Testing the lepton PDFs of the proton at lepton-proton colliders

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#### Lepton PDFs

- At zeroth order in perturbation theory the lepton carries all the momentum of the beam
- At high energies, collinear radiation emitted by splitting of the initial state must be taken into account.
- The case of collinear photon emission from an electron gives the Equivalent Photon Approximation (EPA): given by the Weizsacker-Williams spectrum  $f_{\gamma,l}(x) \approx Q_{\ell}^2 \frac{\alpha}{2\pi} P_{\gamma,\ell}(x) \ln \frac{E^2}{m_{\ell}^2}$

Weizsiieker, Kopenhagen

https://doi.org/10.1007/BF01333110

Williams https://doi.org/10.1103/PhysRev.45.729

Splitting func:  $P_{\gamma,\ell}(x) = \frac{1+(1-x)^2}{x}$ 

the probability of the lepton to emit a photon with a fraction x of its longitudinal momentum.

These radiated photons can furthermore provide other fermion parton content to the colliding lepton

via the  $\gamma \to f\bar{f}$  splitting function  $P_{f,\gamma}(z_f) = 1 - 2z_f + 2z_f^2$   $z_f \to z_f \to z_f$  the fraction of momentum of the fermion with respect to the photon.

The EPA has been generalized to describe EW gauge bosons in high-energy collisions, in what is now known as Effective Vector Approximation (EVA)

## Lepton PDFs

This previous description becomes inadequate at some high scale

At high energies  $E >> m_{\ell}$  the collinear logarithmic term can become sizeable and needs to be resummed for reliable predictions.

Leads to solving QED analog of Dokshitzer-Gribov-Lipatov-Altarelli-Parisi (DGLAP) equations

Generalization of Parton Distribution Functions (PDFs) for charged leptons

Recently, it has been computed for electron and muons



Garosi, et al. https://arxiv.org/abs/2303.16964 Han, et al. https://arxiv.org/abs/2007.14300 https://arxiv.org/abs/2103.09844

### Lepton PDFs of proton

A similar description applies to the QED interactions within the proton.

When estimating proton PDFs involving QED-induced processes, hadronization effects lead to a modification of the logarithmic factor to  $\ln(\mu_F^2/\Lambda^2)$   $\mu_F \rightarrow \text{factorization scale}$  $\Lambda \rightarrow \text{typical hadronic scale}$  arXiv:2005.06477 Buonocore, et al.

In a pure QCD model, protons consist of valence quarks, with abundant soft and collinear QCD radiation.

Order of quark and gluon PDFs  $(lpha_S L)^n$  ,  $L=\ln(\mu_F/\Lambda)$ 

If  $\ L \sim 1/lpha_S$  , all the contributions becomes relevant

#### Notes:

- DGLAP equations account for perturbative evolution (renormalization scale dependence).
- Non-perturbative initial conditions (momentum fraction dependence) are extracted through global fits to DIS and collider data
- Naive expection from power counting that the lepton density in the proton very small  $f_\ell \sim lpha^2 f_q$ But it could be large due to large logarithimic factors

#### Power counting schemes

quark and gluon PDF  $(lpha_S L)^n\,,\,L=\ln(\mu_F/\Lambda)$ 

The photon PDF is reduced by a factor of  $\alpha L$  (collinear-enhanced term)at the LO compared to the quark PDF.

$$\alpha \sim \alpha_s^2 \qquad L \sim 1/\alpha_s \qquad \rightarrow \alpha L \sim \alpha_s$$

- Similarly, leptons will be suppressed by a factor of  $(\alpha L)^2 \sim lpha_s^2$  at the LO
- So,  $f_{p/q}$  and  $f_{p/g} \sim (\alpha_s L)^n \sim \mathcal{O}(1)$   $f_{p/\gamma} \sim \alpha_s$   $f_{p/\ell} \sim \alpha_s^2 C_\ell$   $C_\ell = 3/20$



 Recently, the lepton PDFs of proton has been computed by Buonocore, et al.

arXiv:2005.06477

 One can think of the LHC as a lepton collider with varying energy. (not in terms of backgrounds)

Claim:

 The LHC can probe lepto-philic Z' up to 600 GeV, leptoquarks and doubly charged scalars up to 2.5--3 TeV with order-one couplings at high luminosity

#### Resonant Z-boson production at proton-lepton colliders:

For HERA and EIC,  $e^+$  from P and the  $e^-$  from  $e^-$  collides and produce Z-boson For MUIC, the same process would take place involving  $\mu$  instead of electron

Another processes: involving quarks coming from the colliding proton and lepton beams



PDF formalism is crucial as leptons and quarks originating from photons are in the collinear regime.

- Why Focus on a Lepton-Proton Collider Instead of a Hadron-Hadron Collider Like the LHC?
- Lepton involving process  $p(\ell^{\pm})p(\ell^{\mp}) \to Z$  cannot be tested at hadron colliders, as the quark initiated processes  $p(q)p(\bar{q}) \to Z$  dominate.
- Making the Proton-Lepton collider a unique machine to study lepton initiated Z-resonant production and test the lepton PDFs of proton.
- There is a need to distinguish lepton-initiated processes from quark-initiated processes at the proton-lepton colliders.

• Useful variable: 
$$E - P_z \equiv \sum_i (E^i - P_z^i)$$

Sum runs over all detected particles

 $P_z \rightarrow$  momentum along the Z-direction  $E \rightarrow$  Energy

#### Kinematics:

Resonance featurs the relation:  $4E_1E_2\sim m_Z^2$ 

$$\rightarrow 4E_{e^-/e^-}E_{p/e^+} \sim m_Z^2$$

$$x_{e^-/e^-}E_e \qquad x_{p/e^+}E_p$$

Energies of the colliding partons

- Electron PDFs of electron peaks at  $~~x_{e^-/e^-}\simeq 1$  (at the LO, its a Dirac-Delta func)
- The largest contribution to the resonant Z-production cross section from the lepton-initiated process occurs for

$$x_{p/e^+} = \frac{m_Z^2}{4E_e E_p}$$

- The lepton-initiated process mostly occurs for a fixed value of  ${\mathcal X}_p/e^+$  ,

while for the quark-initiated process, a range of  $x_{e^-/q}$  and  $x_{p/q}$  exists for which significant contributions



[arXiv:0911.0884]

Operated from 1992 to 2007

Electrons or positrons colliding with protons

Proton ~ 920 GeV and Electrons ~ 27.5 GeV

Integrated luminosity: 500  $Pb^{-1}$  per collider experiment (H1 and ZEUS)

$$x_{p/e^+} = \frac{m_Z^2}{4E_e E_p}$$

Lepton-initiated Z production

$$E_{p/e^+} \sim 75 \,\mathrm{GeV}$$
  
 $x_{p/e^+} \sim 0.082$ 



#### Energy sum of decaying objects coming from the produced Z-boson

Focus to the di-jet decay mode of Z



- These distributions suggest that the quark-initiated processes occur for a wide range of longituidinal momentum fraction of the quarks coming from the proton and the lepton.
- Meanwhile, the lepton-initiated process occurs at approximately a fixed value of the longitudinal momentum fraction of the lepton originating from the proton.

#### $\eta\text{-distributions}$ of the leading two jets





 $\eta$  -distributions of jets from Z, produced through the lepton-initited process at MuIC

 $P(\mu^+)\mu^- \to Z \to jj$ 



- $E_{\mu} \sim 1 \,\mathrm{TeV}$   $E_p \sim 275 \,\mathrm{GeV}$
- $\bullet \to E_{p/\mu^+} \sim 2 \,\mathrm{GeV}$

Decaying objects carry longituidinal boosts along the muon-beam direction (  $-\hat{z}$ -direction)

Dominant contribution to Z-production through lepton-initited process

$$E - P_z \equiv \sum_i (E^i - P_z^i)$$

sum runs over all detected particles



 $E \rightarrow \text{Energy}$  $P_z \rightarrow \text{momentum along the z-direction}$ (proton beam direction) •According to longitudinal momentum conservation this quantity should be equal to  $2x_{e/j}E_e$  $j \rightarrow \text{colliding parton within the electron}$ 

e

 $\hat{z}$ 

• cut of  $E - P_z$  close to  $2E_e$  leads to enhancing the purely electron initiated process due to the almost Dirac delta-like behaviour of the electron PDFs of electron near  $x_{e/e} \sim 1$ 

- kinematical distributions of the quark-initiated and photon-initiated processes predominantly peak at relatively small  $x_{e^-/q}$  and  $x_{e^-/\gamma} \rightarrow \text{significantly smaller } E P_z$
- cut of  $E P_z$  close to  $2E_e$  can separate these two contributions

## Other relevant SM processes:

 Di-jet production processes involving photon partons from the lepton beam and gluon/quark partons from the proton

• Selection cuts of  $E - P_z > 52 \,\text{GeV}$   $\longrightarrow$  Only processes with incoming photons or quarks from the electron carrying more than 26 GeV. i.e.,  $x_{e/\gamma}$ ,  $x_{e/q} > 0.945$  contribute.

Photon PDFs of the electron is not too small,
 even at relatively large x-values.

Photon from electron involving processes dominate.

• A di-jet background with no lepton can arise from processes such as  $pe^- \rightarrow je^-$ , j j e<sup>-</sup> and  $je^-\gamma$  where the final state electron is misidentified as a jet.

 $\blacktriangleright$  Negligible number of events in the signal region considering  $10^{-3}$  fake rate

Other production processes are also suppressed.



# $M_{j_1j_2}$ distribution



- Perform a cut-based analysis selection cuts:
  - $E P_Z > 52 \,\mathrm{GeV}$
  - $P_T(j) > 30 \, {\rm GeV}$
  - $|\eta(j)| < 2$
  - 5 GeV selection window of  $M_{j_1j_2}$  and  $E_{j_1j_2}$  around 91 GeV and 103 GeV, respectively

Applying the selection cuts

- Obtain 73 signal events and 267 background events considering the combined integrated luminosity of  $0.98~{\rm fb}^{-1}$  of H1 and ZEUS detectors
- Signal significance of around  $4.5\sigma$  at HERA
- Since our calculations are at LO in QCD and QED, and no showering, hadronization, or detector effects are included, it's better to extract the di-jet background from data by extrapolating from sidebands ---- similar to the approach used for the di-photon background in the Higgs discovery at the LHC.
- Similar analysis can be done in EIC. We obtain 4.3  $\sigma$  signal significance there for the same process.
- Due to the larger luminosity at the EIC, leptonic decay modes (particularly involving muons) of the Z boson can be studied.

#### NLO corrections

There are virtual corrections at one-loop level and real corrections from radiation to the processes that we discuss. For resonant Z production,

NLO QED corrections can be obtained via the photon initiated processes:

 $\gamma e^- \rightarrow Z e^- \quad \hat{\sigma}_{\gamma e^-} \sim \alpha^2 / m_Z^2 \sim \alpha_s^4 / m_Z^2 \quad \text{After convolution with the PDFs} \quad \sigma \sim \alpha_s^5 C_\ell / m_Z^2$ 

Corrections from processes with virtual particles in a loop, as well as those with a final photon

 $e^+e^- \to Z\gamma \qquad \sigma \sim \alpha_s^6 C_\ell / m_Z^2$ 

• NLO QCD corrections in Z + j production

m gq 
ightarrow Zq ~
m and~ q ar q 
ightarrow Zg (initial quark or gluon coming from the colliding lepton)  $\sigma \sim lpha_s^5/m_Z^2$ 

- The lepton-initiated resonant process has  $x_{e^-/e^-} \simeq 1$ , while the NLO QCD correction processes peak at low  $x_{e^-/q,g}$
- A cut on  $E P_Z$  can reject these corrections.
- Finally, processes with a final photon:

 $q \bar{q} 
ightarrow Z \gamma \,$  are suppressed by  $lpha_s^2$  compared with the resonant production.

## Another possibility:

- Study of non-resonant di-leptonic final state
- Signal is dominated by photon mediated t/u and s-channel processes involving the lepton PDFs of the proton
- Concentrate on the ee,  $e\mu$  and  $\mu\mu$  final states. (interesting: same sign different flavor)
  - $E P_Z$  of processes involing electron PDFs of electron beam peaks at  $2E_e \sim 55 \,\mathrm{GeV}$  at HERA



involves positron, electron PDFs of proton

•  $e^-e^+$ ,  $\mu^+\mu^-$  final states can produce via the so-called photoproduction processes Involing photon PDFs of proton and electron.

 $E - P_Z$  would peak relatively smaller value

One can compare with this with the existing data of HERA

#### Comparison with data

https://arxiv.org/pdf/0907.3627

Multi-Leptons at HERA  $(0.94 \text{ fb}^{-1})$ 

Multi-Leptons HERA analysis

Selection cuts:

- at least two central leptons (electrons or muons)  $(20^\circ < \theta < 150^\circ)$
- Leading lepton  $P_T^\ell > 10 \, {\rm GeV}$  , other leptons  $P_T^\ell > 5 \, {\rm GeV}$

 ${\mbox{ }}{\mbox{ }}$  Isolated leptons  $\Delta R>0.5$ 

- • $E P_Z < 45 \,\text{GeV}$  cut to isolate the di-leptonic  $(e^+e^-, \mu^+\mu^-)$  events from photo-production regime.
- The detector efficiencies for e and  $\mu$  in the central region: e - 80% (HI) and 90% (ZEUS),  $\mu - 90\%$  (HI) and 55% (ZEUS).
  - $0.005 \lesssim x_{p/e,\mu} \lesssim 0.03$  region contributes under these cuts
- Check cross-section uncertainties by varying the renormalization scale by a factor of 2 up and down.
   Obtain large uncertainty which reflects that our results are only accurate at LO.

Sample	Data SM		Pair Production (GRAPE)	NC DIS + QEDC	
ee	873	$895\pm57$	$724 \pm 41$	$171 \pm 28$	
$\mu\mu$	298	$320\pm36$	$320 \pm 36$	< 0.5	
$e\mu$	173	$167\pm10$	$152 \pm 9$	$15 \pm 3$	
eee	116	$119\pm7$	$117\pm 6$	< 4	
$e\mu\mu$	140	$147\pm15$	$147\pm15$	< 0.5	
$(\gamma\gamma)_e$	284	$293 \pm 18$	$289 \pm 18$	$4\pm1$	
$(\gamma\gamma)_{\mu}$	235	$247\pm26$	$247\pm26$	< 0.5	

#### Our results:

final state	ℓ-init	$\gamma$ -init	$\gamma$ -init <sub><math>E-P_z&gt;45 \text{GeV}</math></sub>
$e^+e^-$	$216^{+34}_{-41}$	$560^{+54}_{-55}$	$70^{+7}_{-8}$
$e^-e^-$	$346^{+65}_{-64}$	_	—
$\mu^+\mu^-$	$12^{+2}_{-2}$	$430^{+41}_{-43}$	$54^{+5}_{-6}$
$e^-\mu^-$	$174^{+39}_{-32}$	_	_

#### Comparison with data

- We estimate more di-leptonic events than observed at HERA.
- Agreement in a few percent is found for ee and  $\mu\mu$  events for E Pz > 45 GeV region where lepton-initiated processes mostly dominates
- Our results show  $632_{-113}^{+105}$  (ee) and  $66_{-8}^{+7}$  ( $\mu\mu$ ) events vs. 418 (ee) and 63 ( $\mu\mu$ ) reported.
- Agreement is worse in the  $e\mu$  channel, with  $348^{+79}_{-64}$  events estimated vs. 173 observed.
- Overestimates are seen in photon-initiated di-lepton pair-production (The rate of this processes involves the photon PDFs of the electron and the proton.) with  $E - P_z < 45$  GeV:  $490^{+47}_{-47}$  (ee) and  $376^{+36}_{-36}$  ( $\mu\mu$ ) events estimated vs. 284 (ee) and 235 ( $\mu\mu$ ) observed.

These discrepancies could originate from various factors:

- higher order corrections
- showering, hadronization
   estimated numbers may drop to lower values
- proper detector simulation
- Alternatively, this may suggest an overestimation of lepton and photon PDFs in the proton.
   Further theoretical and experimental analysis is needed to understand these discrepancies

#### EIC

Like at HERA, di-leptonic final states from non-resonant processes can be studied at the EIC to test these PDFs.

 $e^+e^- \rightarrow 15067$ 

Lepton-initiated proceses:

 $E - P_Z \sim 2E_e \sim 40 \,\mathrm{GeV}$ 

- Additionally, same-flavor, opposite-charge di-lepton final states can arise from photo-production predominantly in the regions of  $0.1 \leq x_{e^-/\gamma} \leq 1.0$  and  $0.04 \leq x_{p/\gamma} \leq 0.2$ , under the same phase space cuts
- 23900 e<sup>+</sup>e<sup>-</sup> total events at the LO under the same cuts, out of which 2780 events occur in the E - P<sub>Z</sub> > 35 GeV region similar numbers for μ<sup>+</sup>μ<sup>-</sup>
- Once future data is available, a detailed study of these processes can be performed at the EIC to test the lepton PDFs of the proton

#### MuIC

Proton ~ 275 GeV and Muons ~ 1 TeV
Integrated luminosity: 400 fb<sup>-1</sup>
• selection criteria:  $P_T(\ell) > 5 \text{ GeV}, |\eta_\ell| < 2.5$ •  $P\mu^- \rightarrow \mu^+ \mu^- \sim 150 \text{ events}$ •  $\mu^-\mu^- \sim 600 \text{ events}$ •  $\mu^-e^{\pm} \sim 720 \text{ events}$ •  $e^+e^- \sim 30 \text{ events}$ 

Furthermore,

opposite-sign same-flavor di-lepton final states can be produced through photoproduction, involving the photon PDFs of the proton and muon, with  $E - P_z$  peaking at relatively low values.

### Consider rare SM signatures at hadron colliders like LHC

Non-resonant production of back-to-back same sign and/or different favour di-leptons

arXiv:2005.06477

#### At the LO

 $p_{T,\ell} > 20 \,\text{GeV}, \quad |\eta| < 2.4, \quad \mu_{\text{F}} = p_{T,\ell}$ 

Buonocore, et al.

	$e^+\mu^-$	$e^+\tau^-$	$\mu^+ \tau^-$	$e^+e^+$	$\mu^+\mu^+$	$\tau^+ \tau^+$
$\sigma_{13 {\rm TeV}}$ [fb]	$0.29^{+0.13}_{-0.10}$	$0.18\substack{+0.11 \\ -0.08}$	$0.16\substack{+0.10 \\ -0.07}$	$0.24_{-0.08}^{+0.10}$	$0.19\substack{+0.09 \\ -0.07}$	$0.08\substack{+0.06 \\ -0.04}$
$\sigma_{27 {\rm TeV}}$ [fb]	$0.53_{-0.18}^{+0.25}$	$0.34_{-0.15}^{+0.21}$	$0.30\substack{+0.19 \\ -0.14}$	$0.440_{-0.14}^{+0.19}$	$0.34_{-0.12}^{+0.16}$	$0.14_{-0.07}^{+0.12}$

Number of events of such processes are estimated using the lepton PDFs of the proton

Measurable number of events are expected at HL-LHC

- For example, at  $3ab^{-1}$  of data,  $850 e^+\mu^-$ ,  $550e^+\tau^-$  and  $500\mu^+\tau^-$  events are expects at HL-LHC, considering these phase space cuts
- Analysis on this at HL-LHC can also verify the distributions of these leptonic PDFs of proton
- This could be even crucial for studying BSM scenarios, like flavor-violating scalars, as precise knowledge of the SM background is required.

#### Conclusions

- PDFs are key to hadron collider phenomenology, describing the momentum distribution of proton constituents.
- The lepton PDFs of the proton have been recently computed (LuxLep).
- This is crucial for LHC searches in BSM scenarios, including hadrophobic and leptophilic Z', doubly charged scalars, leptoquarks, flavor-violating leptophilic scalars, etc.
- We propose to validate the lepton content of the proton by analyzing both non-resonant di-lepton production and di-jets from resonantly produced Z boson, Using existing HERA data and future measurements from other lepton-proton colliders
  - For Z-resonant production, we propose conducting this study at pe colliders rather than at a hadron collider like the LHC.
  - Focus to di-jet decay modes of Z, due to large branching fractions
  - Feasible to study the leptonic decay modes of the Z boson, particularly in the di-muon final state, at the EIC
  - For non-resonant di-lepton final states, consider all flavor combinations and charge pairings.
  - These can explore a smaller  $x_{p/e,\mu}$  region of the proton lepton PDFs compared to resonant Z-boson production.
  - Comparison with HERA data indicates a slight overestimation of di-lepton production rates by a few tens of percent.
  - A dedicated theoretical and experimental study could clarify these discrepancies.





# Extra

### NLO corrections in non-resonant dilepton processes

- For the non-resonant di-lepton production there are real corrections from photon emission of initial and final particles, with the same initial state as the leading processes, that are suppressed by  $\alpha_s^2$  compared with them.
- Di-lepton plus jet final states can be obtained from  $\gamma q \to \ell \ell q$ , being suppressed by  $\alpha_s$  only, compared with the leading processes.
- When considering the quark and gluon content of the electron one can also obtain di-lepton plus jet final states which are suppressed by  $\alpha_s$
- The kinematic distributions of processes initiated by partons of the electron, other than the electron itself, are peaked towards low values of  $E P_z$ , thus the rates of these processes can be controlled by cuts on this variable.
- Virtual QED corrections are expected to be suppressed by  $\alpha_s^2$  compared with the real ones.
- A dedicated analysis containing these corrections, as well as parton showering, is needed for a precise calculation.