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QUANTUM EFFECTS ON NEUTRINO PARAMETERS FROM A FLAVORED GAUGE BOSON

GENERATING NEUTRINO MASSES VIA ONE-LOOP RGE RUNNING



2nd Hokkaido Workshop on Particle Physics at Crossroads

2026/03/06



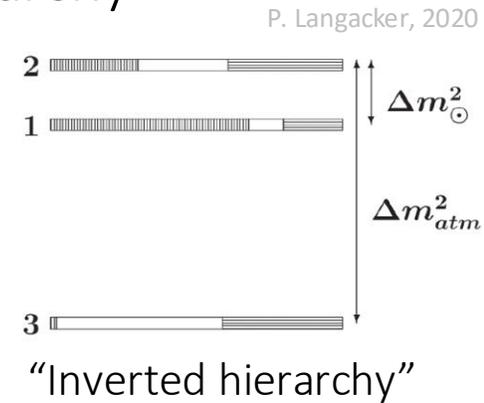
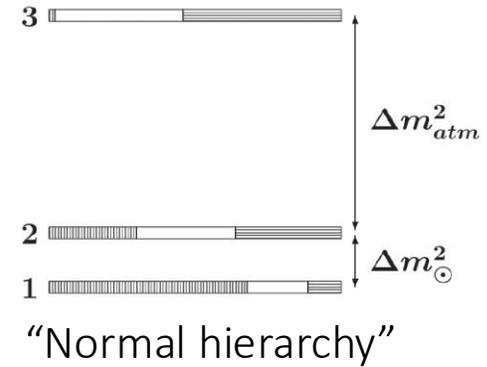
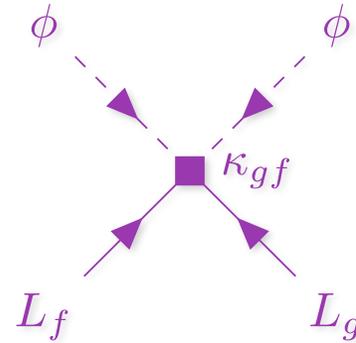
INTRODUCTION

- We know neutrinos have mass from oscillation data, don't know their hierarchy
- Describe using dim. 5 effective SM interaction

$$\mathcal{L}_\kappa \sim \kappa_{gf} L_g L_f \phi \phi \rightsquigarrow \kappa_{gf} v^2 \nu_g \nu_f$$

Weinberg, 1979

$$\underbrace{\kappa_{gf} v^2}_{\sim m_\nu^{gf}}$$



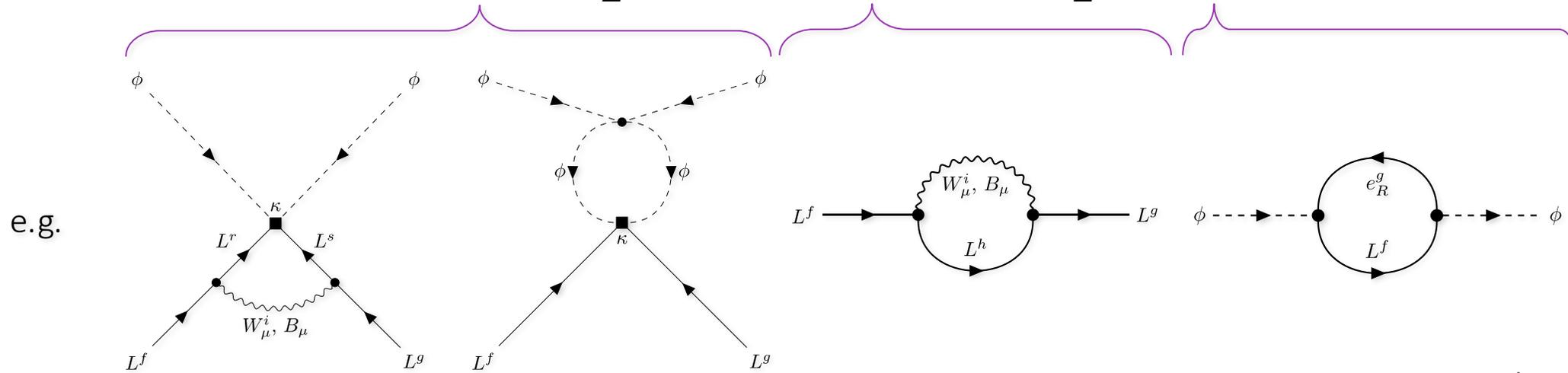
P. Langacker, 2020

- Loop effects may be important; different scales involved in mass generation, creation, detection → Running!

BETA FUNCTION THE WEINBERG OPERATOR

- General formula: $\beta_{\kappa} = \delta\kappa_{,1} - \frac{1}{2}(\delta Z_{\phi,1} + \delta Z_{l,1})^T \kappa - \frac{1}{2}\kappa(\delta Z_{\phi,1} + \delta Z_{l,1})$

Babu+, 1992; Casas+, 2000; Antusch+, 2001



$$\alpha = \lambda - 3g_2^2 + 2\text{Tr}(3Y_u^\dagger Y_u + 3Y_d^\dagger Y_d + Y_e^\dagger Y_e)$$

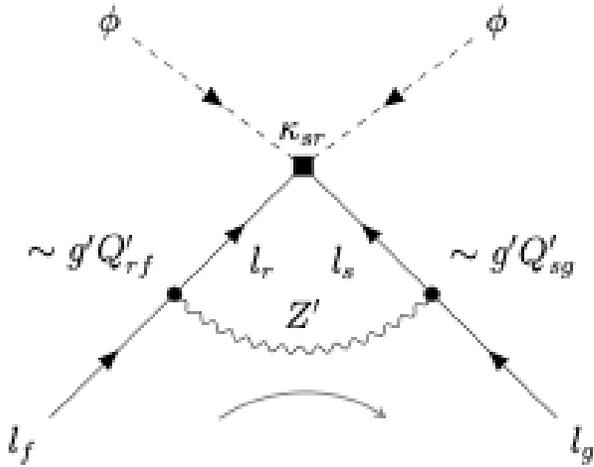
- Has structure $\beta_{\kappa}^{(\text{SM})} = \alpha \kappa + P^T \kappa + \kappa P$ $P = -\frac{3}{2}(Y_e^\dagger Y_e)$

- We consider effects of flavor-dependent gauge couplings, such as $U(1)_{L_\mu - L_\tau}$

- Motivations: oscillation data, $g_\mu - 2$, dark matter, anomaly freedom,...

GENERATING MASSES VIA RGE RUNNING

Vertex renormalization



$$\frac{d\kappa}{dt} = \alpha\kappa + P^T \kappa + \kappa P + G^T \kappa G$$

$$G = \sqrt{6}g'Q'$$

SM

$U(1)_{L_\mu - L_\tau}$

$$\frac{d\kappa_i}{dt} = \alpha\kappa_i + 2(U P U^T)_{ii} \kappa_i + \sum_{k=1}^n [(U G U^T)_{ki}]^2 \kappa_k$$

If initial $\kappa_i = 0$, stays 0

If one initial $\kappa_k \neq 0$,
generates n hierarchical
masses

At one loop, flavor-dep. gauge interaction only way
to get $\beta_\kappa \supset G^T \kappa G$ and raise rank of mass matrix

MIXING ANGLE CONDITION FOR DEGEN. EIGENVALUES

$$\frac{d\kappa_i}{dt} = \alpha \kappa_i + 2(U P U^T)_{ii} \kappa_i + \sum_{k=1}^n [(U G U^T)_{ki}]^2 \kappa_k \quad \tilde{A} = U A U^\dagger$$

$$\frac{dU}{dt} = T U \quad 16\pi^2 T_{ij} = \begin{cases} i \sum_{k=1}^n \text{Im}[\tilde{G}_{ki}^2] \frac{\kappa_k}{2 \kappa_i} & \text{for } i = j, \\ \frac{\kappa_i + \kappa_j}{\kappa_i - \kappa_j} \text{Re}[\tilde{P}_{ij}] + \sum_{k=1}^3 \frac{\kappa_k}{\kappa_i - \kappa_j} \text{Re}[\tilde{G}_{ki} \tilde{G}_{kj}] \\ + i \left(\frac{\kappa_i - \kappa_j}{\kappa_i + \kappa_j} \text{Im}[\tilde{P}_{ij}] + \sum_{k=1}^3 \frac{\kappa_k}{\kappa_i + \kappa_j} \text{Im}[\tilde{G}_{ki} \tilde{G}_{kj}] \right) & \text{for } i \neq j, \end{cases}$$

EXAMPLE: NORMAL HIERARCHY

For $\kappa_1 = \kappa_2 \ll \kappa_3$ we find

$$|\kappa_1|_{M_{Z'}} \simeq \min_{i=1,2} \left\{ \left| (a - bW_{1i}^2 - bW_{2i}^2)\kappa_2 - bW_{3i}^2\kappa_3 \right| \right\},$$

$$|\kappa_2|_{M_{Z'}} \simeq \max_{i=1,2} \left\{ \left| (a - bW_{1i}^2 - bW_{2i}^2)\kappa_2 - bW_{3i}^2\kappa_3 \right| \right\},$$

$$|\kappa_3|_{M_{Z'}} \simeq |\kappa_3|$$

$$W_{31} \simeq 0 : \quad \sin \theta_{13} \simeq -\frac{\sin \theta_{12} \cos \theta_{12} \tan \theta_{23}}{1 + \sin^2 \theta_{12}}$$

$$a = 1 - \alpha / (16\pi^2) \log(M_1/M_{Z'})$$

$$b = 6g'^2 / (16\pi^2) \log(M_1/M_{Z'})$$

$$W_{ij} = \tilde{G}_{ij} / (\sqrt{6}g')$$

For $\kappa_1 = \kappa_2 = 0$, at leading log

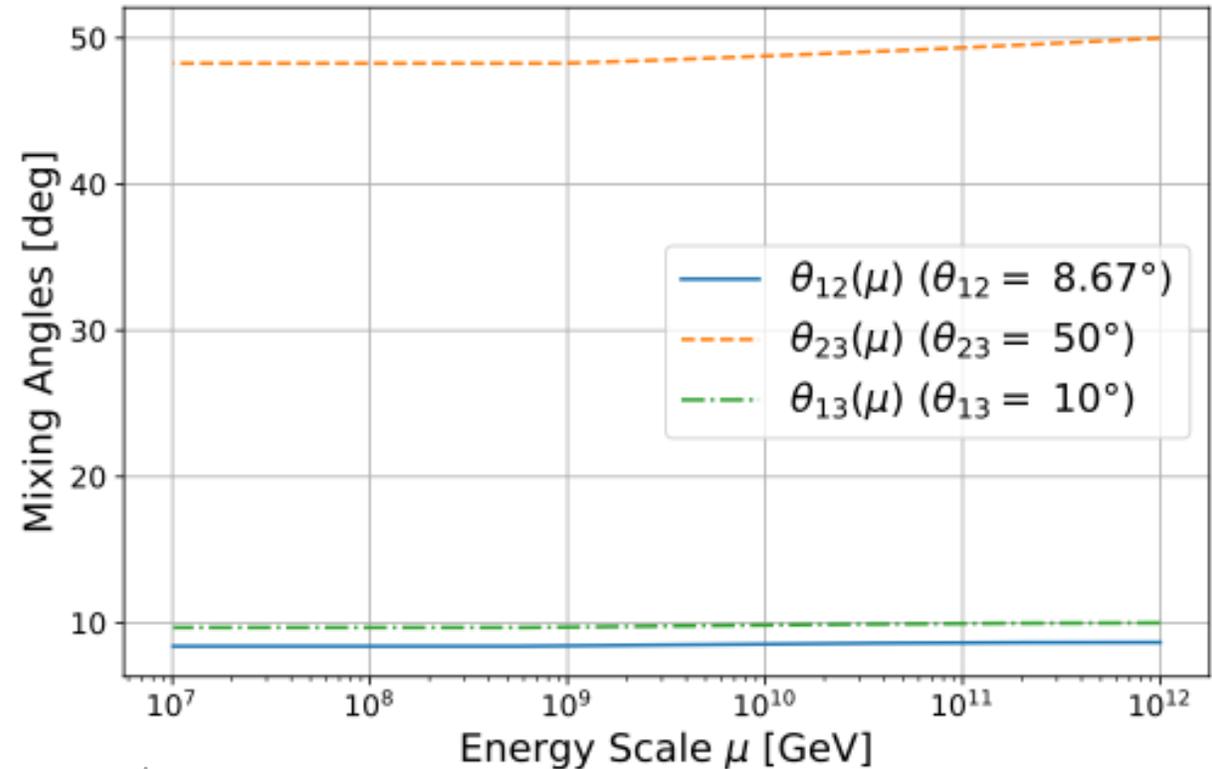
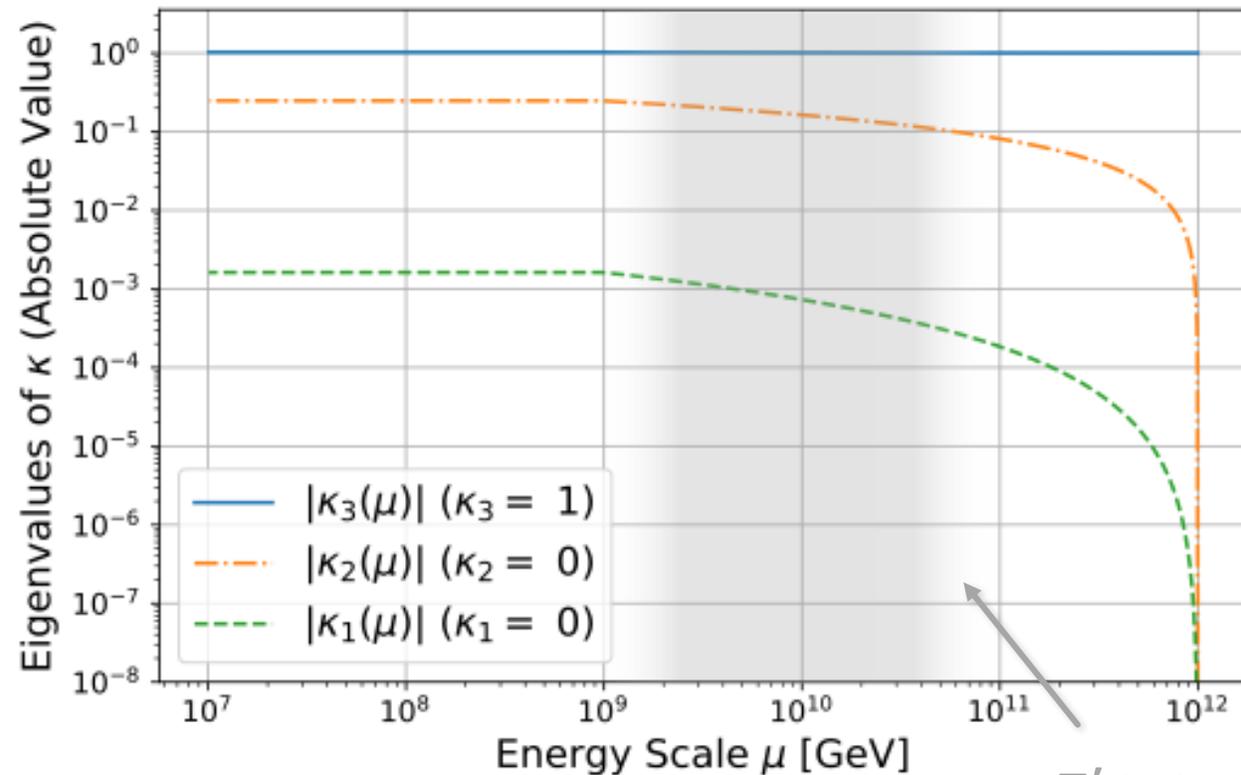
$$\kappa_1|_{M_{Z'}} \simeq 0$$

$$\kappa_2|_{M_{Z'}} \simeq -bW_{32}^2\kappa_3$$

$$\kappa_3|_{M_{Z'}} \simeq \kappa_3$$

$\kappa_1|_{M_{Z'}}$ generated at
subleading log level $\sim b^2$

EXAMPLE: NORMAL HIERARCHY



Generate two new masses via
one-loop RGE
Can reproduce exp. splittings

Z' is integrated out
→ running stops

Including CP, Majorana
phases can reproduce
exp. angles

BACKUP

UV COMPLETION: EXAMPLE

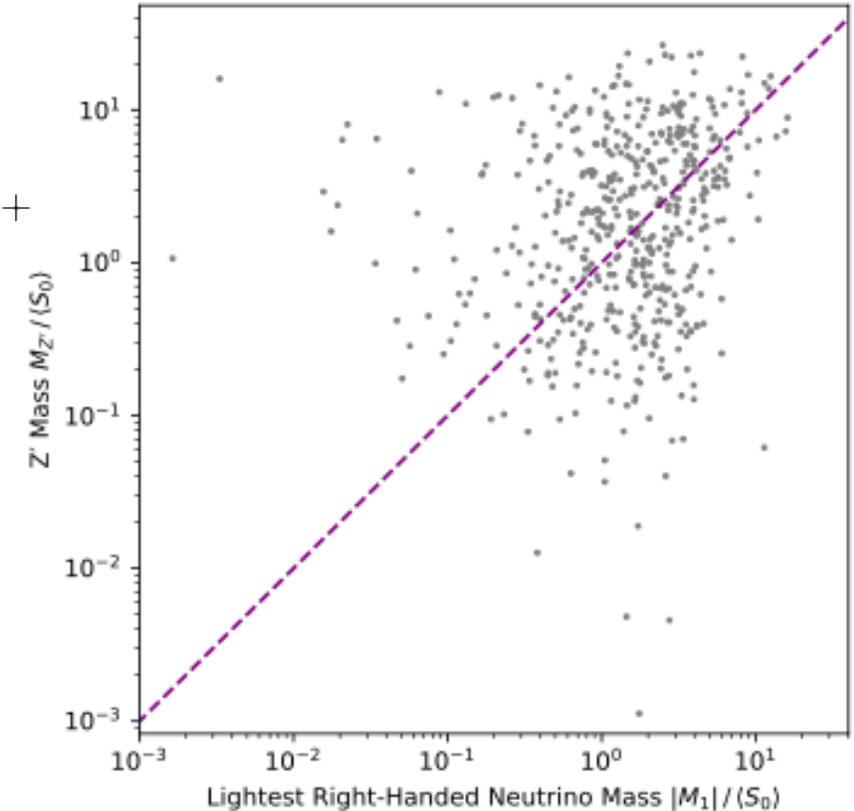
Field	$U(1)_{L_\mu-L_\tau}$ charge	Field	$U(1)_{L_\mu-L_\tau}$ charge
N_1	0	S_0	0
N_2	1	S_1	-1
N_3	-1	S_2	-2

$$2\mathcal{L}_{\text{mass}} = M_{11}\overline{N_1^C}N_1 + M_{23}\overline{N_2^C}N_3 + \text{h.c.},$$

$$2\mathcal{L}_{\text{Yuk}} = 2Y_{\nu,e1}\bar{l}_e\epsilon\phi^*N_1 + 2Y_{\nu,\mu2}\bar{l}_\mu\epsilon\phi^*N_2 + 2Y_{\nu,\tau3}\bar{l}_\tau\epsilon\phi^*N_3 + \lambda_{11}S_0\overline{N_1^C}N_1 + \lambda_{23}S_0\overline{N_2^C}N_3 + \\ + \lambda_{12}S_1\overline{N_1^C}N_2 + \lambda_{13}S_1^\dagger\overline{N_1^C}N_3 + \lambda_{22}S_2\overline{N_2^C}N_2 + \lambda_{33}S_2^\dagger\overline{N_3^C}N_3 + \text{h.c.}$$

$$M_R = \begin{pmatrix} M_{11} + \lambda_{11}\langle S_0 \rangle & \lambda_{12}\langle S_1 \rangle & \lambda_{13}\langle S_1 \rangle \\ \lambda_{12}\langle S_1 \rangle & \lambda_{22}\langle S_2 \rangle & M_{23} + \lambda_{23}\langle S_0 \rangle \\ \lambda_{13}\langle S_1 \rangle & M_{23} + \lambda_{23}\langle S_0 \rangle & \lambda_{33}\langle S_2 \rangle \end{pmatrix}$$

$$M_{Z'} = \sqrt{2}g' \cdot \sqrt{\langle S_1 \rangle^2 + 4\langle S_2 \rangle^2}$$



NON-ABELIAN GAUGE GROUPS

$$\beta_{\kappa} = \alpha \kappa + P^T \kappa + \kappa P + G^T \kappa G + \frac{1}{2} [G_+^T \kappa G_- + G_-^T \kappa G_+]$$

$$\frac{d\kappa_i}{dt} = \alpha \kappa_i + 2\tilde{P}_{ii} \kappa_i + \sum_{k=1}^n \text{Re} \left\{ [\tilde{G}_{ki}]^2 + [\tilde{G}_{+,ki} \tilde{G}_{-,ki}] \right\} \kappa_k$$

Enhanced running!

