

# Thermal precondensation in gauge-fermion theories

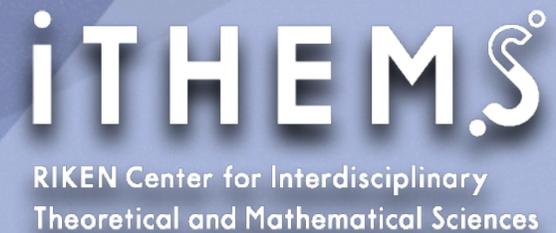
Álvaro Pastor Gutiérrez

KEK Theory Meeting on Particle Physics Phenomenology

17th of February 2026

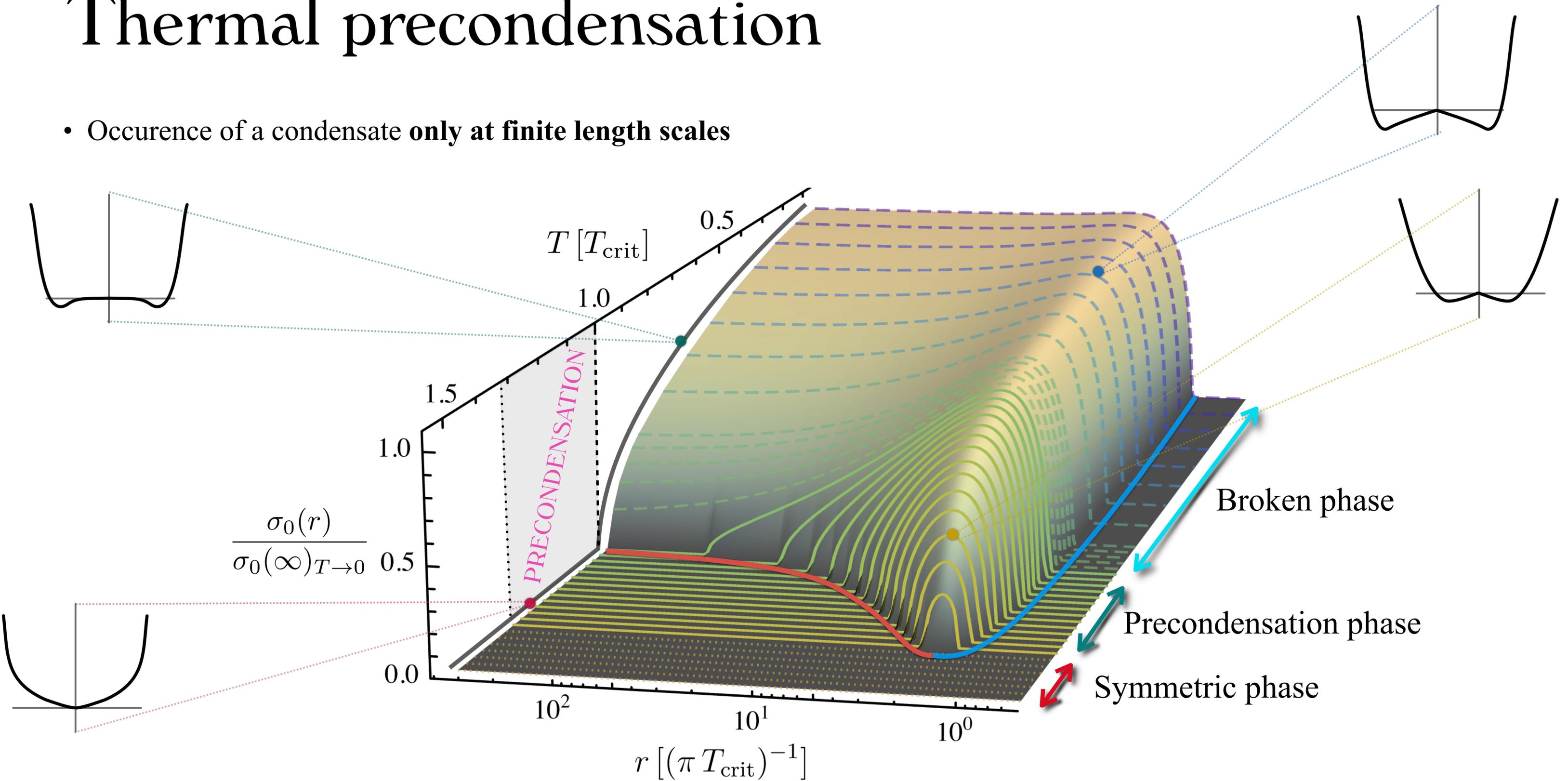


with Jan M. Pawłowski and Franz Sattler



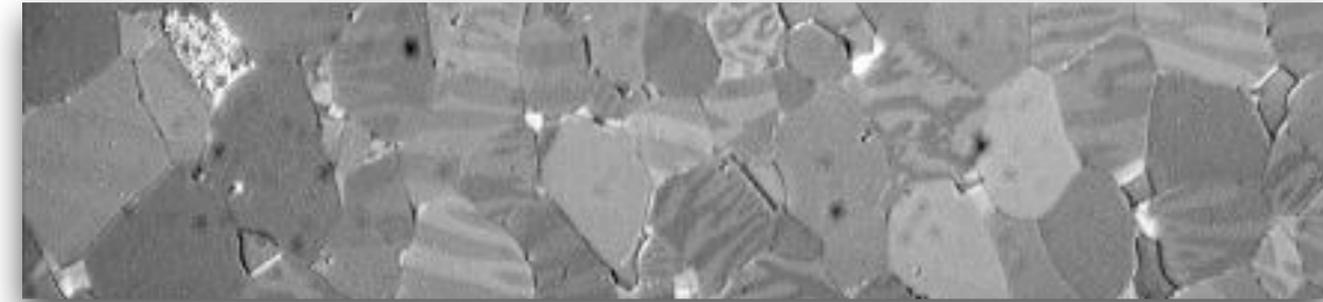
# Thermal precondensation

- Occurrence of a condensate **only at finite length scales**

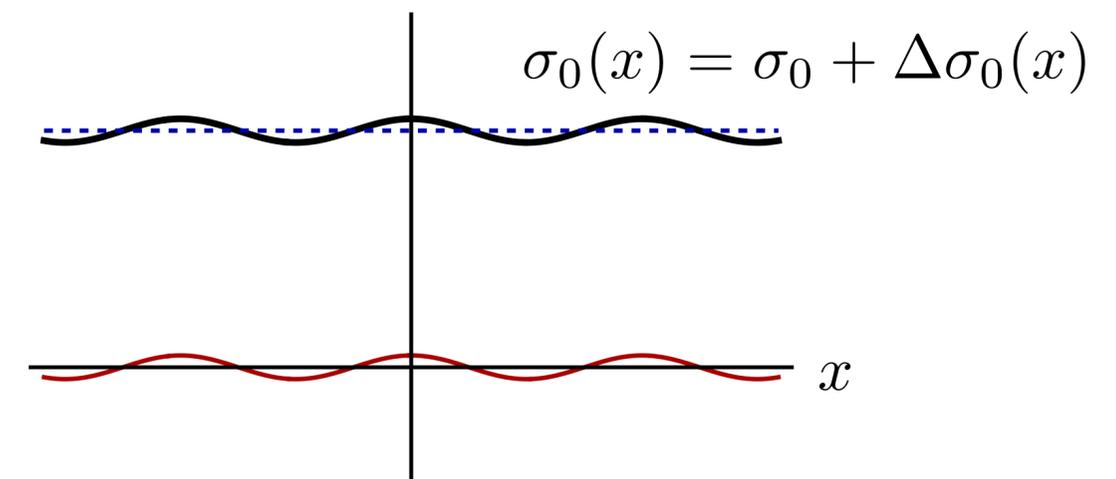
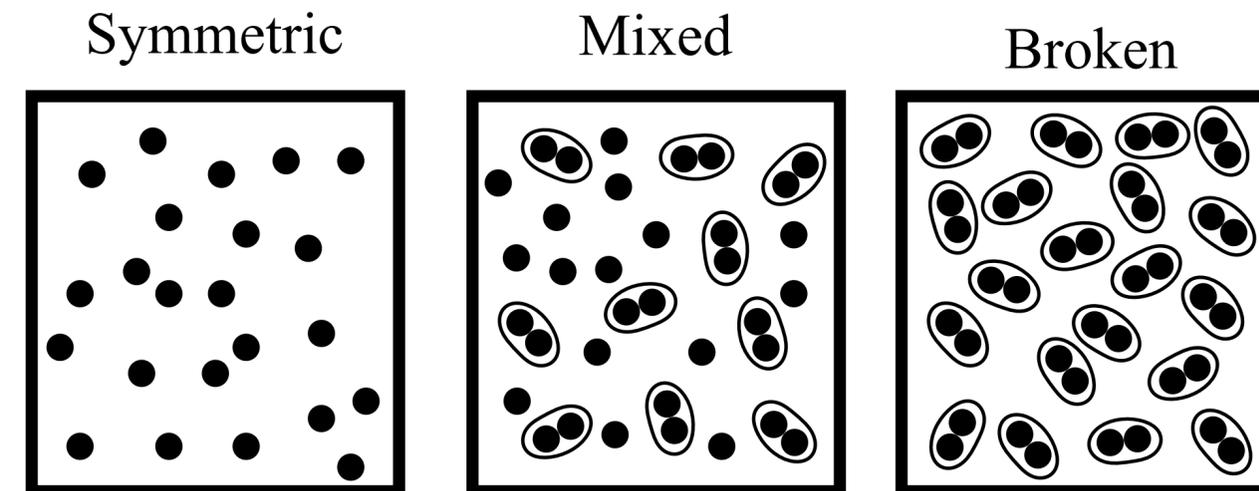


# Thermal precondensation

- Discussed in very different contexts
  - cold atoms: BCE-BCS crossover  
[Boettcher,Pawłowski,Diehl\[1204.4394\]](#)
  - condensed matter: Dirac semi-metals  
[Tolosa-Simeón,Classen,Scherer \[2503.04911\]](#)
  - gauge-fermion: diquark condensate in 2-colour QCD  
[Khan,Pawłowski,Rennecke,Scherer \[1512.03673\]](#)
- Features:
  - **domain formation** in equilibrium
  - **pseudo-gapped** or **mixed phases**
  - **inhomogeneities** and **spatial modulations**



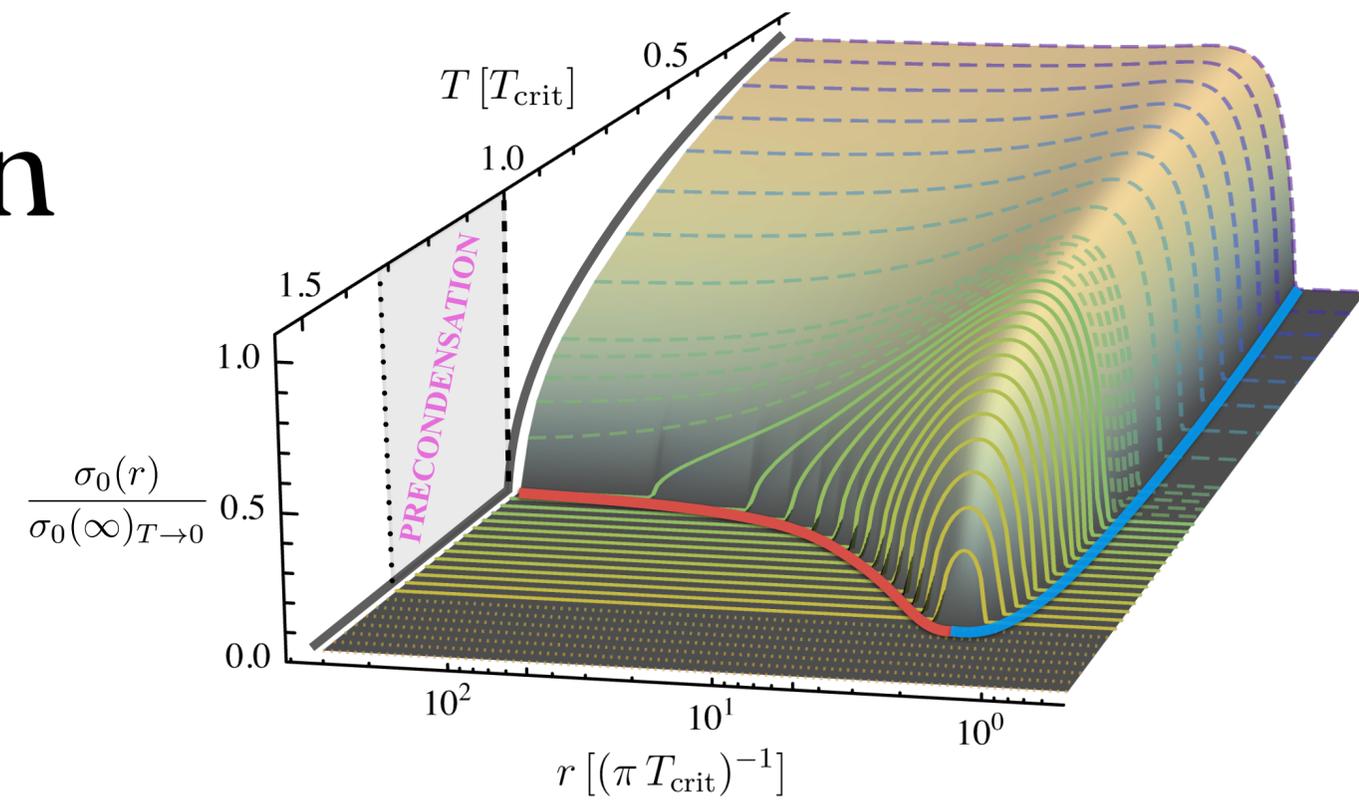
Weiss domains [Wikipedia]



# Hallmarks of precondensation

$$\sigma_0(x) = \sigma_0 + \Delta\sigma_0(x)$$

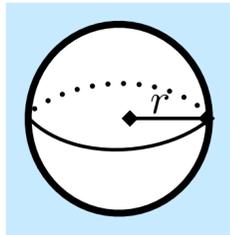
Macroscopic limit  $\sigma_0 = \lim_{\mathcal{V} \rightarrow \infty} \frac{1}{\mathcal{V}} \int d^3x \sigma_0(x) \quad \sigma_0(x) = \langle \hat{\sigma}(x) \rangle$



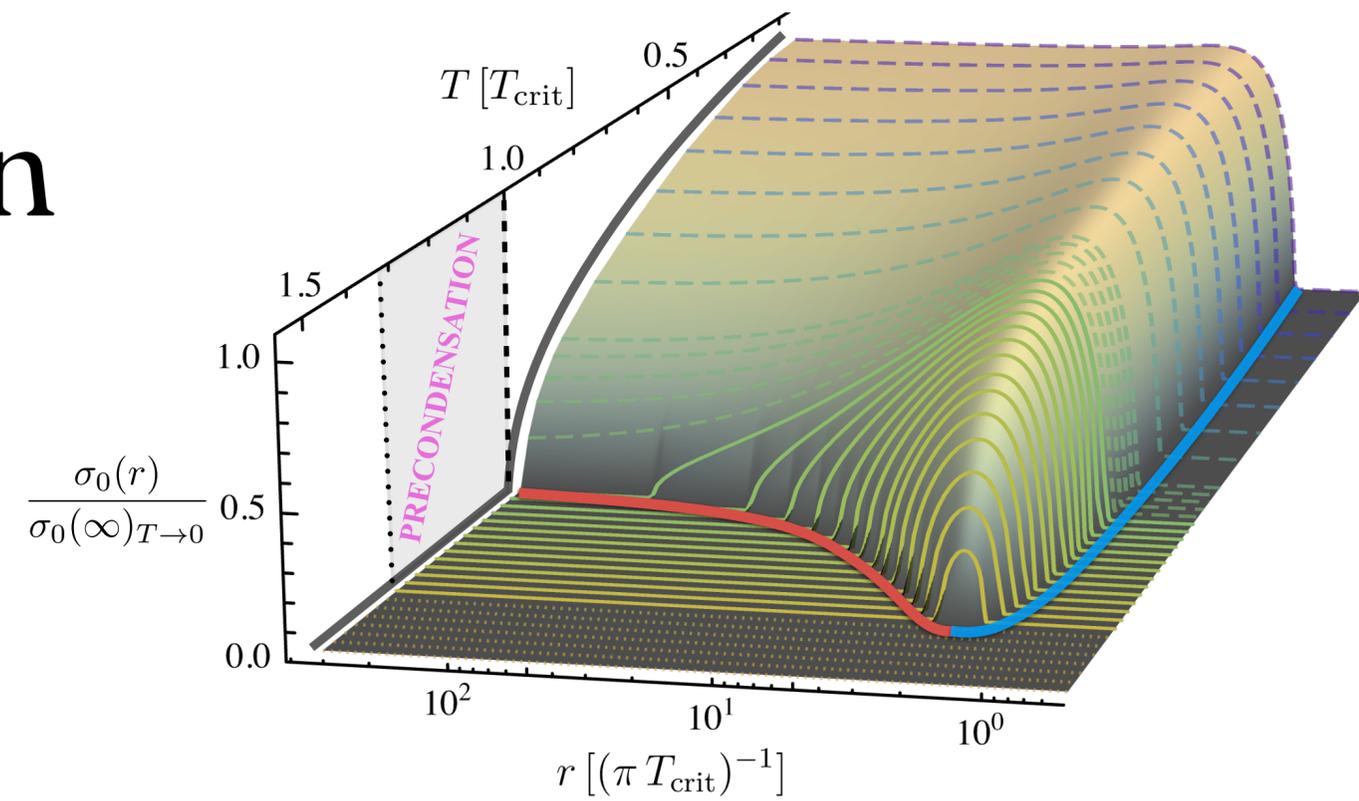
# Hallmarks of precondensation

$$\sigma_0(x) = \sigma_0 + \Delta\sigma_0(x)$$

Macroscopic limit  $\sigma_0 = \lim_{\mathcal{V} \rightarrow \infty} \frac{1}{\mathcal{V}} \int d^3x \sigma_0(x) \quad \sigma_0(x) = \langle \hat{\sigma}(x) \rangle$



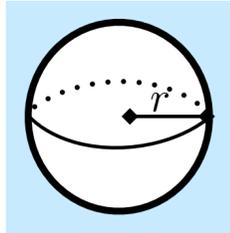
$$\sigma_0^2(\mathbf{x}; r) = \frac{1}{\mathcal{V}_r} \int d^3y \langle \hat{\sigma}(x) \hat{\sigma}(y) \rangle \theta(r^2 - (\mathbf{x} - \mathbf{y})^2)$$



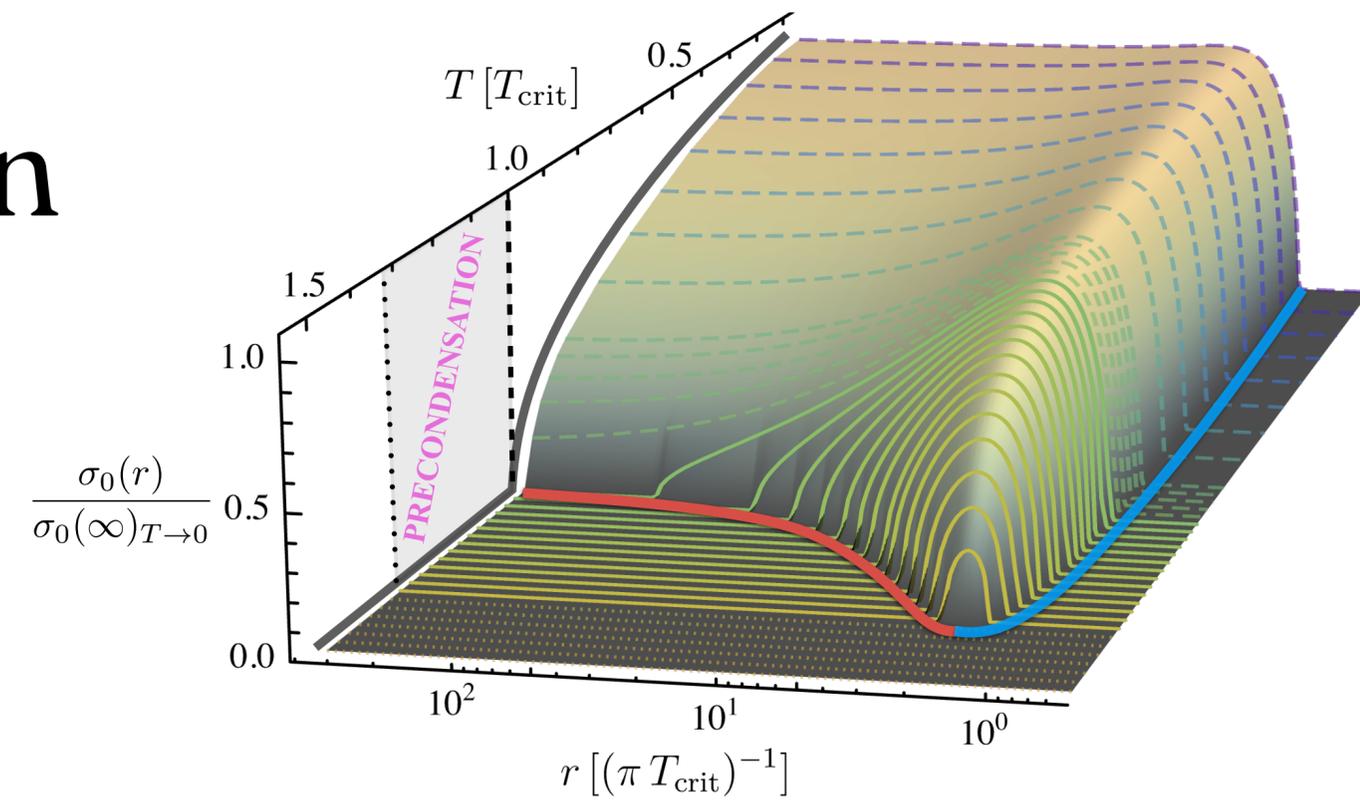
# Hallmarks of precondensation

$$\sigma_0(\mathbf{p}) = \sigma_0 (2\pi)^3 \delta(\mathbf{p}) + \Delta\sigma_0(\mathbf{p})$$

Macroscopic limit  $\sigma_0 = \lim_{\mathcal{V} \rightarrow \infty} \frac{1}{\mathcal{V}} \int d^3x \sigma_0(x) \quad \sigma_0(x) = \langle \hat{\sigma}(x) \rangle$



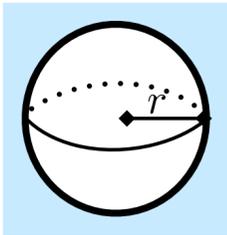
$$\sigma_0^2(\mathbf{p}; r) = \frac{1}{\mathcal{V}_r} \int d^3q \langle \hat{\sigma}(\mathbf{p} - \mathbf{q}) \hat{\sigma}(\mathbf{q}) \rangle \delta_r(\mathbf{q})$$



# Hallmarks of precondensation

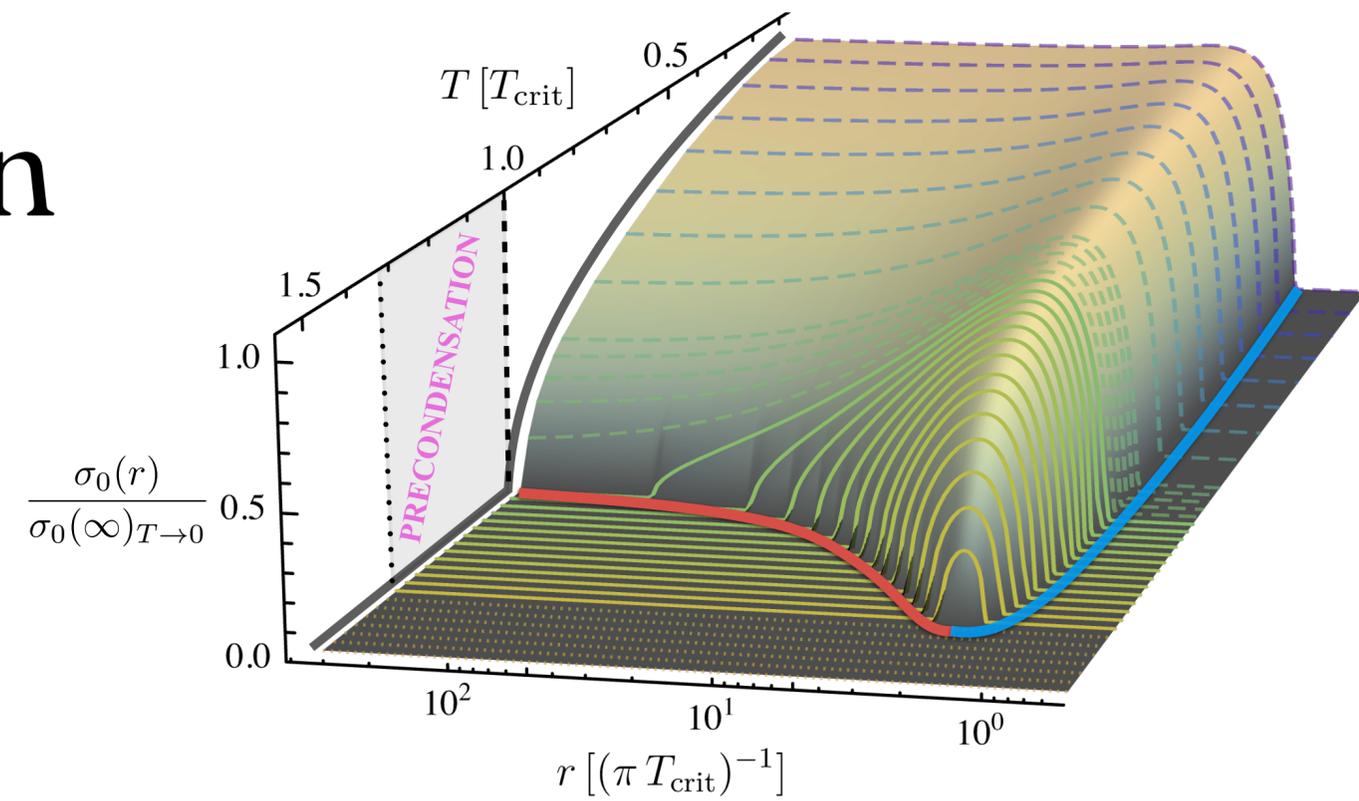
$$\sigma_0(\mathbf{p}) = \sigma_0 (2\pi)^3 \delta(\mathbf{p}) + \Delta\sigma_0(\mathbf{p})$$

Macroscopic limit  $\sigma_0 = \lim_{\mathcal{V} \rightarrow \infty} \frac{1}{\mathcal{V}} \int d^3x \sigma_0(x) \quad \sigma_0(x) = \langle \hat{\sigma}(x) \rangle$



$$\sigma_0^2(\mathbf{p}; r) = \frac{1}{\mathcal{V}_r} \int d^3q \langle \hat{\sigma}(\mathbf{p} - \mathbf{q}) \hat{\sigma}(\mathbf{q}) \rangle \delta_r(\mathbf{q})$$

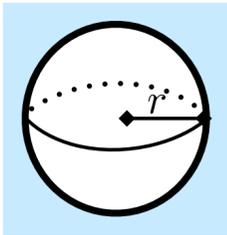
$$\begin{aligned} \sigma_0^2(\mathbf{p}; r) = & \sigma_0^2 (2\pi)^3 \delta(\mathbf{p}) + \sigma_0 \Delta\sigma_0(\mathbf{p}) \left( 1 + \frac{(2\pi)^3}{\mathcal{V}_r} \delta_r(\mathbf{p}) \right) \\ & + \frac{1}{\mathcal{V}_r} \int d^3q \Delta\sigma_0(\mathbf{p} - \mathbf{q}) \Delta\sigma_0(\mathbf{q}) \delta_r(\mathbf{q}) \end{aligned}$$



# Hallmarks of precondensation

$$\sigma_0(\mathbf{p}) = \sigma_0 (2\pi)^3 \delta(\mathbf{p}) + \Delta\sigma_0(\mathbf{p})$$

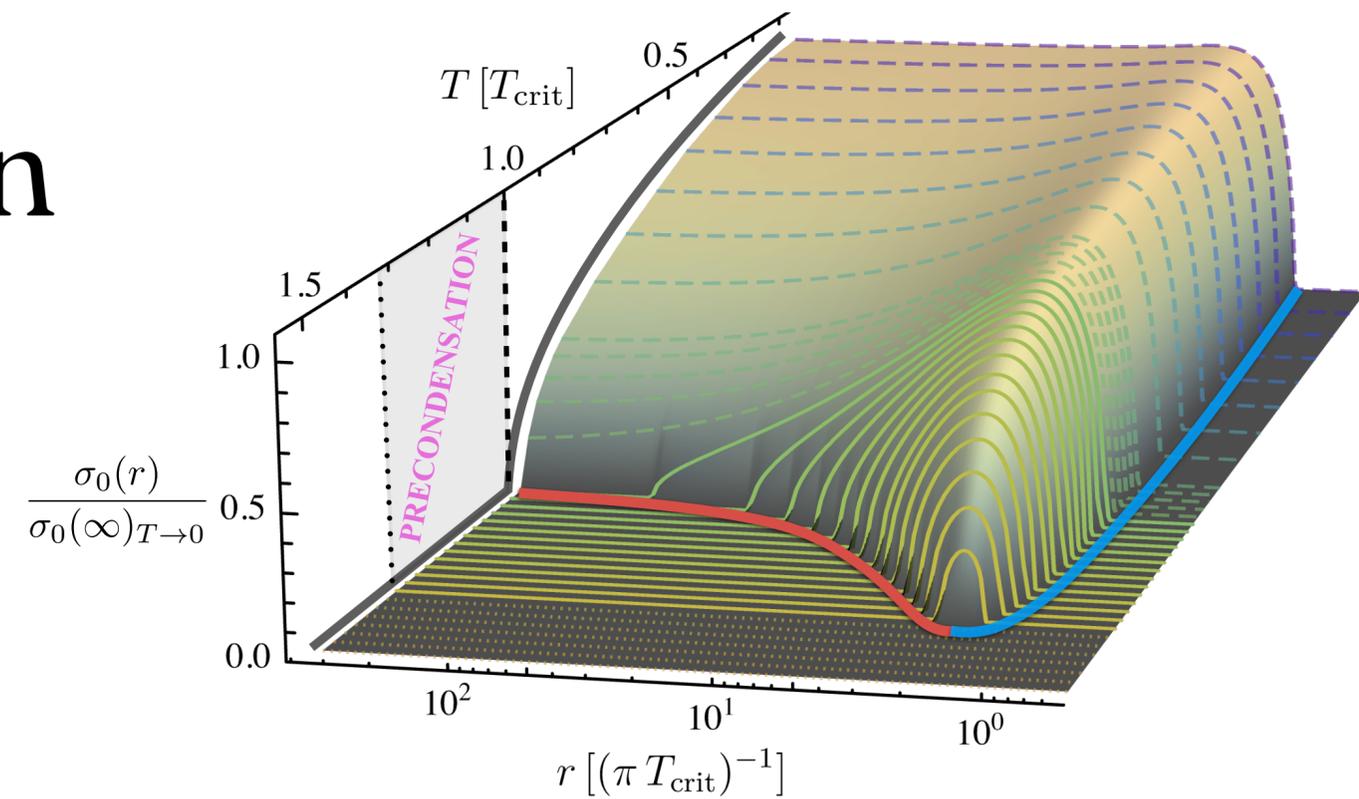
Macroscopic limit  $\sigma_0 = \lim_{\mathcal{V} \rightarrow \infty} \frac{1}{\mathcal{V}} \int d^3x \sigma_0(x) \quad \sigma_0(x) = \langle \hat{\sigma}(x) \rangle$



$$\sigma_0^2(\mathbf{p}; r) = \frac{1}{\mathcal{V}_r} \int d^3q \langle \hat{\sigma}(\mathbf{p} - \mathbf{q}) \hat{\sigma}(\mathbf{q}) \rangle \delta_r(\mathbf{q})$$

$$\begin{aligned} \sigma_0^2(\mathbf{p}; r) = & \sigma_0^2 (2\pi)^3 \delta(\mathbf{p}) + \sigma_0 \Delta\sigma_0(\mathbf{p}) \left( 1 + \frac{(2\pi)^3}{\mathcal{V}_r} \delta_r(\mathbf{p}) \right) \\ & + \frac{1}{\mathcal{V}_r} \int d^3q \Delta\sigma_0(\mathbf{p} - \mathbf{q}) \Delta\sigma_0(\mathbf{q}) \delta_r(\mathbf{q}) \end{aligned}$$

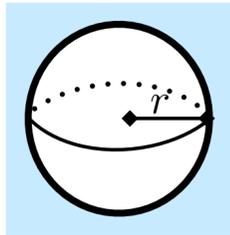
(i) *Symmetric regime:*  $\sigma_0^2(\mathbf{p}; r) \equiv 0$ .



# Hallmarks of precondensation

$$\sigma_0(\mathbf{p}) = \sigma_0 (2\pi)^3 \delta(\mathbf{p}) + \Delta\sigma_0(\mathbf{p})$$

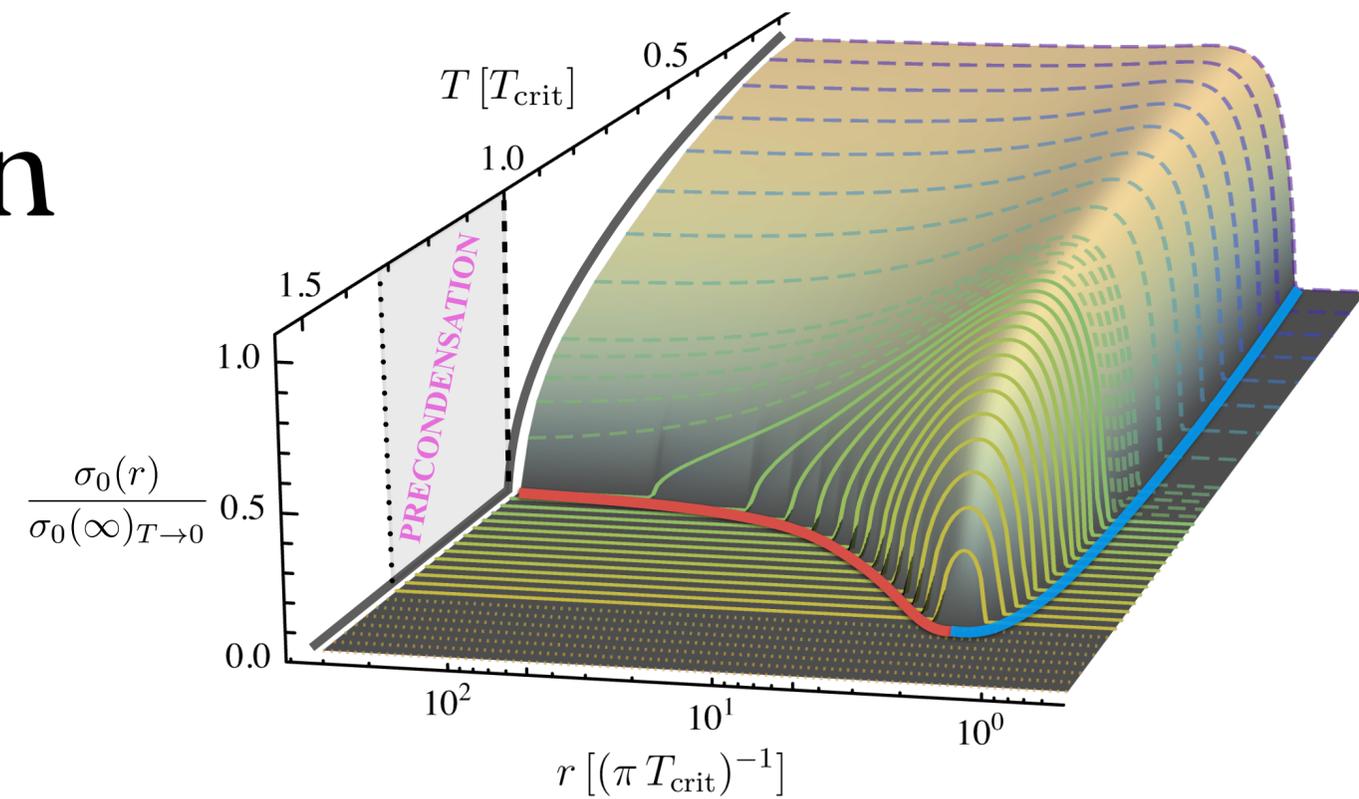
Macroscopic limit  $\sigma_0 = \lim_{\mathcal{V} \rightarrow \infty} \frac{1}{\mathcal{V}} \int d^3x \sigma_0(x) \quad \sigma_0(x) = \langle \hat{\sigma}(x) \rangle$



$$\sigma_0^2(\mathbf{p}; r) = \frac{1}{\mathcal{V}_r} \int d^3q \langle \hat{\sigma}(\mathbf{p} - \mathbf{q}) \hat{\sigma}(\mathbf{q}) \rangle \delta_r(\mathbf{q})$$

$$\begin{aligned} \sigma_0^2(\mathbf{p}; r) = & \sigma_0^2 (2\pi)^3 \delta(\mathbf{p}) + \sigma_0 \Delta\sigma_0(\mathbf{p}) \left( 1 + \frac{(2\pi)^3}{\mathcal{V}_r} \delta_r(\mathbf{p}) \right) \\ & + \frac{1}{\mathcal{V}_r} \int d^3q \Delta\sigma_0(\mathbf{p} - \mathbf{q}) \Delta\sigma_0(\mathbf{q}) \delta_r(\mathbf{q}) \end{aligned}$$

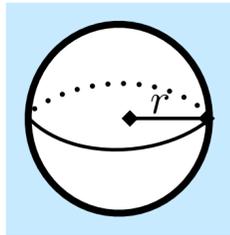
- (i) *Symmetric regime:*  $\sigma_0^2(\mathbf{p}; r) \equiv 0$ .
- (ii) *Precondensation regime:*  $\sigma_0^2(\mathbf{p}; r) \neq 0$  and  $\sigma_0 = 0$ .



# Hallmarks of precondensation

$$\sigma_0(\mathbf{p}) = \sigma_0 (2\pi)^3 \delta(\mathbf{p}) + \Delta\sigma_0(\mathbf{p})$$

Macroscopic limit  $\sigma_0 = \lim_{\mathcal{V} \rightarrow \infty} \frac{1}{\mathcal{V}} \int d^3x \sigma_0(x) \quad \sigma_0(x) = \langle \hat{\sigma}(x) \rangle$



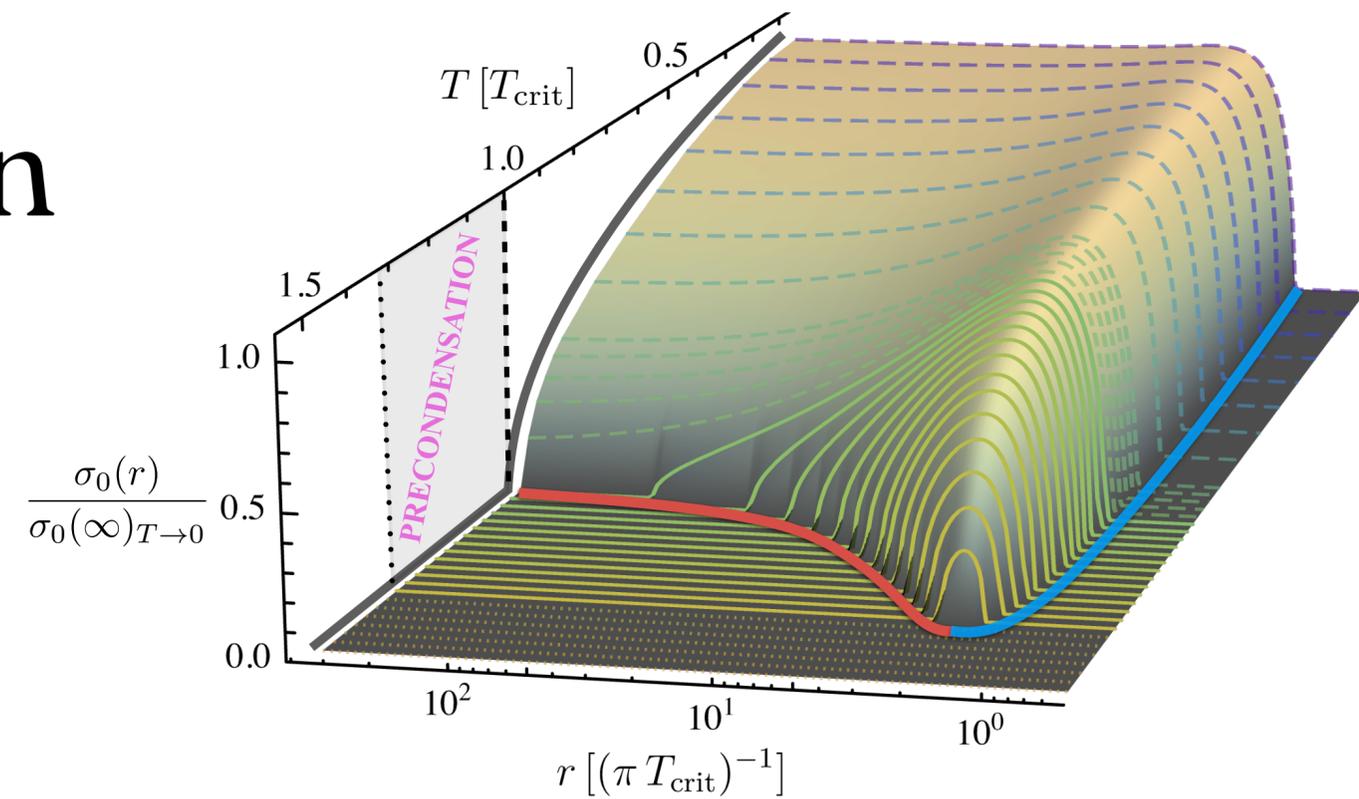
$$\sigma_0^2(\mathbf{p}; r) = \frac{1}{\mathcal{V}_r} \int d^3q \langle \hat{\sigma}(\mathbf{p} - \mathbf{q}) \hat{\sigma}(\mathbf{q}) \rangle \delta_r(\mathbf{q})$$

$$\begin{aligned} \sigma_0^2(\mathbf{p}; r) = & \sigma_0^2 (2\pi)^3 \delta(\mathbf{p}) + \sigma_0 \Delta\sigma_0(\mathbf{p}) \left( 1 + \frac{(2\pi)^3}{\mathcal{V}_r} \delta_r(\mathbf{p}) \right) \\ & + \frac{1}{\mathcal{V}_r} \int d^3q \Delta\sigma_0(\mathbf{p} - \mathbf{q}) \Delta\sigma_0(\mathbf{q}) \delta_r(\mathbf{q}) \end{aligned}$$

(i) *Symmetric regime:*  $\sigma_0^2(\mathbf{p}; r) \equiv 0$ .

(ii) *Precondensation regime:*  $\sigma_0^2(\mathbf{p}; r) \neq 0$  and  $\sigma_0 = 0$ .

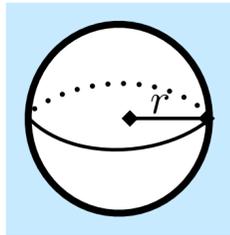
$$\sigma_0^2(\mathbf{p}; r) = \begin{cases} 0, & r < r_{\text{UV}}, \\ \neq 0, & r_{\text{UV}} < r < \xi, \\ 0, & r > \xi. \end{cases}$$



# Hallmarks of precondensation

$$\sigma_0(\mathbf{p}) = \sigma_0 (2\pi)^3 \delta(\mathbf{p}) + \Delta\sigma_0(\mathbf{p})$$

Macroscopic limit  $\sigma_0 = \lim_{\mathcal{V} \rightarrow \infty} \frac{1}{\mathcal{V}} \int d^3x \sigma_0(x) \quad \sigma_0(x) \equiv \langle \hat{\sigma}(x) \rangle$



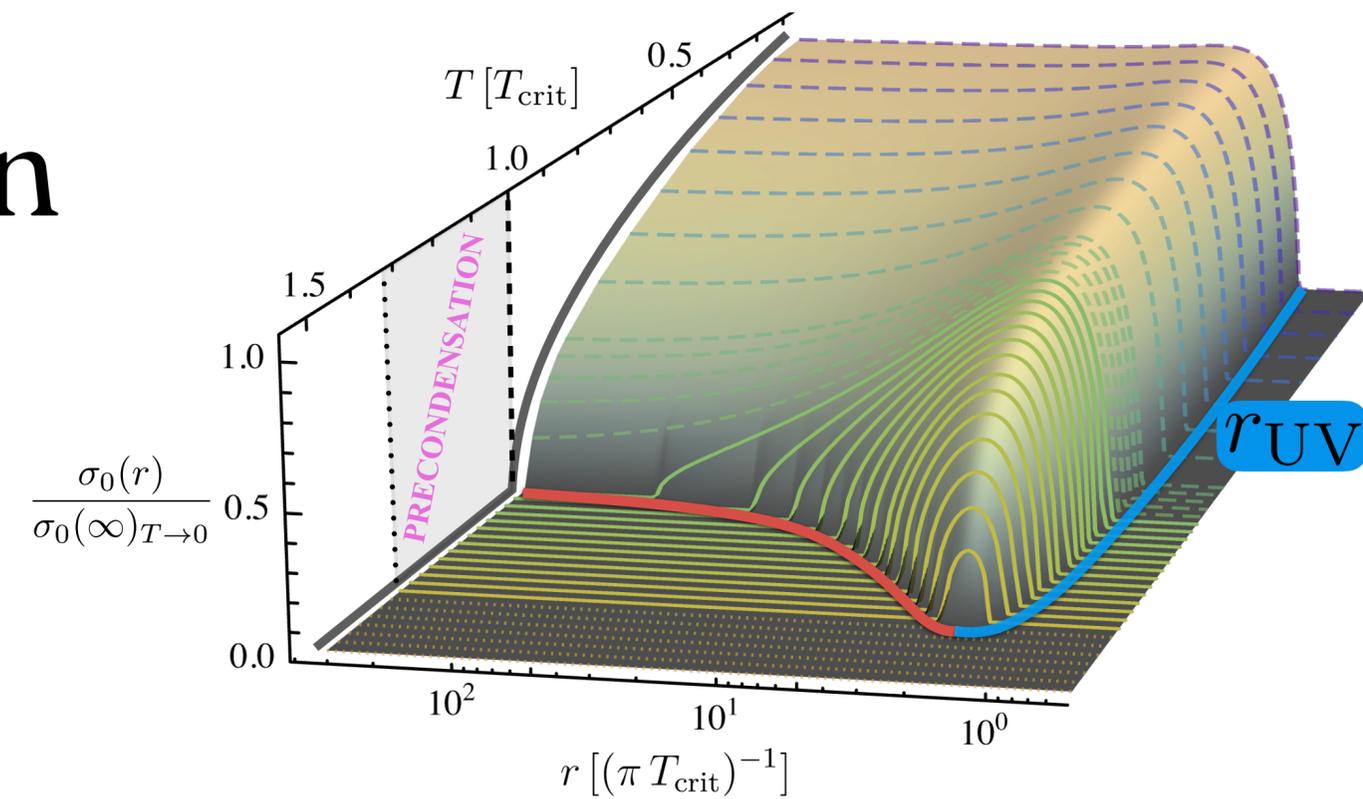
$$\sigma_0^2(\mathbf{p}; r) = \frac{1}{\mathcal{V}_r} \int d^3q \langle \hat{\sigma}(\mathbf{p} - \mathbf{q}) \hat{\sigma}(\mathbf{q}) \rangle \delta_r(\mathbf{q})$$

$$\begin{aligned} \sigma_0^2(\mathbf{p}; r) &= \sigma_0^2 (2\pi)^3 \delta(\mathbf{p}) + \sigma_0 \Delta\sigma_0(\mathbf{p}) \left( 1 + \frac{(2\pi)^3}{\mathcal{V}_r} \delta_r(\mathbf{p}) \right) \\ &+ \frac{1}{\mathcal{V}_r} \int d^3q \Delta\sigma_0(\mathbf{p} - \mathbf{q}) \Delta\sigma_0(\mathbf{q}) \delta_r(\mathbf{q}) \end{aligned}$$

(i) *Symmetric regime:*  $\sigma_0^2(\mathbf{p}; r) \equiv 0$ .

(ii) *Precondensation regime:*  $\sigma_0^2(\mathbf{p}; r) \neq 0$  and  $\sigma_0 = 0$ .

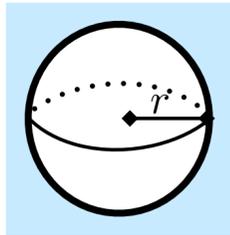
$$\sigma_0^2(\mathbf{p}; r) = \begin{cases} 0, & r < r_{\text{UV}}, \\ \neq 0, & r_{\text{UV}} < r < \xi, \\ 0, & r > \xi. \end{cases}$$



# Hallmarks of precondensation

$$\sigma_0(\mathbf{p}) = \sigma_0 (2\pi)^3 \delta(\mathbf{p}) + \Delta\sigma_0(\mathbf{p})$$

Macroscopic limit  $\sigma_0 = \lim_{\mathcal{V} \rightarrow \infty} \frac{1}{\mathcal{V}} \int d^3x \sigma_0(x) \quad \sigma_0(x) \equiv \langle \hat{\sigma}(x) \rangle$



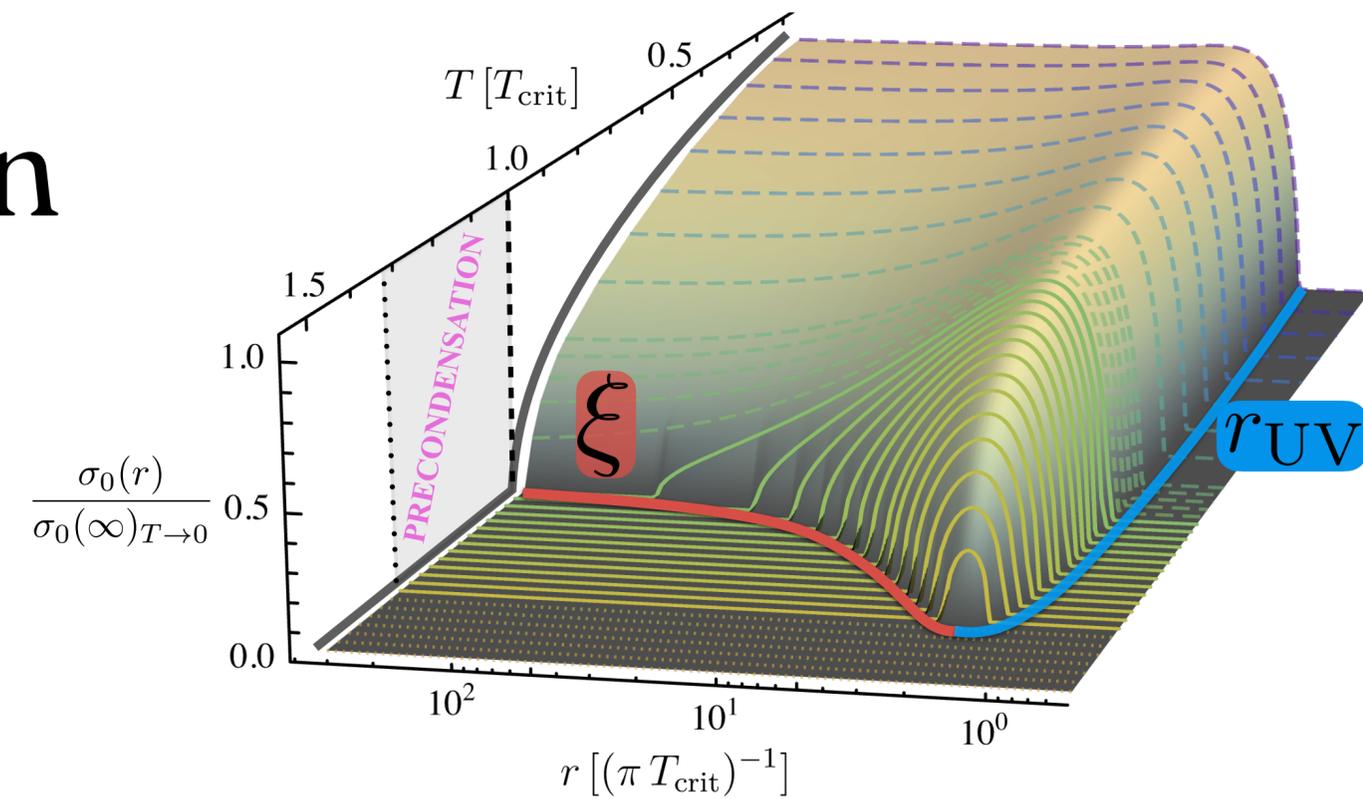
$$\sigma_0^2(\mathbf{p}; r) = \frac{1}{\mathcal{V}_r} \int d^3q \langle \hat{\sigma}(\mathbf{p} - \mathbf{q}) \hat{\sigma}(\mathbf{q}) \rangle \delta_r(\mathbf{q})$$

$$\begin{aligned} \sigma_0^2(\mathbf{p}; r) &= \sigma_0^2 (2\pi)^3 \delta(\mathbf{p}) + \sigma_0 \Delta\sigma_0(\mathbf{p}) \left( 1 + \frac{(2\pi)^3}{\mathcal{V}_r} \delta_r(\mathbf{p}) \right) \\ &+ \frac{1}{\mathcal{V}_r} \int d^3q \Delta\sigma_0(\mathbf{p} - \mathbf{q}) \Delta\sigma_0(\mathbf{q}) \delta_r(\mathbf{q}) \end{aligned}$$

(i) *Symmetric regime:*  $\sigma_0^2(\mathbf{p}; r) \equiv 0$ .

(ii) *Precondensation regime:*  $\sigma_0^2(\mathbf{p}; r) \neq 0$  and  $\sigma_0 = 0$ .

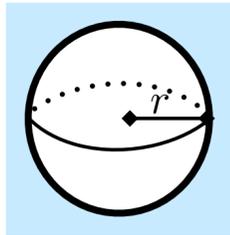
$$\sigma_0^2(\mathbf{p}; r) = \begin{cases} 0, & r < r_{\text{UV}}, \\ \neq 0, & r_{\text{UV}} < r < \xi, \\ 0, & r > \xi. \end{cases}$$



# Hallmarks of precondensation

$$\sigma_0(\mathbf{p}) = \sigma_0 (2\pi)^3 \delta(\mathbf{p}) + \Delta\sigma_0(\mathbf{p})$$

Macroscopic limit  $\sigma_0 = \lim_{\mathcal{V} \rightarrow \infty} \frac{1}{\mathcal{V}} \int d^3x \sigma_0(x) \quad \sigma_0(x) \equiv \langle \hat{\sigma}(x) \rangle$

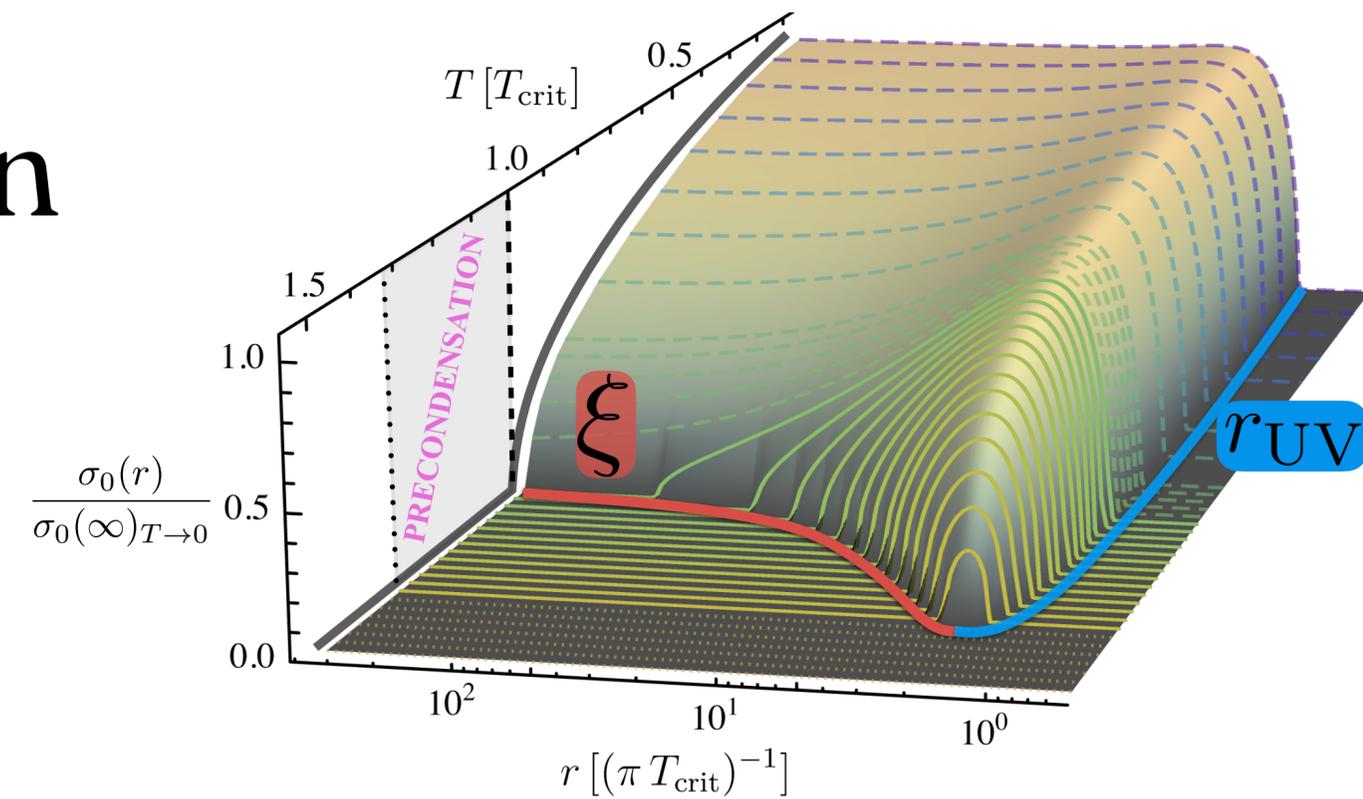


$$\sigma_0^2(\mathbf{p}; r) = \frac{1}{\mathcal{V}_r} \int d^3q \langle \hat{\sigma}(\mathbf{p} - \mathbf{q}) \hat{\sigma}(\mathbf{q}) \rangle \delta_r(\mathbf{q})$$

$$\begin{aligned} \sigma_0^2(\mathbf{p}; r) = & \sigma_0^2 (2\pi)^3 \delta(\mathbf{p}) + \sigma_0 \Delta\sigma_0(\mathbf{p}) \left( 1 + \frac{(2\pi)^3}{\mathcal{V}_r} \delta_r(\mathbf{p}) \right) \\ & + \frac{1}{\mathcal{V}_r} \int d^3q \Delta\sigma_0(\mathbf{p} - \mathbf{q}) \Delta\sigma_0(\mathbf{q}) \delta_r(\mathbf{q}) \end{aligned}$$

- (i) *Symmetric regime:*  $\sigma_0^2(\mathbf{p}; r) \equiv 0$ .
- (ii) *Precondensation regime:*  $\sigma_0^2(\mathbf{p}; r) \neq 0$  and  $\sigma_0 = 0$ .
- (iii) *Broken phase:*  $\sigma_0 \neq 0$ .

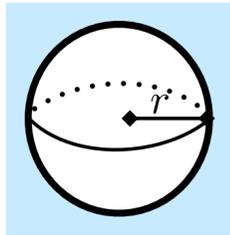
$$\sigma_0^2(\mathbf{p}; r) = \begin{cases} 0, & r < r_{\text{UV}}, \\ \neq 0, & r_{\text{UV}} < r < \xi, \\ 0, & r > \xi. \end{cases}$$



# Hallmarks of precondensation

$$\sigma_0(\mathbf{p}) = \sigma_0 (2\pi)^3 \delta(\mathbf{p}) + \Delta\sigma_0(\mathbf{p})$$

Macroscopic limit  $\sigma_0 = \lim_{V \rightarrow \infty} \frac{1}{V} \int d^3x \sigma_0(x) \quad \sigma_0(x) \equiv \langle \hat{\sigma}(x) \rangle$

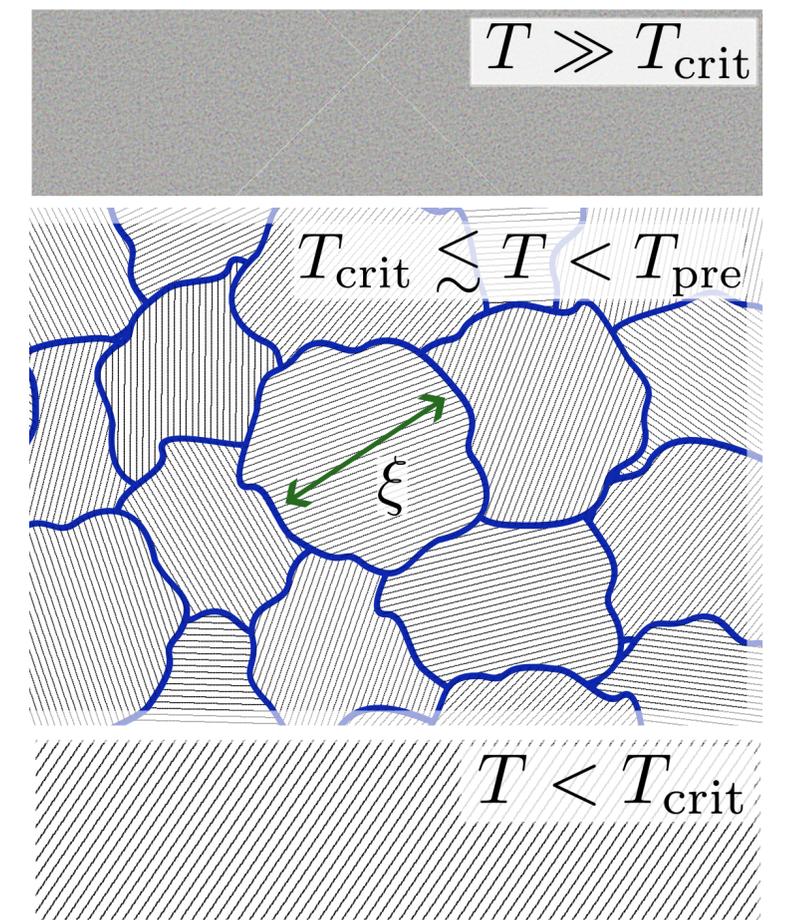
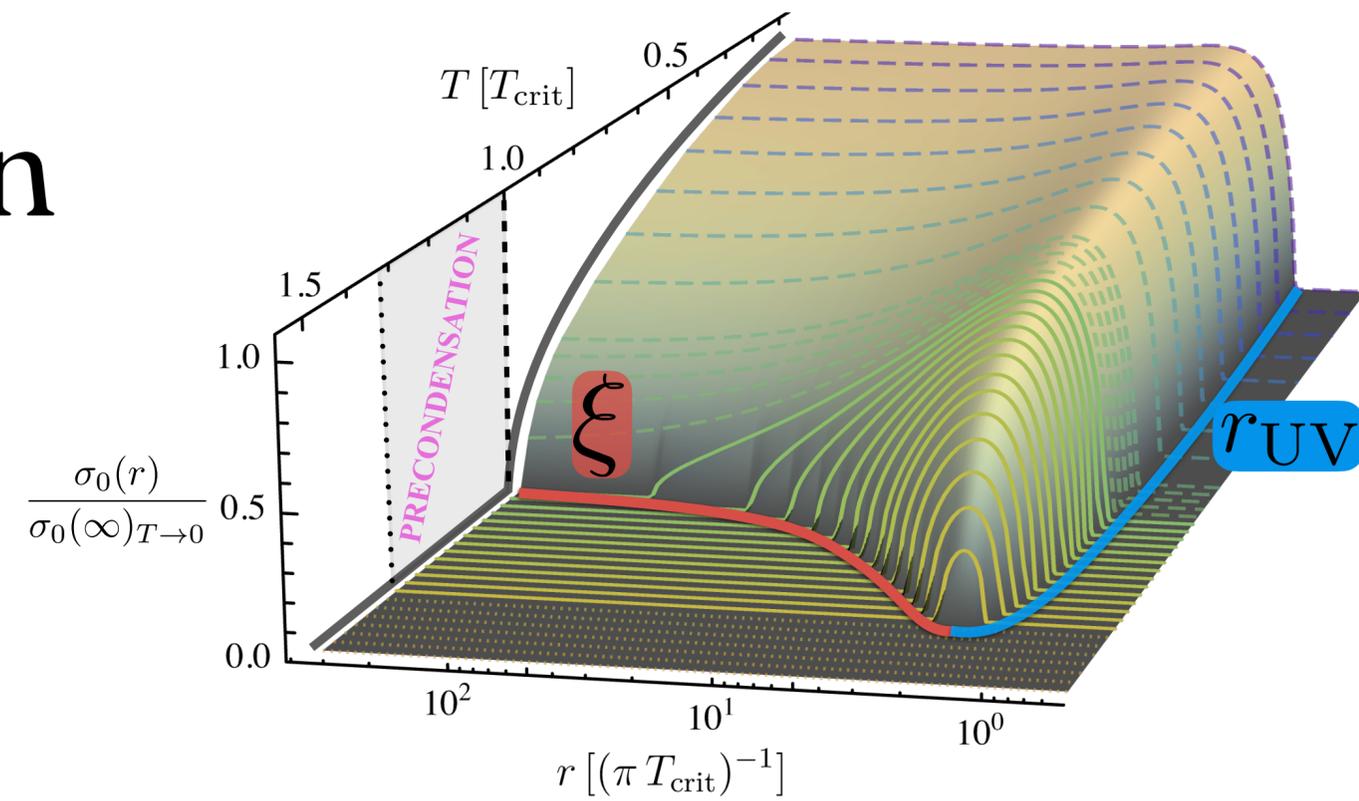


$$\sigma_0^2(\mathbf{p}; r) = \frac{1}{V_r} \int d^3q \langle \hat{\sigma}(\mathbf{p} - \mathbf{q}) \hat{\sigma}(\mathbf{q}) \rangle \delta_r(\mathbf{q})$$

$$\begin{aligned} \sigma_0^2(\mathbf{p}; r) = & \sigma_0^2 (2\pi)^3 \delta(\mathbf{p}) + \sigma_0 \Delta\sigma_0(\mathbf{p}) \left( 1 + \frac{(2\pi)^3}{V_r} \delta_r(\mathbf{p}) \right) \\ & + \frac{1}{V_r} \int d^3q \Delta\sigma_0(\mathbf{p} - \mathbf{q}) \Delta\sigma_0(\mathbf{q}) \delta_r(\mathbf{q}) \end{aligned}$$

- (i) *Symmetric regime:*  $\sigma_0^2(\mathbf{p}; r) \equiv 0$ .
- (ii) *Precondensation regime:*  $\sigma_0^2(\mathbf{p}; r) \neq 0$  and  $\sigma_0 = 0$ .
- (iii) *Broken phase:*  $\sigma_0 \neq 0$ .

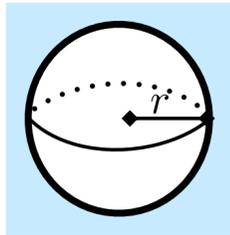
$$\sigma_0^2(\mathbf{p}; r) = \begin{cases} 0, & r < r_{UV}, \\ \neq 0, & r_{UV} < r < \xi, \\ 0, & r > \xi. \end{cases}$$



# Hallmarks of precondensation

$$\sigma_0(\mathbf{p}) = \sigma_0 (2\pi)^3 \delta(\mathbf{p}) + \Delta\sigma_0(\mathbf{p})$$

Macroscopic limit  $\sigma_0 = \lim_{V \rightarrow \infty} \frac{1}{V} \int d^3x \sigma_0(x) \quad \sigma_0(x) \equiv \langle \hat{\sigma}(x) \rangle$

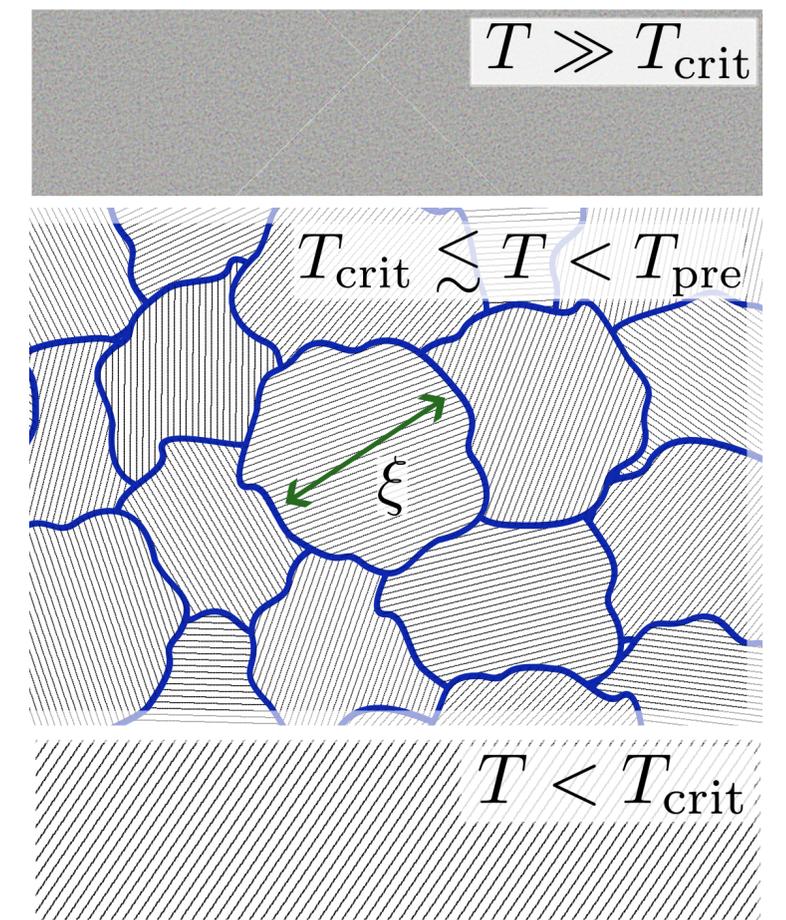
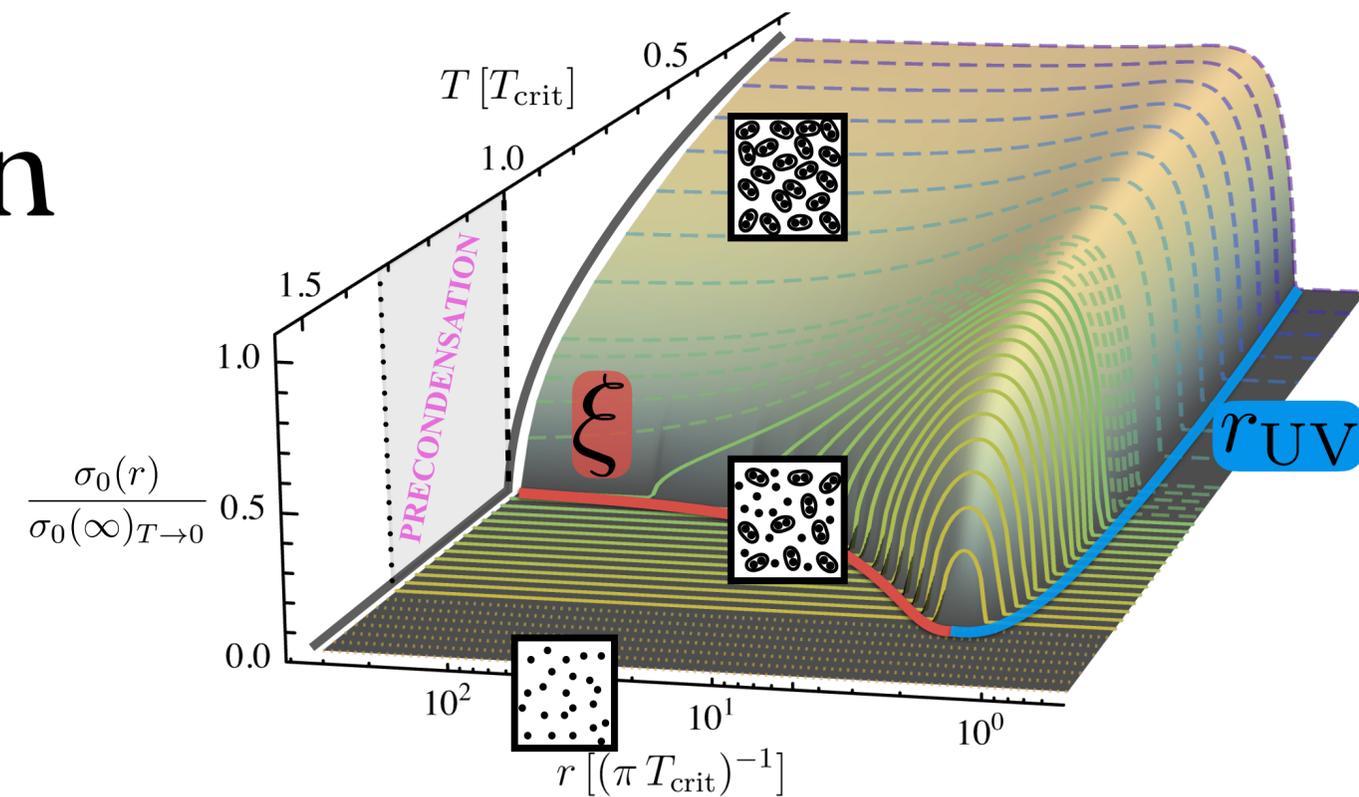


$$\sigma_0^2(\mathbf{p}; r) = \frac{1}{V_r} \int d^3q \langle \hat{\sigma}(\mathbf{p} - \mathbf{q}) \hat{\sigma}(\mathbf{q}) \rangle \delta_r(\mathbf{q})$$

$$\begin{aligned} \sigma_0^2(\mathbf{p}; r) = & \sigma_0^2 (2\pi)^3 \delta(\mathbf{p}) + \sigma_0 \Delta\sigma_0(\mathbf{p}) \left( 1 + \frac{(2\pi)^3}{V_r} \delta_r(\mathbf{p}) \right) \\ & + \frac{1}{V_r} \int d^3q \Delta\sigma_0(\mathbf{p} - \mathbf{q}) \Delta\sigma_0(\mathbf{q}) \delta_r(\mathbf{q}) \end{aligned}$$

- (i) *Symmetric regime:*  $\sigma_0^2(\mathbf{p}; r) \equiv 0$ .
- (ii) *Precondensation regime:*  $\sigma_0^2(\mathbf{p}; r) \neq 0$  and  $\sigma_0 = 0$ .
- (iii) *Broken phase:*  $\sigma_0 \neq 0$ .

$$\sigma_0^2(\mathbf{p}; r) = \begin{cases} 0, & r < r_{UV}, \\ \neq 0, & r_{UV} < r < \xi, \\ 0, & r > \xi. \end{cases}$$



# Functional Renormalisation Group

- **Scale-dependent effective action:**  $\Gamma_k[\phi]$   $k \sim 1/r$

$$\int [\mathcal{D}\phi]_{p>k} = \int \mathcal{D}\phi \exp(-\Delta S_k[\phi]) \quad \Delta S_k[\phi] = \int_p \phi(p) R_k \phi(-p)$$

$$\Gamma_k[\phi] = \int_x J(x)\phi(x) - \mathcal{W}_k[J] - \Delta S_k[\phi]$$

Wetterich '89

Flow equation:

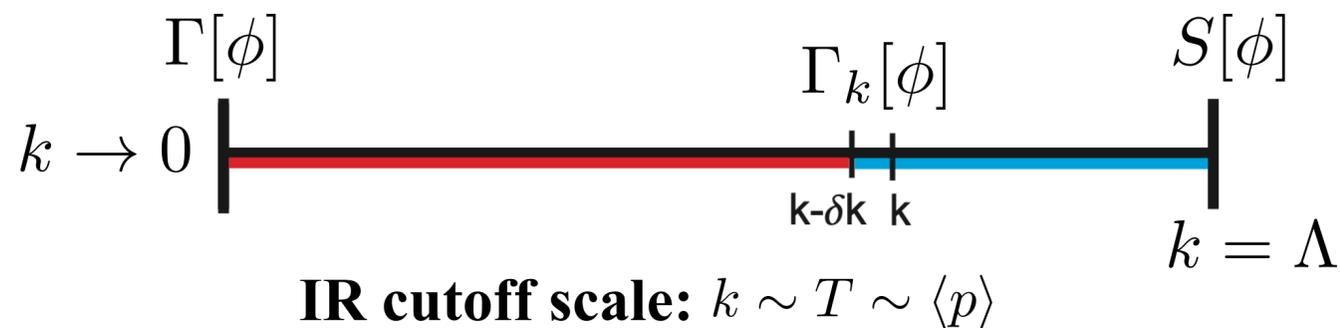
$$\partial_t \Gamma_k[\phi] = \frac{1}{2} \text{Tr} \left[ \frac{1}{\Gamma_k^{(2)} + R_k} \partial_t R_k \right] = \frac{1}{2} \text{Diagram}$$

$$\partial_t \equiv k \partial_k$$

Wetterich '93

- Average action of fields over a  $k^{-d}$  space-time volume
- Kadanoff's block-spinning idea in the continuum limit

- Exact derivation
- One loop
- **Non-perturbative**
- Mass-dependent
- Analytic regulators
- Versatile
- Systematic expansion schemes
- UV-IR finite
- Diagrammatic
- Real-time formulation
- ...



# Gauge-fermion theories

$$S = \int_x \frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a + \mathcal{L}_{\text{gf}} + \mathcal{L}_{\text{gh}} + \bar{\psi} \not{D} \psi$$

- Gauge-fermion theories in the chiral limit
  - Rich phenomena (confinement, dχSB, walking, ...)
  - Relevance for BSM physics
    - ★ Composite Higgs & Technicolour (HP)
    - ★ Strong Dark Sectors (DM)
    - ★ Cosmological evolution
    - ★ 1st order chiral phase transition, GWs

[Reichert,Sannino,Wang,Zhang\[2109.11552\]](#)      [Iso,Okada,Orikasa\[0902.4050\]](#)

[Pashechnik,Reichert,Sannino,Wang\[2309.16755\]](#)      [Iso,Serpico,Shimada\[1704.04955\]](#)

...

- Precondensation opens path to **new observational test**

- Hanbury Brown-Twiss spectroscopy

[Rennecke,Pisarski,Rischke\[2301.11484\]](#)

- ... ?

[Fukushima,Hidaka,Inoue,Shigaki,Yamaguchi\[2306.17619\]](#)

# Gauge-fermion theories

$$S = \int_x \frac{1}{4} F_{\mu\nu}^a F_{\mu\nu}^a + \mathcal{L}_{\text{gf}} + \mathcal{L}_{\text{gh}} + \bar{\psi} \not{D} \psi$$

- Gauge-fermion theories in the chiral limit

- Rich phenomena (confinement,  $d\chi$ SB, walking, ...)
- Relevance for BSM physics

- ★ Composite Higgs & Technicolour (HP)

- ★ Strong Dark Sectors (DM)

- ★ Cosmological evolution

- ★ 1st order chiral phase transition, GWs

Reichert,Sannino,Wang,Zhang[2109.11552]      Iso,Okada,Orikasa[0902.4050]

Pashechnik,Reichert,Sannino,Wang[2309.16755]      Iso,Serpico,Shimada[1704.04955]

- Precondensation opens path to **new observational test**

- Hanbury Brown-Twiss spectroscopy

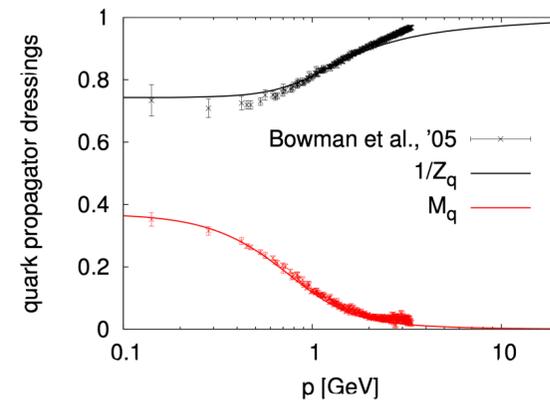
Rennecke,Pisarski,Rischke[2301.11484]

- ... ?

Fukushima,Hidaka,Inoue,Shigaki,Yamaguchi[2306.17619]

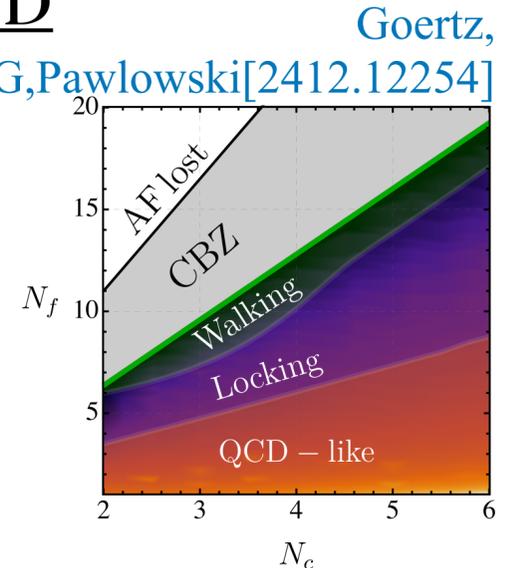
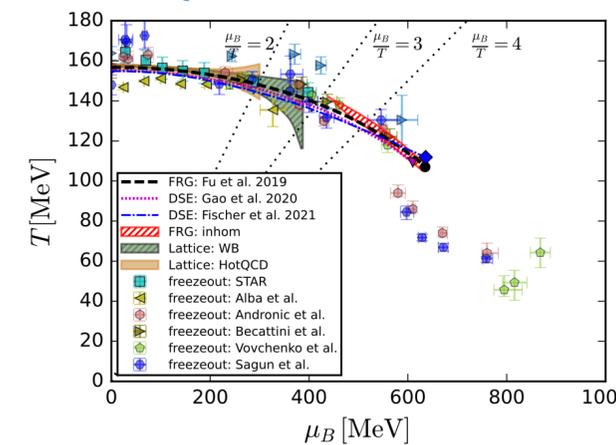
## fRG approach to QCD

Mitter,Pawlowski,Strodthoff[1411.7978]



fQCD collaboration

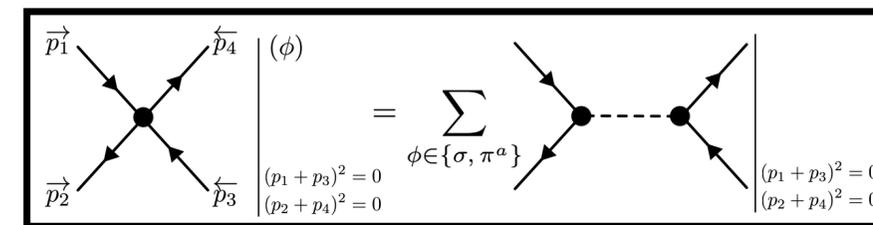
APG,Pawlowski[2412.12254]



- $d\chi$ SB encoded in higher dimensional fermionic operators

$$\Gamma_k[\Phi] \supset - \int_x \bar{\lambda}_\sigma (\bar{\psi} \mathcal{T}_{(S-P)} \psi)^2 + \dots$$

- Exact field transformation along the RG flow

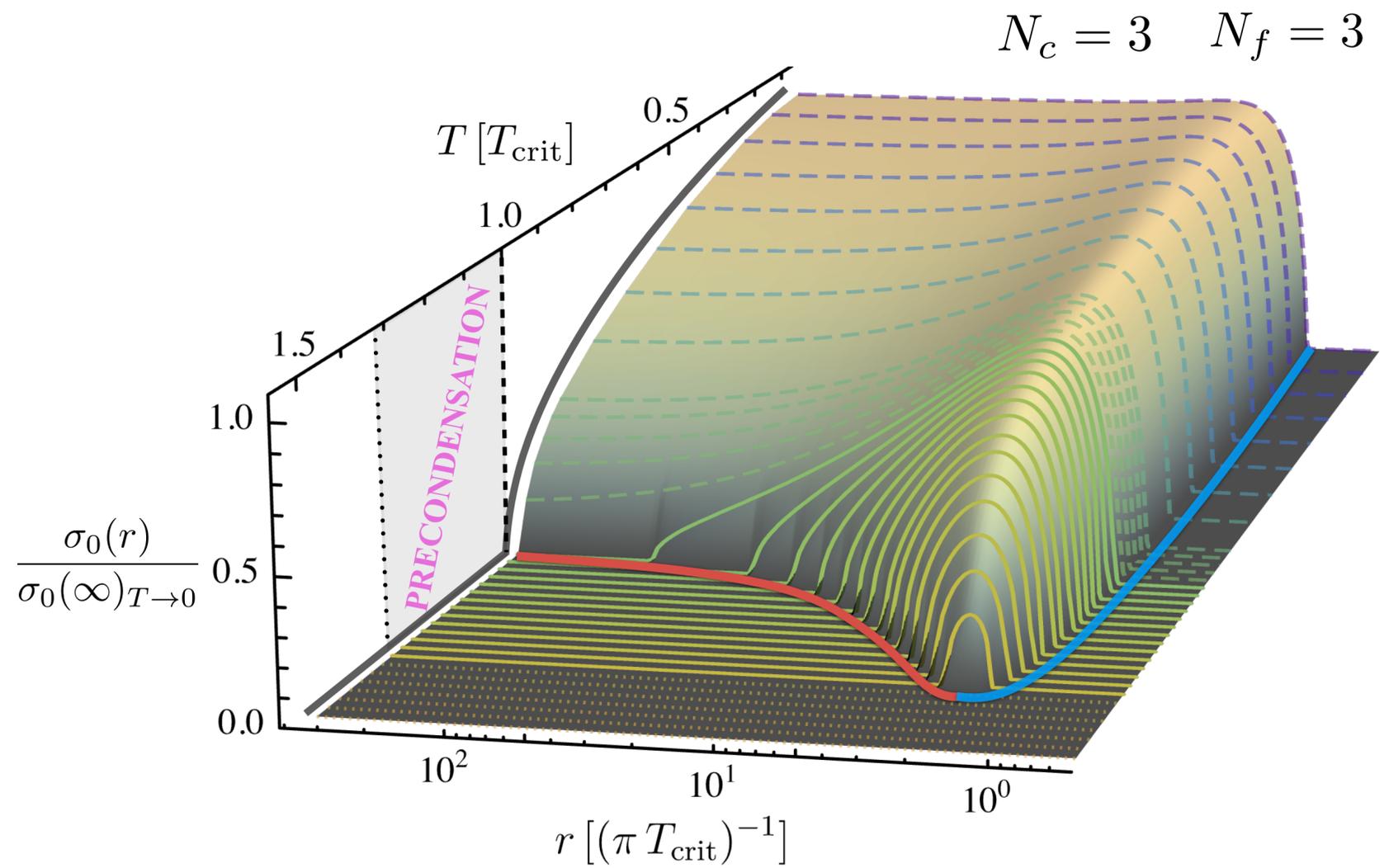
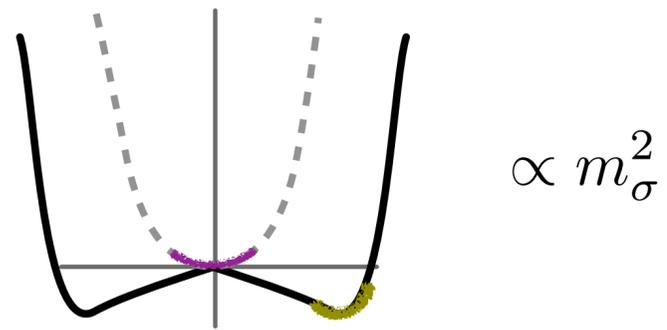


Stratonovich'57 Hubbard'59

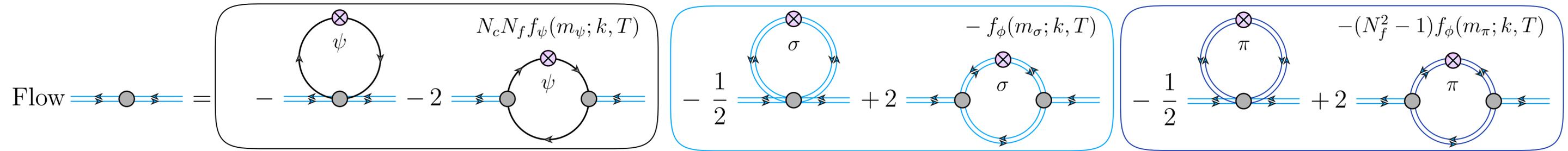
$$\Gamma_k[\Phi] \supset \int_x \bar{h} \bar{\psi} (T_f^0 \sigma + i\gamma_5 T_f^a \pi^a) \psi + Z_\phi (\partial_\mu \phi)^2 + V(\phi^2) + \dots$$

# Microscopic origin

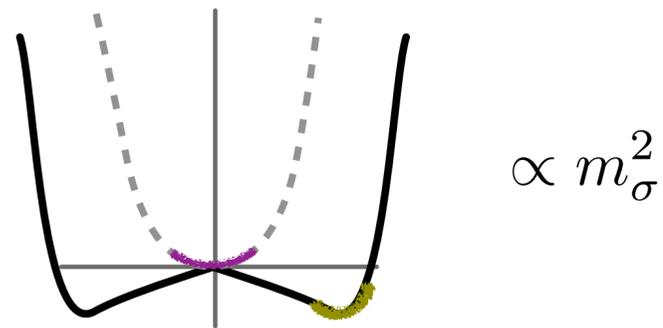
Flow of the effective chiral potential:



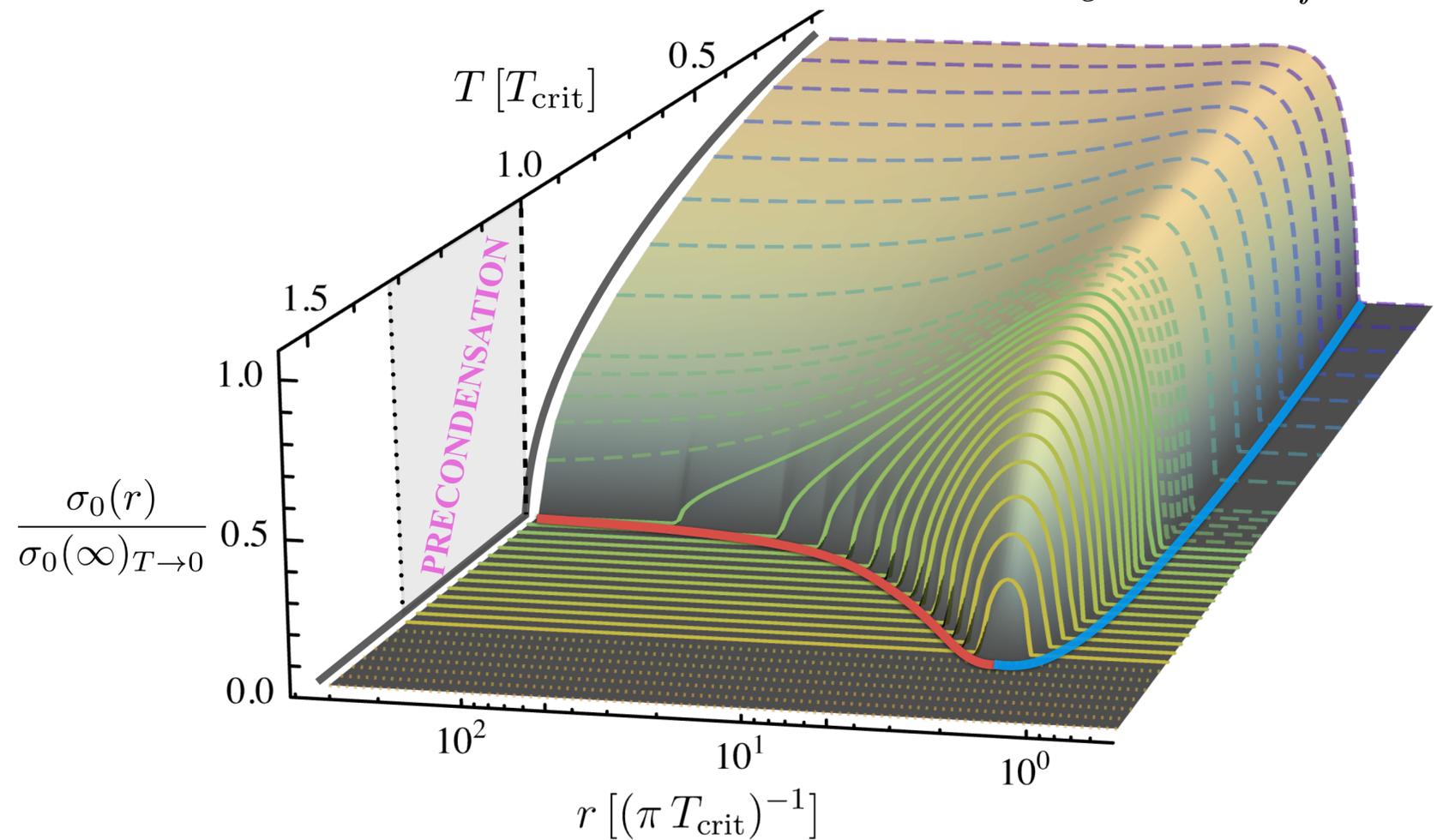
# Microscopic origin



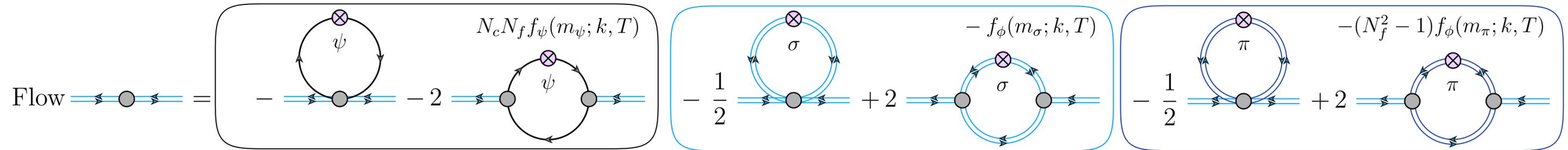
Flow of the effective chiral potential:



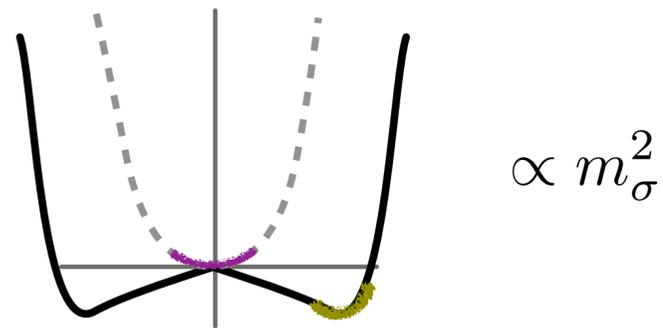
$N_c = 3 \quad N_f = 3$



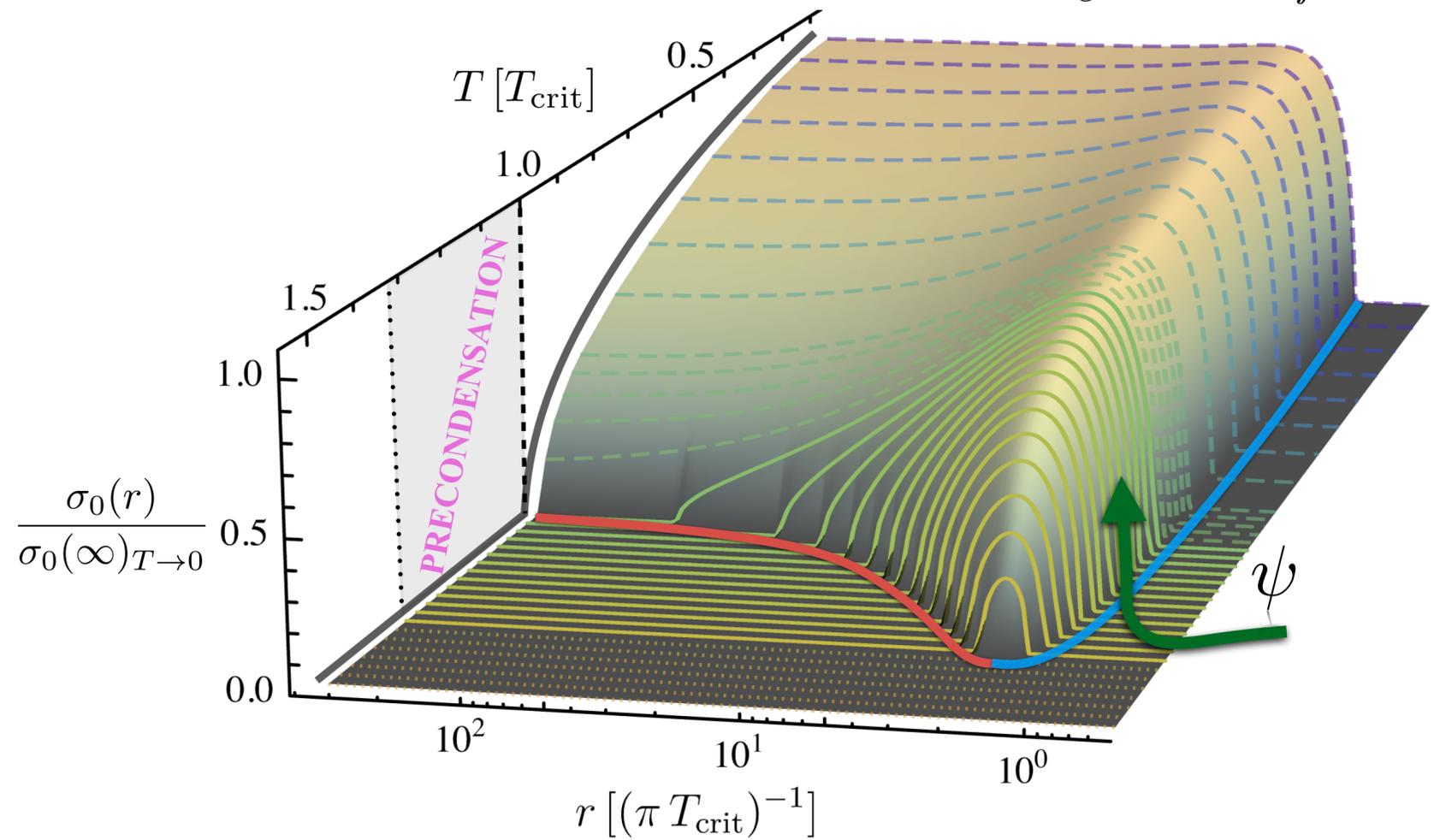
# Microscopic origin



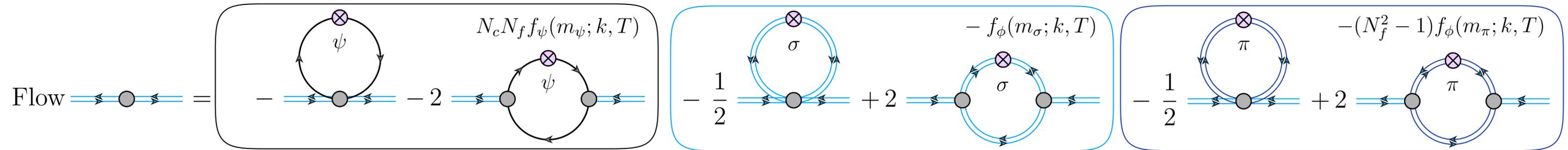
Flow of the effective chiral potential:



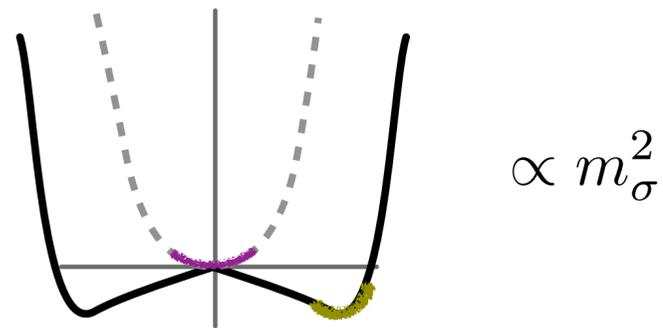
$N_c = 3 \quad N_f = 3$



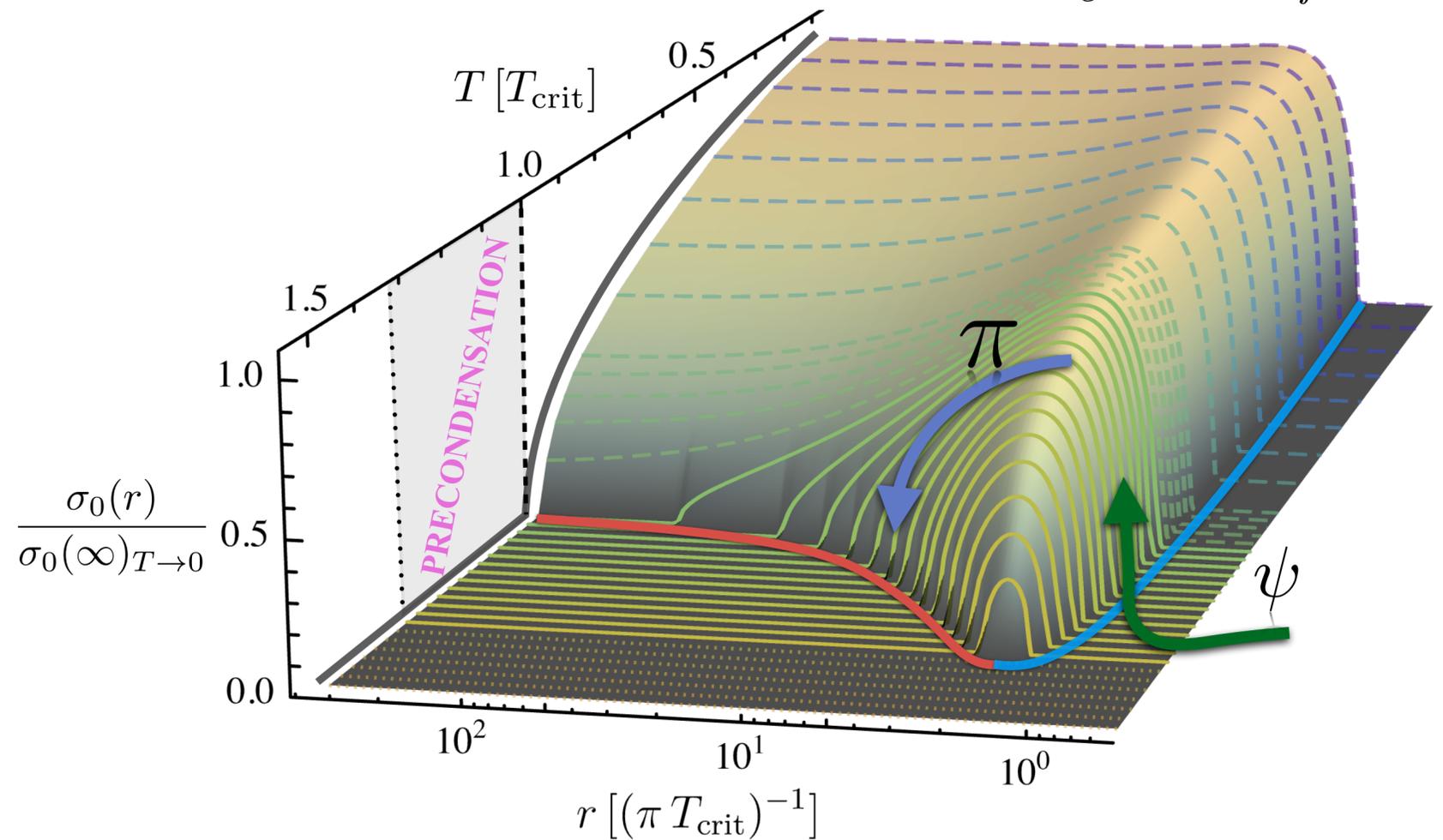
# Microscopic origin



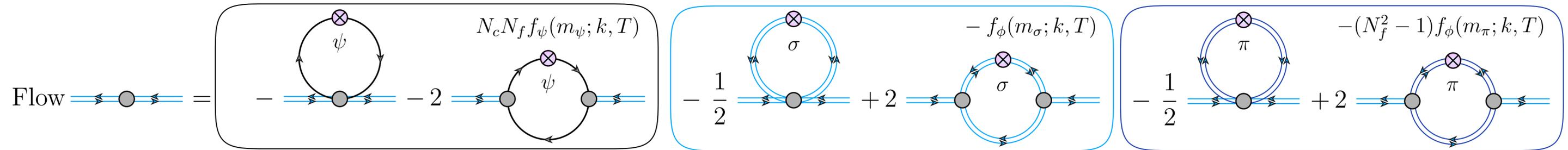
Flow of the effective chiral potential:



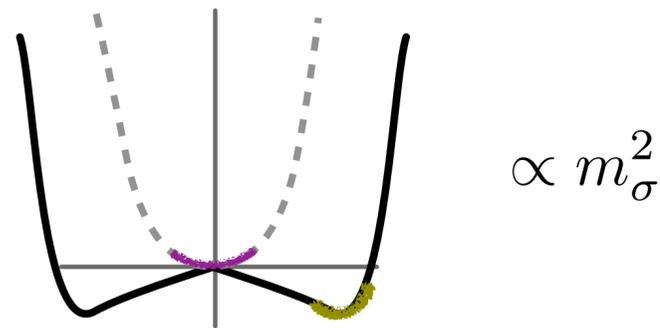
$N_c = 3 \quad N_f = 3$



# Microscopic origin



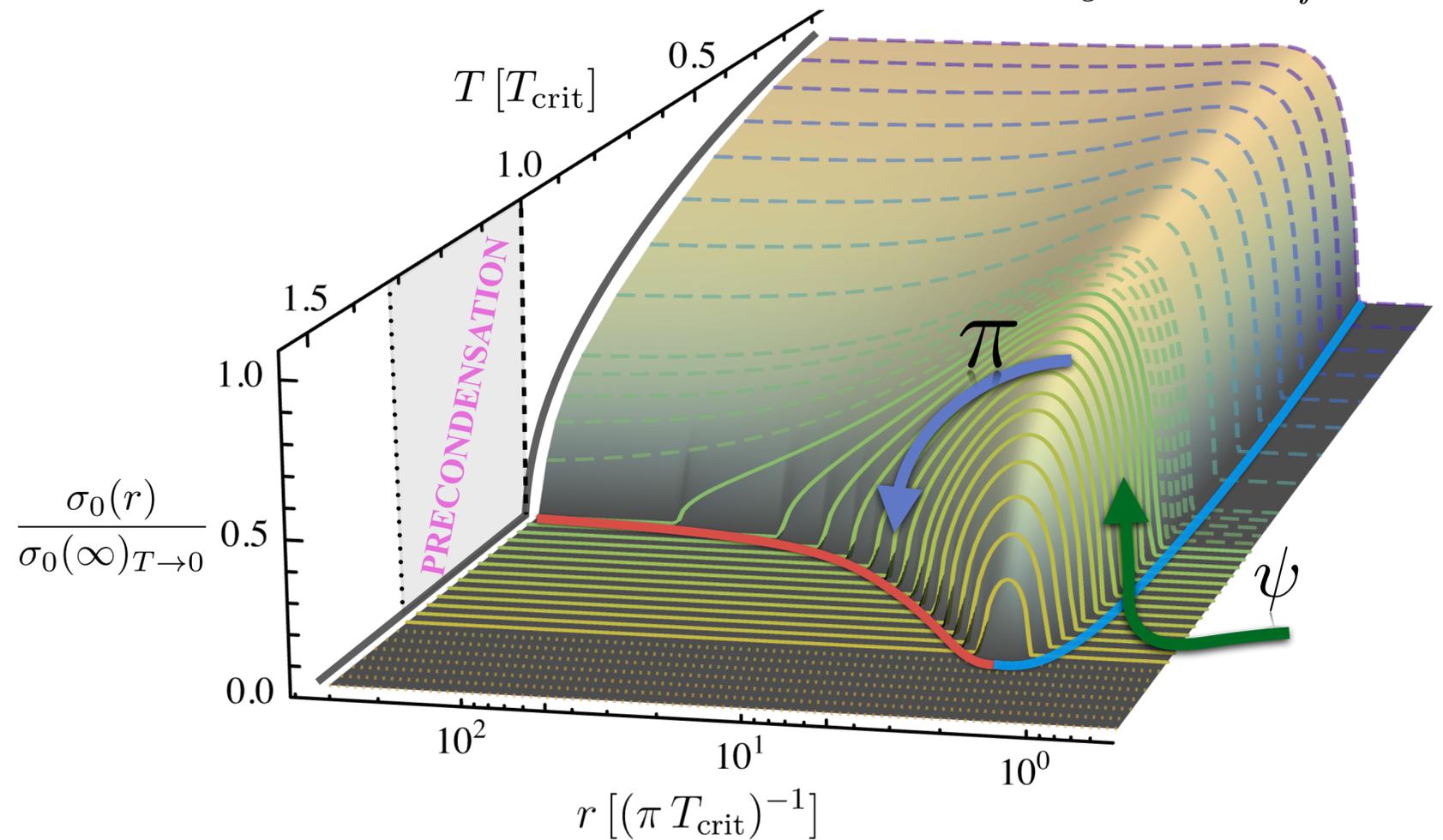
Flow of the effective chiral potential:



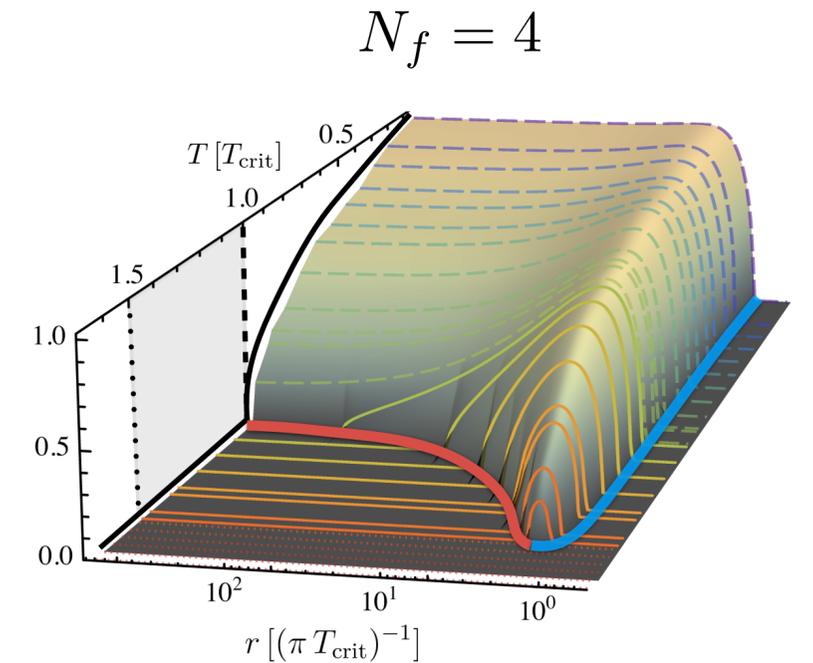
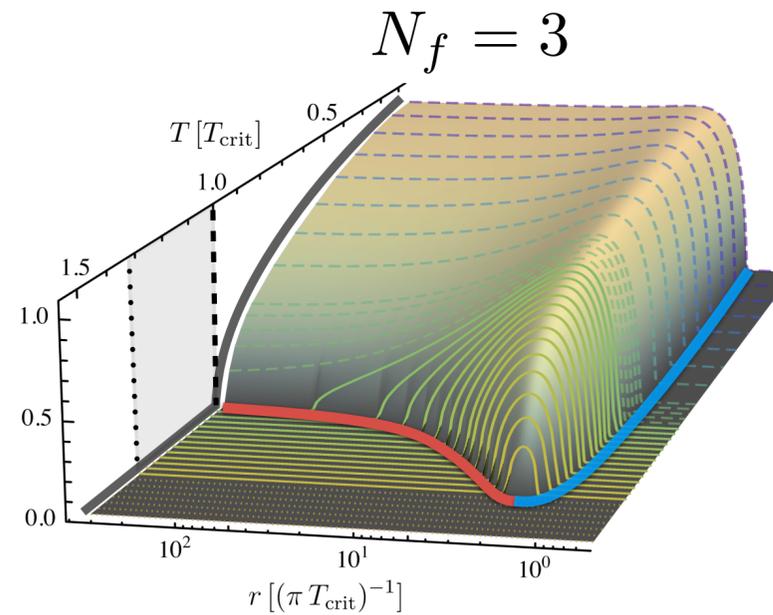
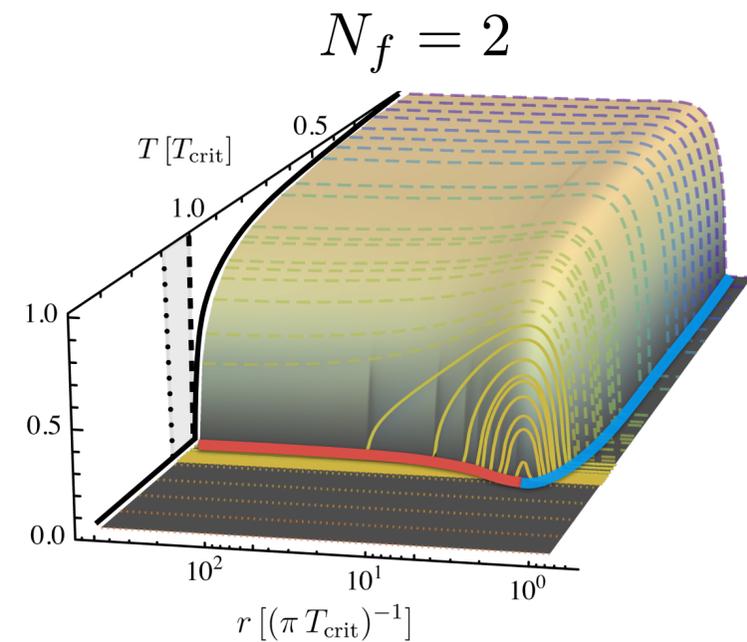
$N_c = 3 \quad N_f = 3$

• Necessary key ingredients:

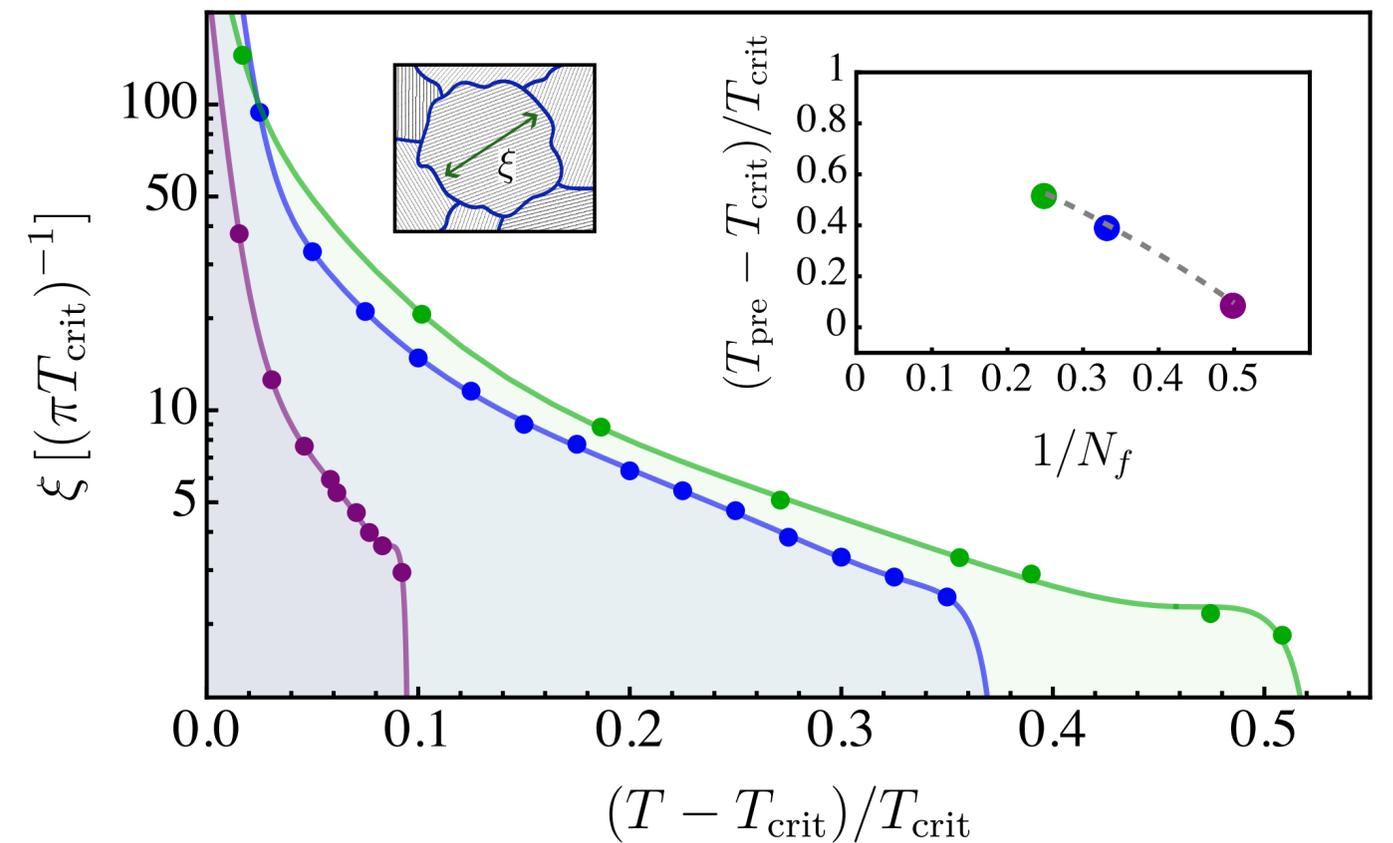
- Competing counteracting effects: fermions and bosons
- External parameter:  $T, \mu, \dots$
- Massless modes: exact Goldstones



# Many flavour scaling

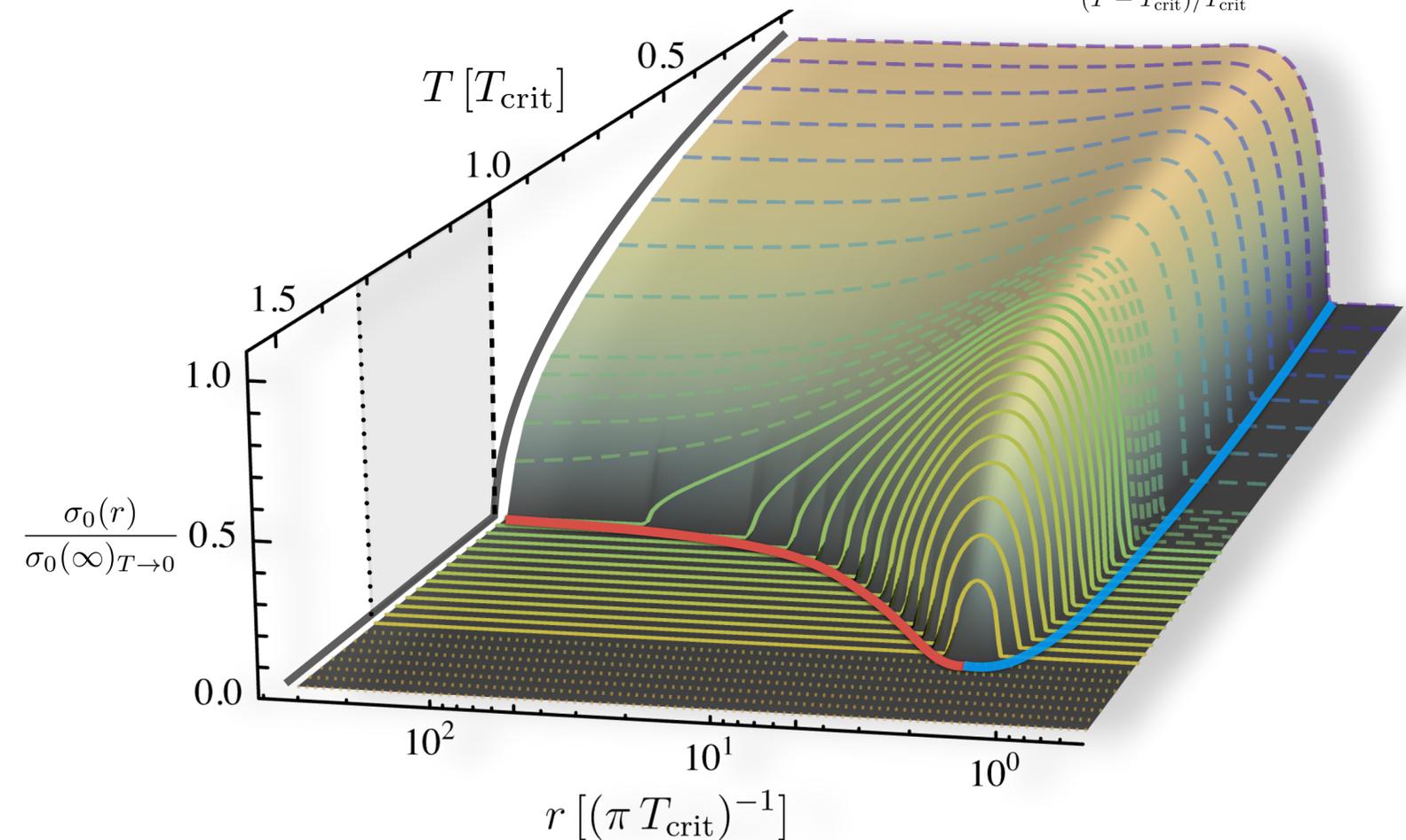
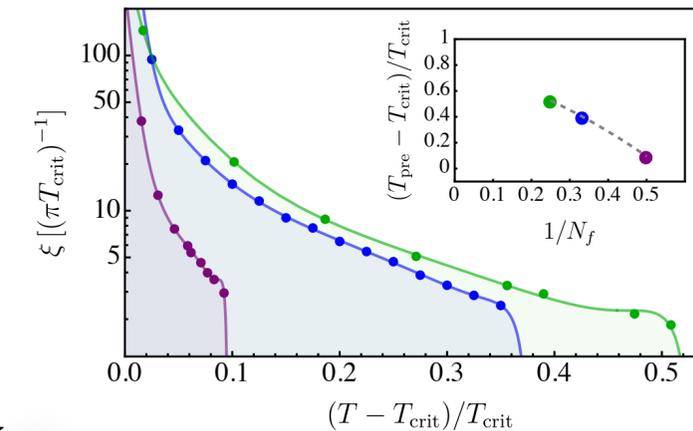
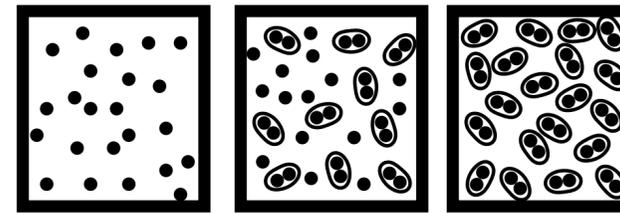
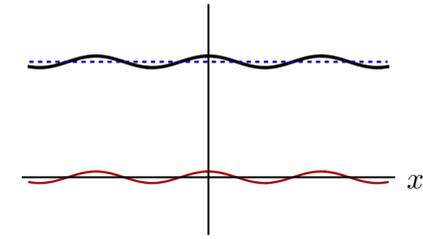
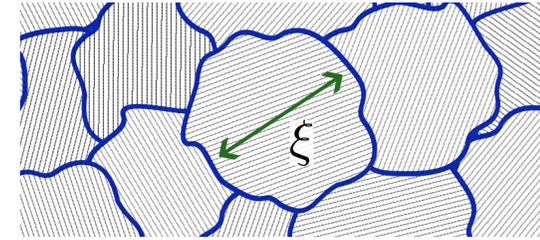


- **Growing size** of precondensation regime in  $T$
- **Growing number of Goldstones**  $N_f^2 - 1$
- **Relevant role** in the near-critical dynamics



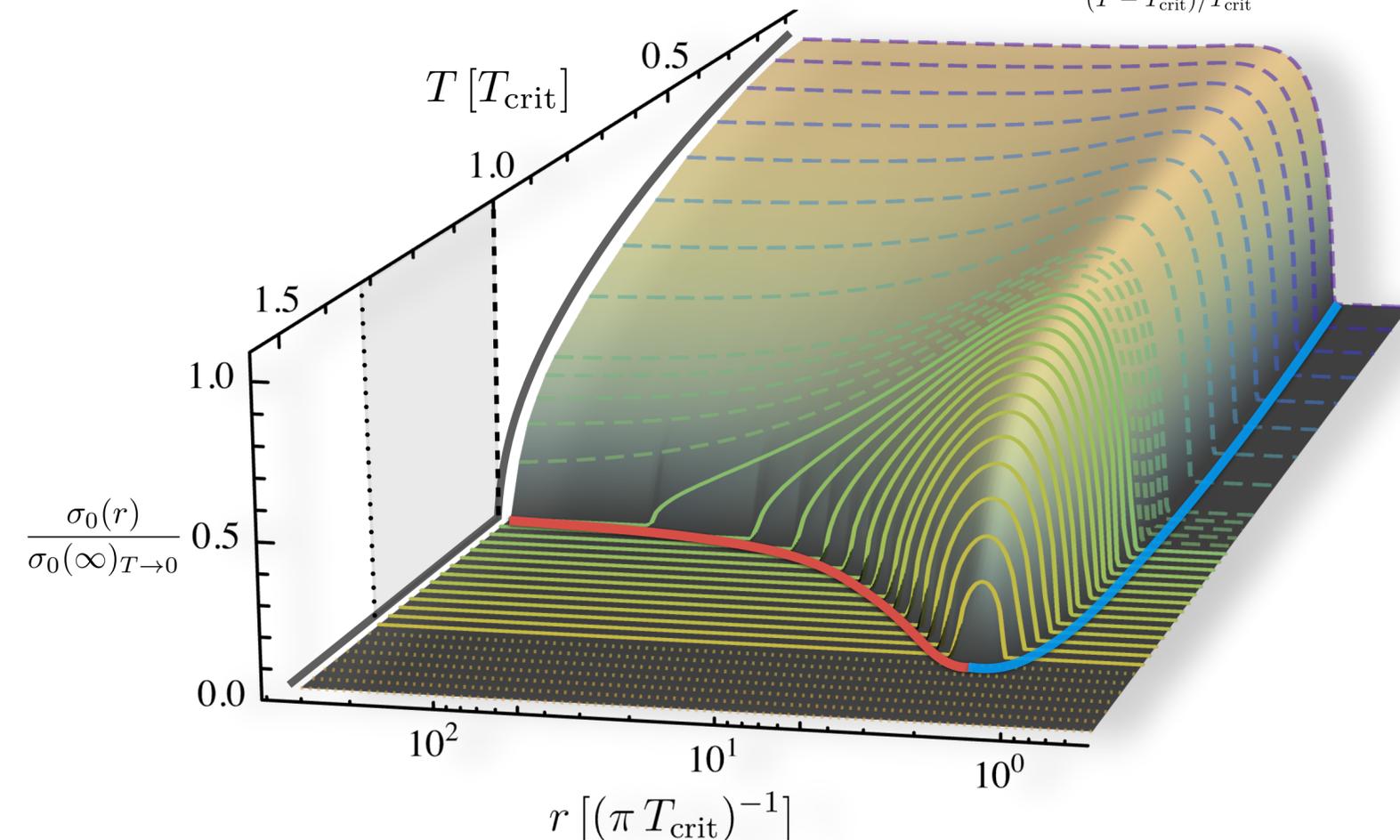
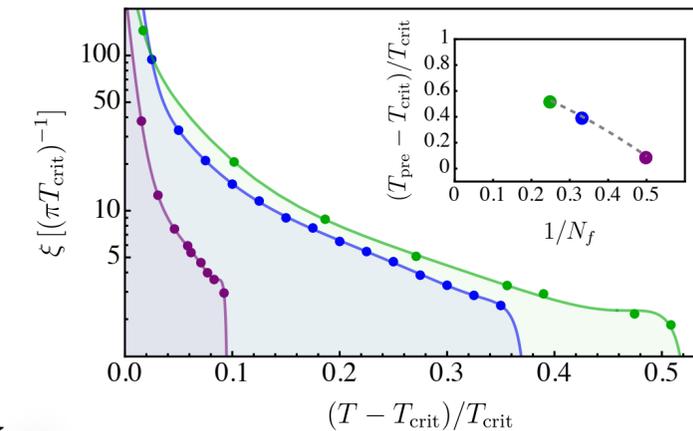
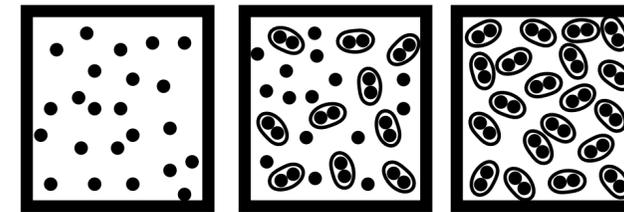
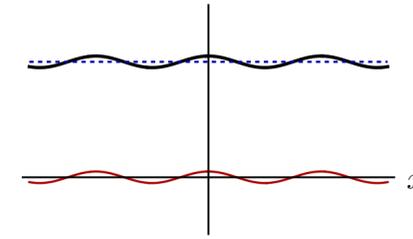
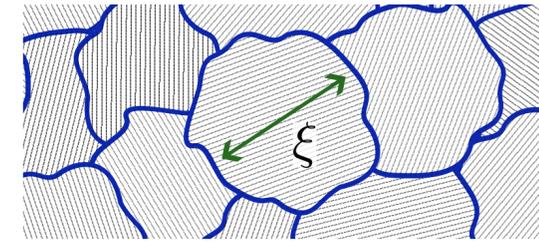
# Conclusions

- **Thermal precondensation** has interesting features
  - **Domains, spatial modulations, mixed phases**
- Occurs in **multiple contexts in nature** (condensed matter, cold atoms, gauge-fermion theories in the chiral limit)
- **Key ingredients:**
  - Competing counteracting effects
  - External parameters
  - Massless modes
- Increasingly relevant in the **many-flavour limit**
- Potentially relevant for **BSM physics** and **observational prospects**



# Conclusions

- **Thermal precondensation** has interesting features
  - **Domains, spatial modulations, mixed phases**
- Occurs in **multiple contexts in nature** (condensed matter, cold atoms, gauge-fermion theories in the chiral limit)
- **Key ingredients:**
  - Competing counteracting effects
  - External parameters
  - Massless modes
- Increasingly relevant in the **many-flavour limit**
- Potentially relevant for **BSM physics** and **observational prospects**



Thank you for your attention!