

Quantification of the flavor diagonal hadronic CP violation

Nodoka Yamanaka
(Tohoku University / Riken)

CP violation

Discrete symmetries : P and C

Parity transformation (P):

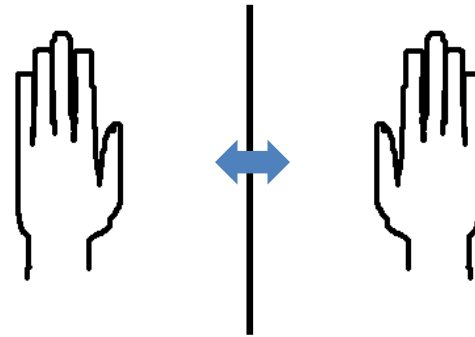
Invert all spatial coordinates

$$\vec{\mathbf{X}} \rightarrow -\vec{\mathbf{X}}$$

⇒ Mirror

For Dirac spinors,

$$\psi_L \leftrightarrow \psi_R$$



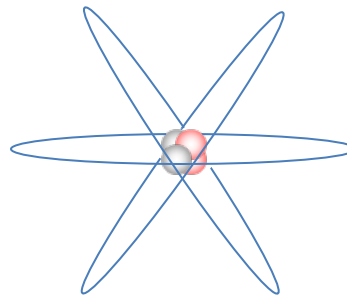
Charge conjugation (C):

Invert sign of all charges

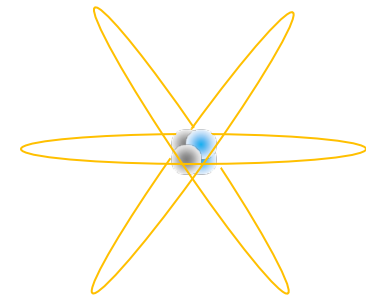
$$\psi = e^{i\theta} |\psi|$$

$$\rightarrow \psi^* = e^{-i\theta} |\psi|$$

⇒ Invert complex phases



Matter (atom)



Antimatter (anti-atom)

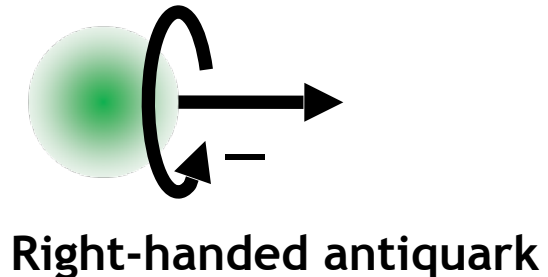
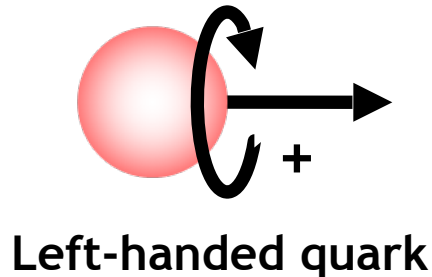
CP symmetry

Fermion $\xleftrightarrow{C}$ Antifermion

BUT !!

Fermions have spin (rotation of charge), opposite for antifermions
(antifermions have opposite P)

⇒ **Left-handed quark** and **right-handed antiquark**
have same properties if no interactions



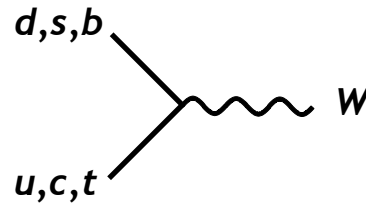
⇒ This is the **genuine symmetry between particles and antiparticles**
for massless and noninteracting fermions !

CP phases of Cabibbo-Kobayashi-Maskawa (CKM) matrix

CP violation arises when **couplings have imaginary part**,
but not all are CP violating

Example of CKM matrix:

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix}$$



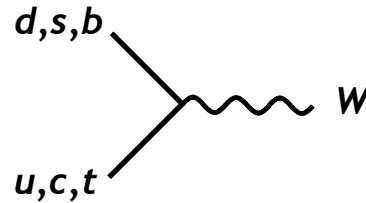
3 x 3 complex elements $\Rightarrow$ Naively, 18 free parameters . But,

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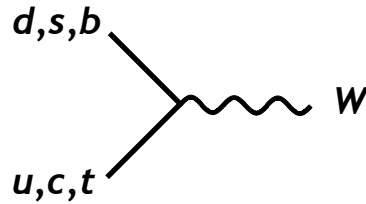
Unitarity : $V_{\text{CKM}} V_{\text{CKM}}^\dagger = \hat{1}$ **9 equations**, reduce to 9 free parameters

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Field redefinition : $\psi_q \rightarrow e^{i\theta_q} \psi_q$ **Absorb 5** complex CP phases

CP phases of Cabibbo-Kobayashi-Maskawa (CKM) matrix

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Example of CKM matrix:

$$\begin{pmatrix} V_{ud} & V_{us} & V_{ub} \\ V_{cd} & V_{cs} & V_{cb} \\ V_{td} & V_{ts} & V_{tb} \end{pmatrix} = \begin{pmatrix} C_{12}C_{13} & S_{12}C_{13} & S_{13}\underline{e^{-i\delta}} \\ -S_{12}C_{23} - C_{12}S_{23}\underline{S_{13}e^{i\delta}} & C_{12}C_{23} - S_{12}S_{23}\underline{S_{13}e^{i\delta}} & S_{23}C_{13} \\ S_{12}S_{23} - C_{12}C_{23}\underline{S_{13}e^{i\delta}} & -C_{12}S_{23} - S_{12}C_{23}\underline{S_{13}e^{i\delta}} & C_{23}C_{13} \end{pmatrix}$$

δ : CP violating phase

3 x 3 complex elements $\Rightarrow$ Naively, 18 free parameters . But,

Unitarity : $V_{\text{CKM}} V_{\text{CKM}}^\dagger = \hat{1}$ **9 equations**, reduce to 9 free parameters

Field redefinition : $\psi_q \rightarrow e^{i\theta_q} \psi_q$ **Absorb 5** complex CP phases

Remaining degree of freedom

 $\left\{ \begin{array}{l} 3 \text{ mixing angles (flavor change)} \\ \mathbf{1 \text{ complex phase : CP violation of SM !}} \end{array} \right.$

A new BSM theory MUST violate CP

In SM, all CP phases have already been fixed (absorbed)

➡ If we add new particles and interactions in the SM,
get nonzero relative phase, must violate CP !

(we call this “NATURALNESS”)

There is no reason for BSM couplings to have the same (aligned) complex phases as SM interactions, unless there is some physics behind

➡ We can look for new physics BSM by
experimentally probing CP violation !

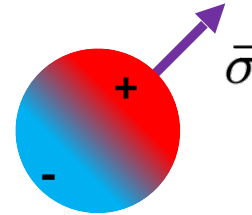
EDM, beta decay, T-violation in neutron scattering ...

Electric dipole moment

Electric dipole moment (EDM)

Permanent polarization of internal charge of a particle.

$$\vec{d}_\psi = \sum_i \langle \psi | Q_i e \vec{r}_i | \psi \rangle$$



(The unit of EDM is **e cm** :
effective distance of internal charge displacement)

● Direction: $\vec{d} \propto \vec{\sigma}$

● Interaction: $H_{\text{EDM}} = -d \langle \vec{\sigma} \rangle \cdot \vec{E}$

● Transformation properties:

● Under parity tr.:

$$\left\{ \begin{array}{ccc} \vec{E} & \xrightarrow{P} & -\vec{E} \\ \vec{\sigma} & \xrightarrow{P} & \vec{\sigma} \end{array} \right. \rightarrow H_{\text{EDM}} \text{ is P-odd}$$

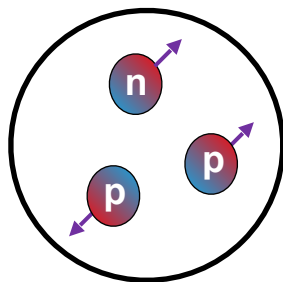
● Under time reversal:

$$\left\{ \begin{array}{ccc} \vec{E} & \xrightarrow{T} & \vec{E} \\ \vec{\sigma} & \xrightarrow{T} & -\vec{\sigma} \end{array} \right. \rightarrow H_{\text{EDM}} \text{ is CP-odd !}$$

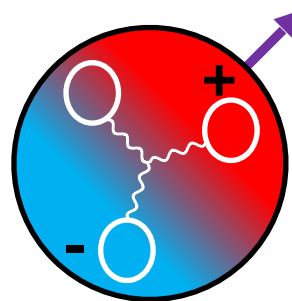
EDM of composite systems

EDM is often measured in composite systems (neutron, atoms, molecules, nuclei)

EDM is not only generated by the EDM of the components, but also by **CP violating many-body interactions**.

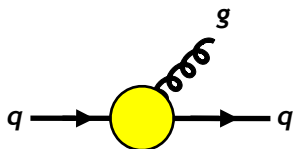


EDM of constituents

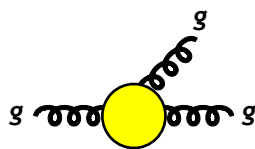


CP-odd many-body interaction

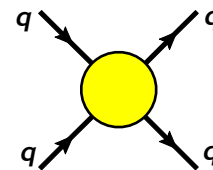
Example of QCD level many-body interactions inducing neutron EDM:



quark chromo-EDM



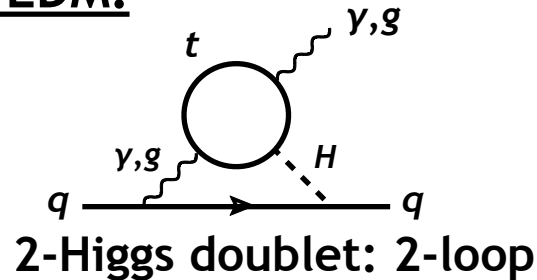
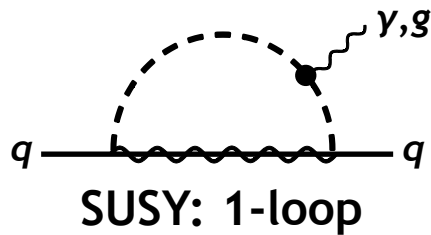
Weinberg operator



P, CP-odd 4-quark interaction

Elementary level CP violation from BSM physics (SMEFT)

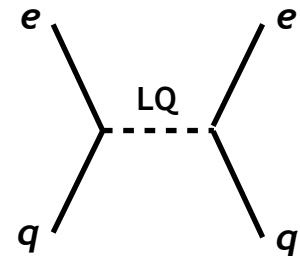
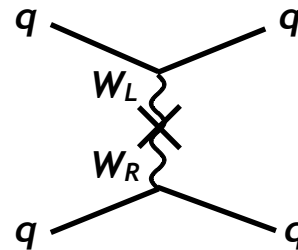
● Fermion EDM, quark chromo-EDM:



● CP-odd 4-fermion interaction:

Tree level:

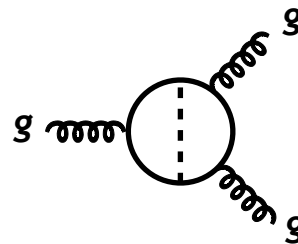
- * Left-right symmetric model
- * Scalar exchange (e.g. Higgs)
- * Leptoquarks



● Weinberg operator:

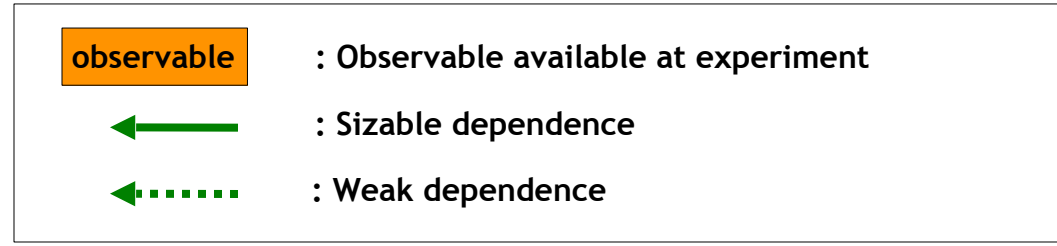
2-loop diagram:

- * 2-Higgs doublet model
- * Vectorlike quark model

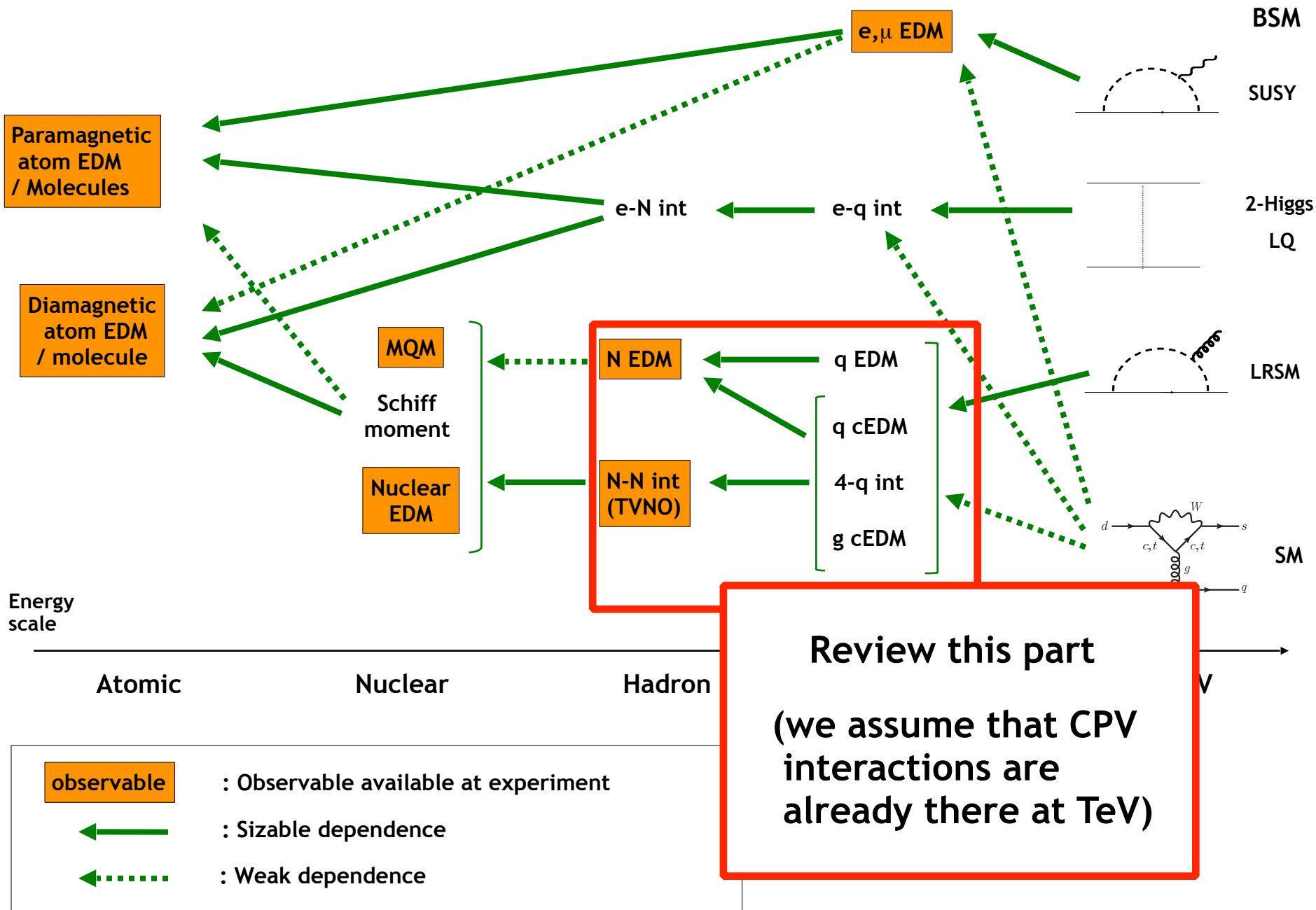


CP violating interactions appear at scale below BSM particles (TeV?)

Energy scale



EDM from CPVBSM



Hadron level CP violation

Hadron level CP violation

Hadronic CP violation has large uncertainty due to QCD nonperturbative physics

Ideally, calculated in lattice QCD, but quantified results are very limited

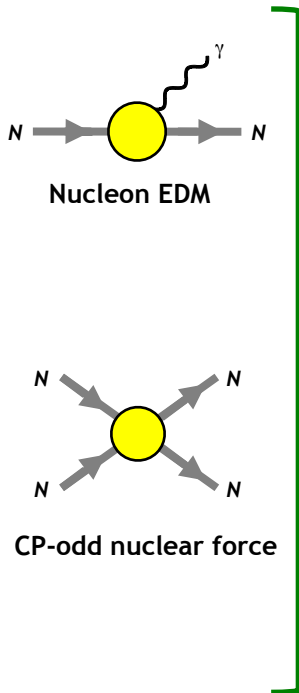
But recently, there were significant progresses in the quantification

In particular, the **leading hadronic contribution has recently been identified**

In this section, we review the hadronic CPV contributing to EDMs

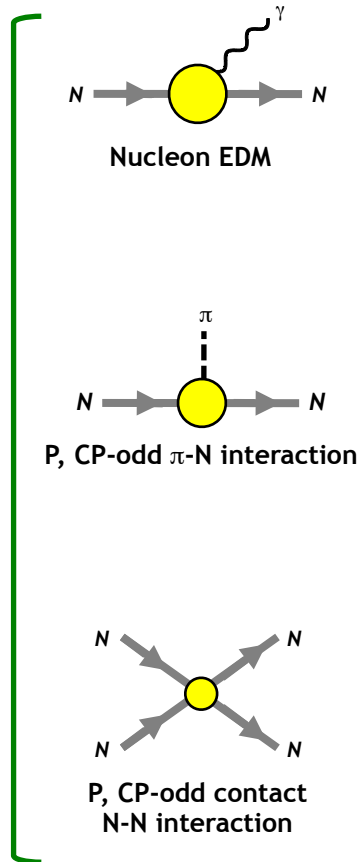
Hadronic CPV : from TeV to hadron level

Nuclear level
inputs

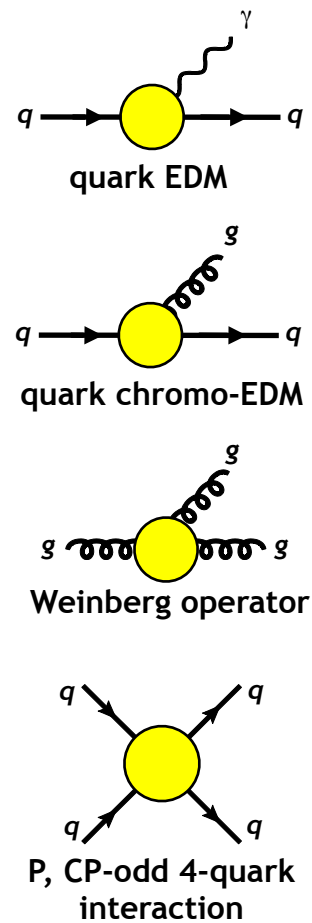


Chiral
EFT

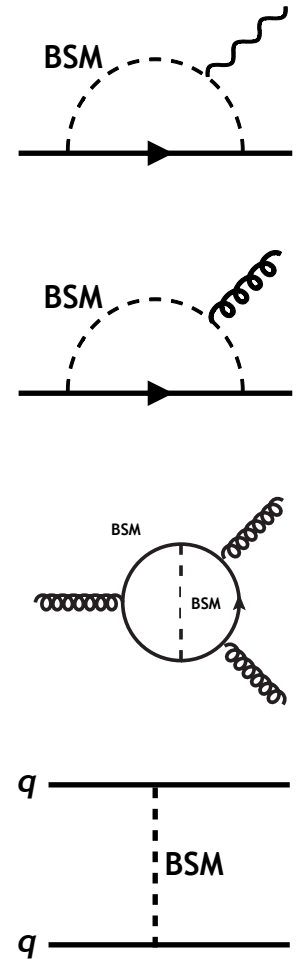
CPV hadron EFT



GeV scale CPV QCD



TeV scale CPV



RGE

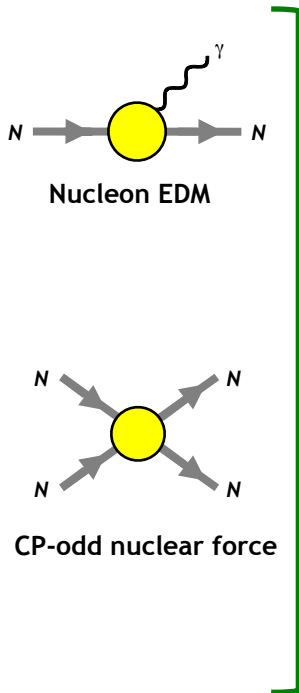
To nuclear level
calculation

QCD calculations

E. Mereghetti et al., Phys. Lett. B **696** (2011) 97;
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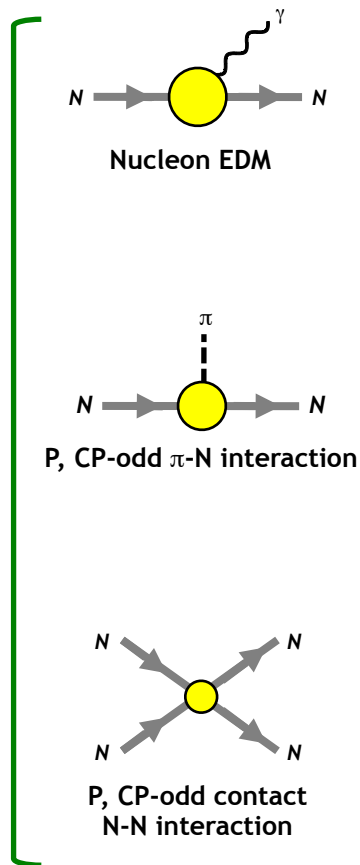
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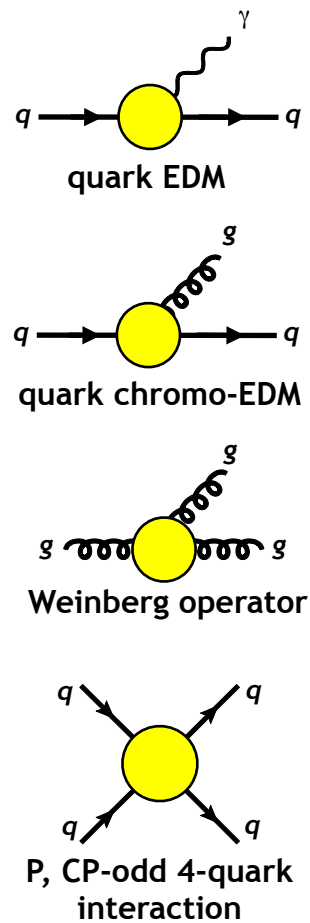


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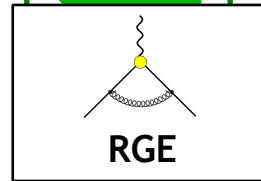
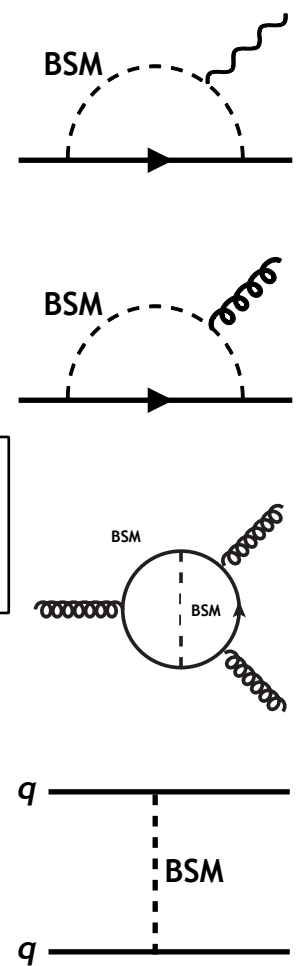
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TeV scale CPV



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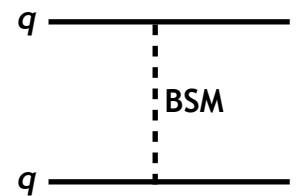
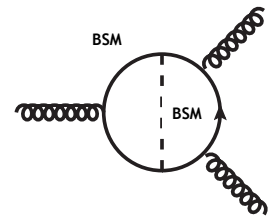
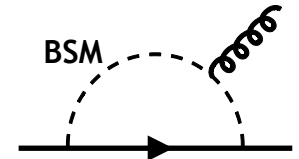
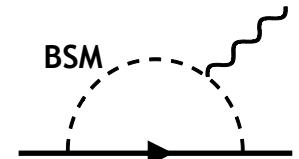
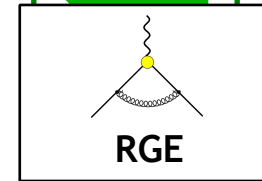
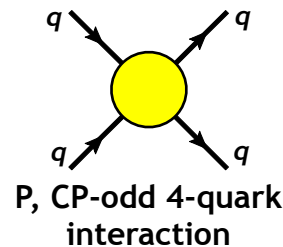
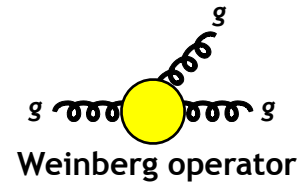
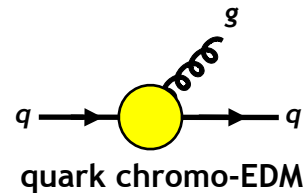
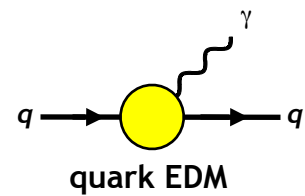
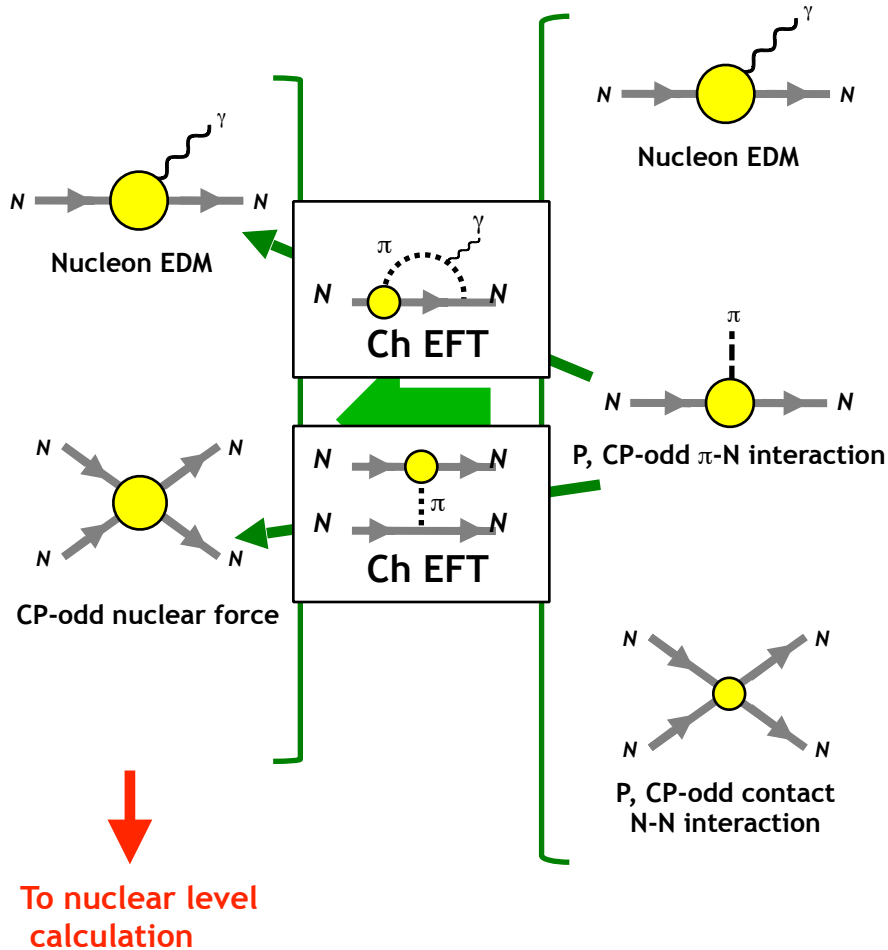
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CPV hadron EFT

GeV scale CPV QCD

TeV scale CPV



QCD calculations

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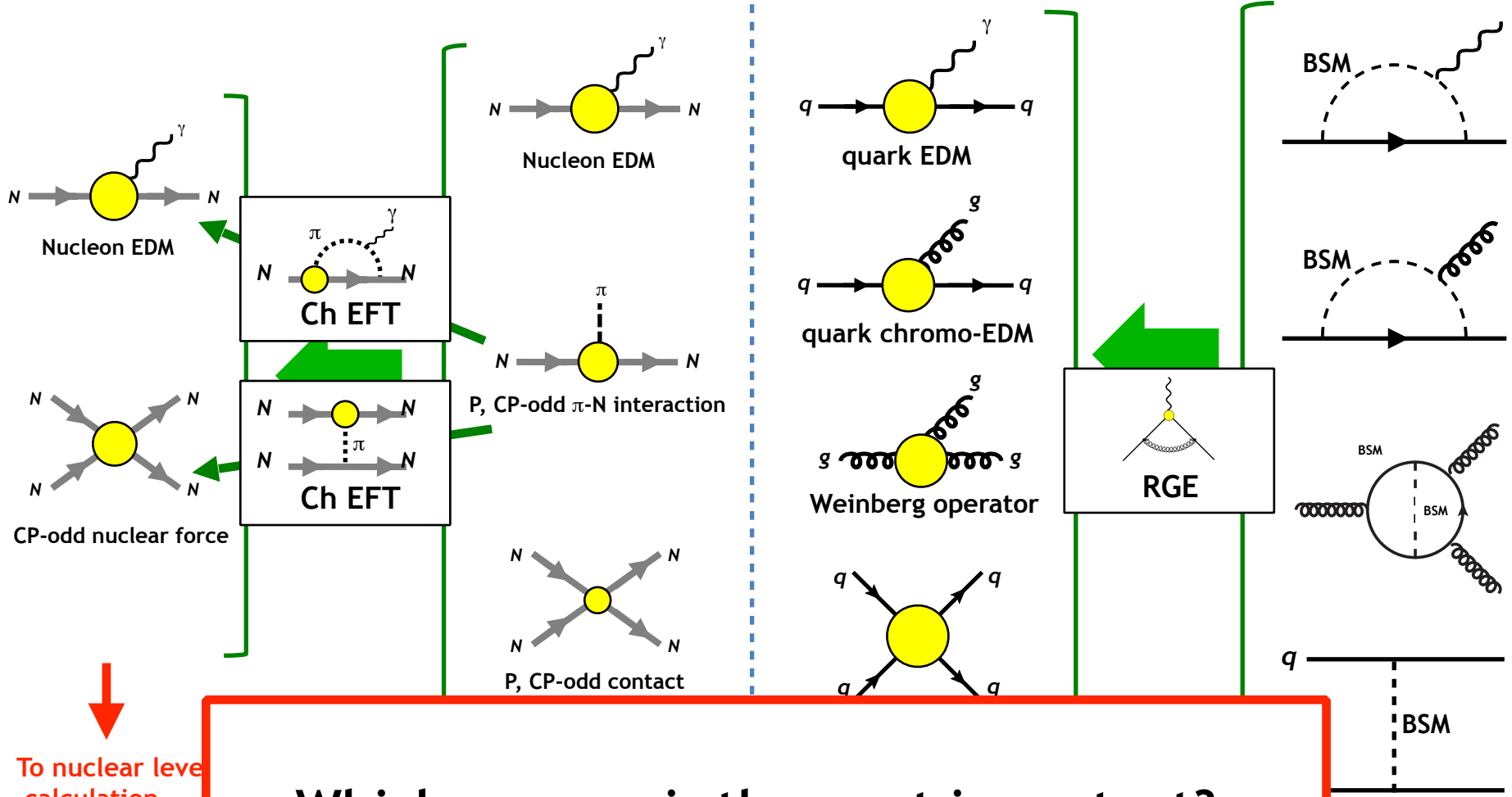
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Which process is the most important?

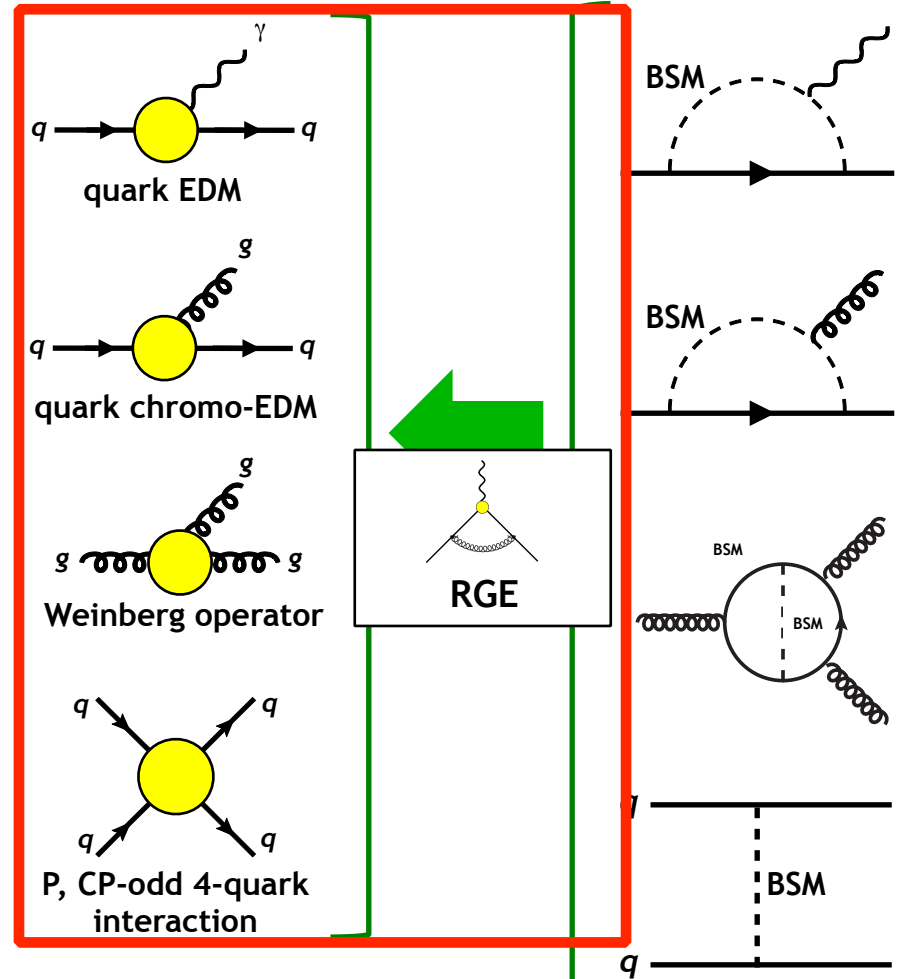
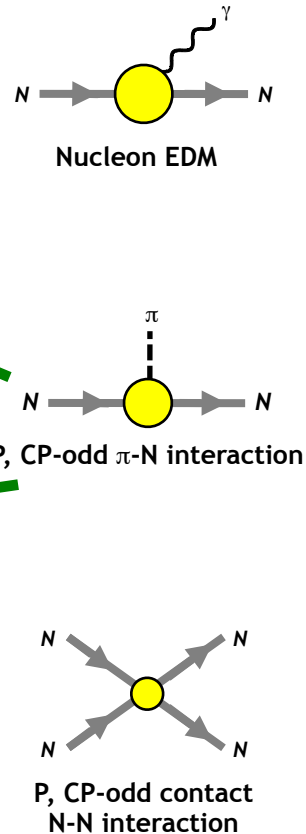
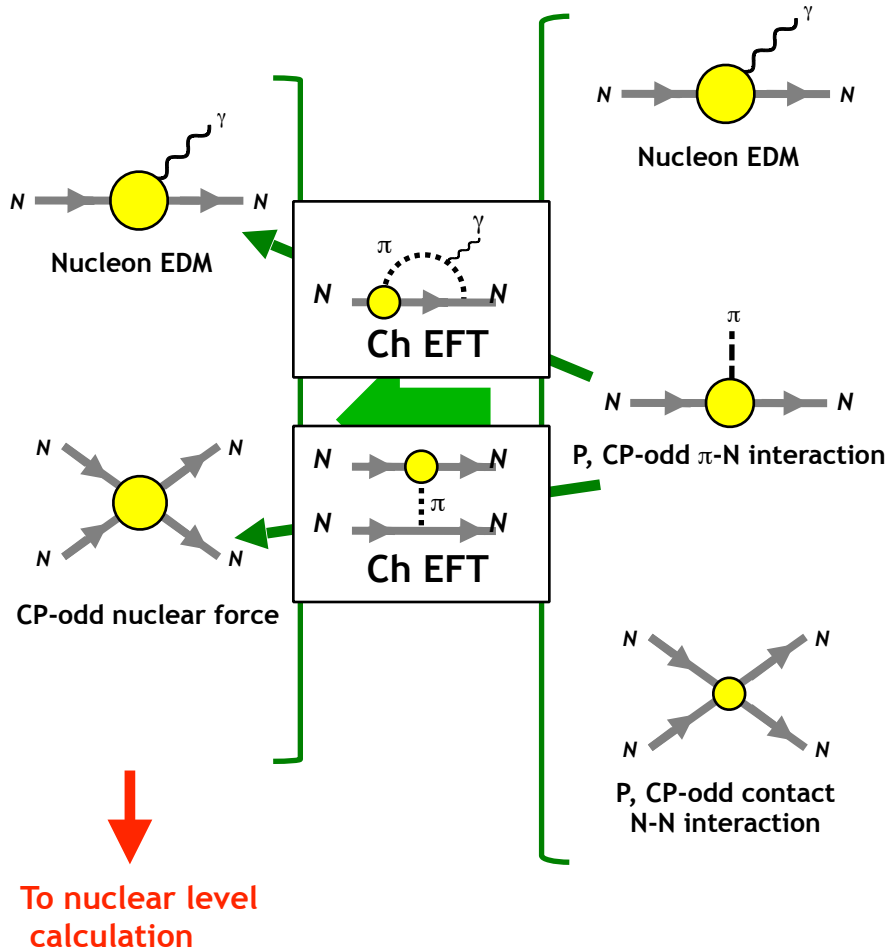
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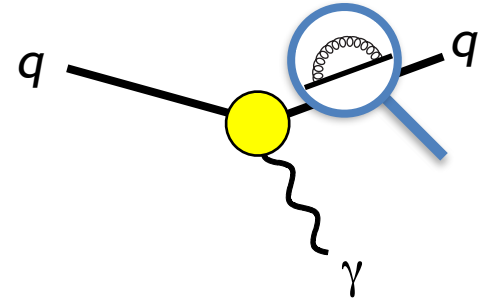
Renormalization group evolution of CPV QCD operators

Quark-gluon interactions change in scale due to quantum corrections

Calculated in perturbative QCD

- Scalar density:

$$\begin{array}{ccc} \bar{q}q & \longrightarrow & 2 \bar{q}q \\ \mu = 1 \text{ TeV} & & \mu = 1 \text{ GeV} \end{array}$$



- Quark EDM:

$$\begin{array}{ccc} d_q & \longrightarrow & 0.8 d_q \\ \mu = 1 \text{ TeV} & & \mu = 1 \text{ GeV} \end{array}$$

- Weinberg operator:

$$\begin{array}{ccc} \text{Diagram 1} & \longrightarrow & 0.34 \text{ Diagram 2} - 0.59 \text{ Diagram 3} + 0.52 \text{ Diagram 4} \\ \mu = 1 \text{ TeV} & & \mu = 1 \text{ GeV} \end{array}$$

Diagram 1: A quark line with a gluon loop (g) and a photon (gamma) attached to the loop.

Diagram 2: A quark line with a gluon loop (g) and a photon (gamma) attached to the loop.

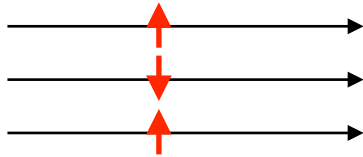
Diagram 3: A quark line with a gluon loop (g) and a photon (gamma) attached to the loop.

Diagram 4: A quark line with a gluon loop (g) and a photon (gamma) attached to the loop.

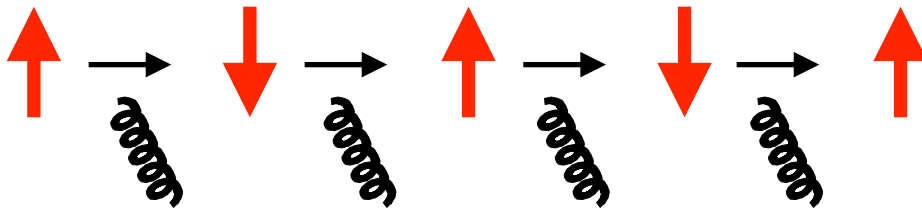
Roughly, **scalar increases** and **spin decreases** when scale goes down

Rough tendency: scalar and spin

● Spin : suppression

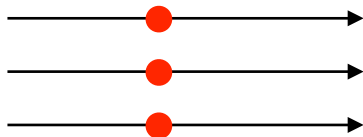


In composite states : **Pairing of spin**
⇒ No enhancement

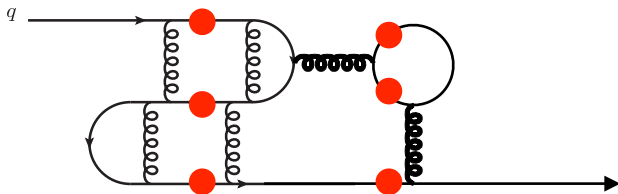


Qu. corrections: **superposition**
of spin flips after
gluon emissions/absorptions
⇒ **Suppressed**

● Scalar density ($\bar{\psi}\psi$) : enhancement



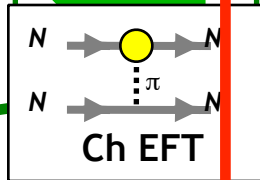
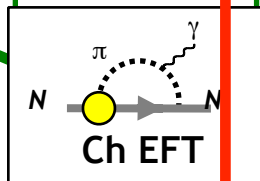
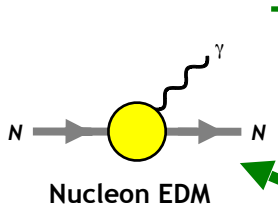
In composite states: coherent sum
⇒ **Enhanced!**



Qu. Corrections: $\bar{\psi}\psi$ of particles and
antiparticles have **same sign**
⇒ **Enhanced** by Z-graphs, loops

Hadronic CPV : from TeV to hadron level

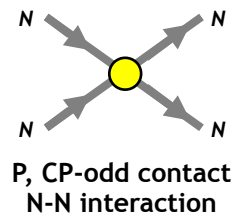
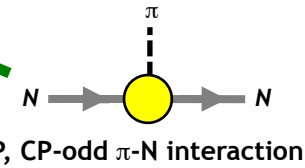
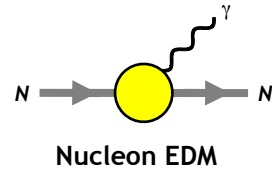
Nuclear level
inputs



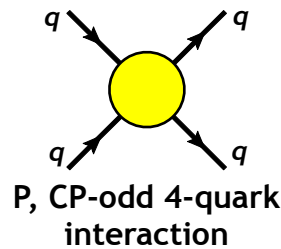
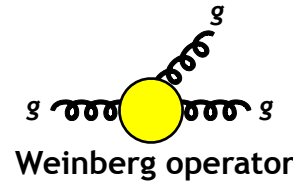
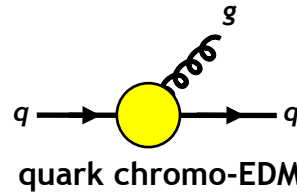
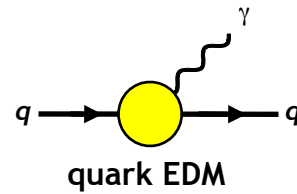
CP-odd nuclear force

To nuclear level
calculation

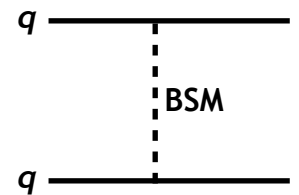
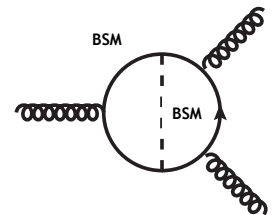
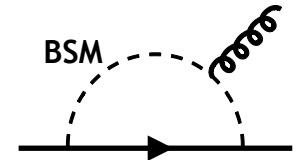
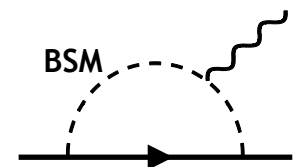
CPV hadron EFT



GeV scale CPV QCD



TeV scale CPV

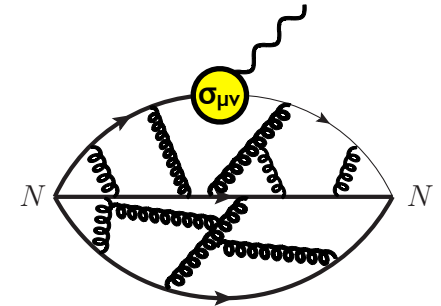


QCD calculations

E. Mereghetti et al., Phys. Lett. B **696** (2011) 97;
J. de Vries et al., Phys. Lett. B **766**, 254 (2017).

Quark EDM to nucleon EDM : lattice QCD

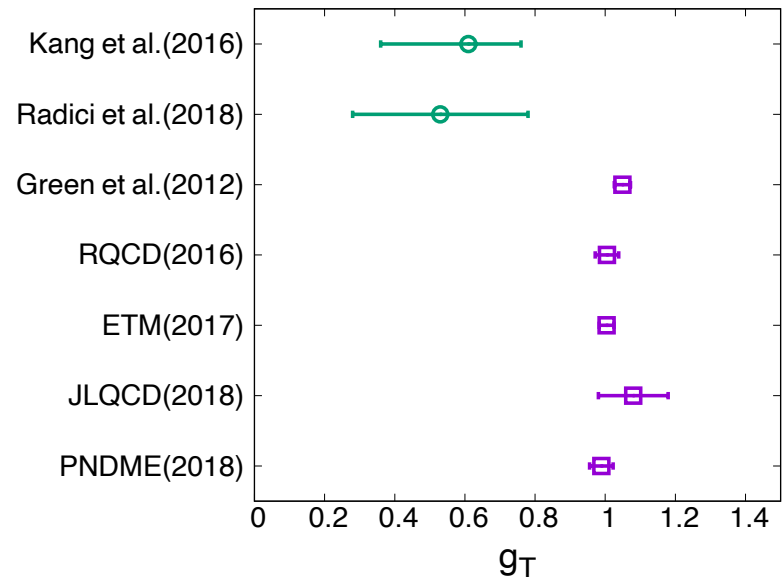
Quark EDM contribution to nucleon EDM is given by
the tensor charge $\langle N | \bar{q} \sigma_{\mu\nu} q | N \rangle$



⇒ Lattice QCD results available

⇒ Accurate up to 10% error
(Little disagreement with pQCD)

⇒ No particular enhancement



Lattice calculations of other interactions are still very difficult
though some first studies started to appear

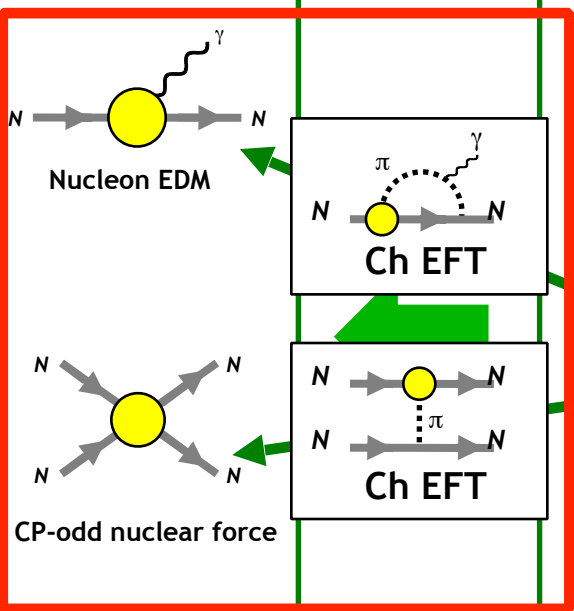
Hadronic CPV : from TeV to hadron level

Nuclear level
inputs

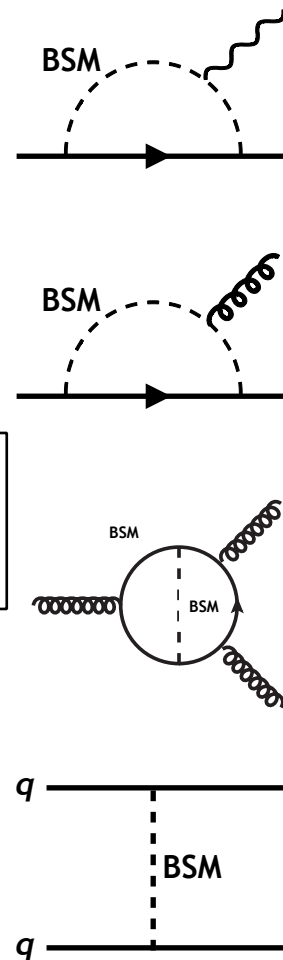
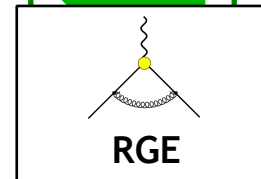
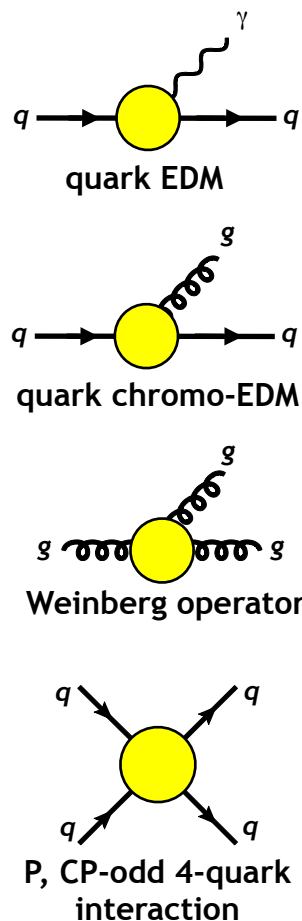
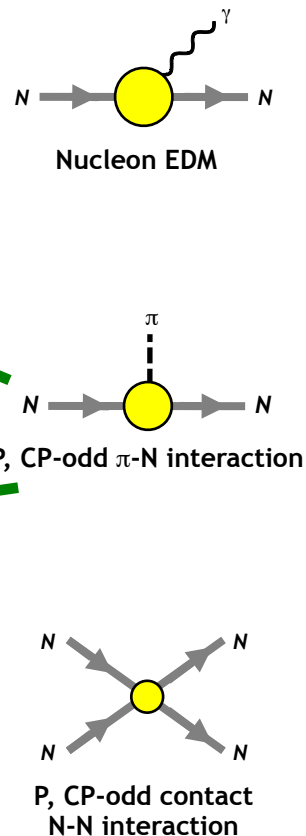
CPV hadron EFT

GeV scale CPV QCD

TeV scale CPV



To nuclear level
calculation



QCD calculations

E. Mereghetti et al., Phys. Lett. B **696** (2011) 97;
J. de Vries et al., Phys. Lett. B **766**, 254 (2017).

Chiral perturbation theory (Chiral EFT)

$$\mathcal{L}_{\text{QCD}} = \underbrace{-\frac{1}{4}F_{\mu\nu,a}F_a^{\mu\nu} + \sum_{\psi=u,d,\dots} \bar{\psi}i\not{D}\psi}_{\text{Chiral symmetric}} - \underbrace{m\bar{\psi}\psi}_{\text{Chiral symmetry breaking}}$$

Chiral symmetric

Chiral symmetry breaking

QCD has small u, d masses ($\sim\text{MeV}$) compared to Λ_{QCD} ($\sim 200\text{ MeV}$)

$\Rightarrow$ Expansion of hadron level interactions in $m_{u,d}/\Lambda_{\text{QCD}} \sim 0.01$

$\Rightarrow$ Can organize and fit hadronic interactions with powers of $m_{u,d}/\Lambda_{\text{QCD}}$

CPV hadronic interactions may also break chiral symmetry

E.g. $\mathcal{L}_{\pi NN} = \bar{g}_{\pi NN}^{(0)}\pi_a\bar{N}\tau_a N + \bar{g}_{\pi NN}^{(1)}\pi_0\bar{N}N \Rightarrow \text{Break chiral sym.}$

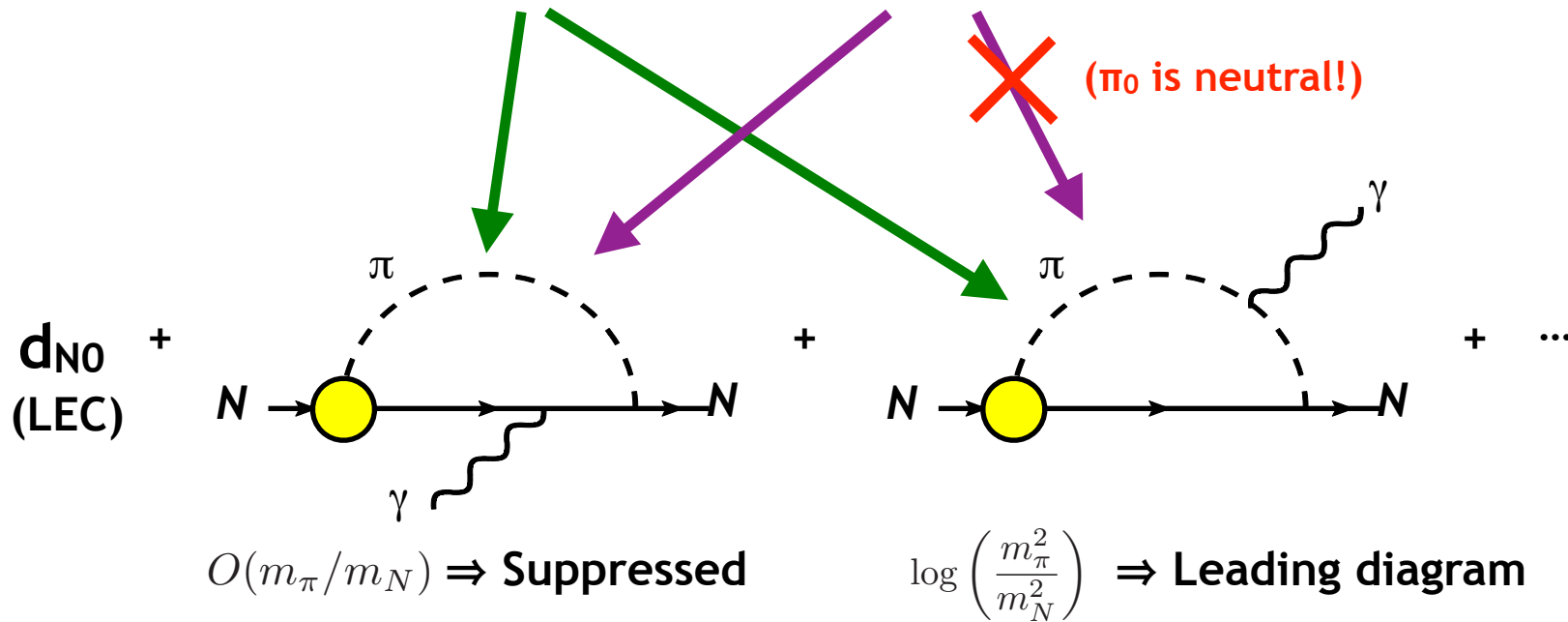
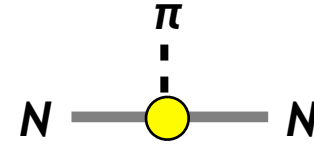
$\mathcal{L}_{NN} = \bar{C}_1\bar{N}i\gamma_5 N \bar{N}N \Rightarrow \text{Conserve chiral sym.}$

 Quantification using diagrams in terms of

$m_{u,d}/\Lambda_{\text{QCD}}$, isospin, CPV, ... is possible ! (Chiral EFT)

Chiral EFT analysis of nucleon EDM

$$\mathcal{L}_{\pi NN} = \underbrace{\bar{g}_{\pi NN}^{(0)} \pi_a \bar{N} \tau_a N}_{\text{Isoscalar}} + \underbrace{\bar{g}_{\pi NN}^{(1)} \pi_0 \bar{N} N}_{\text{Isovector}}$$

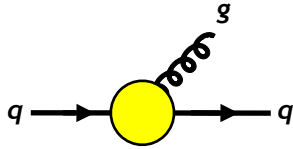


$\Rightarrow$ Essentially loop effect, no notable enhancement

$\Rightarrow$ Only $\bar{g}^{(0)}$ is important (sensitivity to specific BSM)

Chromo-EDM $\rightarrow$ CPV π NN interaction in chiral EFT

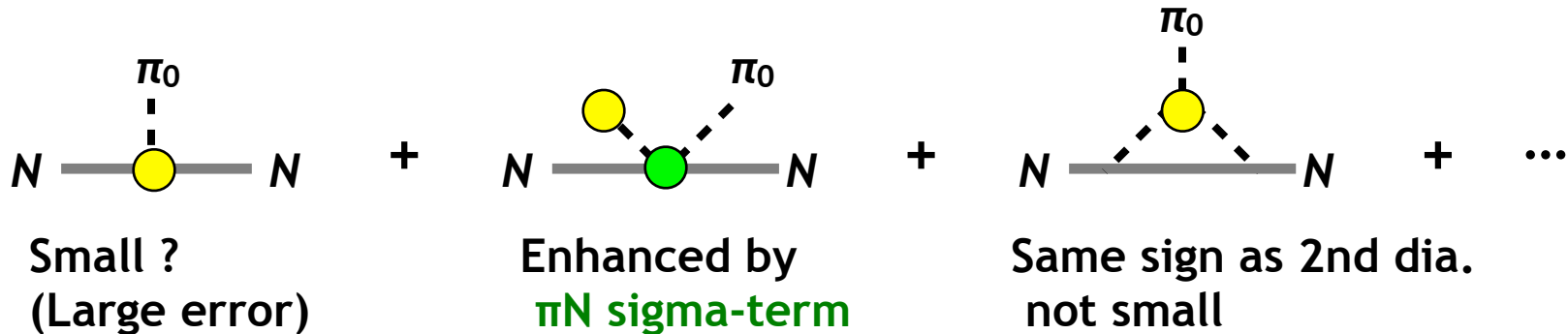
cEDM



is chiral sym. breaking, may contribute to both π NN

$$\mathcal{L}_{\pi NN} = \bar{g}_{\pi NN}^{(0)} \pi_a \bar{N} \tau_a N + \underline{\bar{g}_{\pi NN}^{(1)} \pi_0 \bar{N} N}$$

Chiral EFT analysis : **isovector** CPV π NN interaction is enhanced



Isovector CPV π NN interaction suppressed for nucleon EDM,
but large for, **atomic/nuclear EDM** (via CPV pion-exchange nuclear force)

$$d_p \approx 1 \times e(d_u^c - d_d^c)$$

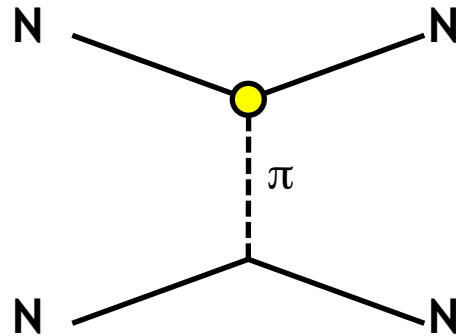
$$d_{\text{He}} \approx 16 \times e(d_u^c - d_d^c)$$



Example of ${}^3\text{He}$ EDM :
nuclear level CPV is much larger

P, CP-odd nuclear force

P, CP-odd nuclear force : one-pion exchange most important



$$\sim \frac{1}{q^2 - m_\pi^2} \bar{N} N \bar{N} i \gamma_5 N$$

● P, CP-odd Hamiltonian (3-types):

$$H_{CPV} = -\frac{g_{\pi NN}}{8\pi m_N} \left[\underbrace{(\bar{g}_{\pi NN}^{(0)} \tau_a \cdot \tau_b)}_{\text{Isoscalar}} + \underbrace{\bar{g}_{\pi NN}^{(2)} (\tau_a \cdot \tau_b - 3\tau_a^z \tau_b^z)}_{\text{Isotensor}} \right] (\vec{\sigma}_a - \vec{\sigma}_b) + \underbrace{\bar{g}_{\pi NN}^{(1)} (\tau_a^z \vec{\sigma}_a - \tau_b^z \vec{\sigma}_b)}_{\text{Isovector}} \cdot \vec{\nabla} \frac{e^{-m_\pi r}}{r}$$

Probed by atomic/nuclear EDM, CPV β decay, TV neutron optics, ...

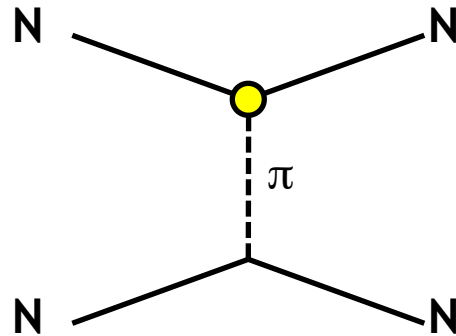
● Isovector CPVNN enhanced by nucleon number (via sigma-term)

⇒ The most sensitive process for quark-gluon CPV

● Enhancement by resonances, deformation at nuclear level

P, CP-odd nuclear force

P, CP-odd nuclear force : one-pion exchange most important



Isoscalar CPV NN-int. :

Axion DM, θ -term, ... $i\gamma_5 N$

● P, CP-odd Hamiltonian (3-types):

$$H_{CPV} = -\frac{g_{\pi NN}}{8\pi m_N} \left[\underbrace{(\bar{g}_{\pi NN}^{(0)} \tau_a \cdot \tau_b)}_{\text{Isoscalar}} + \underbrace{\bar{g}_{\pi NN}^{(2)} (\tau_a \cdot \tau_b - 3\tau_a^z \tau_b^z)}_{\text{Isotensor}} \right] (\vec{\sigma}_a - \vec{\sigma}_b) + \underbrace{\bar{g}_{\pi NN}^{(1)} (\tau_a^z \vec{\sigma}_a - \tau_b^z \vec{\sigma}_b)}_{\text{Isovector}} \cdot \vec{\nabla} \frac{e^{-m_\pi r}}{r}$$

Probed by atomic/nuclear EDM, CPV β decay, TV neutron optics, ...

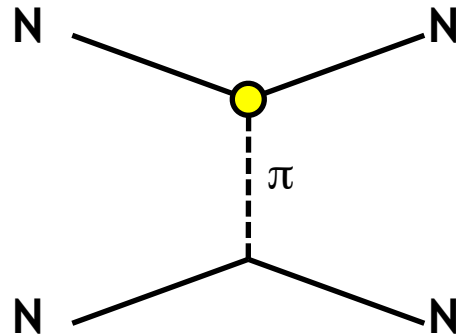
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Isoscalar CPV NN-int. :
Axion DM, θ -term, ...

Isovector CPV NN-int. :
Higgs, SUSY, ...

● P, CP-odd Hamiltonian (3-types):

$$H_{CPV} = -\frac{g_{\pi NN}}{8\pi m_N} \left[\underbrace{\bar{g}_{\pi NN}^{(0)}}_{\text{Isoscalar}} \tau_a \cdot \tau_b + \underbrace{\bar{g}_{\pi NN}^{(2)}}_{\text{Isotensor}} (\tau_a \cdot \tau_b - 3\tau_a^z \tau_b^z) \right] (\vec{\sigma}_a - \vec{\sigma}_b) + \underbrace{\bar{g}_{\pi NN}^{(1)}}_{\text{Isovector}} (\tau_a^z \vec{\sigma}_a - \tau_b^z \vec{\sigma}_b) \cdot \vec{\nabla} \frac{e^{-m_\pi r}}{r}$$

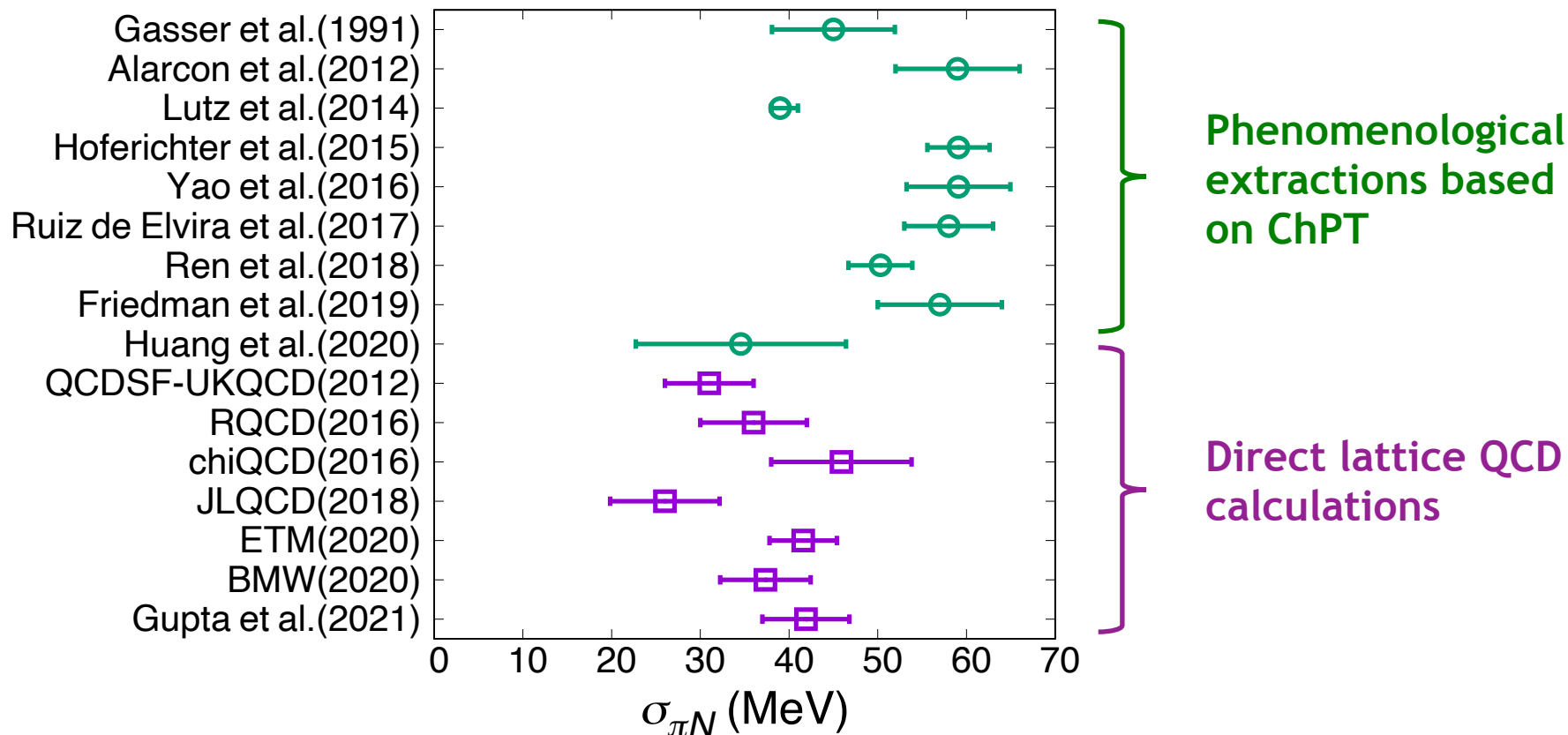
Probed by atomic/nuclear EDM, CPV β decay, TV neutron optics, ...

● Isovector CPVNN enhanced by nucleon number (via sigma-term)

$\Rightarrow$ The most sensitive process for quark-gluon CPV

● Enhancement by resonances, deformation at nuclear level

Present status of the calculations of πN sigma-term

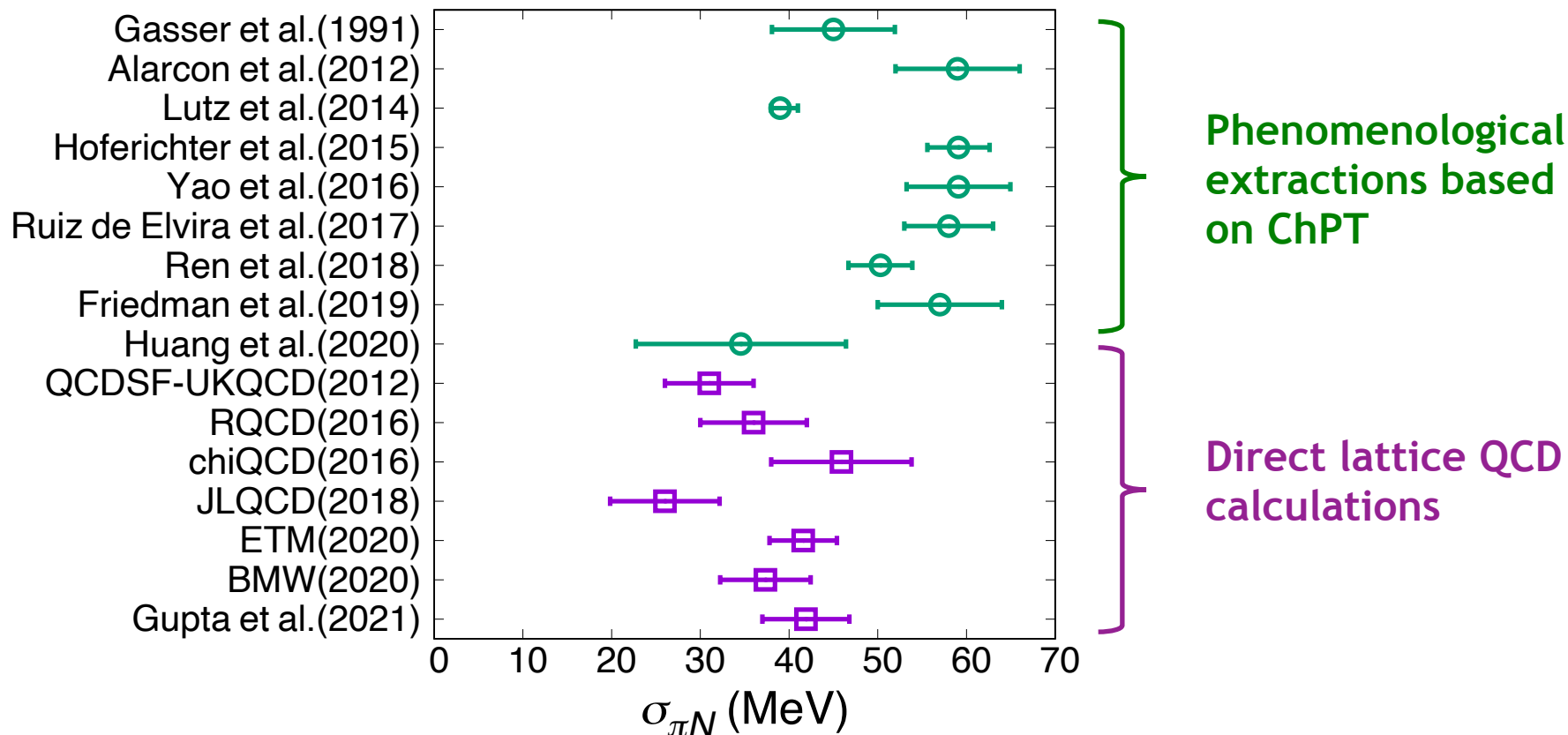


Extractions of $\sigma_{\pi N} \equiv \frac{m_u + m_d}{2} \langle N | \bar{u}u + \bar{d}d | N \rangle$

suggests $\langle N | \bar{u}u + \bar{d}d | N \rangle \sim 10$ (c.f. scalar charge)

⇒ Enhancement of CPV nuclear force via $\sigma_{\pi N}$!

Present status of the calculations of πN sigma-term

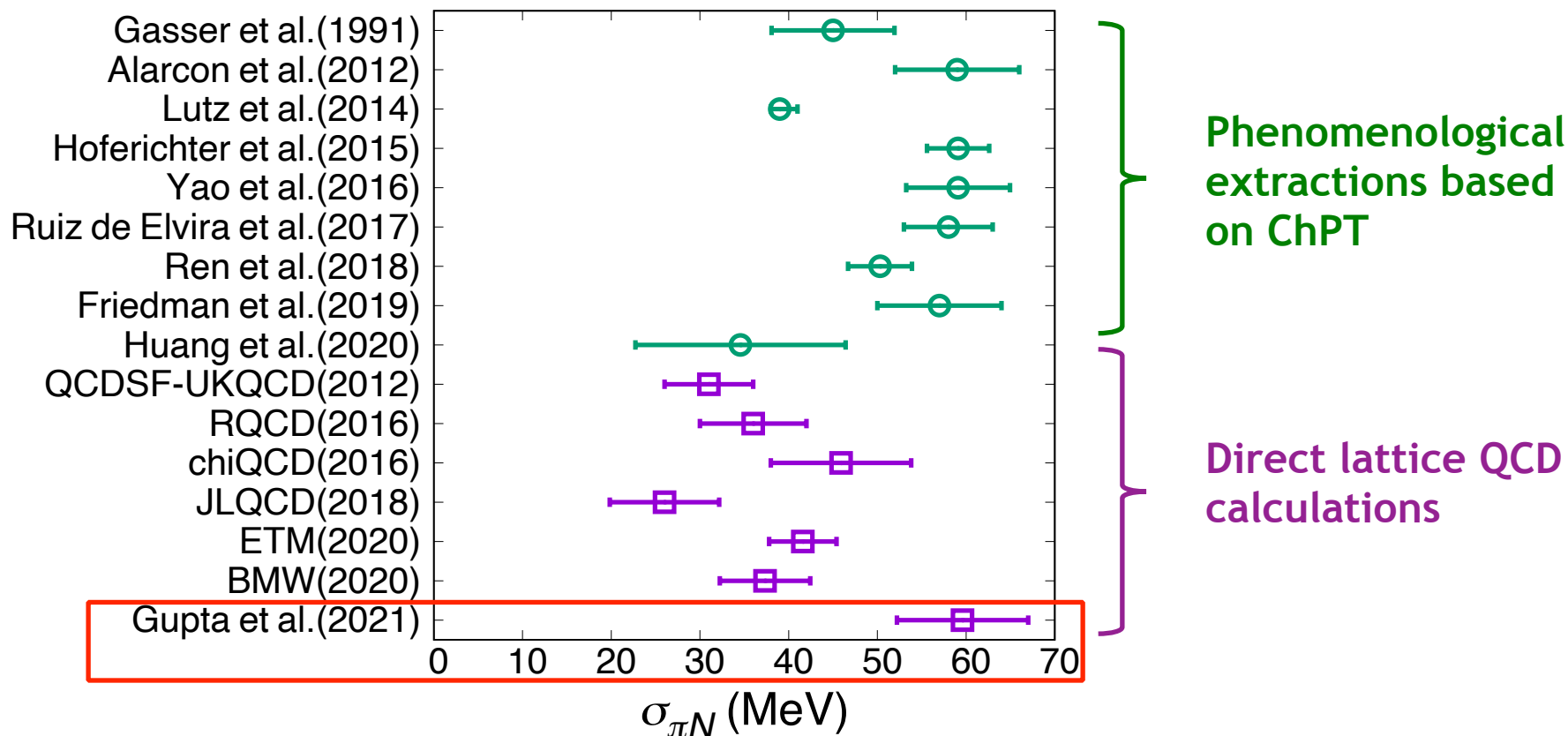


Phenomenological extractions : $\sigma_{\pi N} \sim 60$ MeV

Lattice QCD calculations : $\sigma_{\pi N} \sim 30$ MeV

⇒ Visible disagreement

Present status of the calculations of πN sigma-term



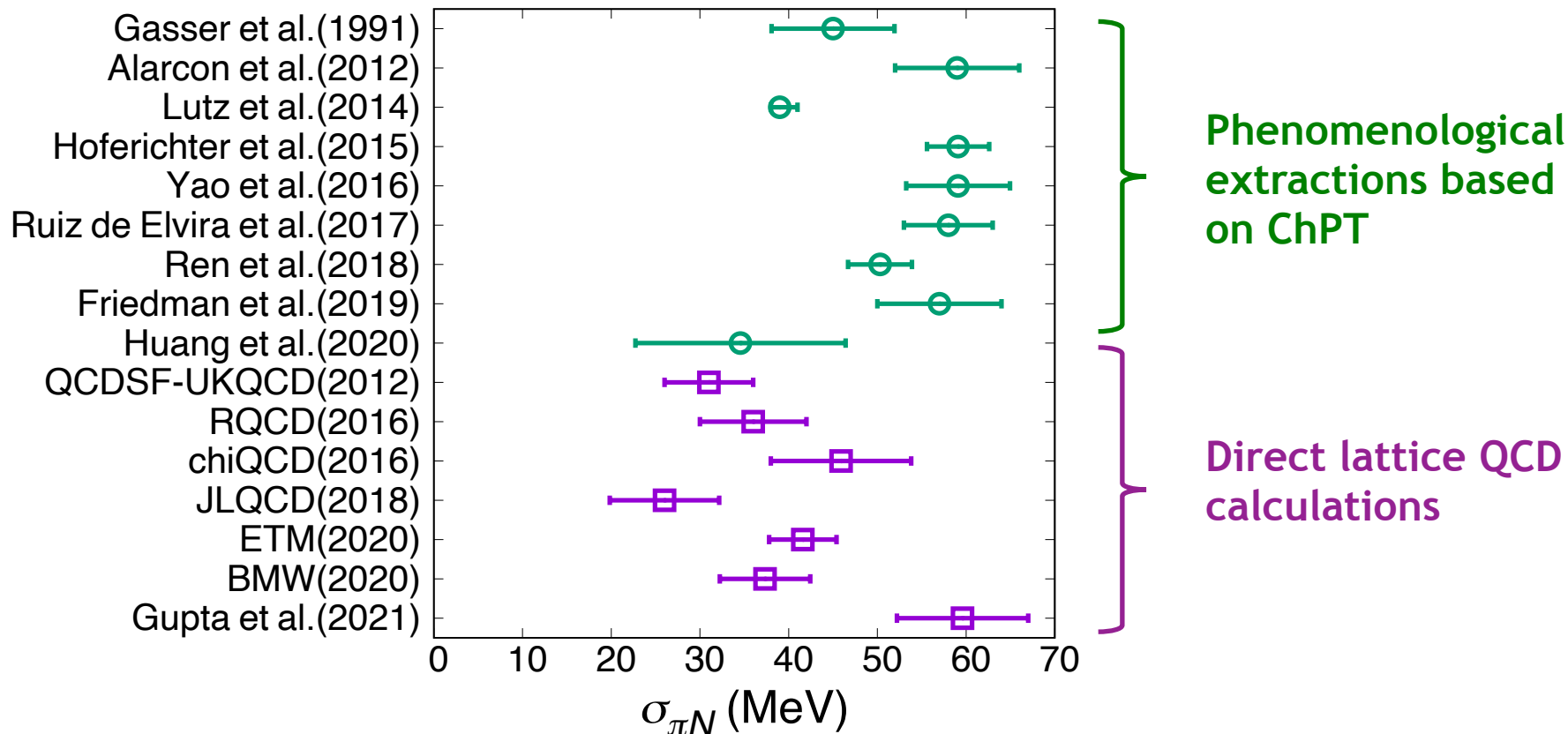
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⇒ Visible disagreement

⇒ Gupta et al. are suggesting some potential resolution??

Present status of the calculations of πN sigma-term



Phenomenological extractions : $\sigma_{\pi N} \sim 60$ MeV

Lattice QCD calculations : $\sigma_{\pi N} \sim 30$ MeV

⇒ Visible disagreement

Homework:

We have to resolve this problem for quantifying EDMs

Summary of hadronic CPV

No notable enhancement for quark EDM

Weinberg operator suppressed by RGE and
hadron matrix elements

 **Chromo-EDM contribution is the most enhanced!**

(Some 4-quark interactions may also be enhanced,
but model specific, like left-right sym. models)

Enhanced by πN sigma-term , still large uncertainty

**Theoretical challenge to
resolve the Strong CP problem**

Theta-term

$$\mathcal{L}_{\text{QCD}} = -\frac{1}{4}F_{\mu\nu,a}F_a^{\mu\nu} + \theta \frac{\alpha_s}{8\pi} F_{\mu\nu,a} \tilde{F}_a^{\mu\nu} \quad \tilde{F}_a^{\mu\nu} \equiv \frac{1}{2}\epsilon^{\mu\nu\rho\sigma}F_{\rho\sigma}^a$$

θ -term is like E.B $\Rightarrow$ Violate P and CP !

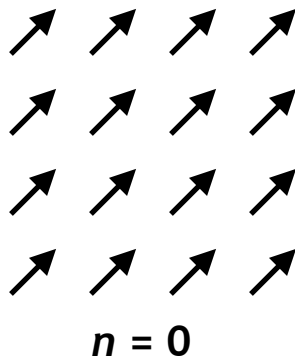
θ -term contains topological information of QCD

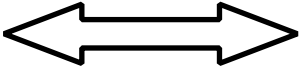
$$\int d^4x F_{\mu\nu,a} \tilde{F}_a^{\mu\nu} = \int d^4x \partial_\mu K^\mu = \Delta n \quad \text{Integer!}$$

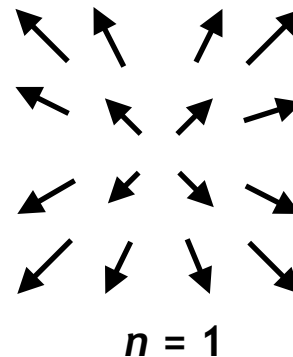
$$K_\mu \equiv \frac{\alpha_s}{8\pi} \epsilon_{\mu\nu\rho\sigma} \left[A_a^\nu F_a^{\rho\sigma} - \frac{g_s}{3} f_{abc} A_a^\nu A_b^\rho A_c^\sigma \right]$$

Integral of total derivative, but nonzero!

$\Rightarrow \Delta n$ = change of **winding number** of gauge configurations




Change of winding #



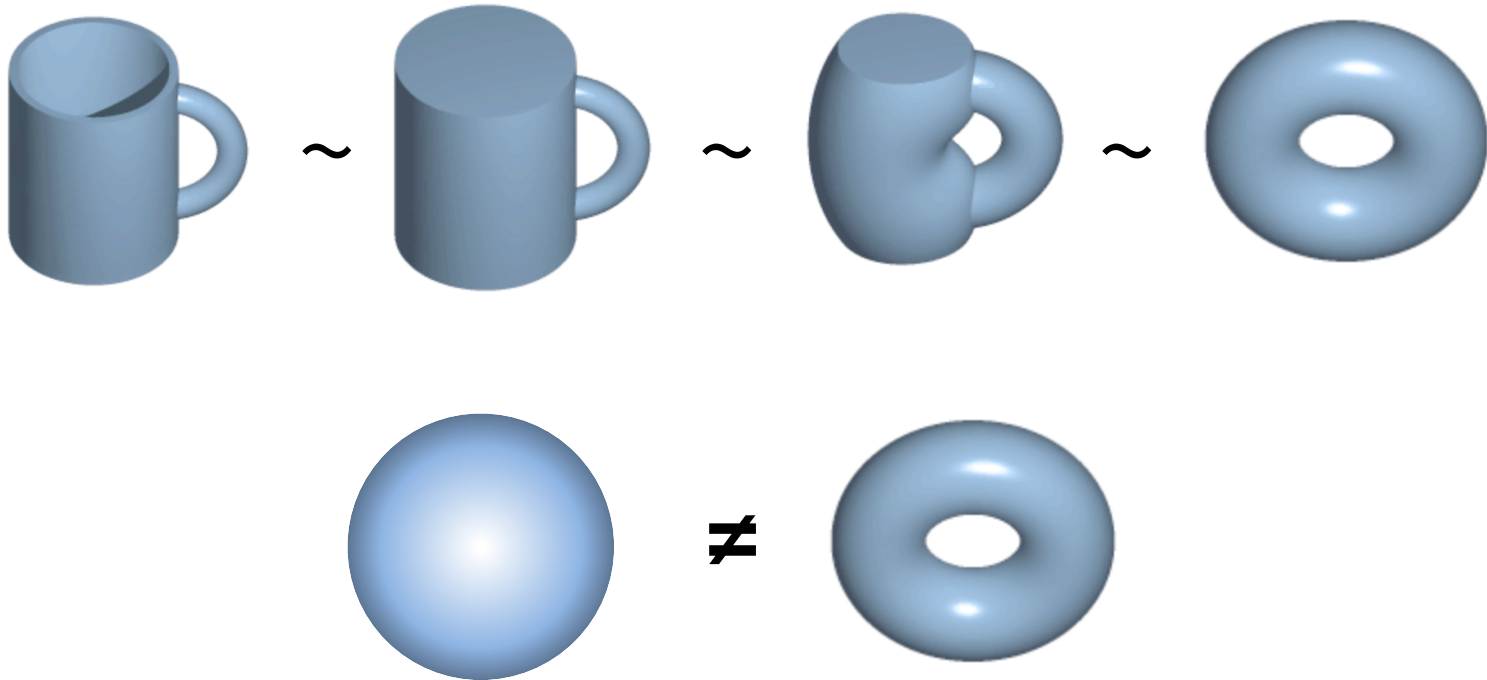
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θ -term is like E.B $\Rightarrow$ Violate P and CP !

θ -
Topology:



$$\left[{}_c A_a^\nu A_b^\rho A_c^\sigma \right]$$

Theta-term

$$\mathcal{L}_{\text{QCD}} = -\frac{1}{4}F_{\mu\nu,a}F_a^{\mu\nu} + \theta \frac{\alpha_s}{8\pi} F_{\mu\nu,a} \tilde{F}_a^{\mu\nu} \quad \tilde{F}_a^{\mu\nu} \equiv \frac{1}{2}\epsilon^{\mu\nu\rho\sigma} F_{\rho\sigma}^a$$

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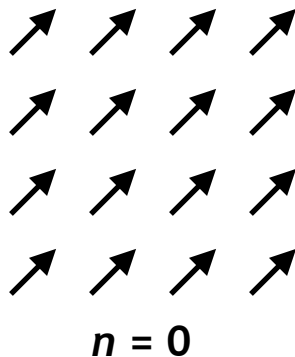
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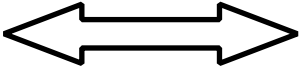
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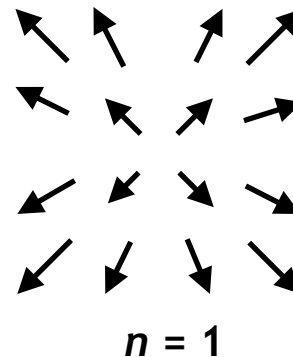
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Integral of total derivative, but nonzero!

$\Rightarrow \Delta n$ = change of **winding number** of gauge configurations




Change of winding #



Strong CP problem

θ -term : $\mathcal{L} = \theta \frac{g^2}{32\pi^2} F_{,a}^{\mu\nu} F_{,a}^{\rho\sigma} \epsilon_{\mu\nu\rho\sigma}$

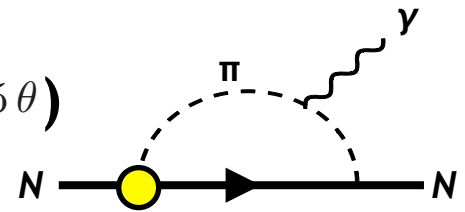
$\Rightarrow$ A “natural” interaction term of QCD, **θ should be O(1)**

Neutron EDM experimental data:

$$|d_n| < 1.8 \times 10^{-26} \text{ e cm}$$

C. Abel et al., Phys. Rev. Lett. 124, 081803 (2020).

Converting to θ via chiral EFT analysis ($\bar{g}_{\pi NN}^{(0)} \approx 0.015 \theta$)



$$|\theta| < 1 \times 10^{-10} \quad \text{...Why so small??}$$

➡ Strong CP problem

Resolutions proposed for the Strong CP problem

Conventional resolutions:

- Massless up quark
- Spontaneous CP breaking
- Axion mechanism

} Need BSM

Resolutions proposed for the Strong CP problem

Conventional resolutions:

- Massless up quark
- Spontaneous CP breaking
- Axion mechanism

} Need BSM

Recently, several resolutions within QCD were also proposed:

- CP conserving correlations under instanton background

W.-Y. Ai, J. S. Cruz, B. Garbrecht, C. Tamarit, Phys. Lett. B **822**, 136616 (2021).

- Color field screening by finite θ

Y. Nakamura and G. Schierholz, Nucl. Phys. B **986**, 116063 (2023).

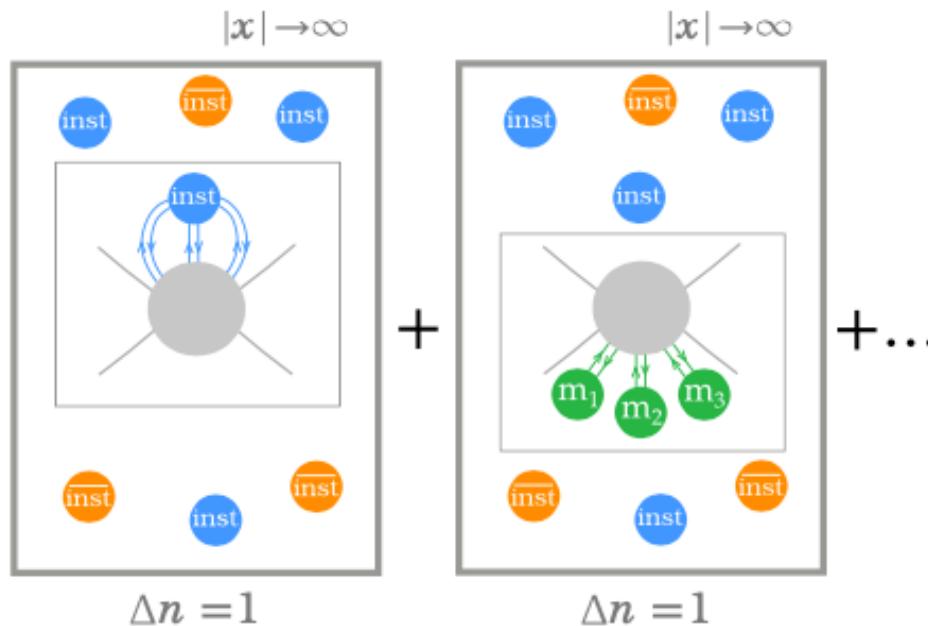
- Unphysical topological charge

NY, arXiv:2212.10994 [hep-th]; arXiv:2212.11820 [hep-ph].

CP conserving correlations under instanton background

Basic idea :

Calculate quark correlators with CP-odd mass within dilute instanton gas



Take first infinite volume
(take all instantons)

Then, sum over topology

⇒ Effect of θ cancels with fermion mass complex phase!

⇒ Effective θ is zero!

Color field screening by finite θ

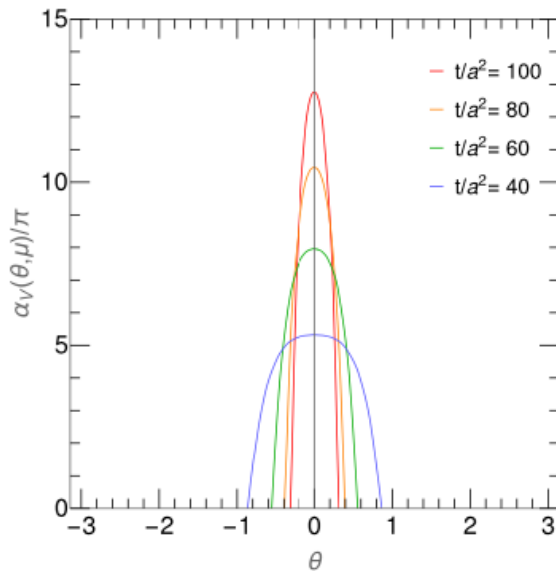
Basic idea : Witten effect

Dual superconductor : magnetic monopole condenses in color confining vacuum

Magnetic monopoles get electric charge under finite θ

E. Witten, Phys. Lett. B 86 (1979) 283.

⇒ Color screening, no monopole condensation,
deconfined colored particles with $|\theta| > 0$!



Calculation of α_s scale evolution
In lattice :

α_s damps for $|\theta| > 0$
with decreasing scale!
(α_s only grows for $\theta = 0$)

⇒ **Confinement only at $\theta = 0$!**

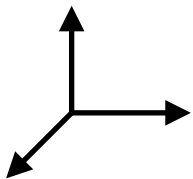
Unphysical topological charge

Topological charge is a triple product of gauge fields

$$\frac{\alpha_s}{8\pi} \int d^4x F_{\mu\nu,a} \tilde{F}_a^{\mu\nu} =$$

$$\frac{ig_s\alpha_s}{24\pi} \int d^3\vec{x} f_{abc} \epsilon_{ijk} A_{ia}(\vec{x}) A_{jb}(\vec{x}) A_{kc}(\vec{x}) \Big|_{t=-\infty}^{t=+\infty}$$

Time is frozen, restricted to 3-dimensional space

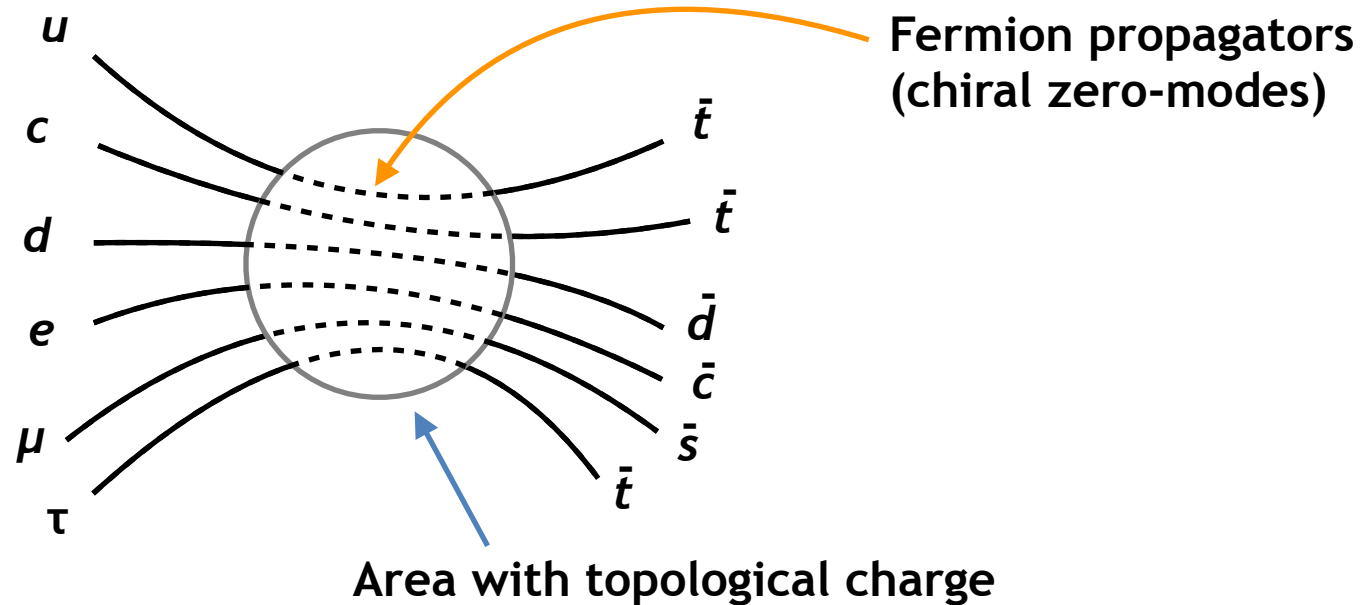
ϵ_{ijk} means  $\Rightarrow$ Covers all 3-dimensional directions!

$\Rightarrow$ Also covers the **unphysical gauge freedom direction!**
 $(\vec{A}_a \rightarrow \vec{A}_a + \vec{\nabla} \chi_a)$

 **Topological charge is unphysical !!**

Implication to sphaleron process

't Hooft vertex : multi-fermion vertex generated by chiral Dirac zero-modes



G. 't Hooft, PRL 37, 8 (1976), PRD 14, 3432 (1976).

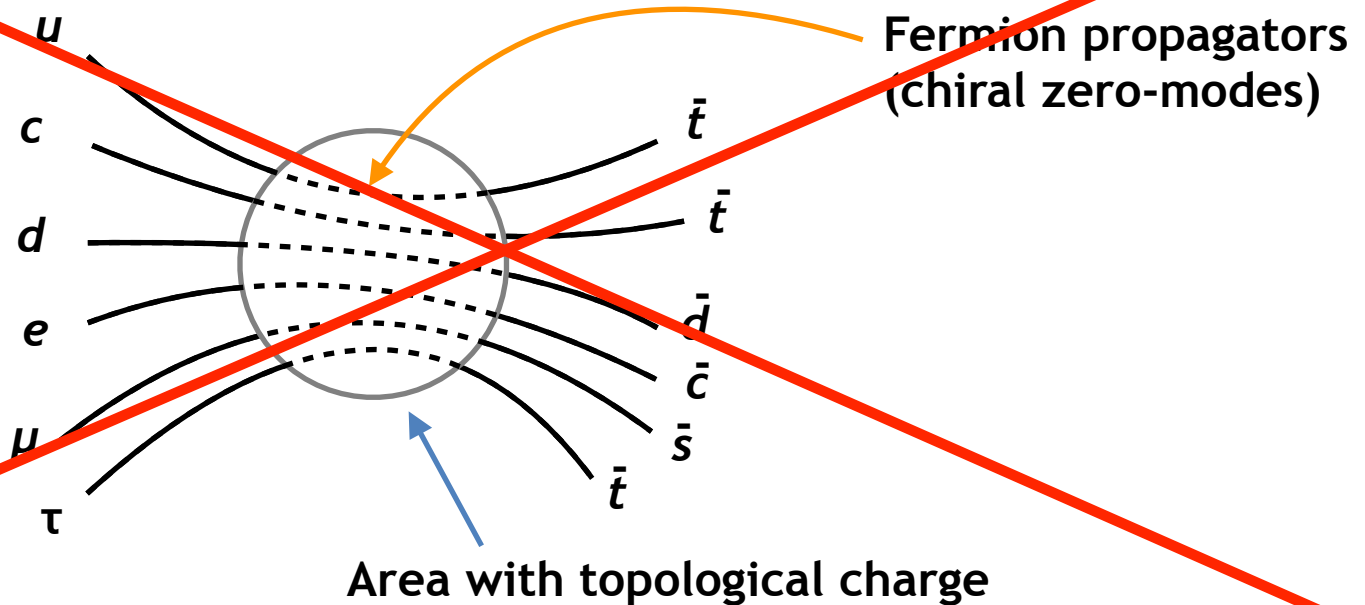
(Nonlocal) B+L violating 12-fermion interaction

⇒ Baryon and lepton numbers may be generated with weak interaction!

V. A. Kuzmin, V. A. Rubakov and M. E. Shaposhnikov, Phys. Lett. B 155 (1985), 36;
M. Fukugita and T. Yanagida, Phys. Lett. B 174 (1986) 45.

Implication to sphaleron process

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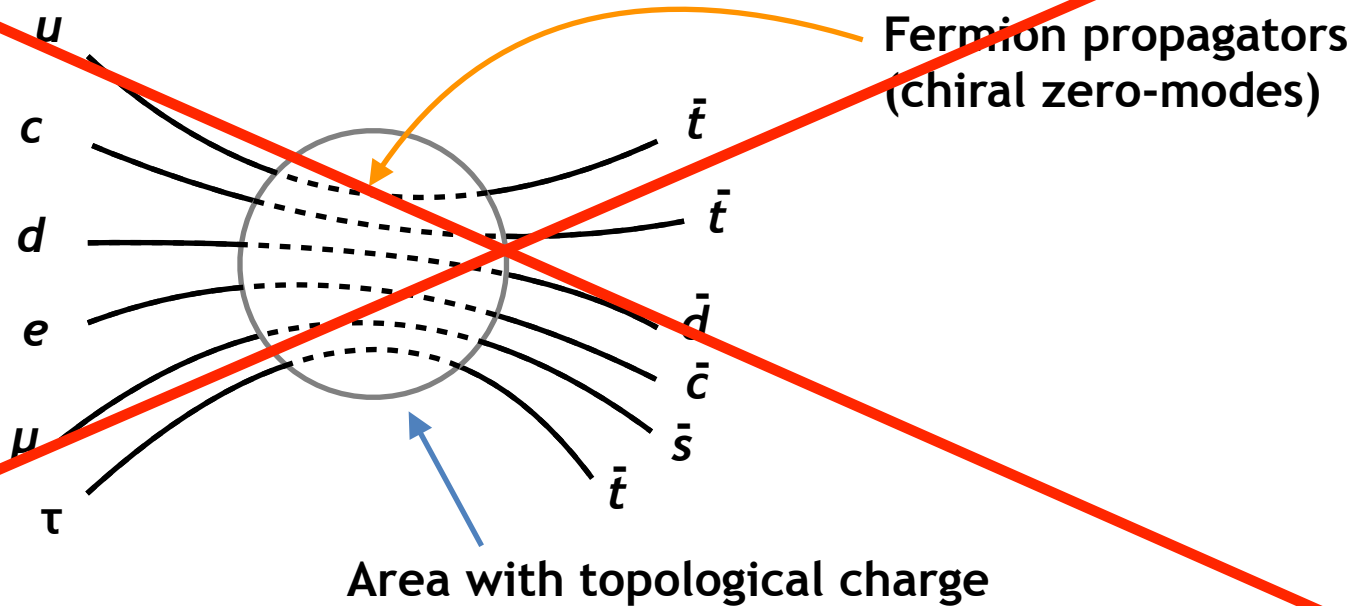
Atiyah-Singer theorem : topological charge = # of chiral Dirac zero-modes

$\Rightarrow$ 't Hooft vertex is unphysical !

since 't Hooft vertex is generated by the propagation of zero-modes

Implication to sphaleron process

't Hooft vertex : multi-fermion vertex generated by chiral Dirac zero-modes



G. 't Hooft, PRL 37, 8 (1976), PRD 14, 3432 (1976).

No matter - antimatter asymmetry generation at TeV scale

⇒ Baryogenesis happened at GUT ?

Quiz to physicists

Topological objects **spread across infinite space(-time)**

(just imagine infinite cosmic string)



(Fig. from BBC Science)

⇒ Their observations provide the **information of infinitely large distance**

⇒ **Zero energy-momentum resolution !**

Very fine low energy detector needed due to uncertainty principle

⇒ Impossible to observe ?

Please try to answer this question

Conjecture:

All topological quantities cannot be observed ?

Summary

- Non-removable phases of theory after field redefinition generates CPV.
- General BSM physics brings new phases: CP violation is a direct probe of BSM physics.
- Review of qualitative mechanisms of hadronic CP violation: chromo-EDM $\rightarrow$ CPV nuclear force is the most sensitive process.
- Review of Strong CP problem: theoretical resolutions without axions?

Future prospects:

- Homework : resolve the systematics of πN sigma-term.
- No sphaleron induced baryogenesis : new scenario in early Universe ?

End