



```
OPENQASM 2.0;
include "qelib1.inc";
qreg q[4];
cx q[0],q[1];
rz(-0.09609732239232643) q[1];
cx q[0],q[1];
cx q[1],q[2];
rz(-0.06088586113564654) q[2];
cx q[1],q[2];
...
```



# **We are open to collaborate!**

Hosting/sending students/researchers is always welcome.

Contact me at [mitarai.kosuke.es@osaka-u.ac.jp](mailto:mitarai.kosuke.es@osaka-u.ac.jp)