

# Gravitational Wave Sources and Origin of Massive Binary Black Holes

Tomoya Kinugawa

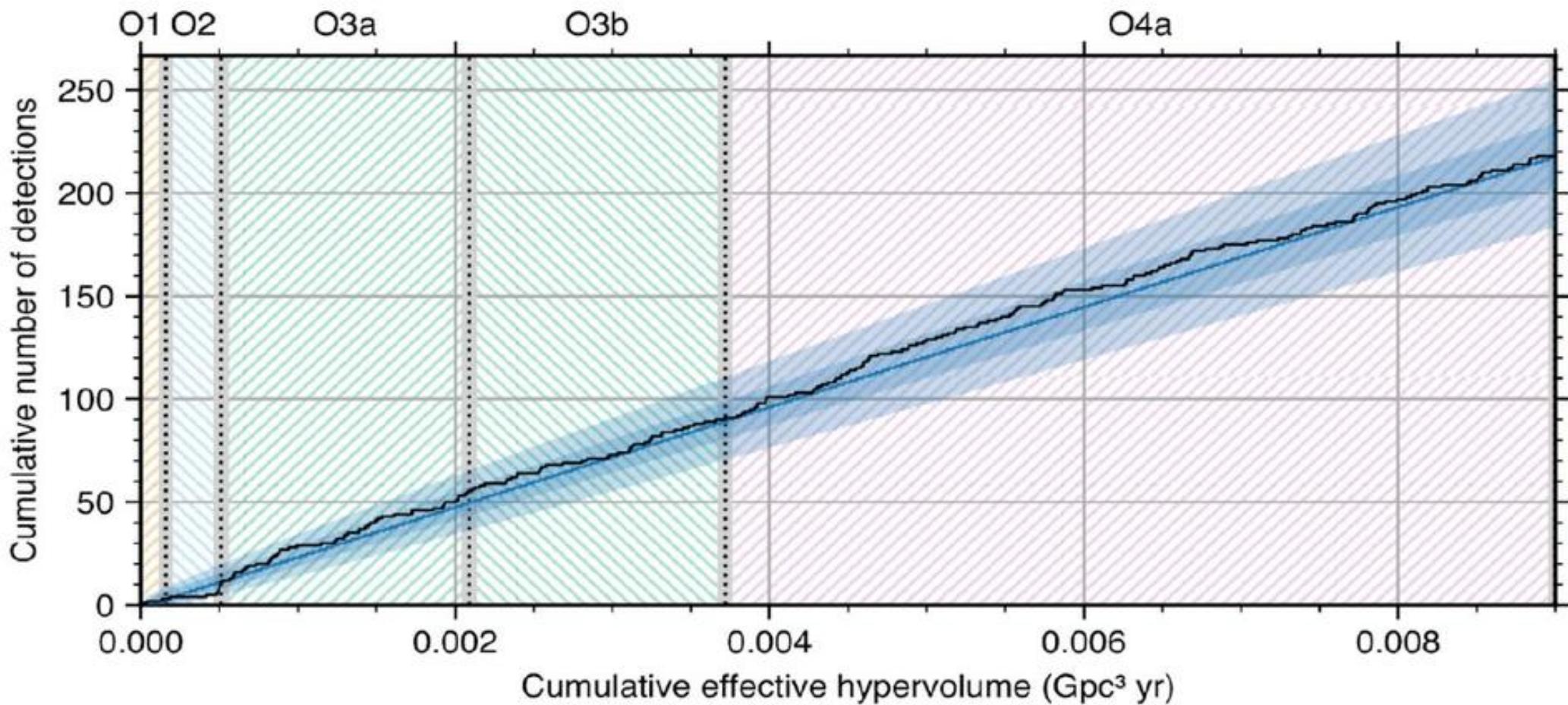
Shinshu University

# GW observatories



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# detections



The cumulative number of detections (candidates found with a probability of being astrophysical greater than 50%) against the approximate space-time hypervolume surveyed by the detectors (source: LVK consortium).

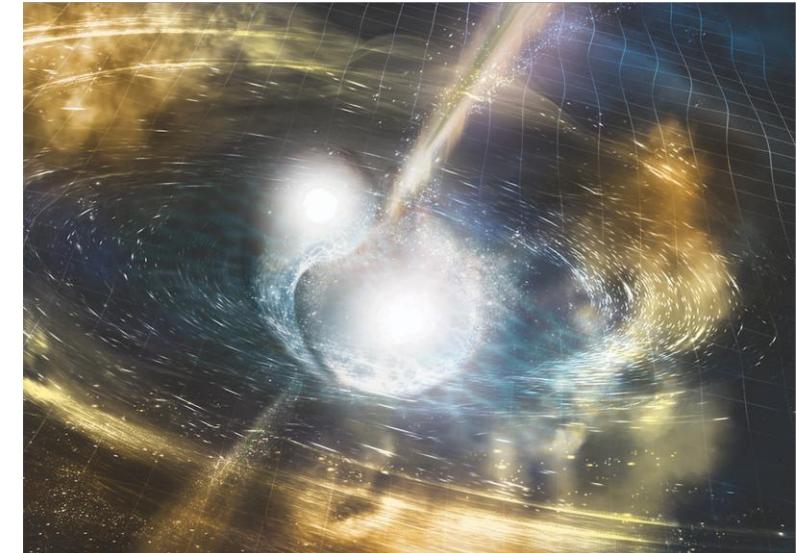
# GW sources



**Binary Black Hole**

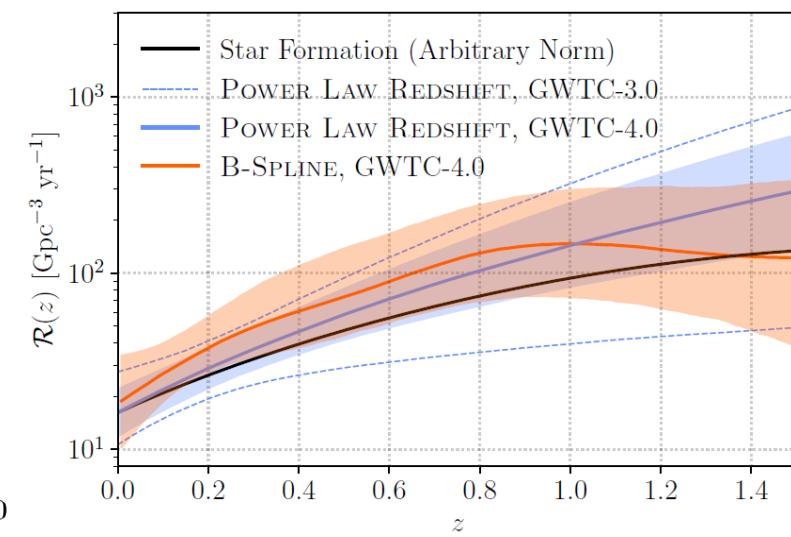
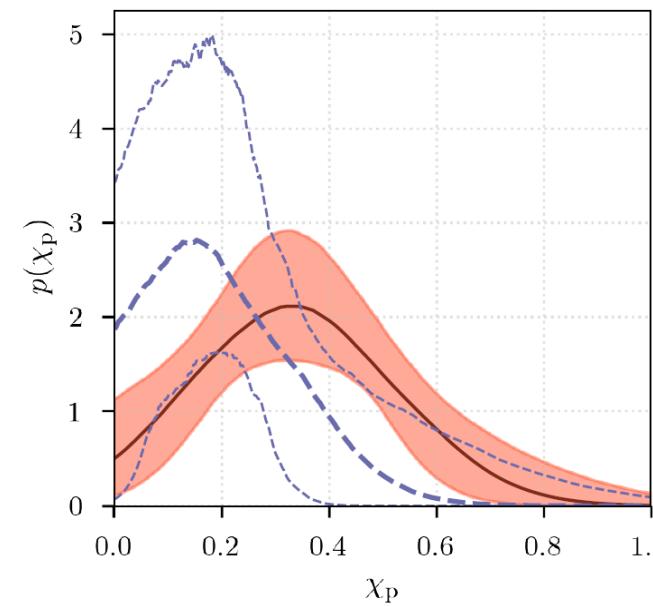
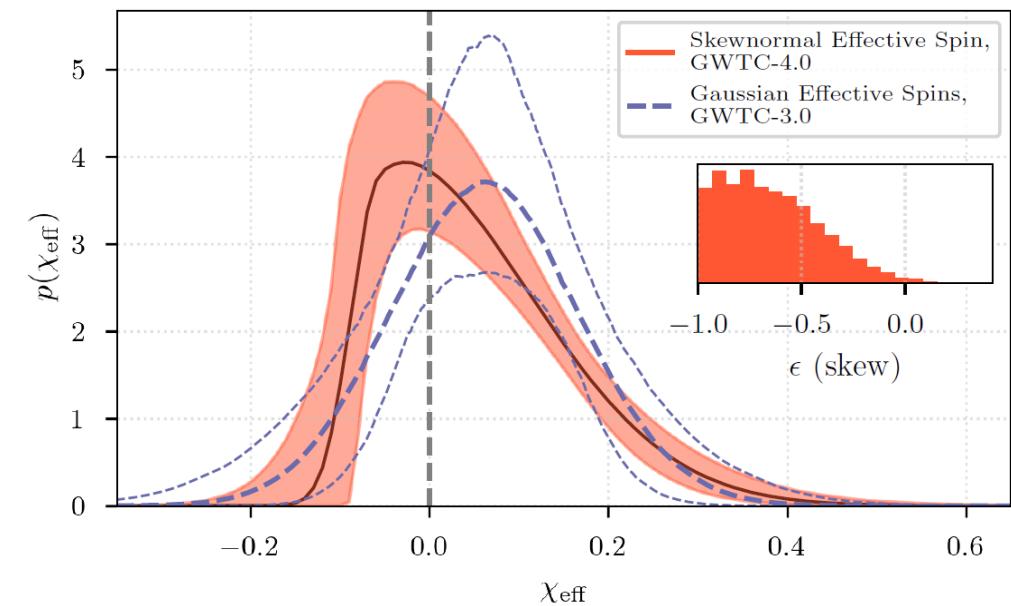
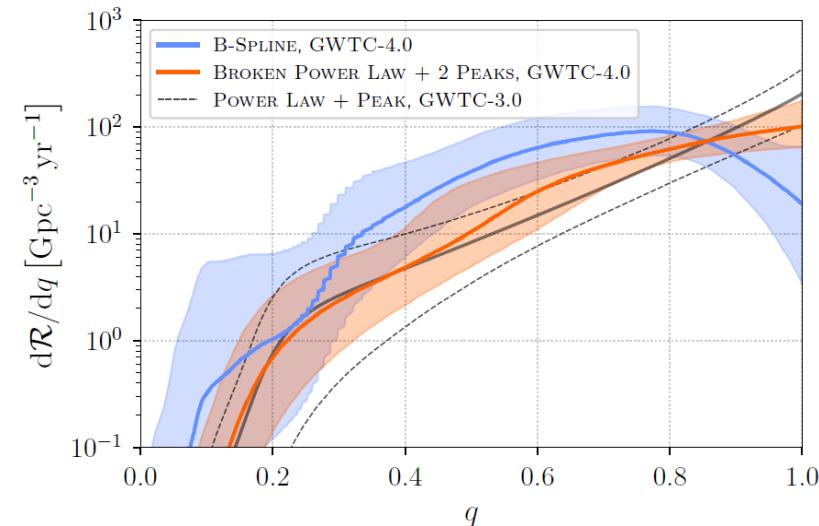
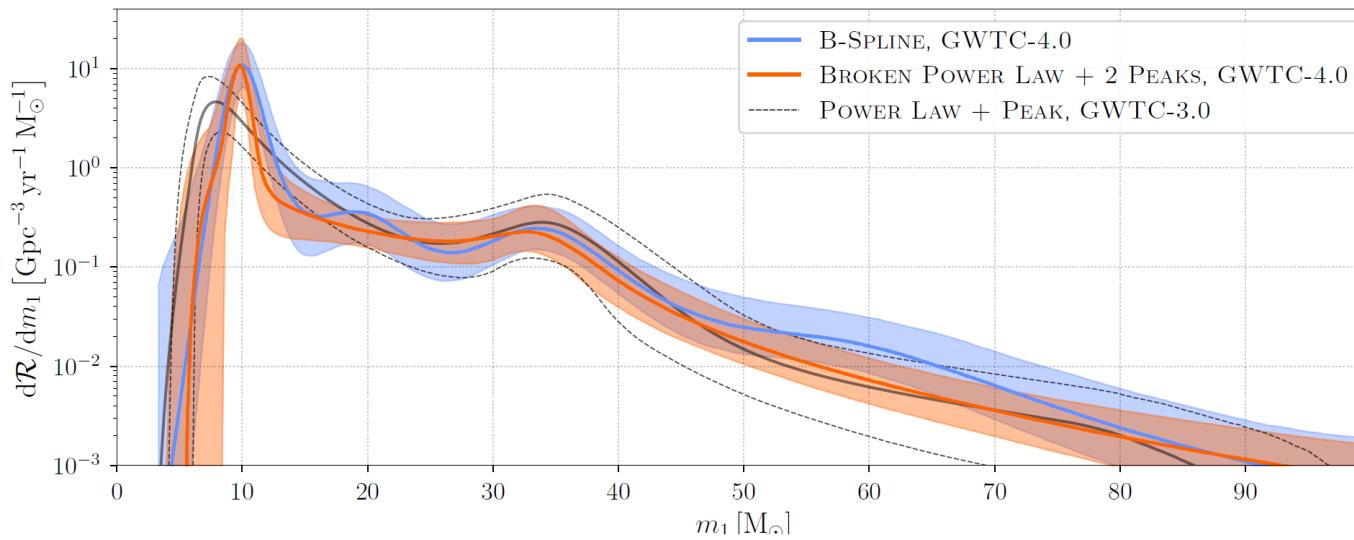


**Neutron Star Black Hole  
Binary**



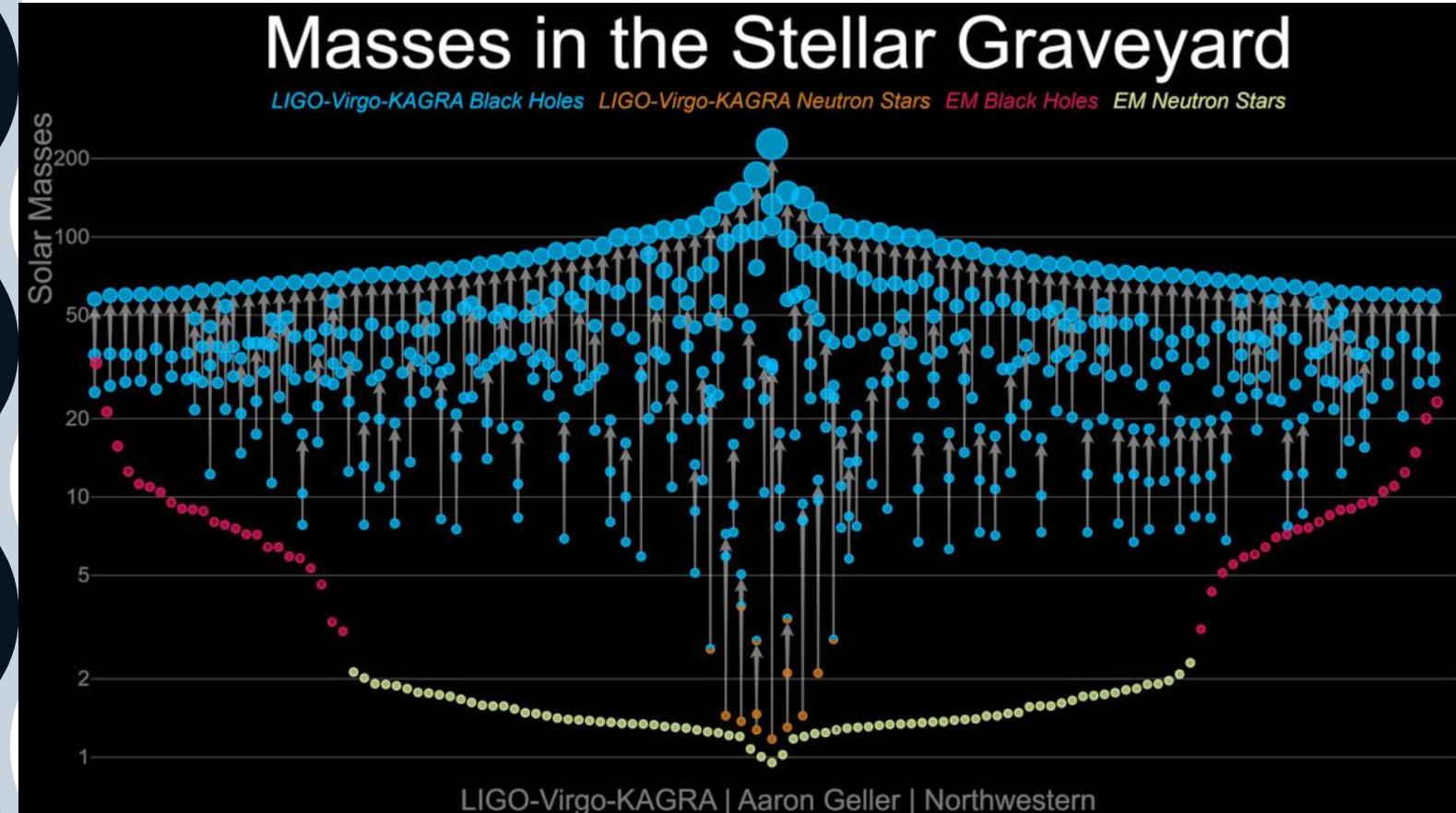
**Binary Neutron Star**

# BBH merger profile



# Mystery of GW events

- GW events show that there are many massive BHs ( $\gtrsim 30$  Msun).
- On the other hand, the typical mass of BHs in X-ray binaries is  $\sim 10$  Msun.



# Origin of massive BBHs

In order to explain the origin of such massive BBHs

Many theories exist such as

- 1) Pop I and Pop II BBH (present day and low metal stars)
- 2) Pop III BBH (First stars)
- 3) Dynamical formation (Dense stellar environment)
- 4) AGN disk
- 5) Primordial BBH
- .....

# Origin of massive BBHs

In order to explain the origin of such massive BBHs

Many theories exist such as

**Stellar origin BH**

- 1) Pop I and Pop II BBH (present day and low metal stars)
- 2) Pop III BBH (First stars)
- 3) Dynamical formation (Dense stellar environment)
- 4) AGN disk
- 5) Primordial BBH
- .....

**Non-stellar origin BH**

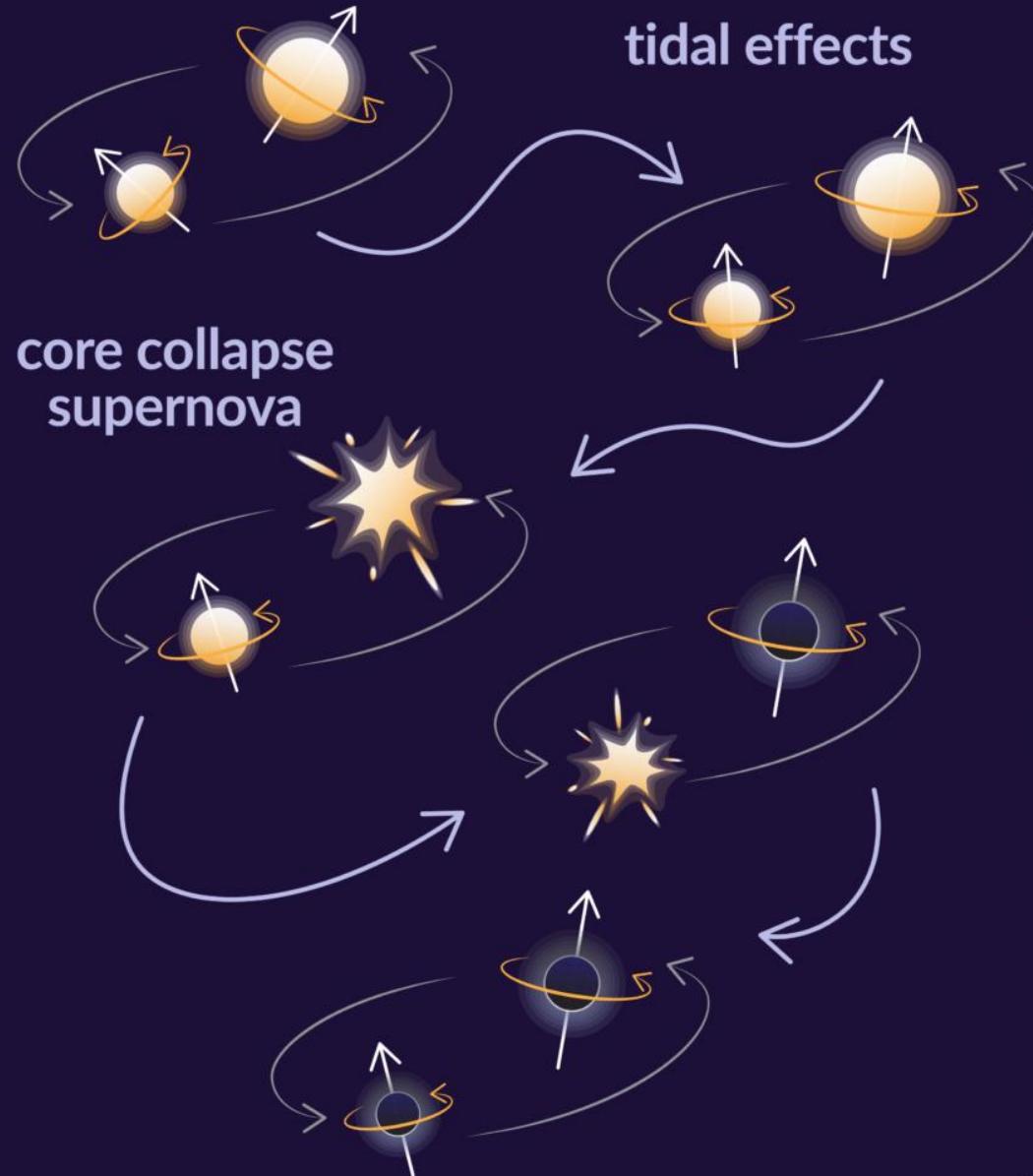
# Origin of massive BBHs

In order to explain the origin of such massive BBHs

Many theories exist such as **Isolated Binary**

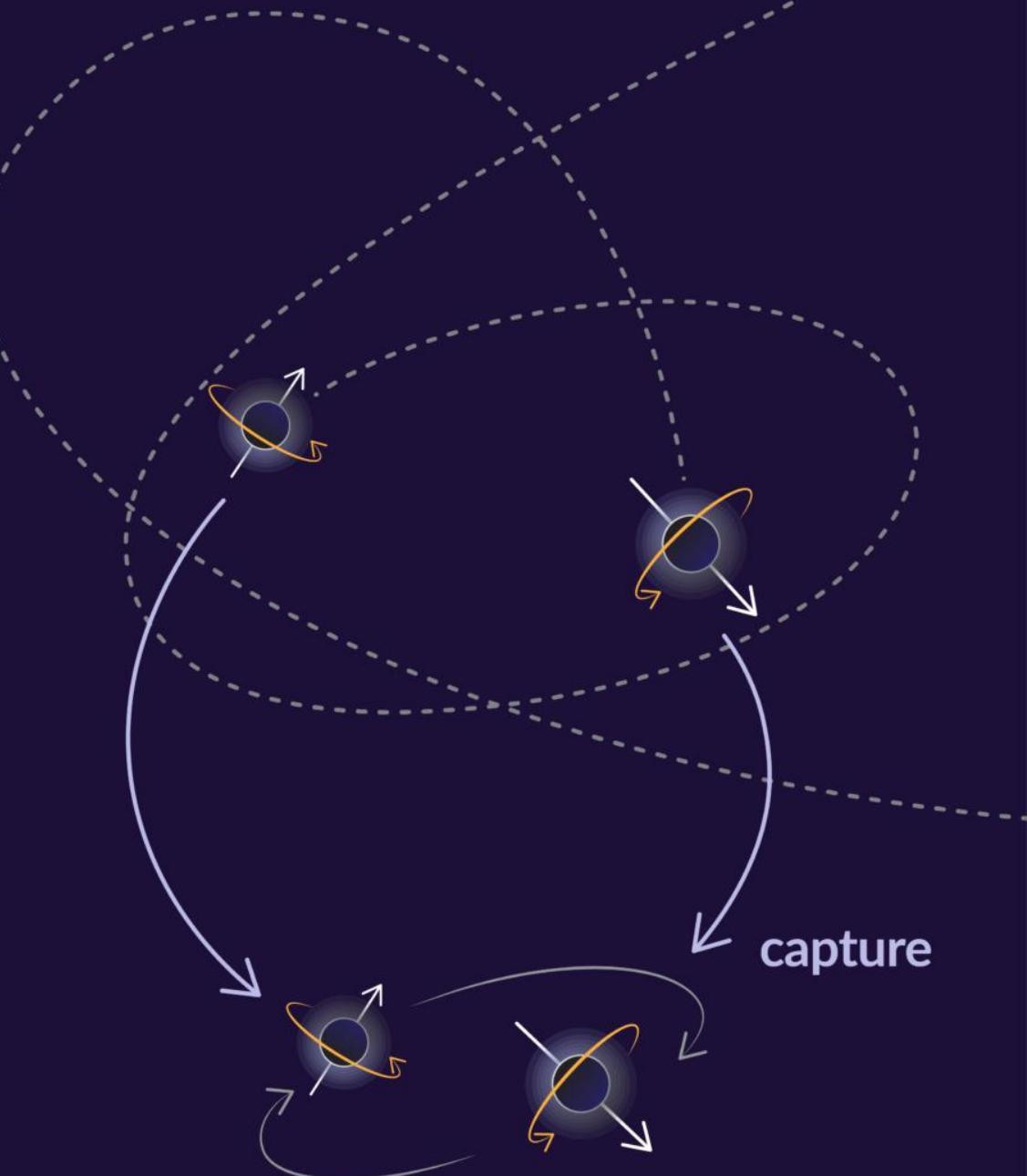
- 1) Pop I and Pop II BBH (present day and low metal stars)
- 2) Pop III BBH (First stars)
- 3) Dynamical formation (Dense stellar environment)
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- 5) Primordial BBH
- .....

**Dynamical**

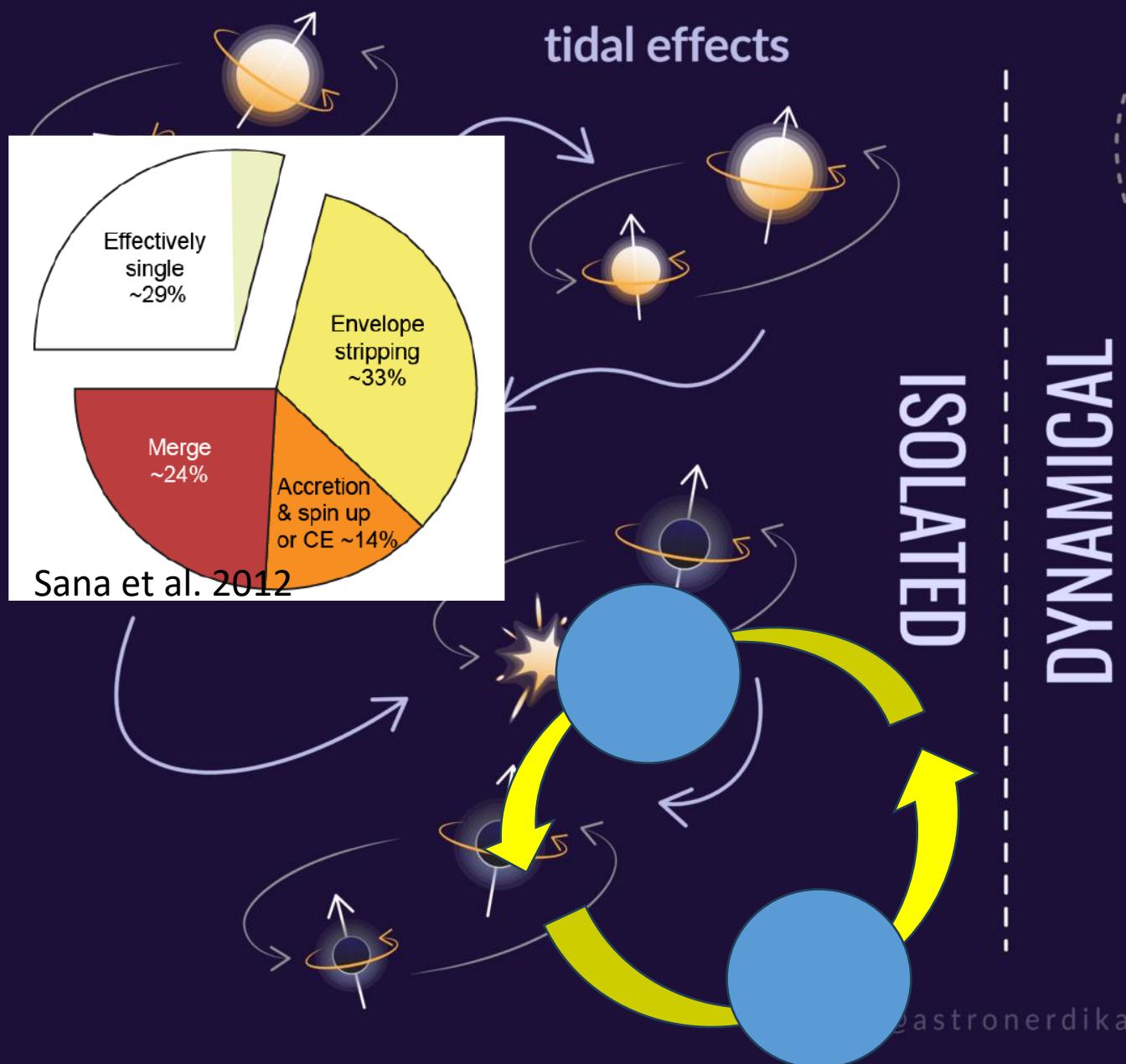


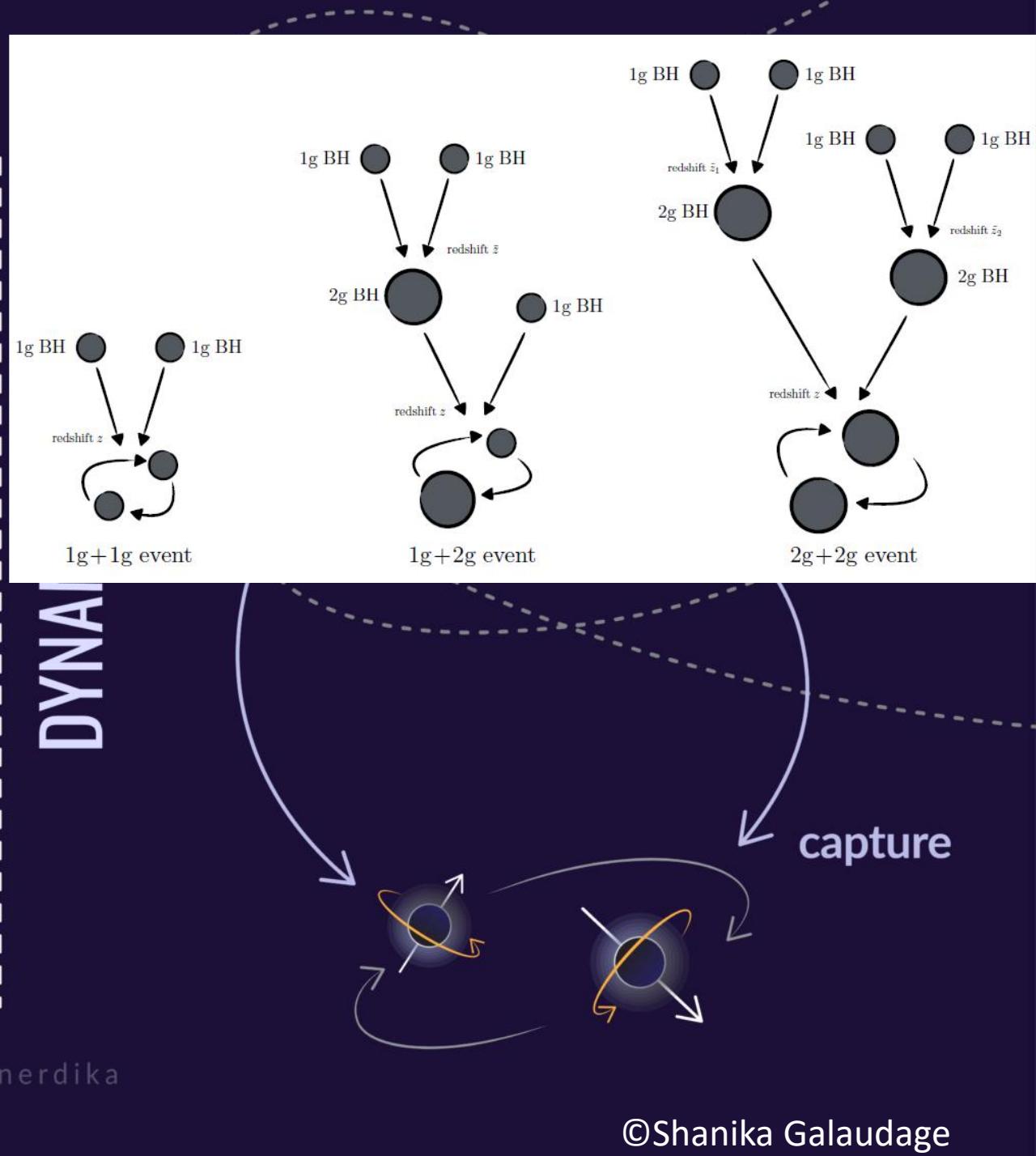
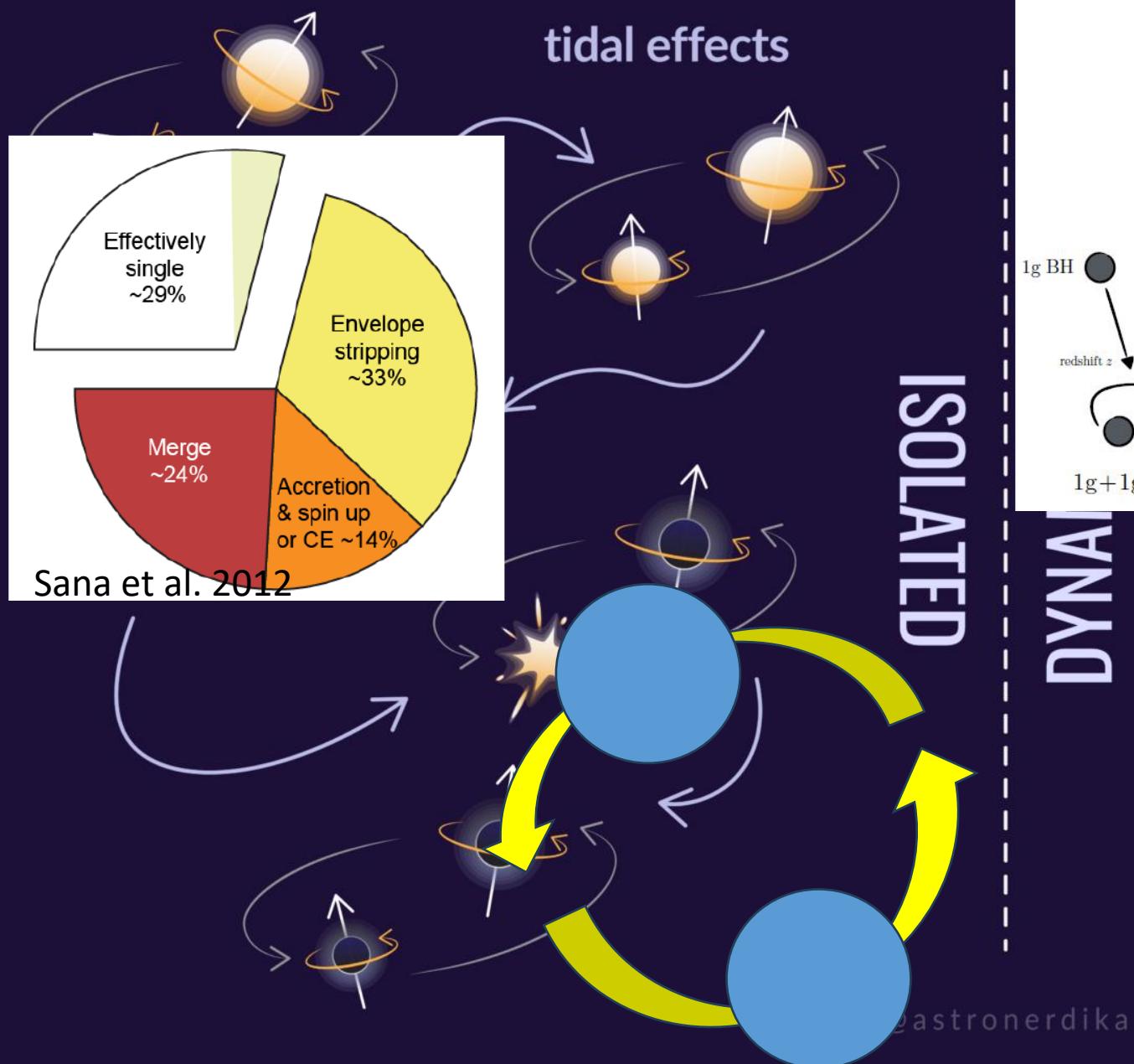
ISOLATED  
DYNAMICAL

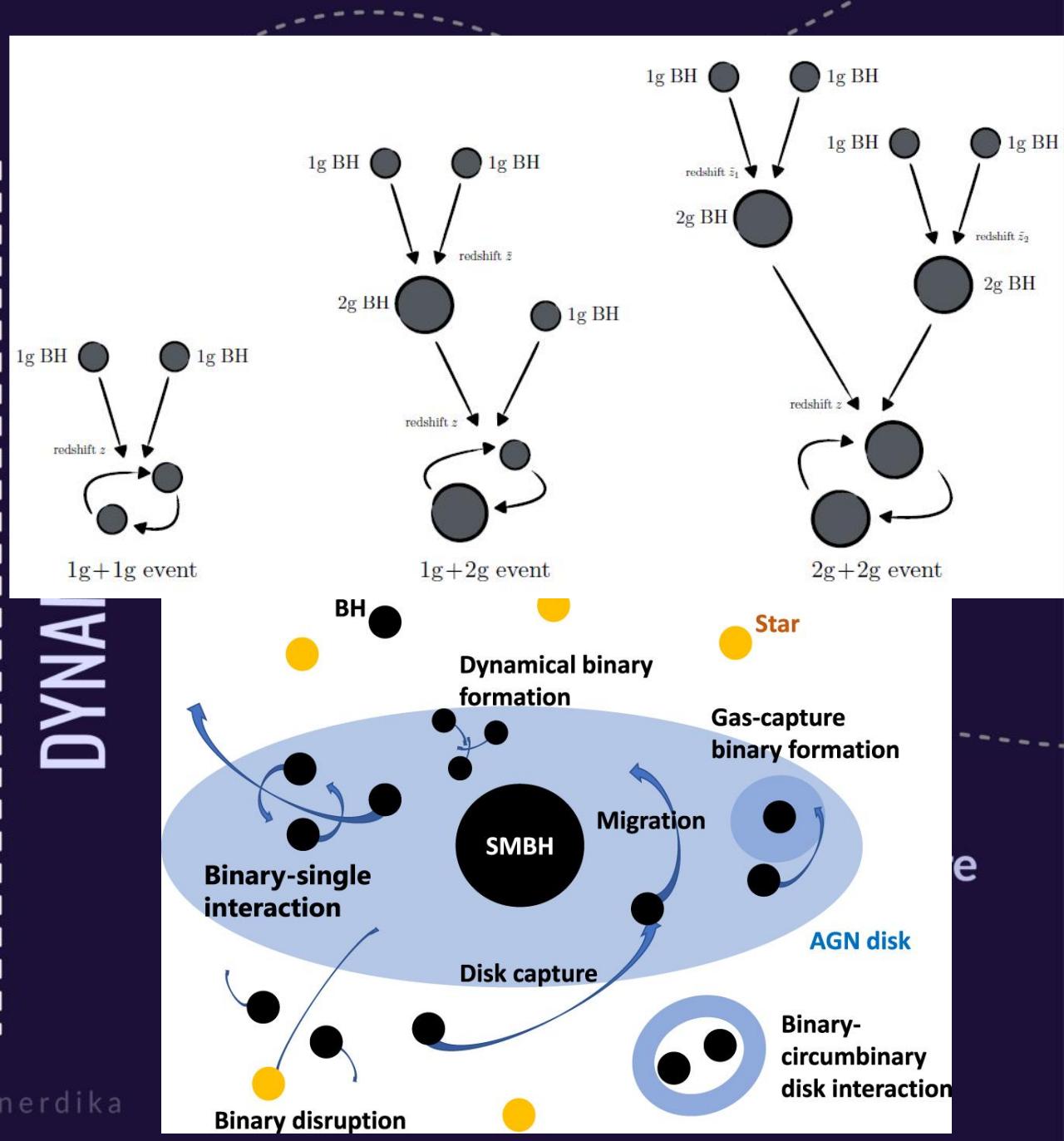
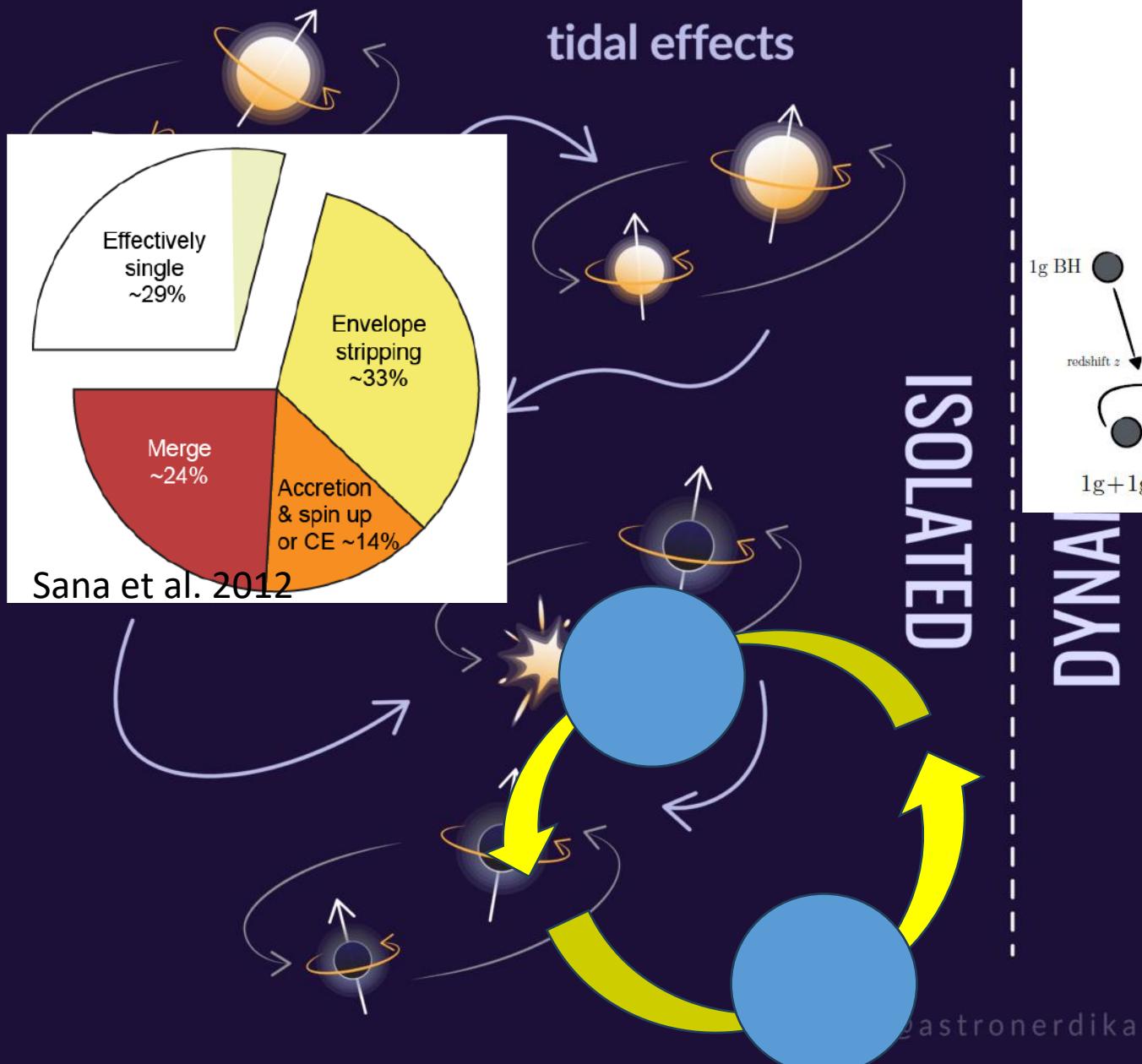
@astronerdika



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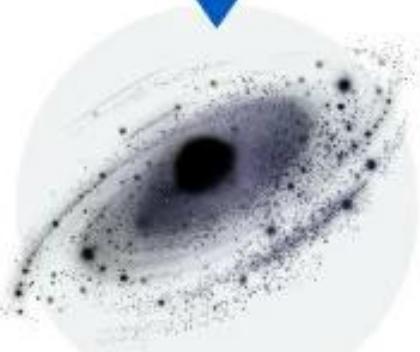


# Primordial Black Hole

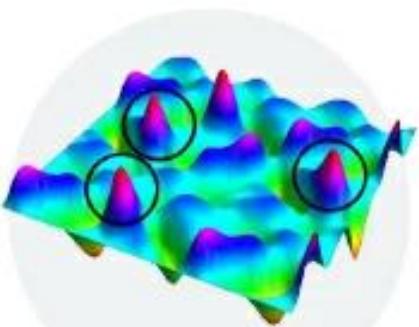


Star

↓  
10 billion years



Black hole



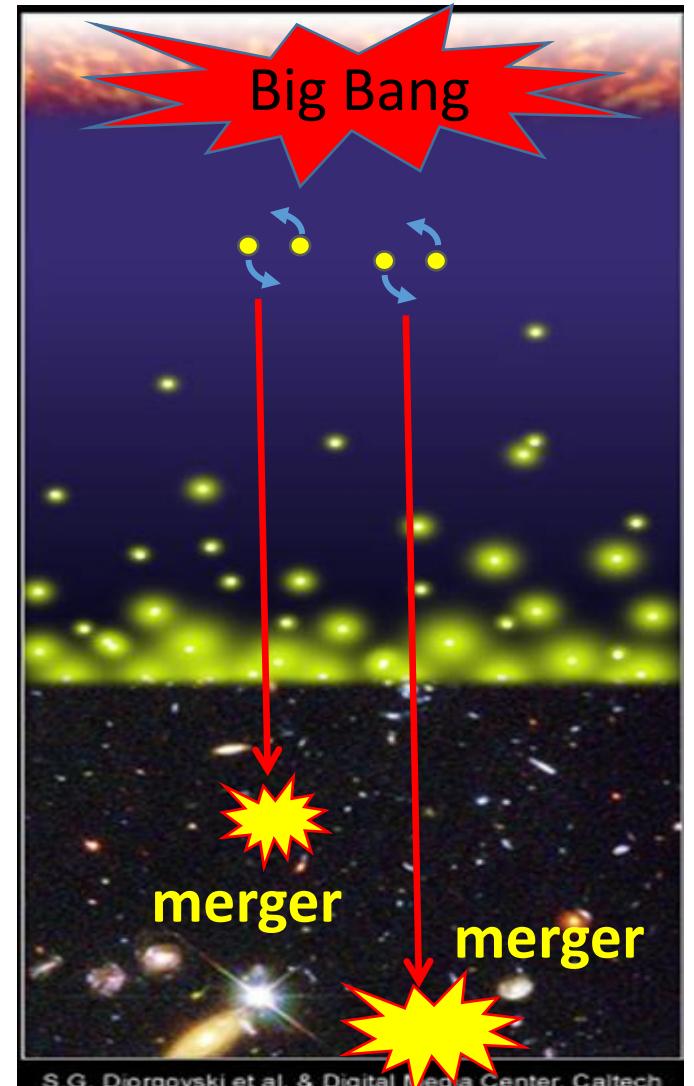
Energy fluctuations immediately  
after the birth of the universe

↓  
0.1 ms



Primordial black hole

time



S.G. Djorgovski et al. & Digital Media Center, Caltech

Djorgovski et al.&Digital Media  
Center

# Origin of massive BBHs

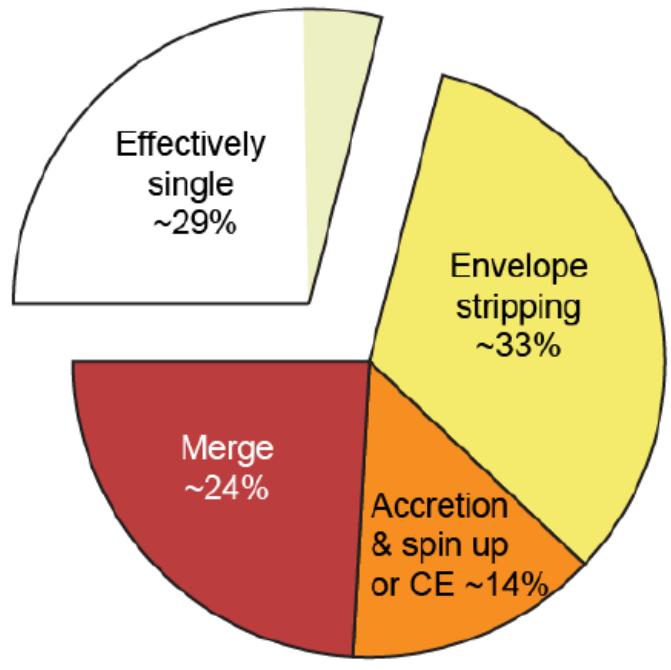
In order to explain the origin of such massive BBHs

Many theories exist such as **Isolated Binary**

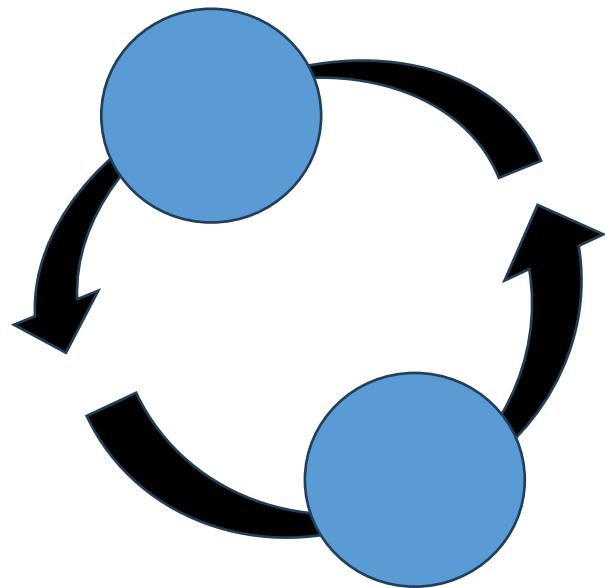
- 1) Pop I and Pop II BBH (present day and low metal stars)
- 2) Pop III BBH (First stars)
- 3) Dynamical formation (Dense stellar environment)
- 4) AGN disk
- 5) Primordial BBH
- .....

# Why is isolated binary important for BBH?

- Binary fraction of massive stars is high (~70% e.g. Sana et al. 2012)
- Almost BH progenitors might evolve in binary systems

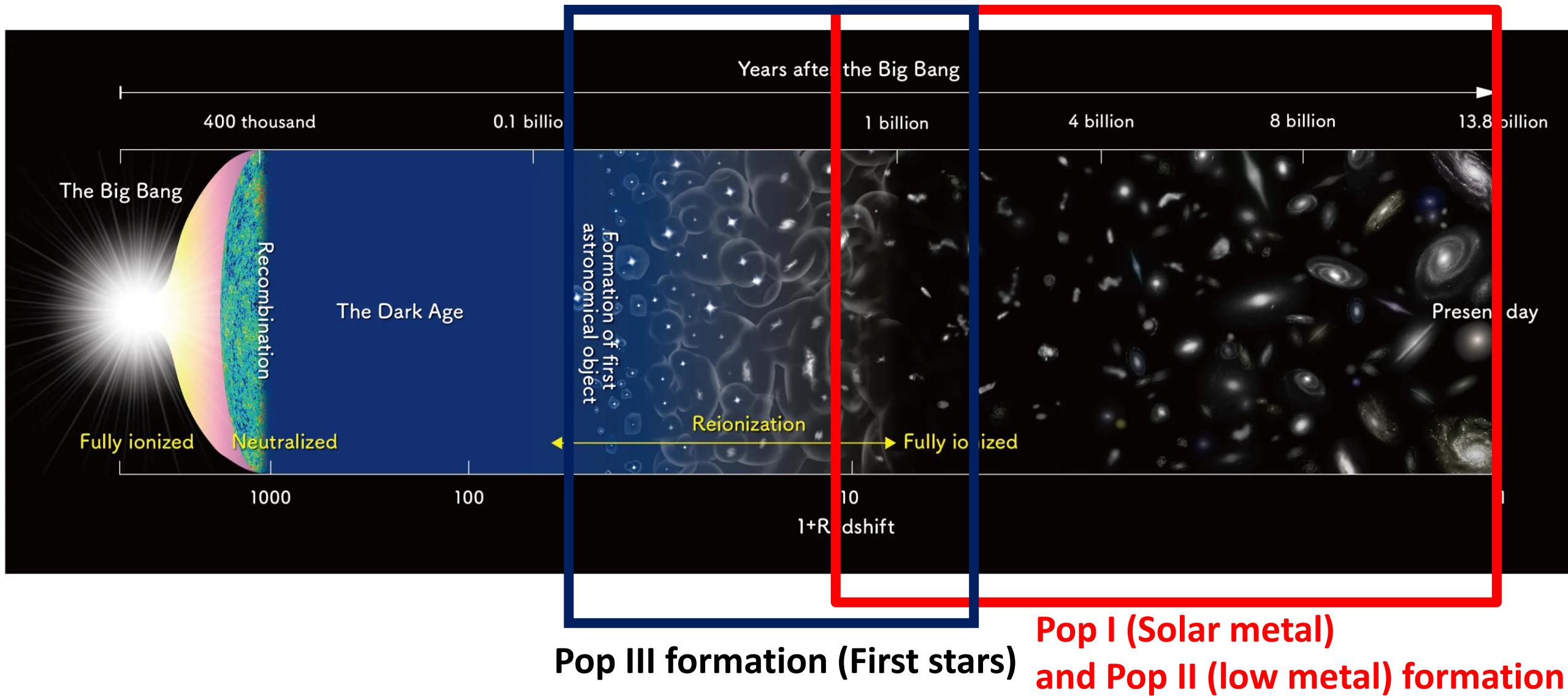


Sana et al. 2012

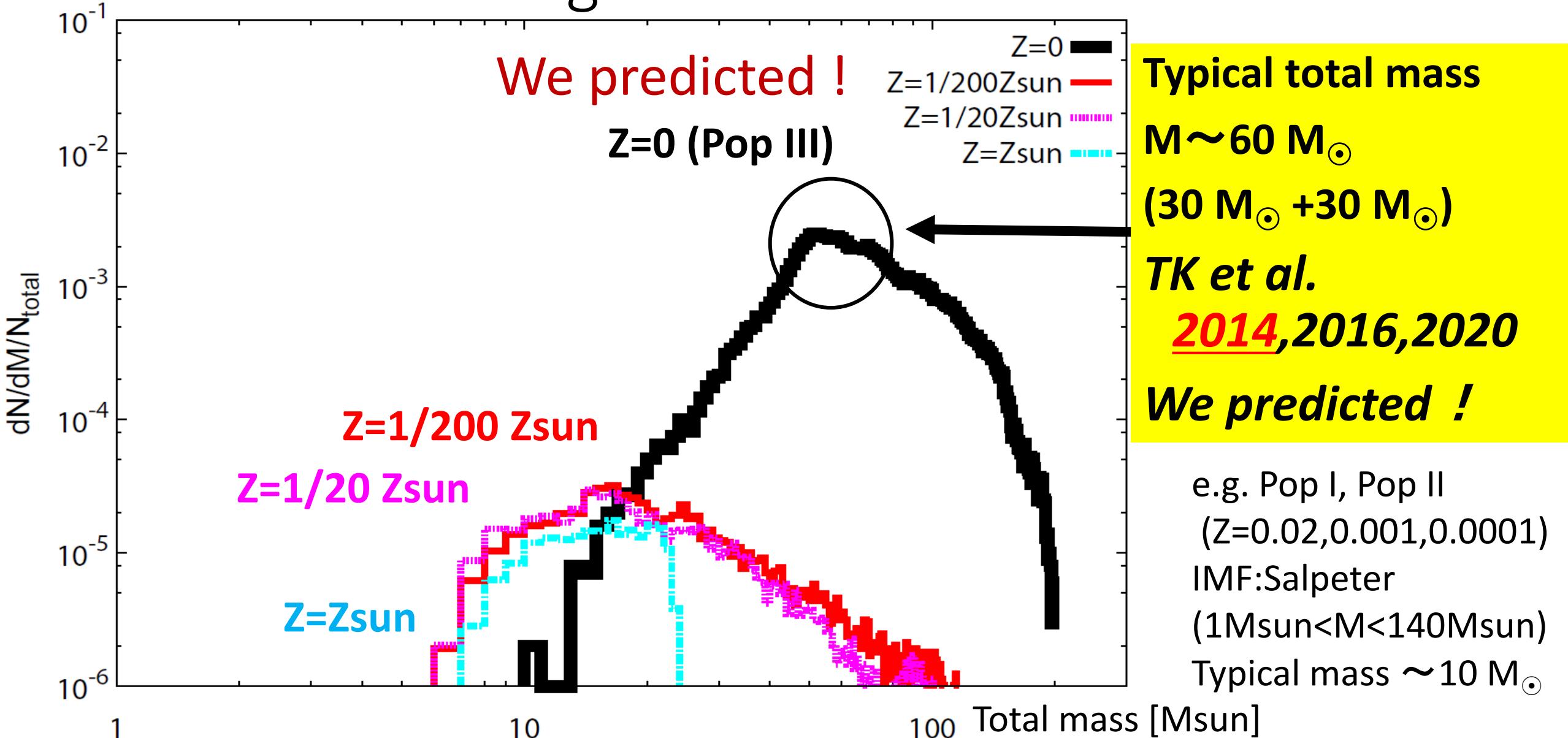


©star wars

# Isolated Binary scenarios

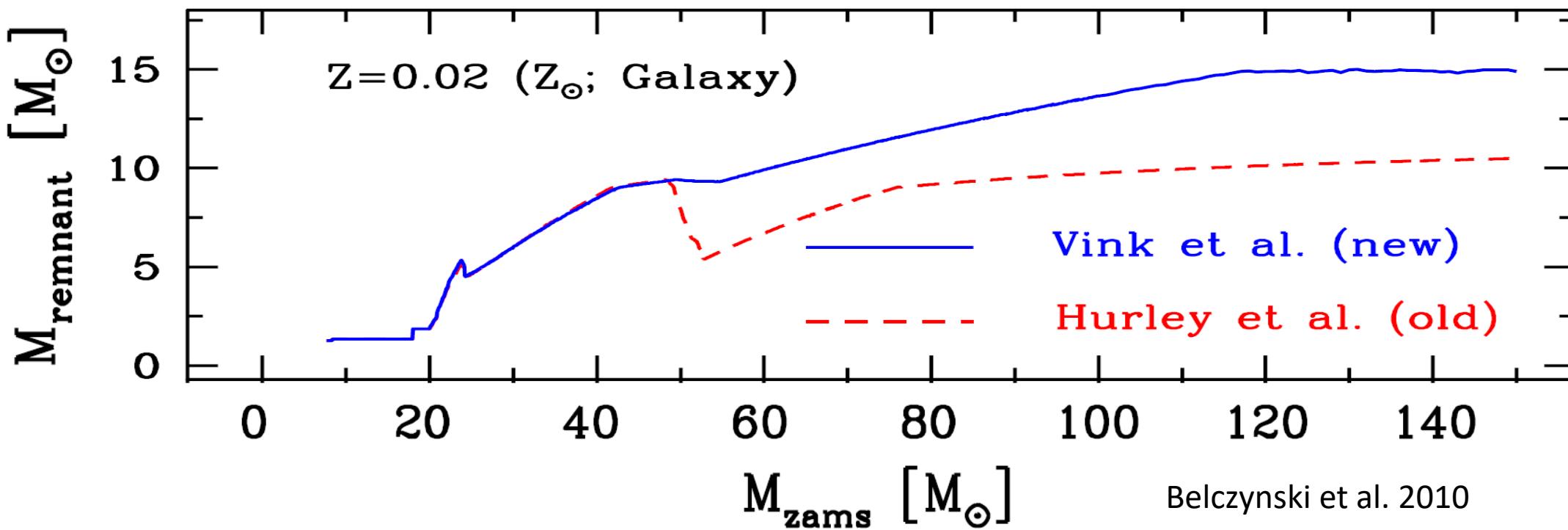


# Formation fraction of BBH which merge within the Hubble time



# Wind mass loss & IMF

- If the progenitor of BH is Pop I (=Solar metal stars)
- Typical mass is small ( $\text{IMF} \propto M^{-2.35}$ ,  $0.1M_{\text{sun}} < M < 100M_{\text{sun}}$ )
- Stars lose a lot of mass due to the strong stellar wind



# Wind mass loss & IMF

If the progenitor is low metal,

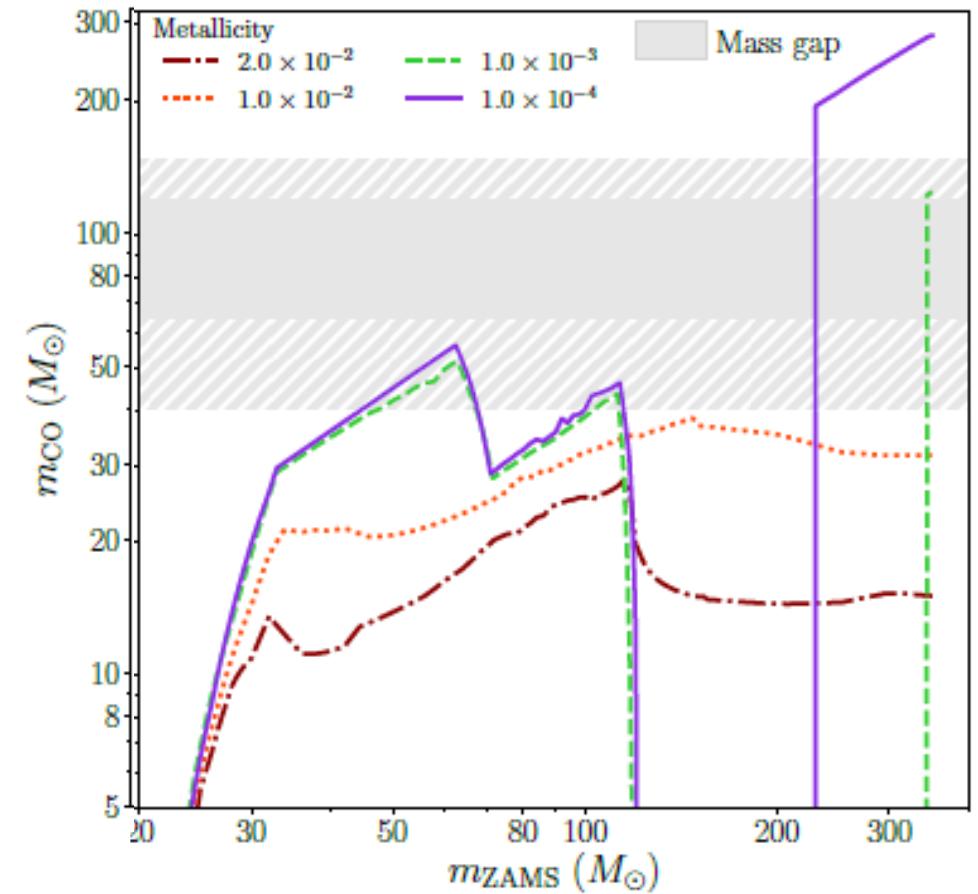
- Pop II (Metal<0.1SolarMetal)  
Typical mass is same as Pop I  
But, weak wind mass loss
- Pop III (No metal)

Pop III stars are ***the first stars*** after the Big Bang.

Typical mass is more massive than Pop I, II

$M_{\text{PopIII}} \sim 10-100 M_{\odot}$

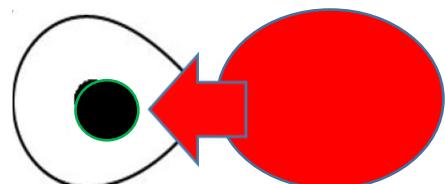
No wind mass loss due to no metal.



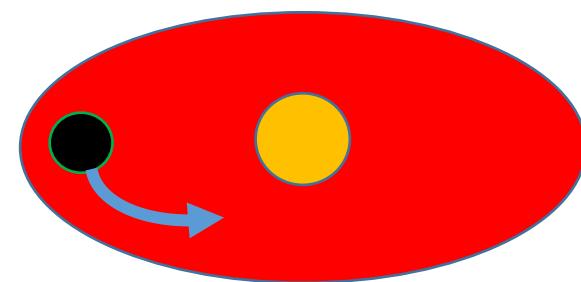
Abbot et al. 2020

# Binary interaction changes progenitor mass

- Mass transfer
- Common envelope



Mass transfer

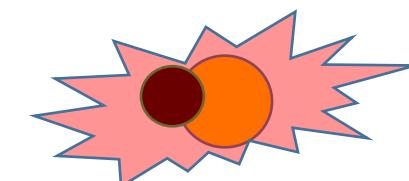


Common envelope

Red Giants tend to become CE



Close binary



or merge

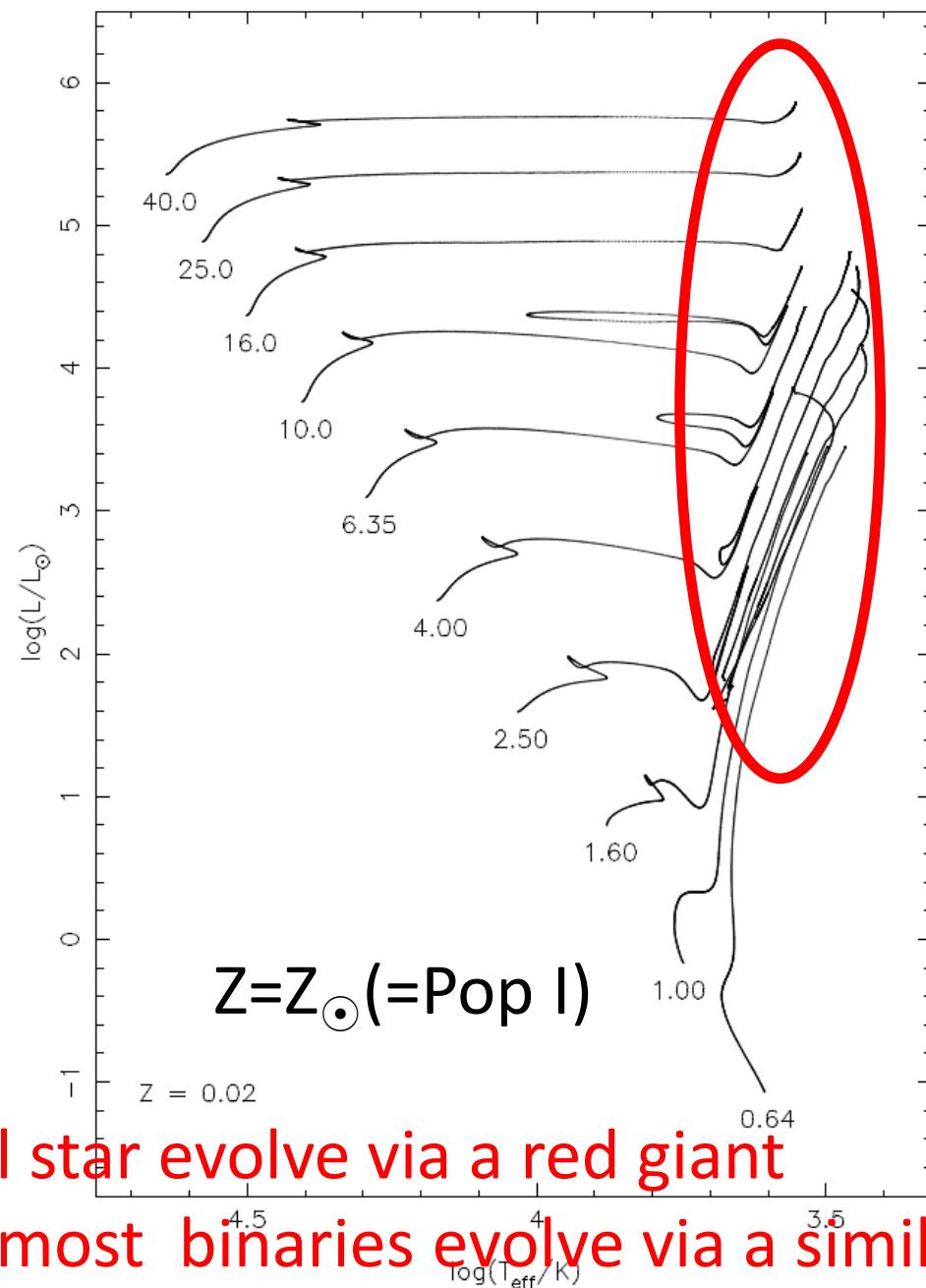


Figure 1. Selected OVS evolution tracks for  $Z = 0.02$ , for masses  $0.64, 1.0, 1.6, 2.5, 4.0, 6.35, 10, 16, 25$  and  $40 M_{\odot}$ .

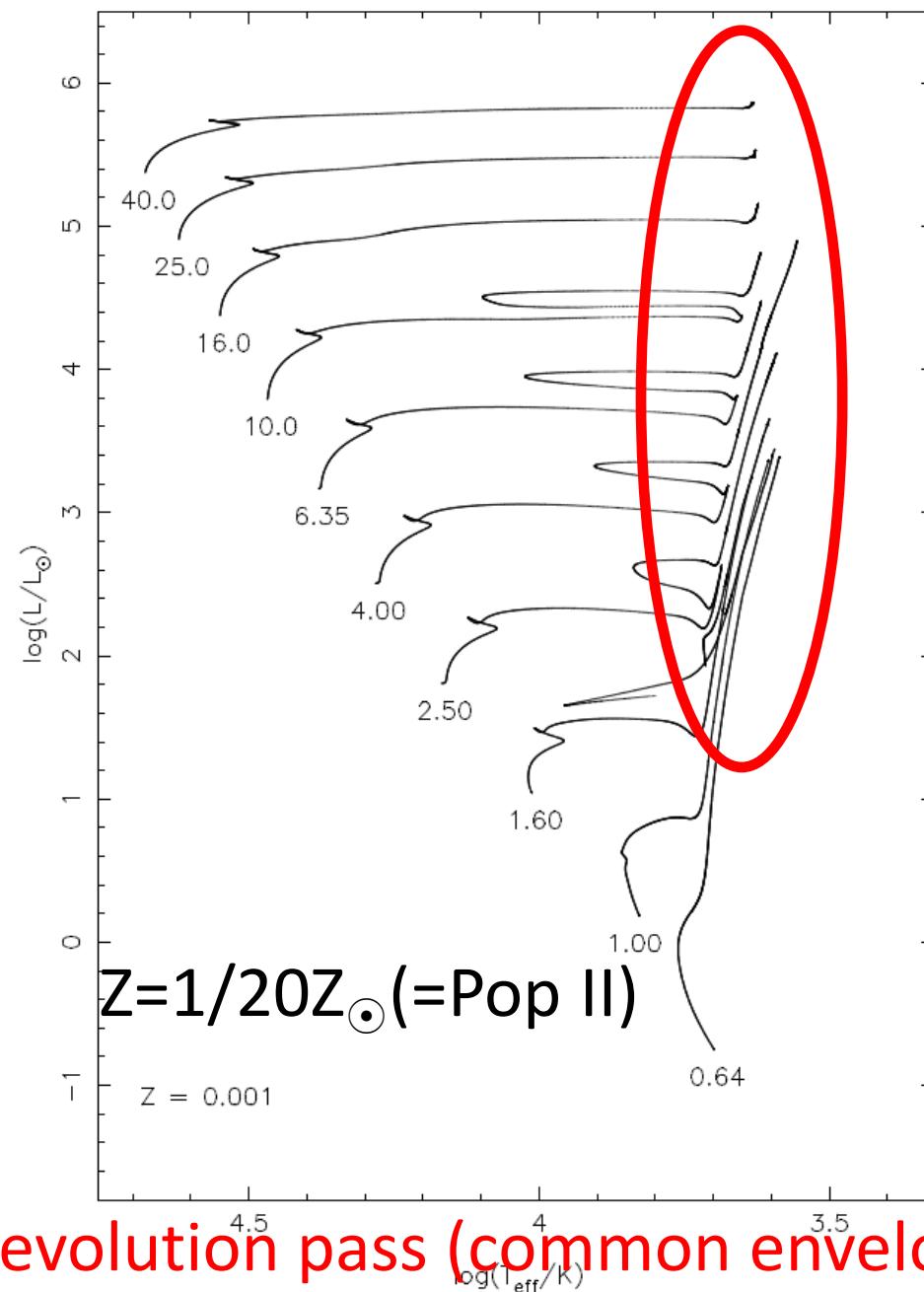
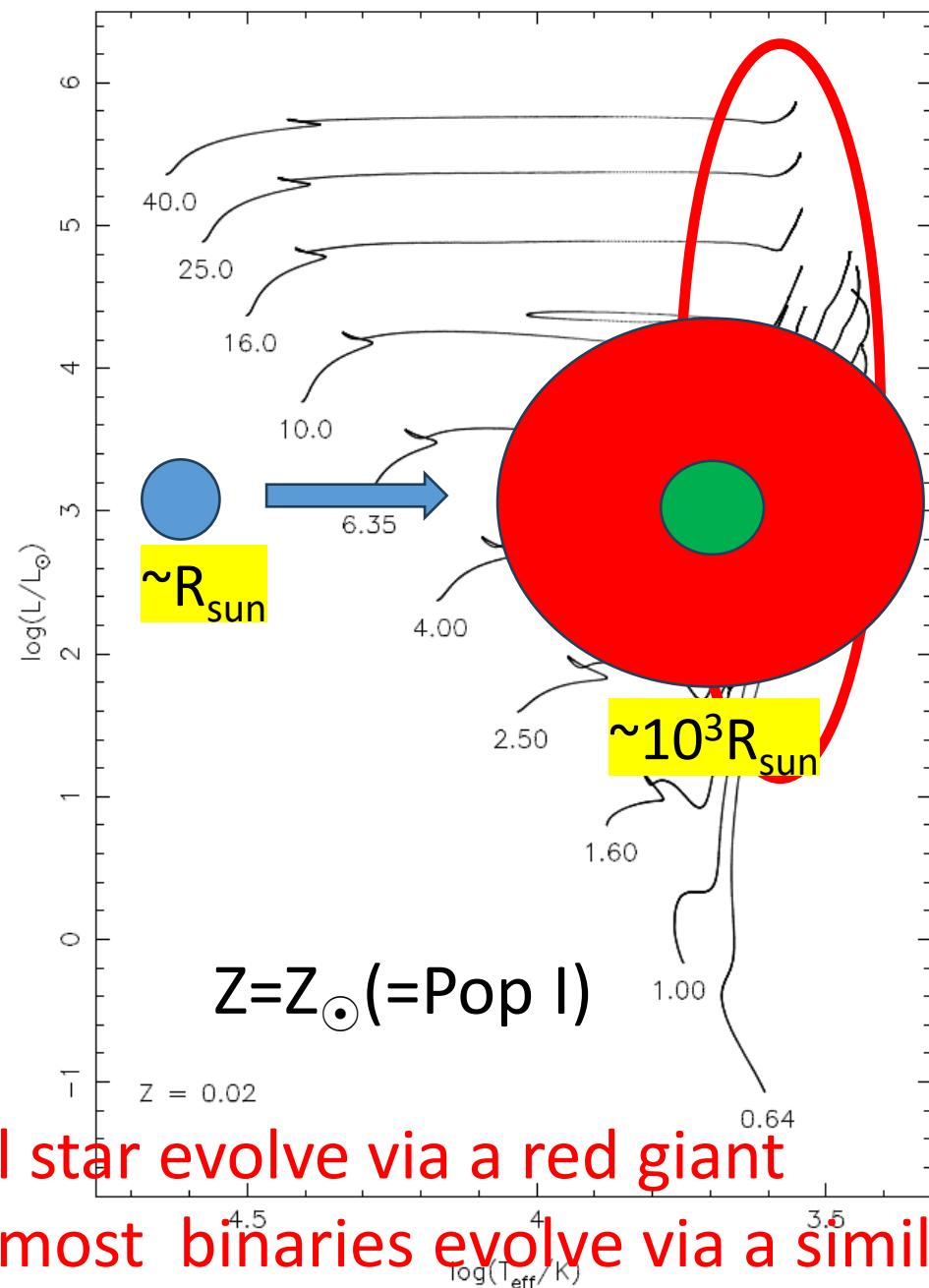


Figure 2. Same as Fig. 1 for  $Z = 0.001$ . The  $1.0 M_{\odot}$  post He flash track has been omitted for clarity.



All star evolve via a red giant

Almost binaries evolve via a similar evolution pass (common envelope)

Figure 1. Selected OVS evolution tracks for  $Z = 0.02$ , for masses  $0.64, 1.0, 1.6, 2.5, 4.0, 6.35, 10, 16, 25$  and  $40 M_\odot$ .

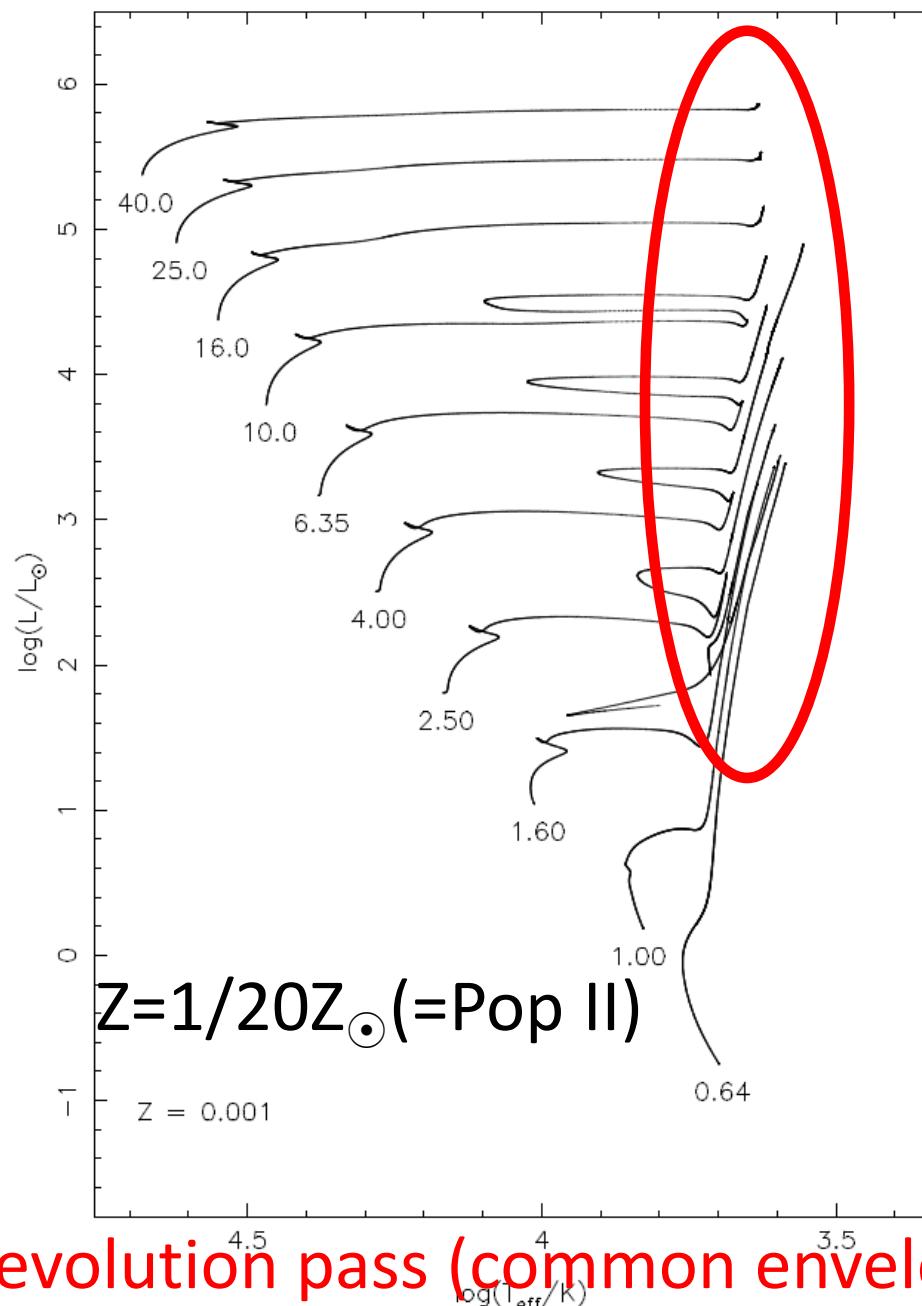
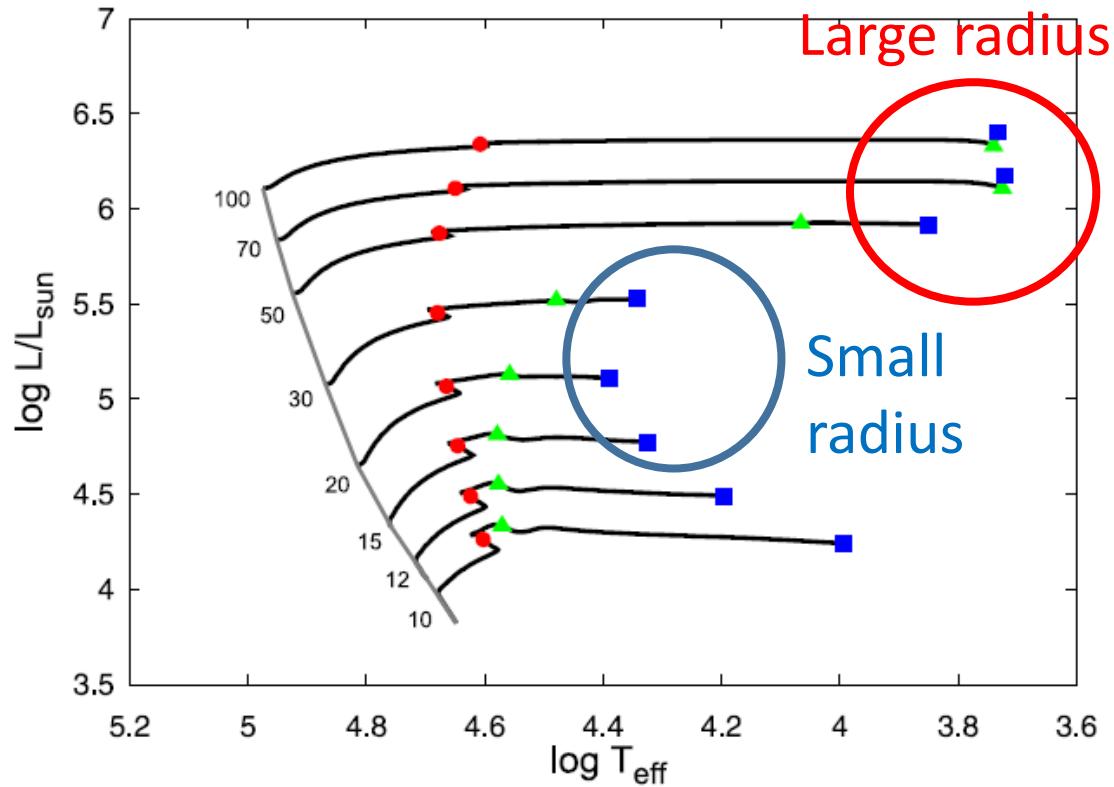


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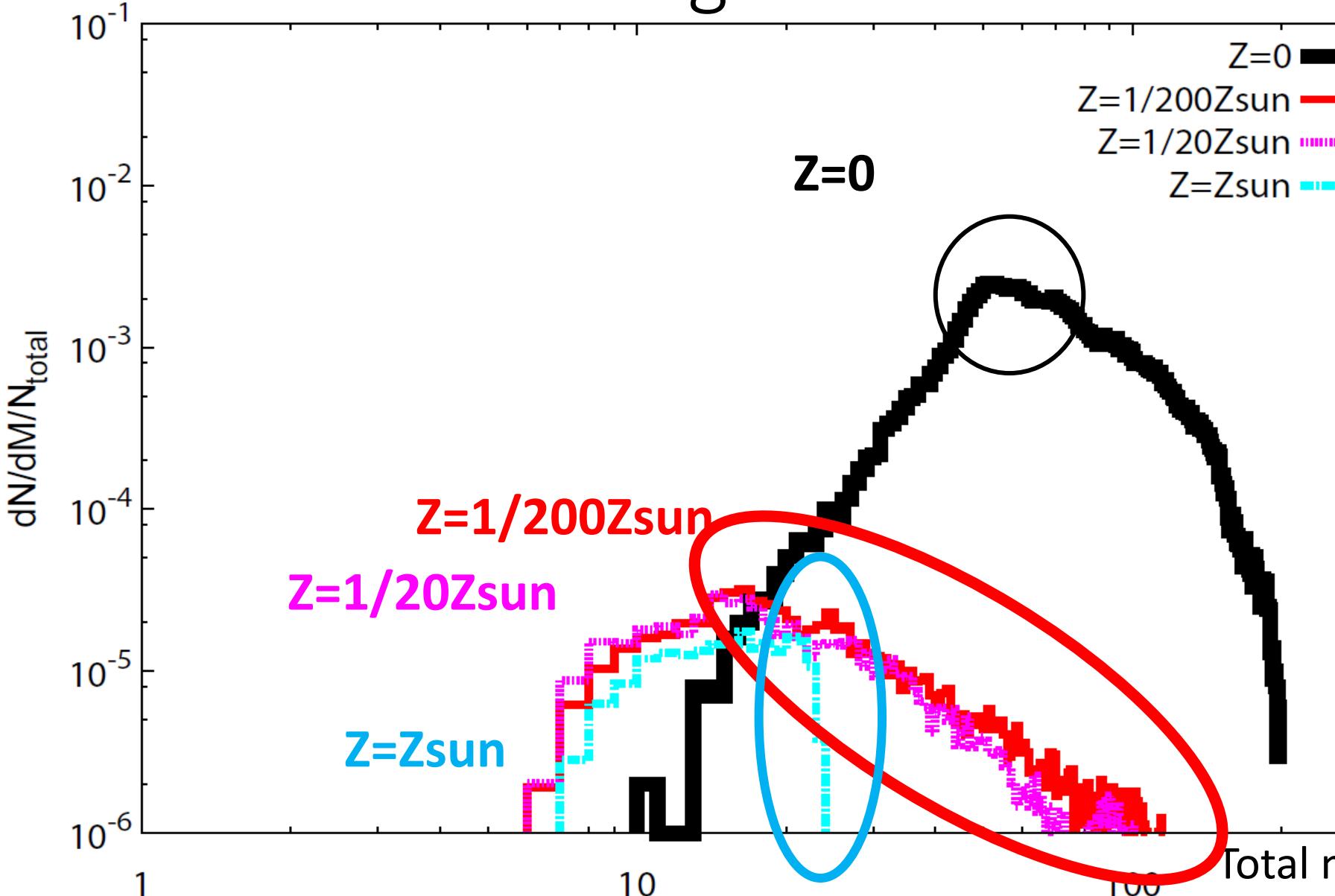
# Why Pop III binaries become 30Msun BH-BH



Marigo et al. 2001

- $M > 50\text{Msun}$  **red giant**
  - Mass transfer tend to be unstable
  - common envelope
  - $1/3 \sim 1/2$  of initial mass ( $\sim 30\text{Msun}$ )
- $M < 50\text{Msun}$  **blue giant**
  - Mass transfer tend to be stable
  - mass loss is not so effective
  - $2/3 \sim 1$  of initial mass ( $30\text{Msun}$ )

# Formation fraction of BBH which merge within the Hubble time

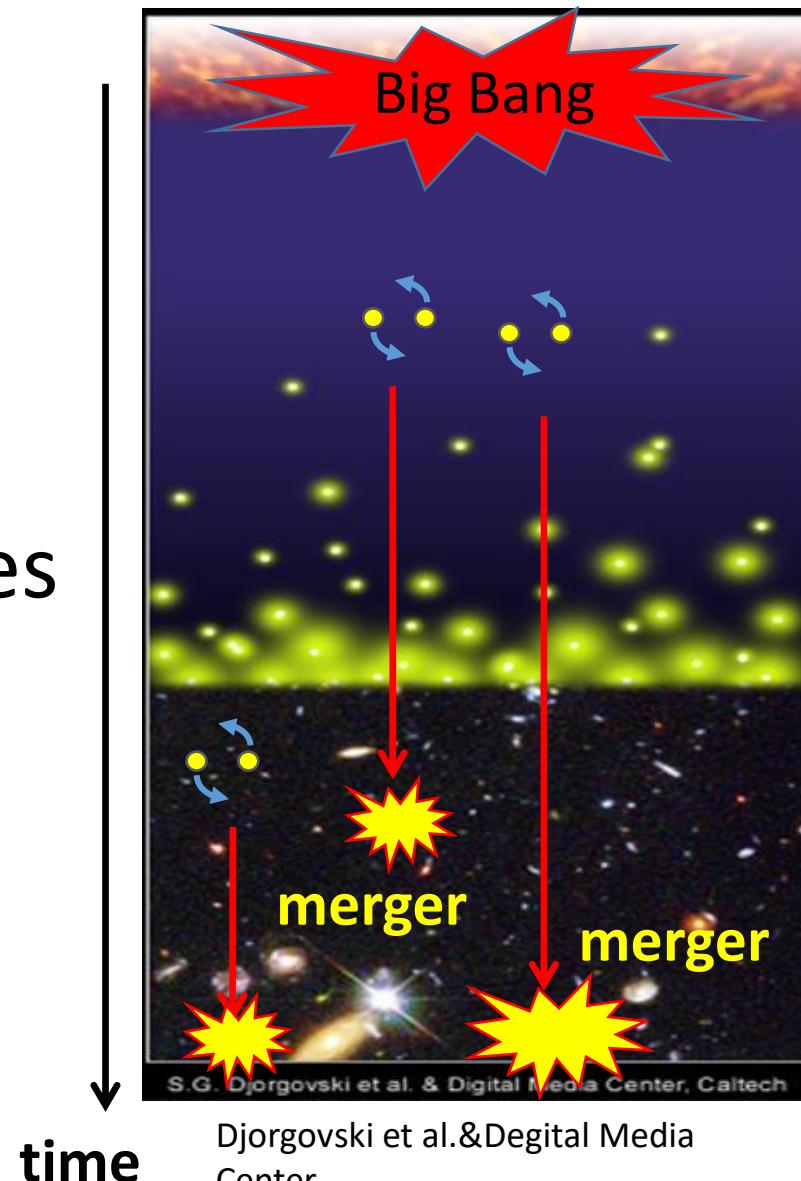


This shape reflects  
the influence of  
Pop III stellar  
evolution

These shapes have  
the influence of IMF  
and the influence of  
stellar wind mass loss

# Pop III BBH remnants for gravitational wave

- Pop III stars were born and died at  $z \gtrsim 10$ .
- The typical merger time of compact binaries  $\sim 10^{8-10}$  yr  $dN/dt \propto t^{-1}$  (Kinugawa et al .2014,2020 Inayoshi et al. 2017)
- We can see Pop III BBH at the present day!

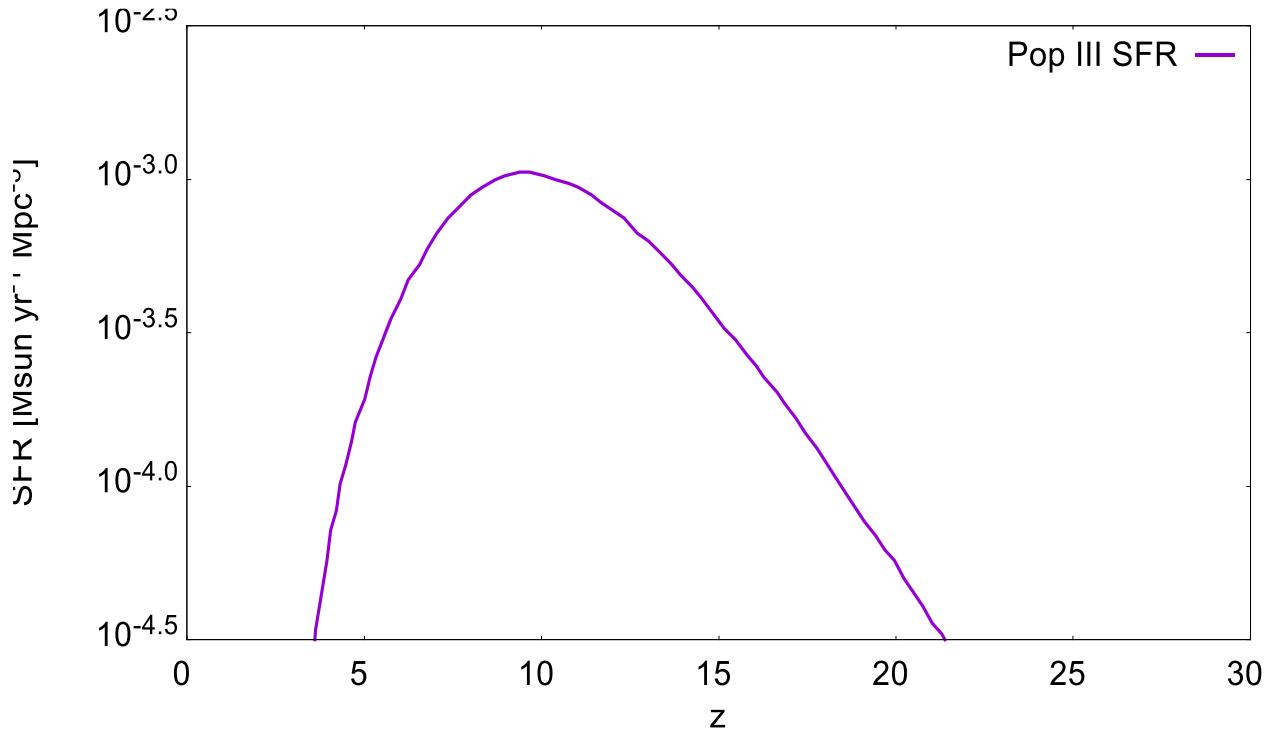


# The star formation rate of Pop III

In order to calculate merger rate,  
we need to know

- When were Pop III stars born?
- How many were Pop III stars born?

⇒ Star formation rate



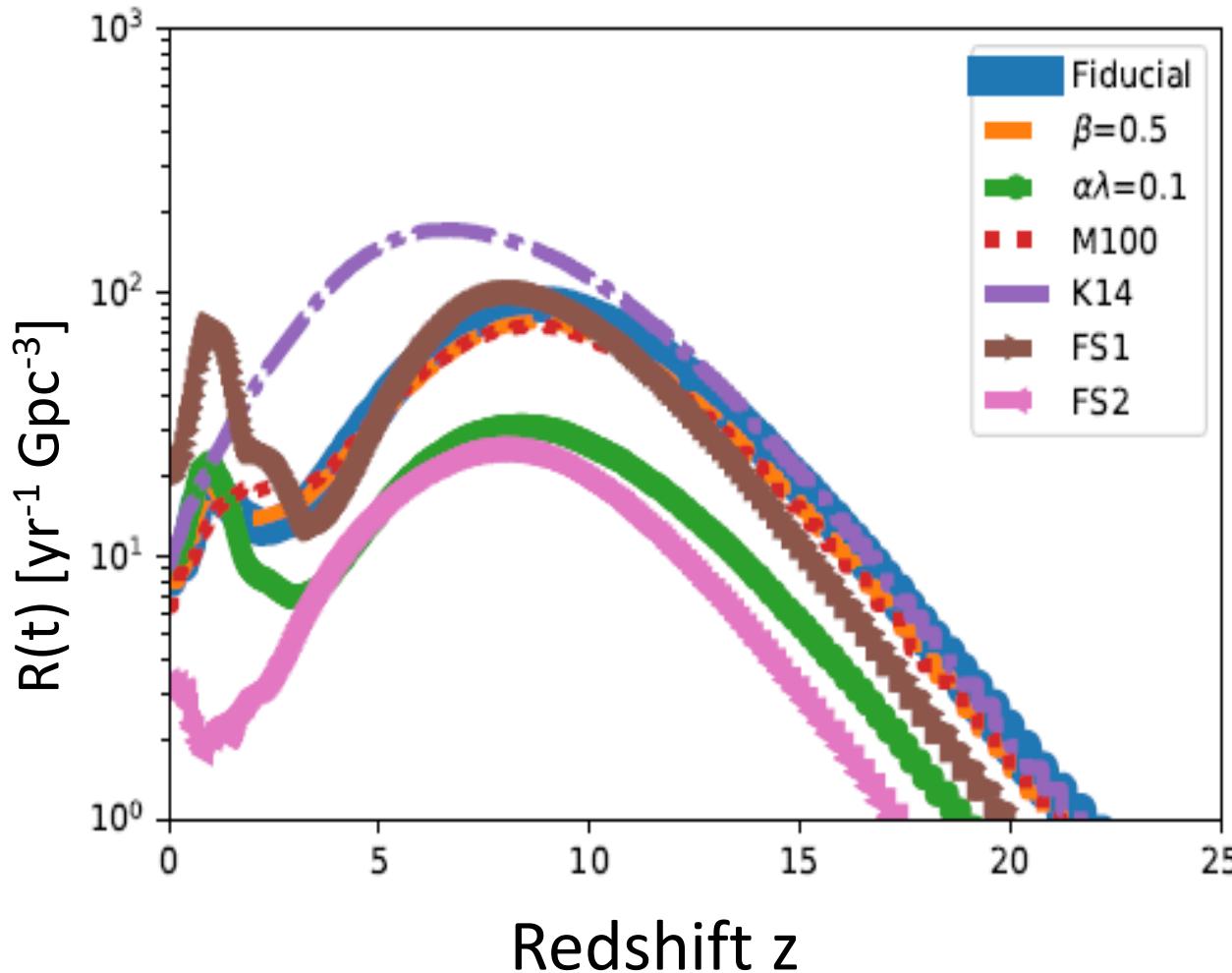
We adopt the total Pop III stellar mass density by Inayoshi et al. 2016

$$\rho = 6 \times 10^5 \text{ Msun/Mpc}^3$$

(SFR peak at  $z=10$ )

We assume the binary fraction  $f_b=0.5$

# The Pop III BH-BH merger rate density

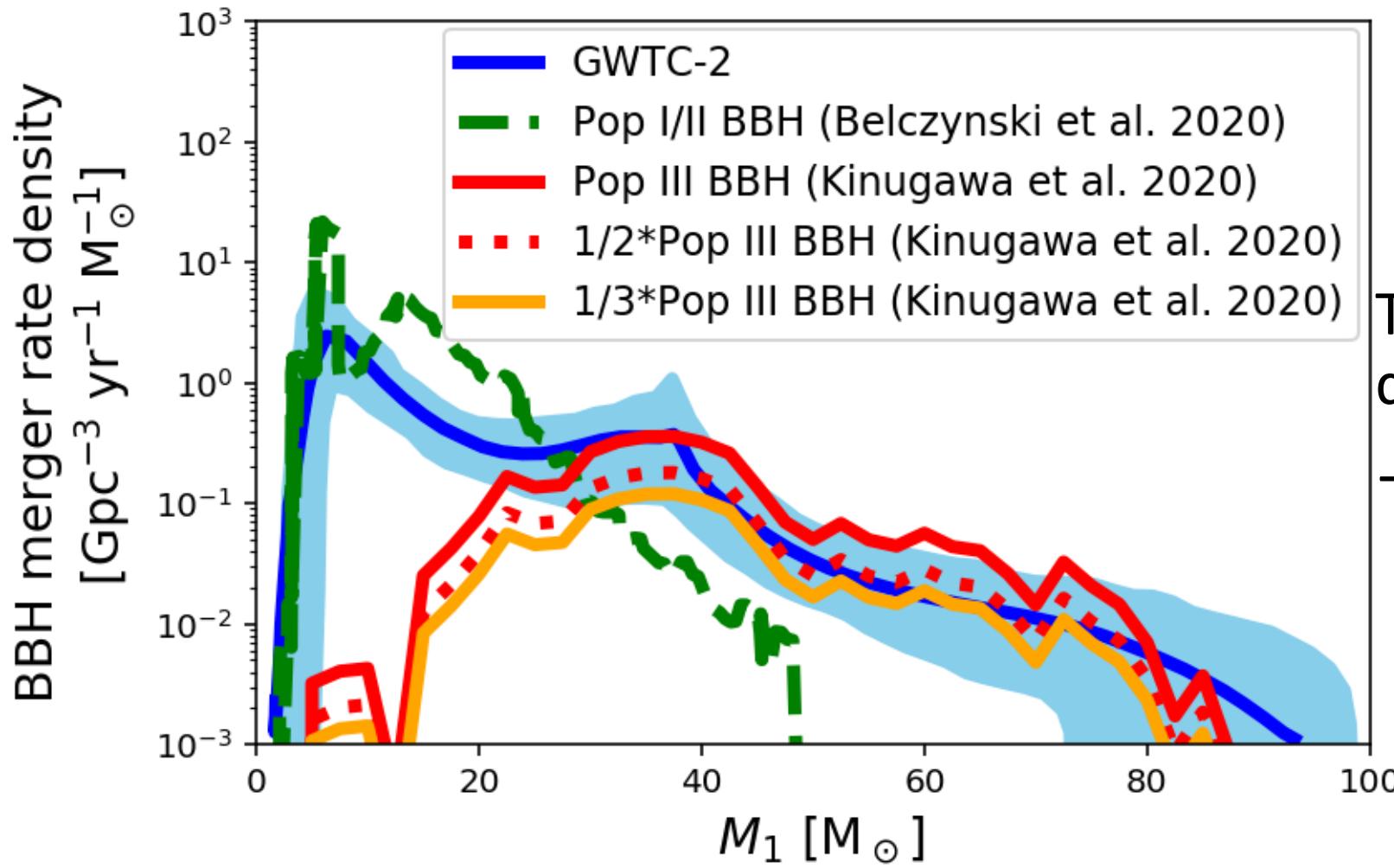


Pop III BH-BH merger rate at  $z=0$   
In our fiducial model

$$R \sim 10 \left( \frac{\rho_{\text{Pop III}}}{6 \times 10^5 M_\odot / \text{Mpc}^3} \right) \left( \frac{f_b / (1+f_b)}{0.33} \right) [\text{yr}^{-1} \text{Gpc}^{-3}]$$

(Kinugawa et al. 2014, 2016, 2020)

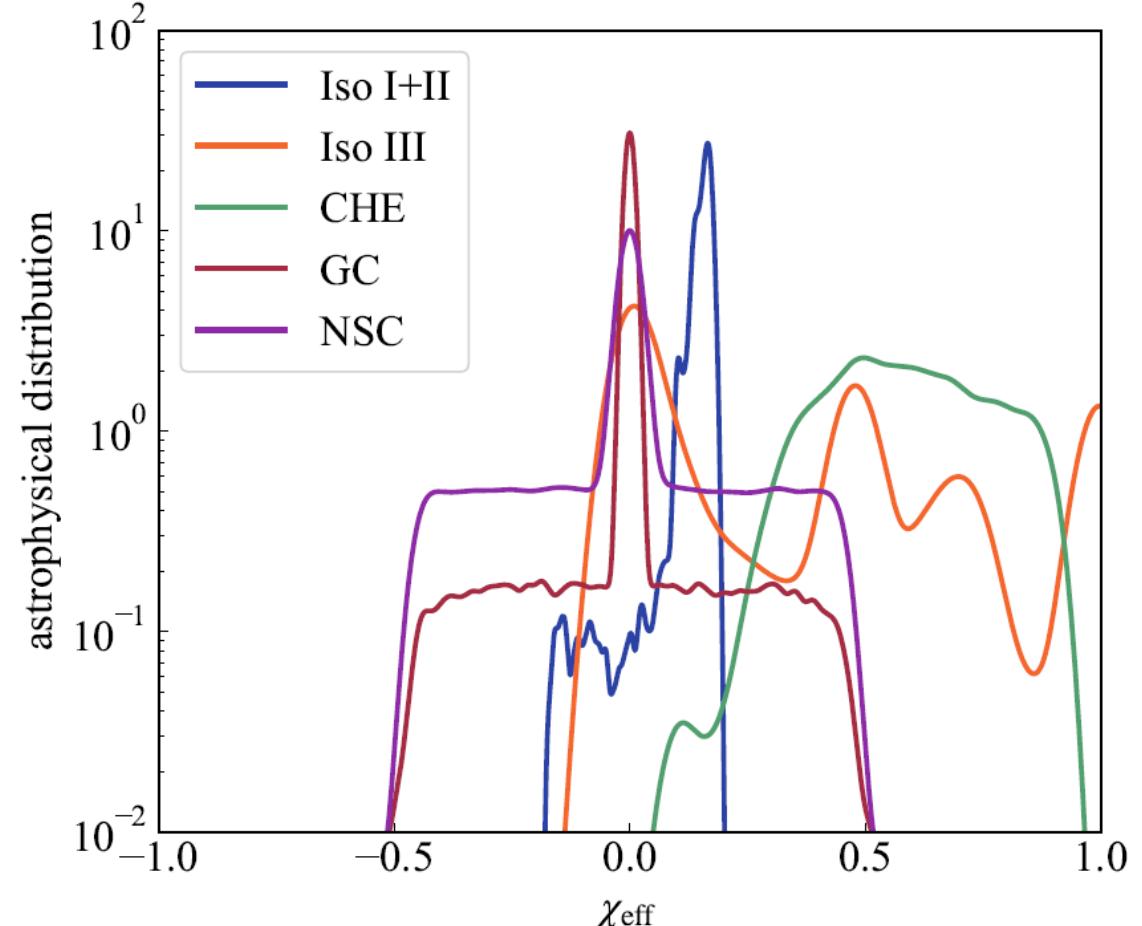
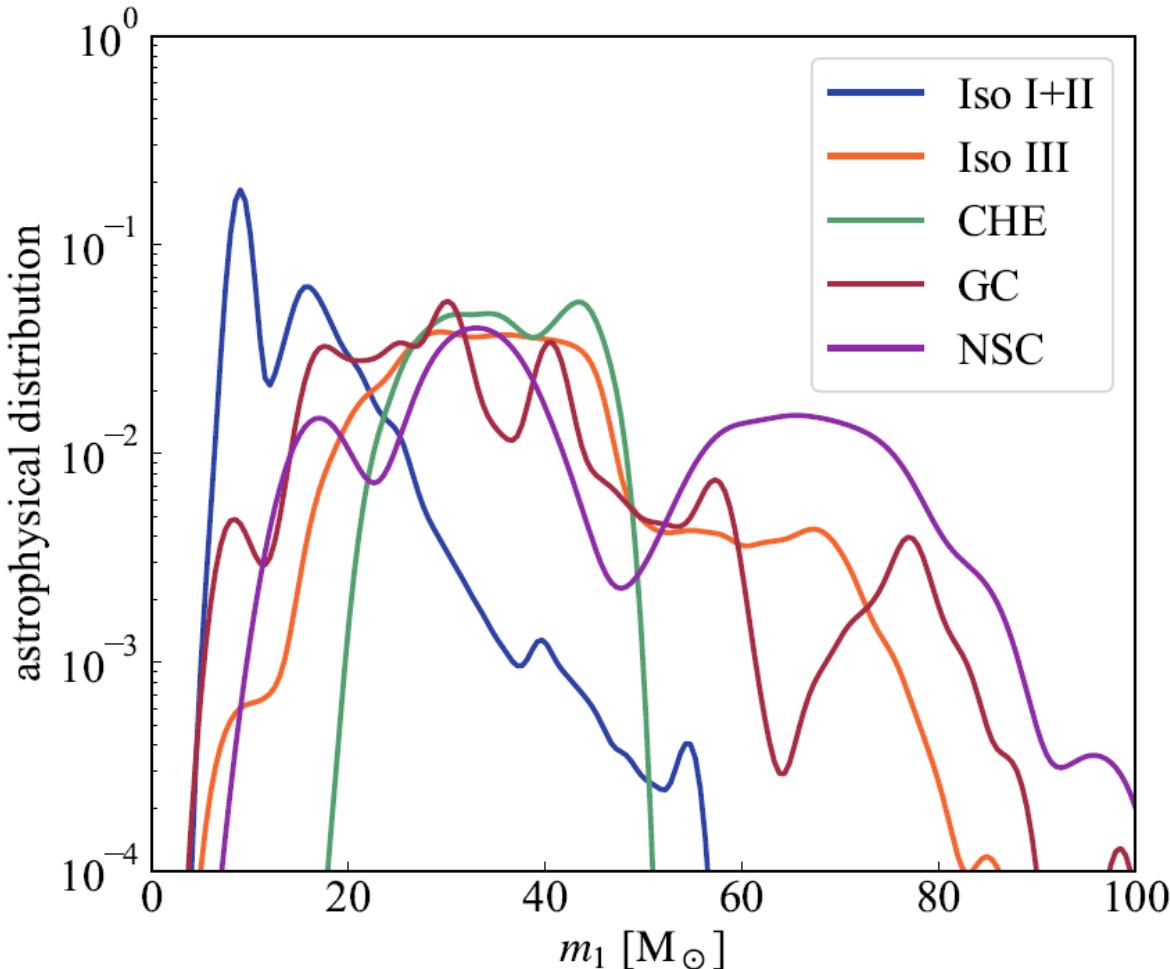
# Comparison with mass distributions of observed BBHs



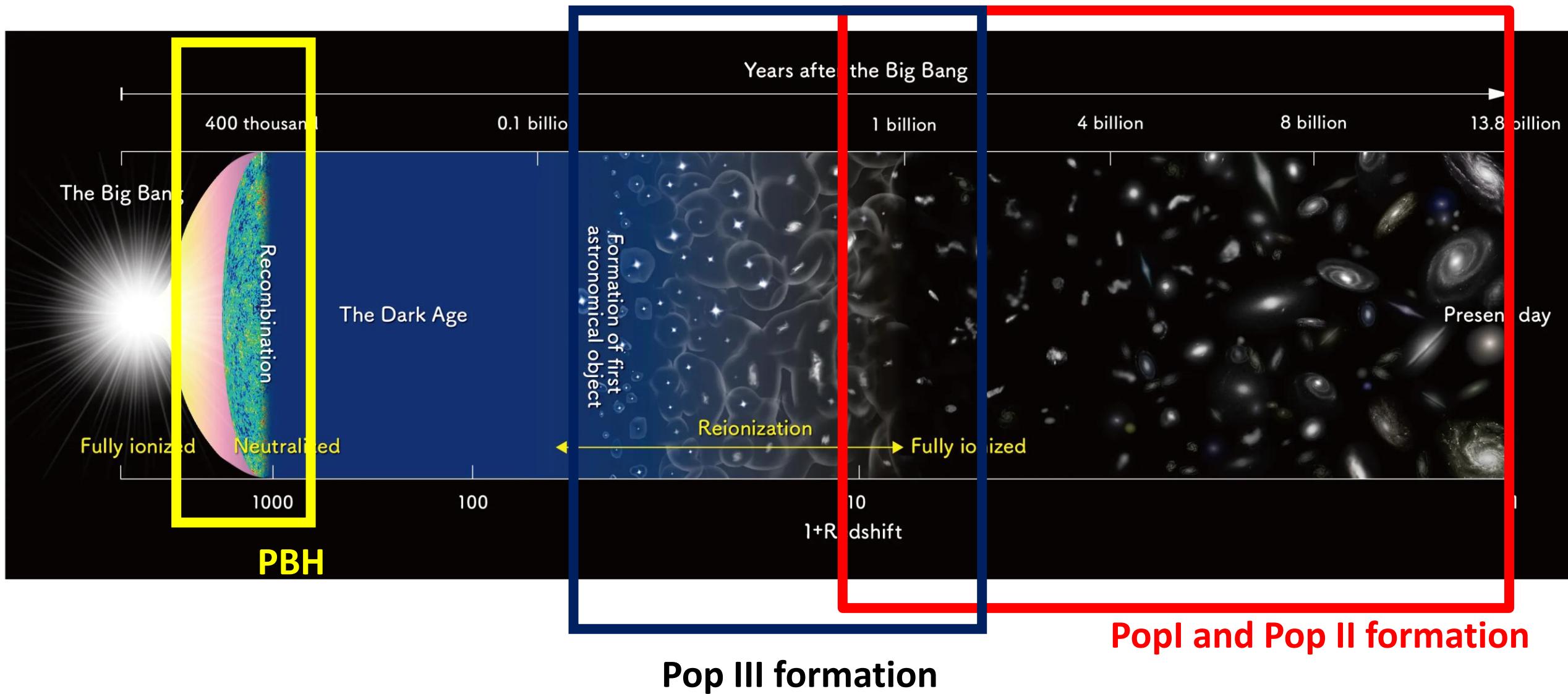
(Kinugawa, Nakamura& Nakano 2021)

The mass distribution might  
distinguish Pop III from Pop I/II  
→ The evidence of Pop III ?  
Not yet

# Best combination model for GWTC-3 (Iwaya, TK, Tagoshi in prep.)



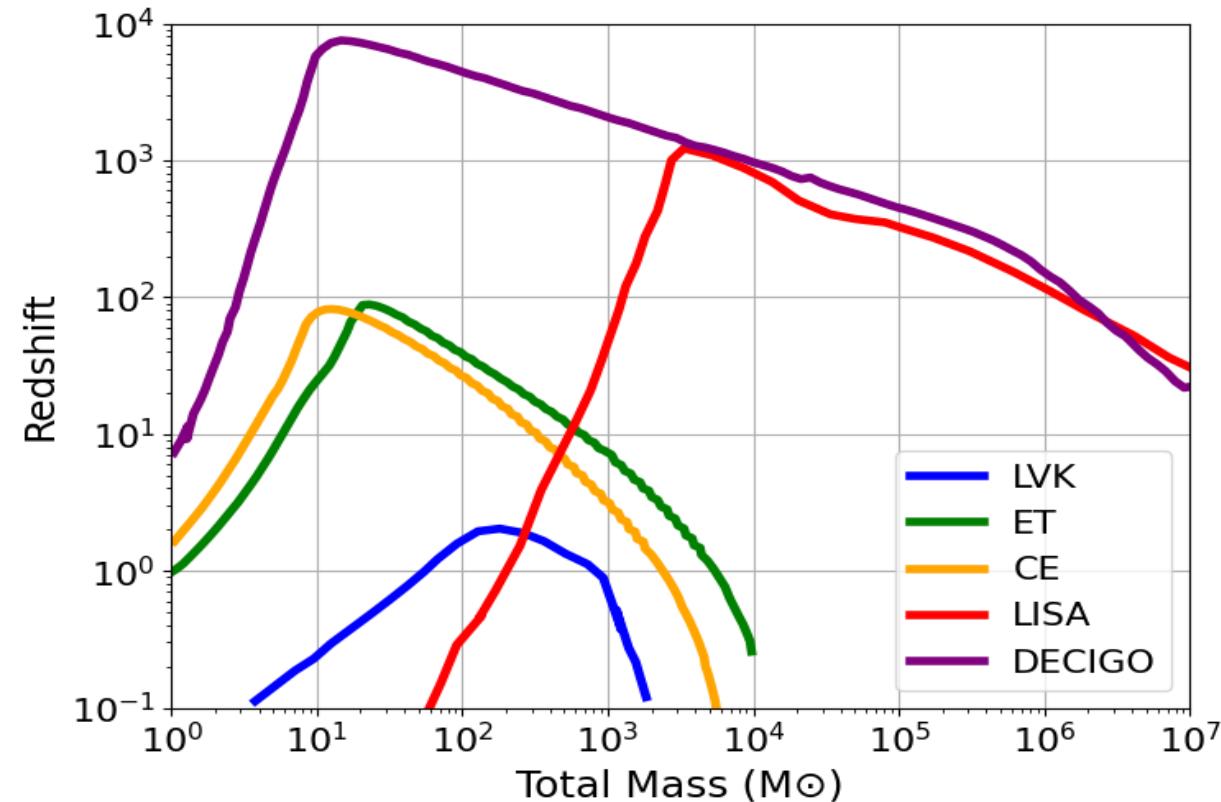
# What is the smoking gun of the origin of BBHs?



# Future plan of GW observer : ET, CE, B-DECIGO and DECIGO

- Einstein telescope (ET): the next generation GW observatory of Europe
- Cosmic explorer (CE) : the next generation GW observatory of US.
- DECIGO: Japanese space gravitational wave observatory project

We can see Pop III BH-BHs  
when Pop III stars were born ( $z>10$ )!  
(Nakamura, Ando, Kinugawa et al. 2016)



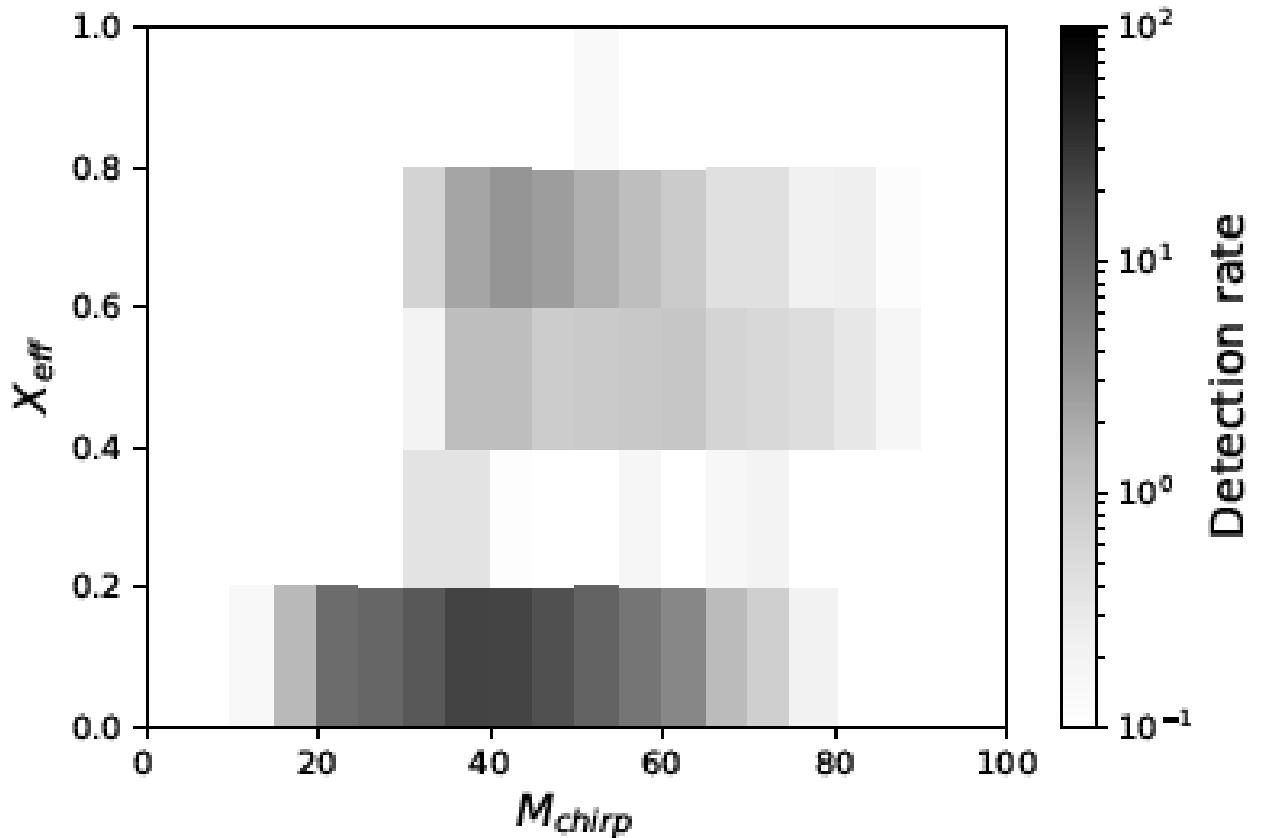
# Merger time dependence of Pop III BBH spin

	$a_1/M_1 < 0.1$ $a_2/M_2 < 0.1$	$a_1/M_1 < 0.1$ $a_2/M_2 > 0.9$	$a_1/M_1 > 0.9$ $a_2/M_2 < 0.1$	$a_1/M_1 > 0.9$ $a_2/M_2 > 0.9$
Merger time <1Gyr	25%	36%	0%	23%
Merger time >10Gyr	70%	0.3%	4%	0%

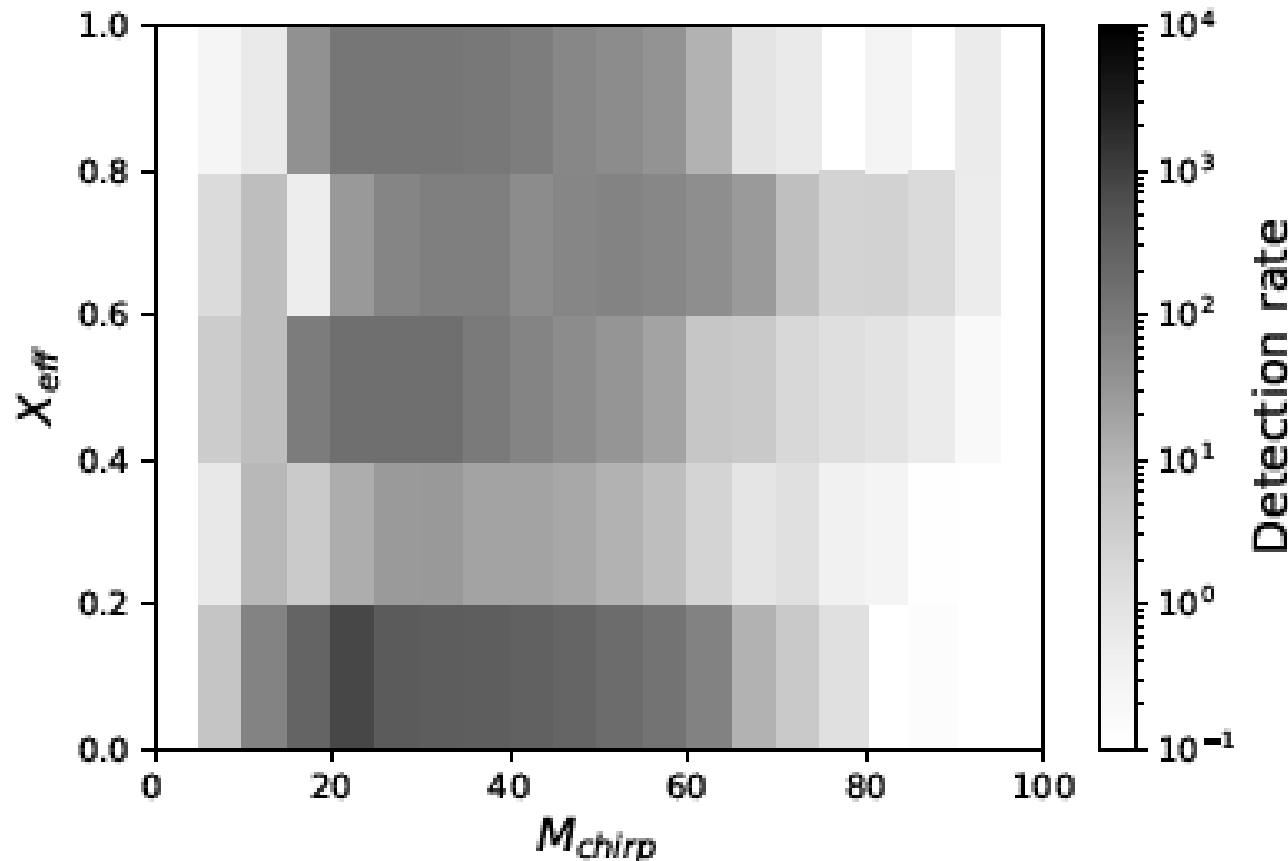
- If the origin of massive BBHs is Pop III,  
high spin BBHs are easier to be detected at high redshift

# Detection rate of Pop III BBH for ET and aLIGO design sensitivity

aLIGO



ET



(Kinugawa et al.2020)

# Summary

- There are many BBH formation theories.
- BBHs detected by LIGO/Virgo/KAGRA might be a mixture of different origins.
- Pop III binaries tend to become **30Msun+30Msun BH-BH**
- Pop III might explain the **GW190521** and **GW231123** like massive BBHs
- **Pop III BBH merger rate density at present day.**

$$R \sim 10 \left( \frac{\rho_{\text{Pop III}}}{6 \times 10^5 M_\odot / \text{Mpc}^3} \right) \left( \frac{f_b / (1 + f_b)}{0.33} \right) [\text{yr}^{-1} \text{ Gpc}^{-3}]$$

- The mass distribution or the redshift dependence might distinguish BH origins.

# Summary

- There are two populations of BBHs
- BBHs can have different mass distributions and origins
- Populations of BBHs are different
- Populations of BBHs are different
- Populations of BBHs are different

Rate density of BBHs

$$R \sim 10^{-3} \text{ yr}^{-1} \text{ Gpc}^{-3} \left( \frac{f_b/(1+f_b)}{0.33} \right) \left( \frac{\text{yr}^{-1} \text{ Gpc}^{-3}}{0.33} \right)$$

- The mass distribution or the redshift dependence might distinguish BH origins.

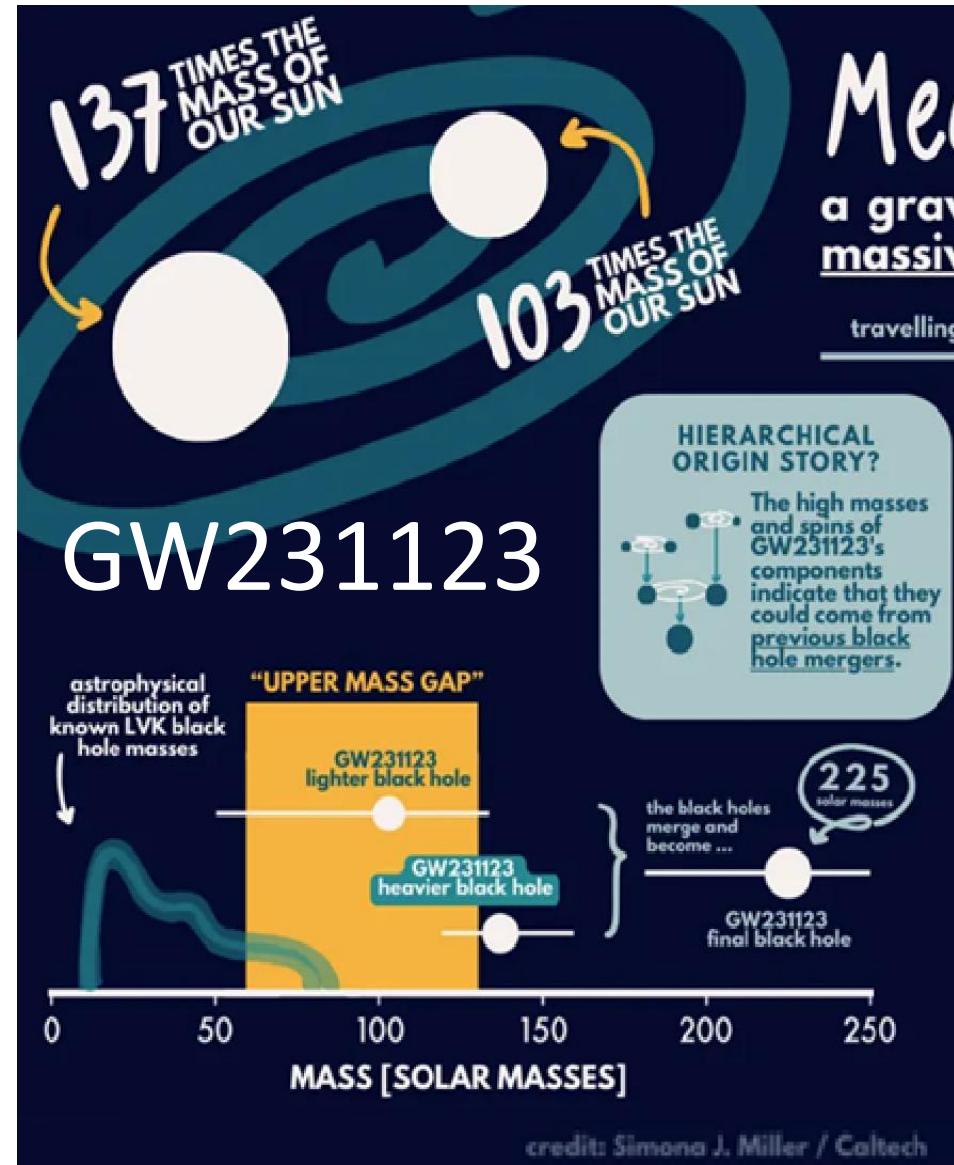
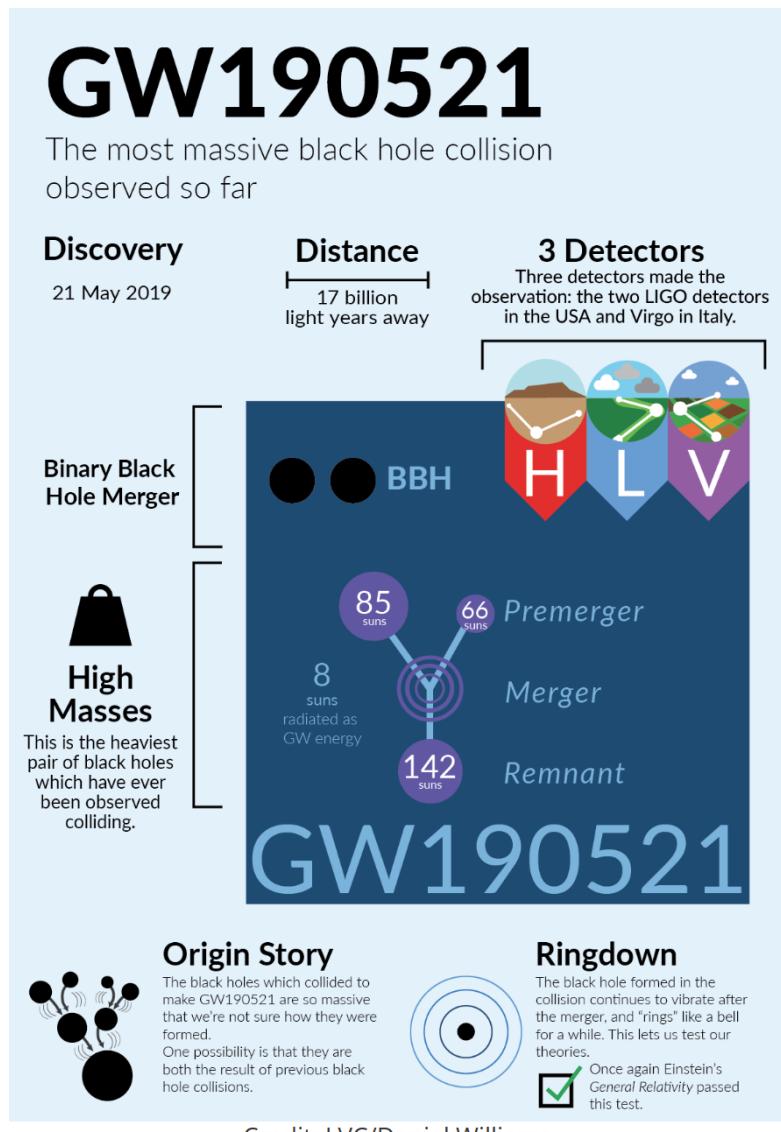


formation rate  
/yr/Gpc<sup>3</sup> might be a mixture of different  
origins. The BBHs  
become 2nd run BBHs.  
e GW190521 and GW231206 like massive BBHs  
**rate density at present day:**

**Massive BBHs = the fossil of ancient BH?**

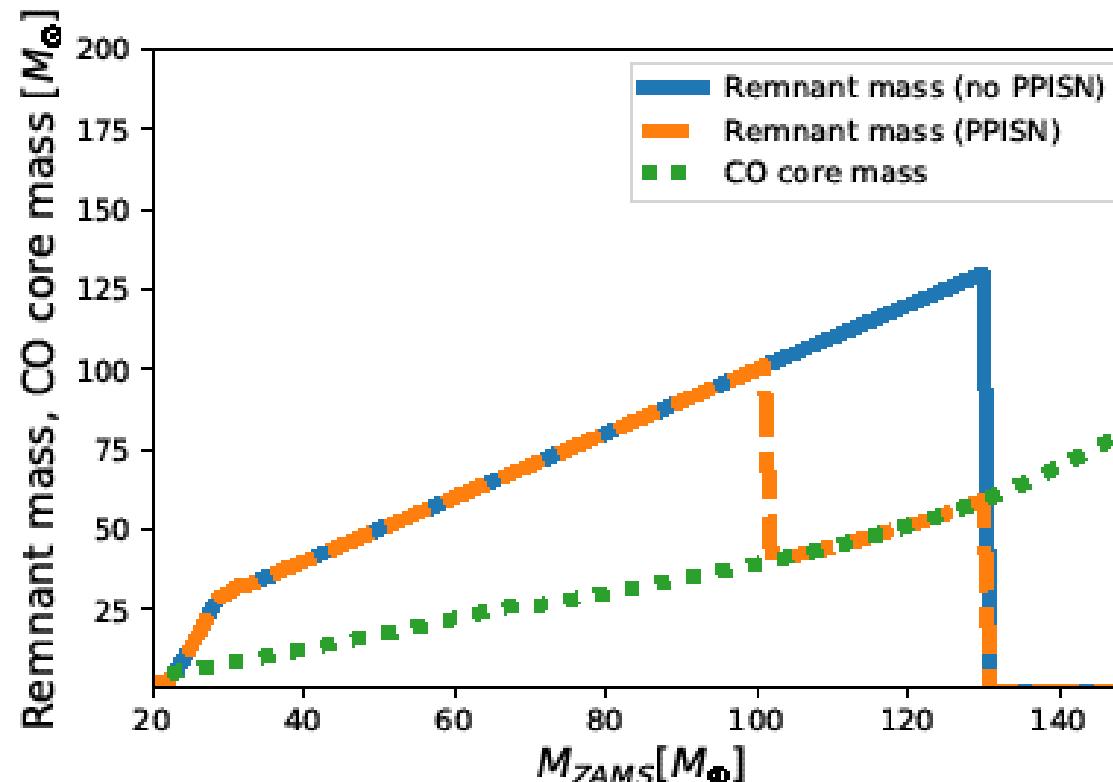


# Can Pop III explain “Mass gap” BBH mergers?



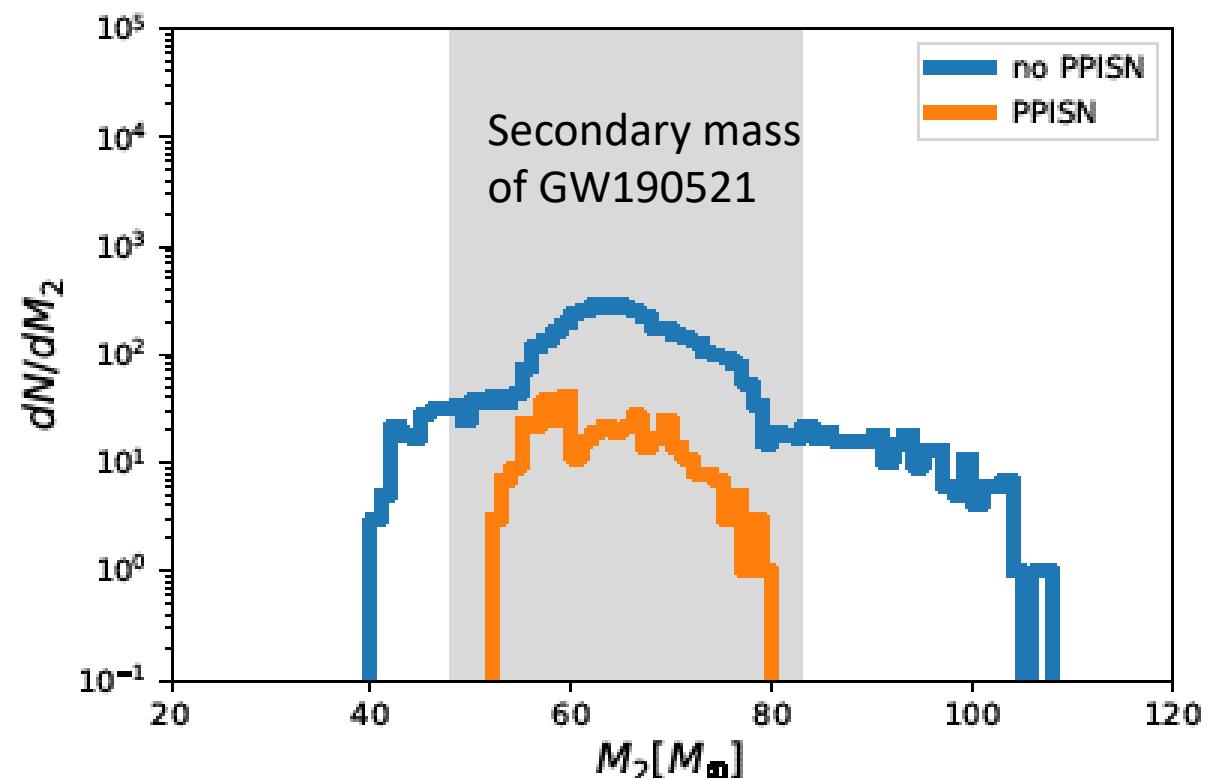
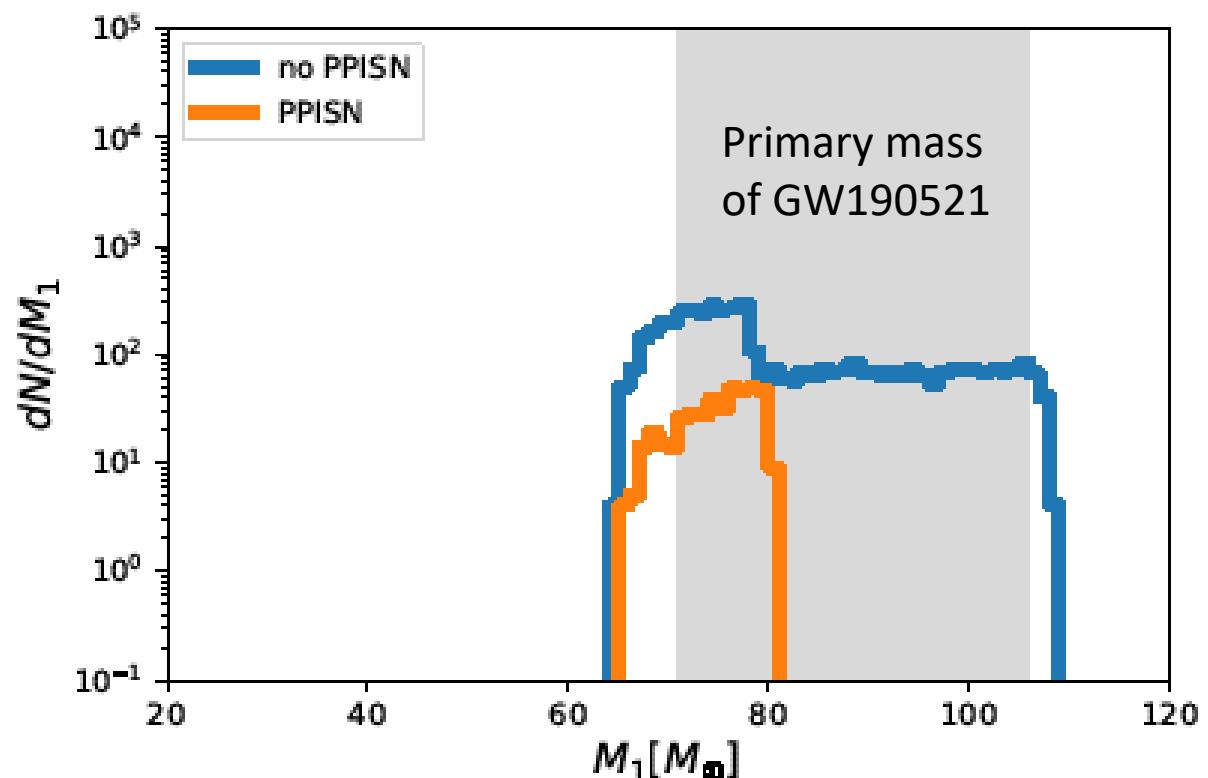
# GW190521

- Pop III can make GW190521 like BBH!  
(Kinugawa et al. 2020, Farrell et al. 2020, Tanikawa et al. 2020)



Kinugawa et al. 2021

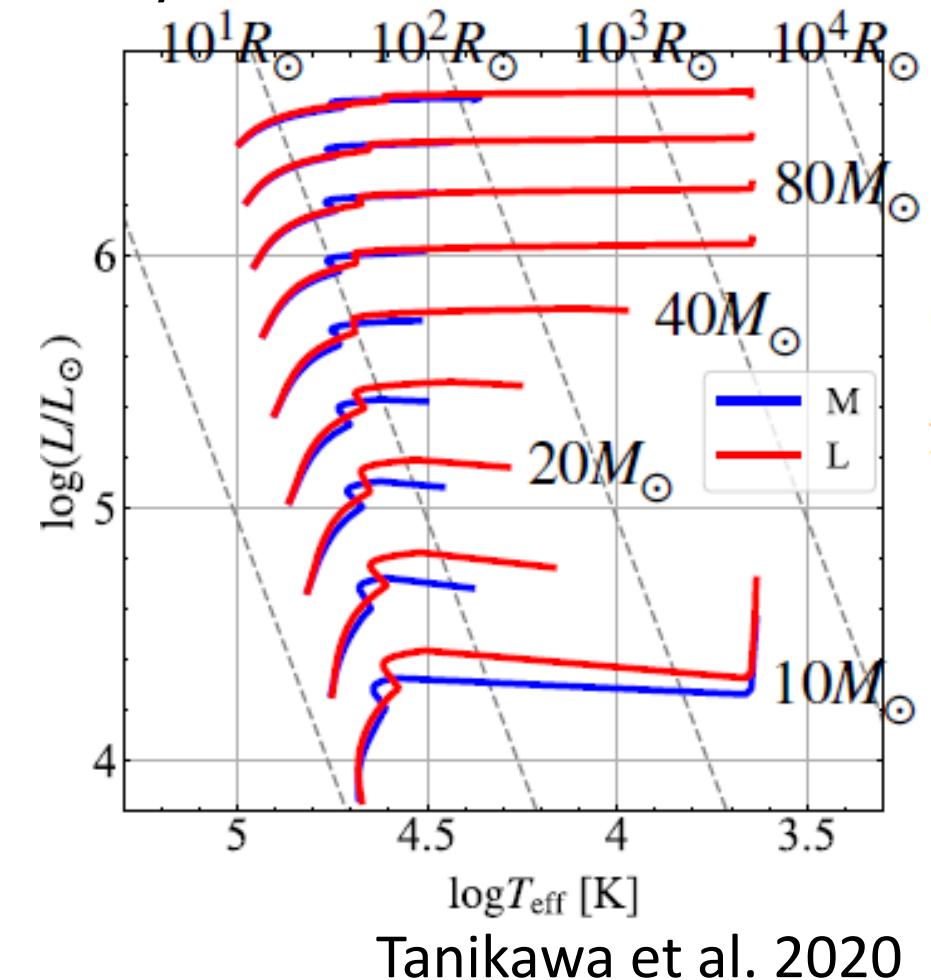
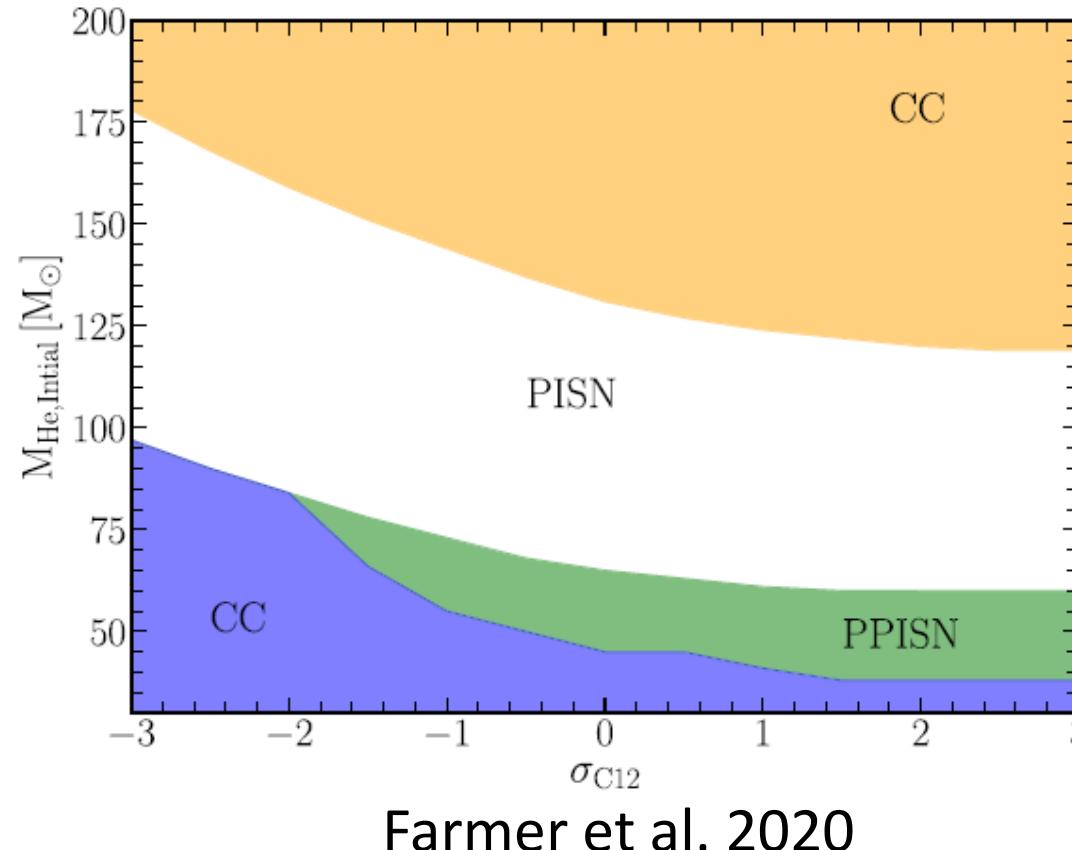
# GW190521



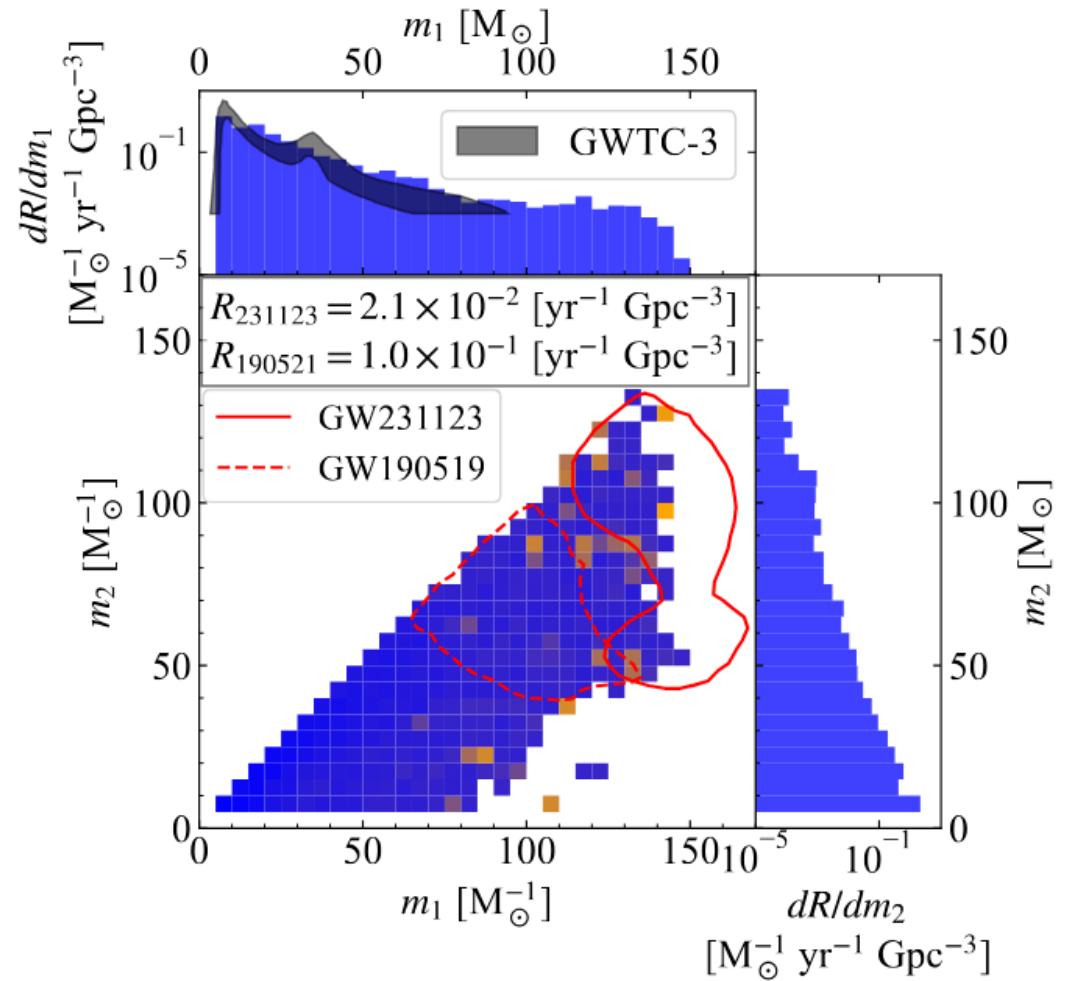
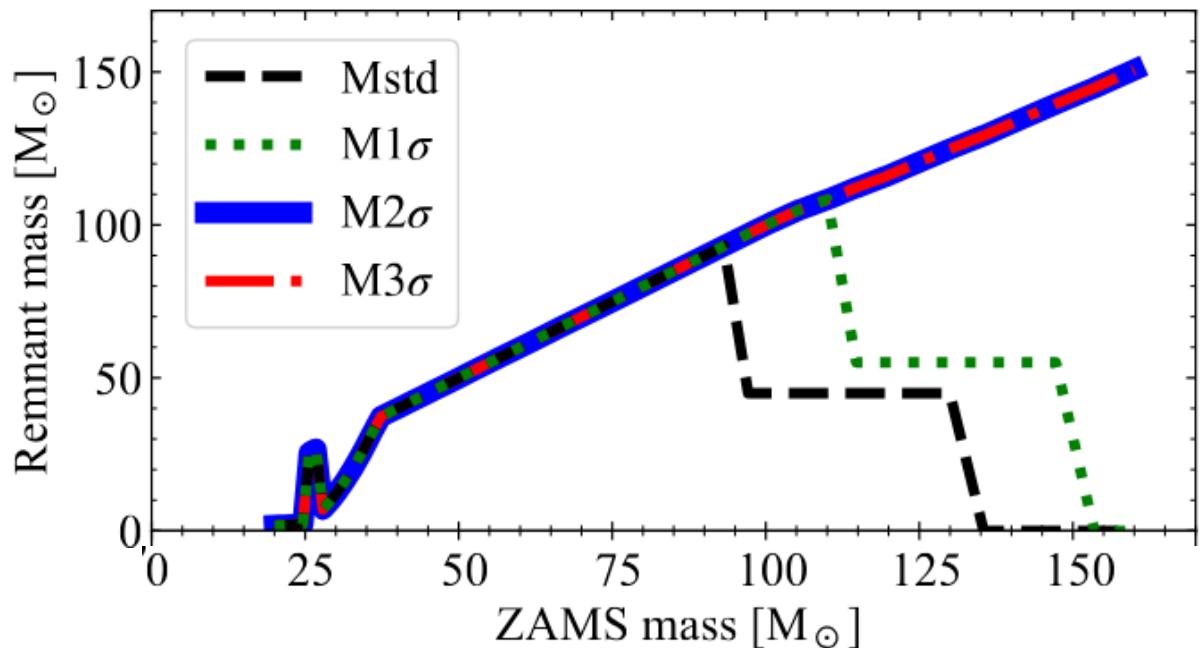
# GW231123

- 100M<sub>sun</sub> BH is **very difficult**,

But it might be explained by uncertainty of  $^{12}\text{C}(\alpha, \gamma)^{16}\text{O}$  and overshooting parameter (Tanikawa et al. 2025)



# GW231123



# Can Pop III BBH explain massive BBHs?

30Msun BBHs	✓
GW190521	✓
GW231123	?

Small  $^{12}\text{C}(\alpha, \gamma)^{16}\text{O}$  reaction rate and  
small overshooting parameter needed

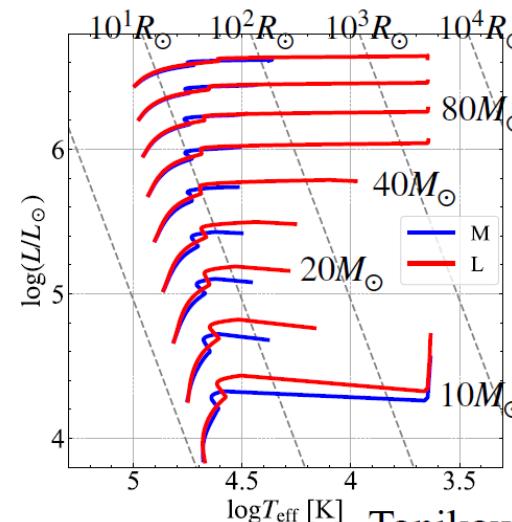
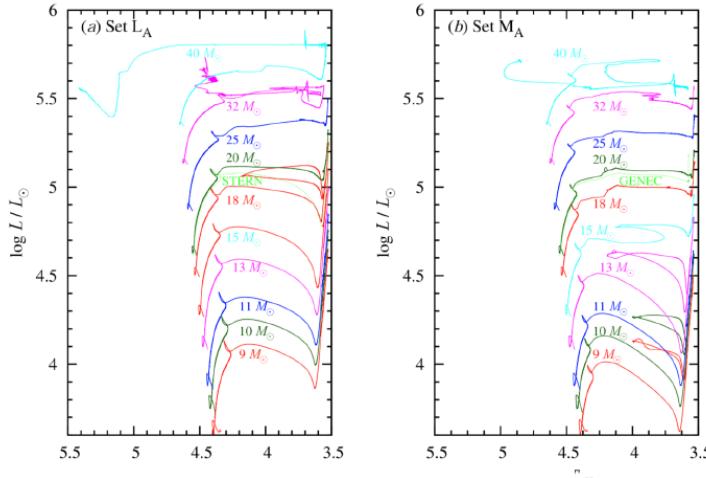
# Uncertainty in Pop. III model

- No massive Pop. III stars discovered so far
- Extrapolation from nearby stars to Pop. III stars
- Nearby star models
  - AB-type stars in MW open clusters, GENEC(Ekstrom et al. 2012), adopted by Farrell et al. (2020)
  - Early B-type stars in LMC, Stern (Brott et al. 2011)
- The maximums radius of a  $80M_{\odot}$  star
  - M model:  $\sim 40R_{\odot}$ , similar to Farrell et al. (2020)
  - L model:  $\sim 3 \times 10^3 R_{\odot}$ , similar to Yoon et al. (2012)

M model

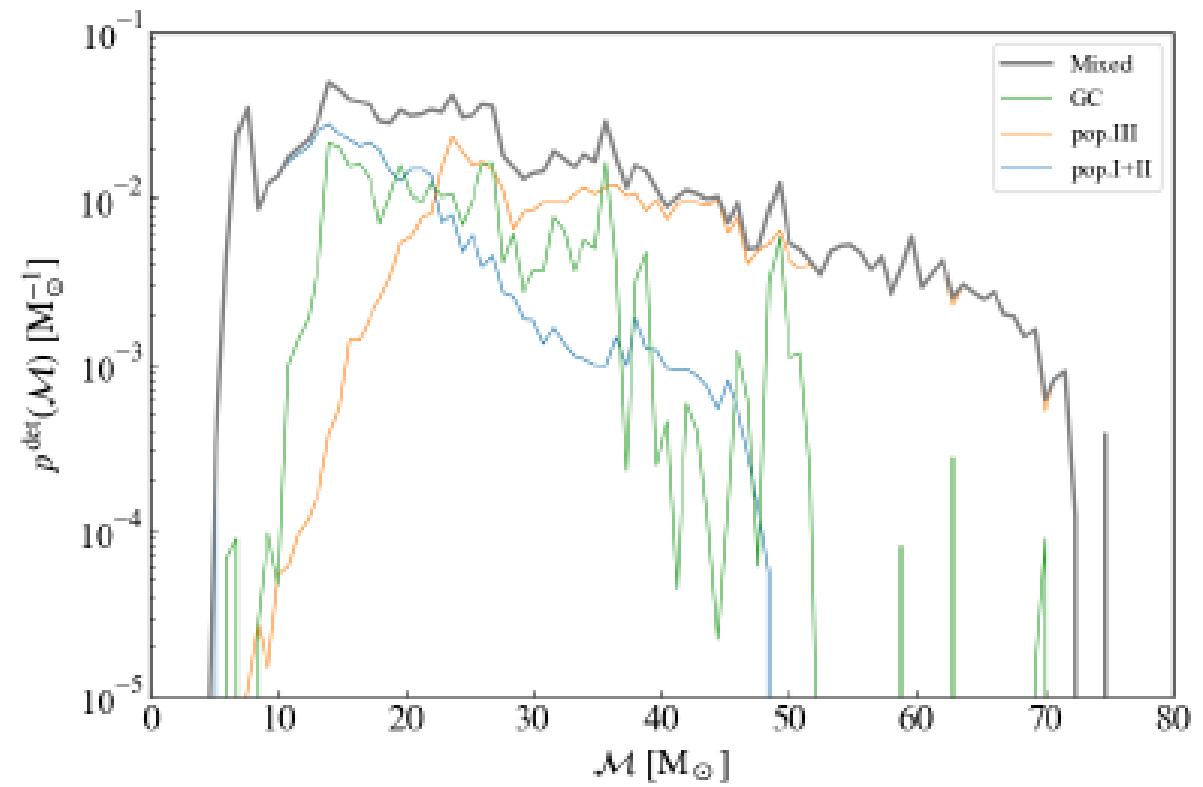
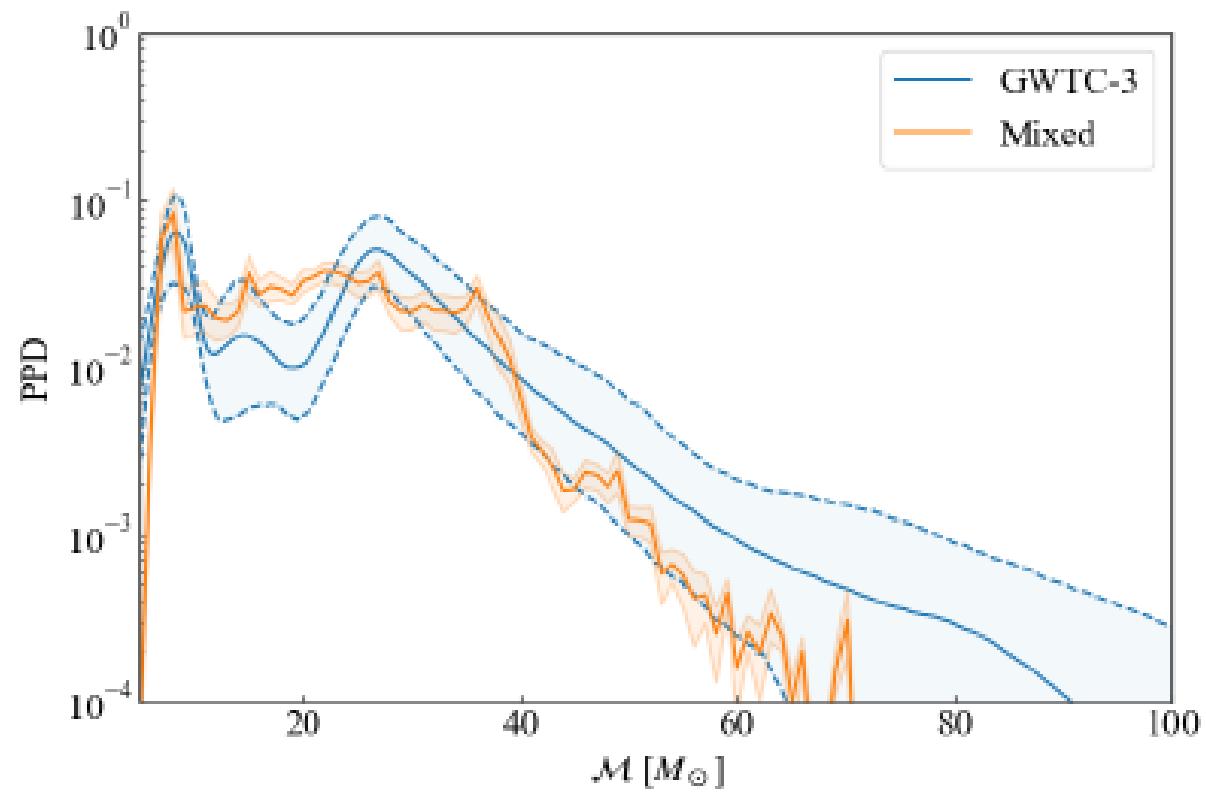
L model

Yoshida et al. (2019)



Two Pop. III models

Tanikawa et al. (2020c)



**Table 1.** Parameter sets and merger rate density of GW231123- and GW190521-like events in each set.

Name	Single star model (Overshoot efficiency)	PPISN/PISN model ( $^{12}\text{C}(\alpha, \gamma)^{16}\text{O}$ rate)	GW231123-like event [yr $^{-1}$ Gpc $^{-3}$ ]	GW190521-like event [yr $^{-1}$ Gpc $^{-3}$ ]
Mstd	M (inefficient)	Standard	0	$4.9 \times 10^{-2}$
M $1\sigma$	M (inefficient)	$1\sigma$ lower	0	$1.0 \times 10^{-1}$
M $2\sigma$	M (inefficient)	$2\sigma$ lower	$2.1 \times 10^{-2}$	$1.0 \times 10^{-1}$
M $3\sigma$	M (inefficient)	$3\sigma$ lower	$2.1 \times 10^{-2}$	$1.0 \times 10^{-1}$
Lstd	L (efficient)	Standard	0	0
L $3\sigma$	L (efficient)	$3\sigma$ lower	0	$8.6 \times 10^{-2}$

NOTE—The Mstd and L $3\sigma$  sets are the same as the fiducial and L- $3\sigma$  sets in [A. Tanikawa et al. \(2022\)](#).

# Merger rate (GWTC-3)

13-1900 /yr/Gpc<sup>3</sup>  
(NS-NS)

7.4-320 /yr/Gpc<sup>3</sup>  
(NS-BH)

17.3-45 /yr/Gpc<sup>3</sup>  
(BH-BH)

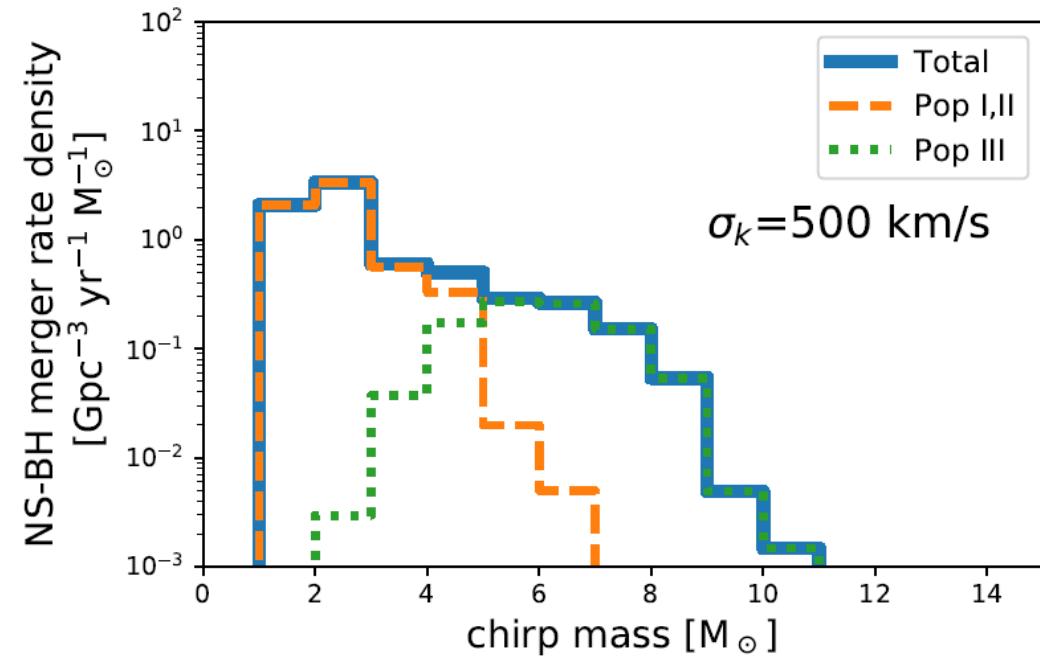
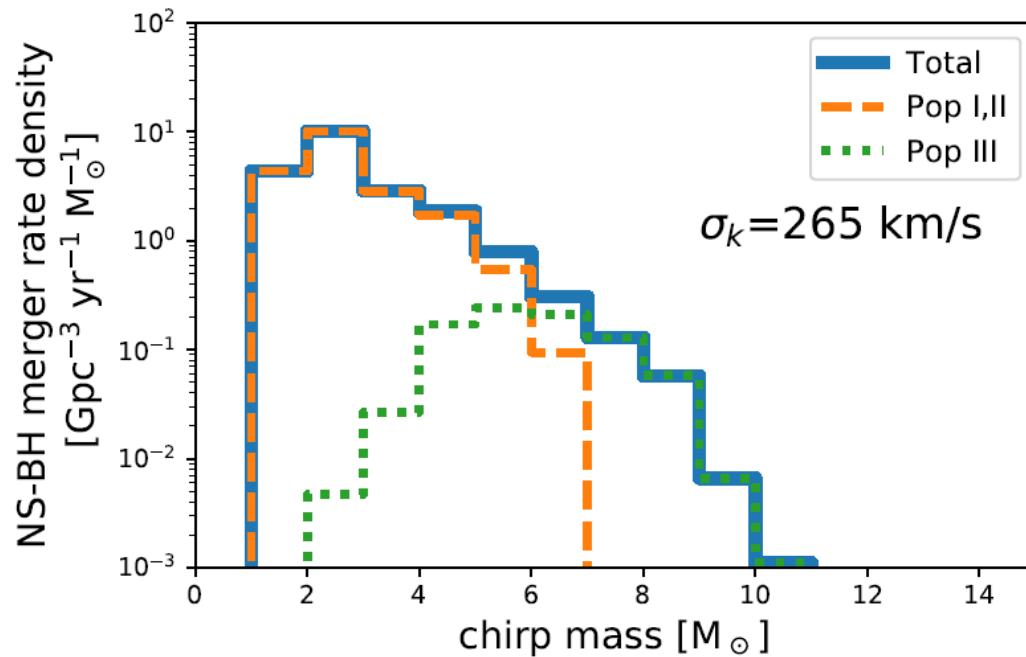
# NS-BH formation (Kinugawa et al. 2017)

- Pop I/II
- Pop III

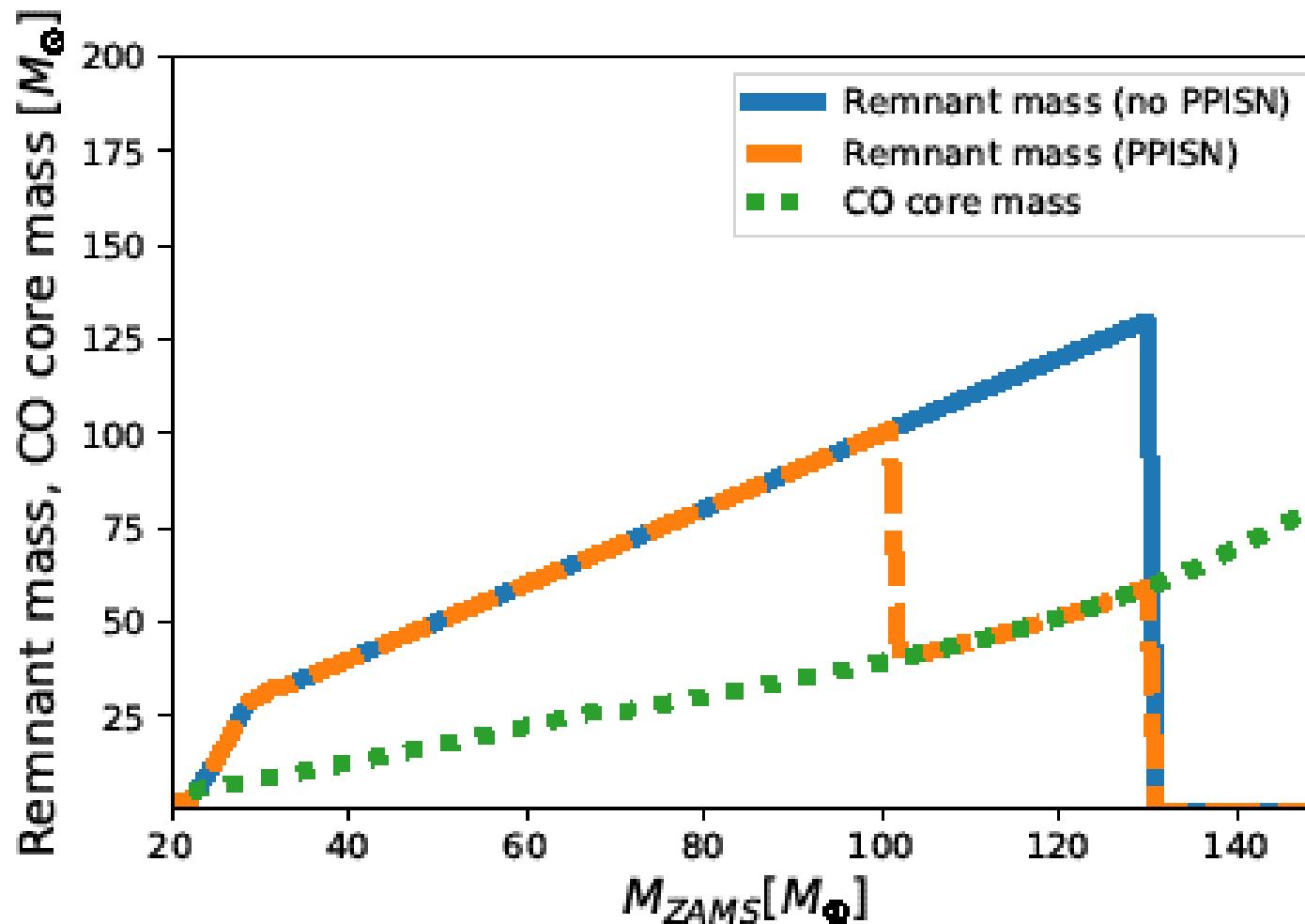
Table 2: The number of NS-BH formations and the number of NS-BHs which merge within 15 Gyrs for each metallicity for the initial  $10^6$  binaries. The numbers are for the  $\sigma_k = 265$  km/s models, while the numbers in the parenthesis are for the  $\sigma_k = 500$  km/s models.

$Z$	$Z_\odot$	$10^{-0.5} Z_\odot$	$10^{-1} Z_\odot$	$10^{-1.5} Z_\odot$	$10^{-2} Z_\odot$	0
NS-BH	148 (32)	598 (169)	1296 (416)	1686 (576)	1896 (617)	22638 (11192)
merging NS-BH	15 (2)	191 (67)	525 (213)	755 (377)	862 (401)	9089 (5856)

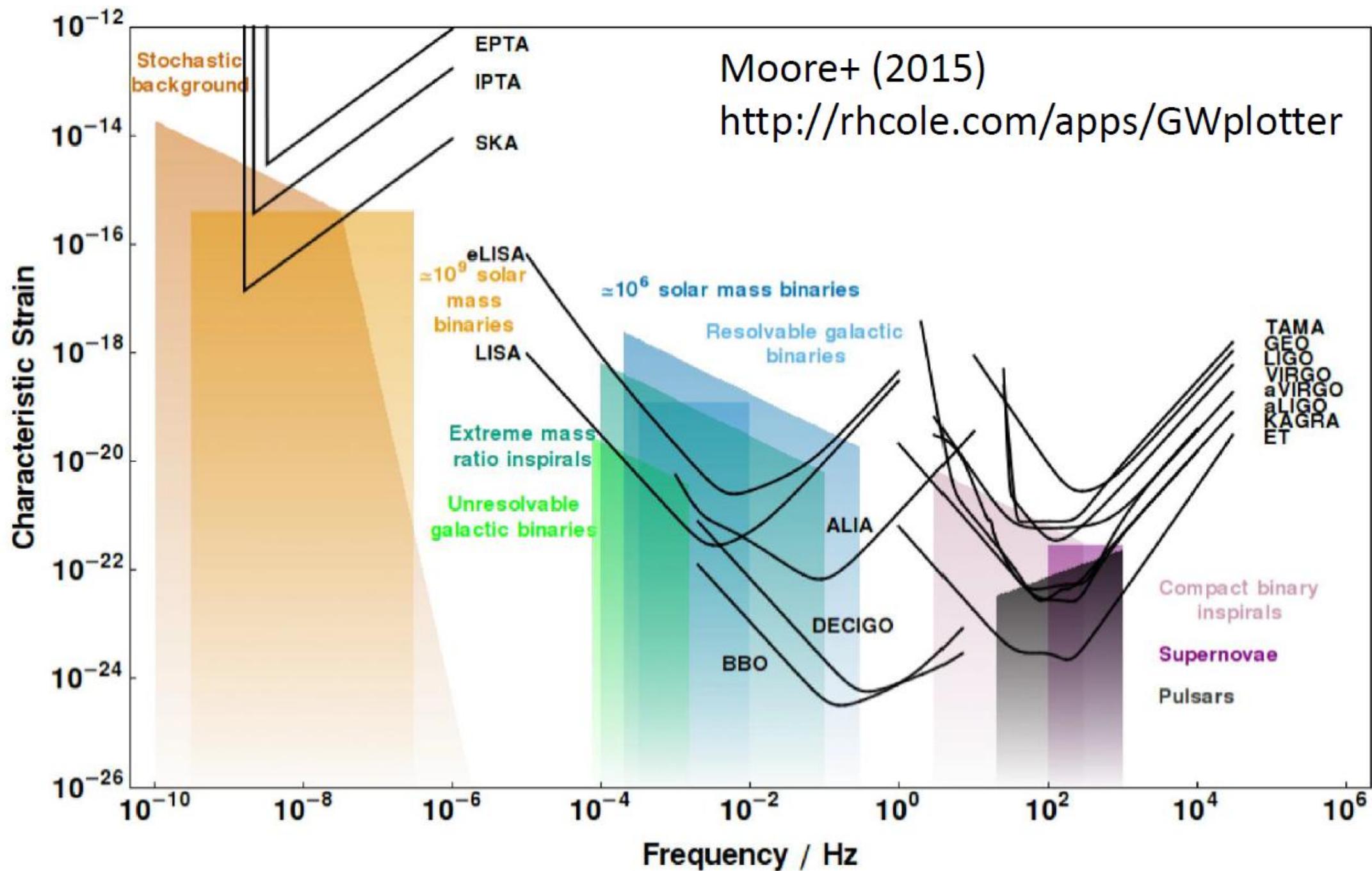
# Chirp mass distribution of observable NS-BH

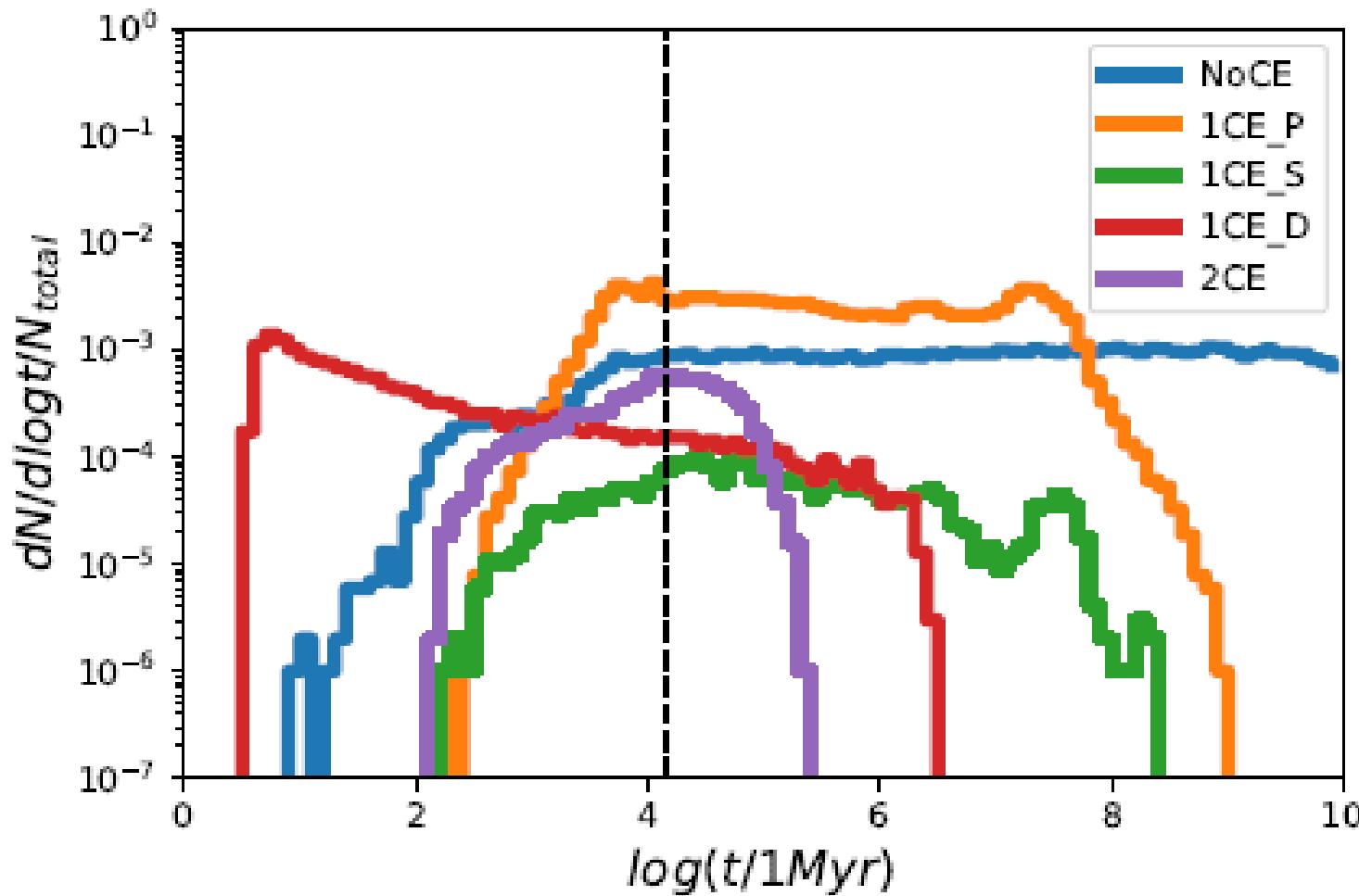


# Pop III remnant mass for single star case



- We assume  $M_{\text{He}} = 40-60 \text{ Msun} \rightarrow \text{PPISN}$





(a) fiducial model

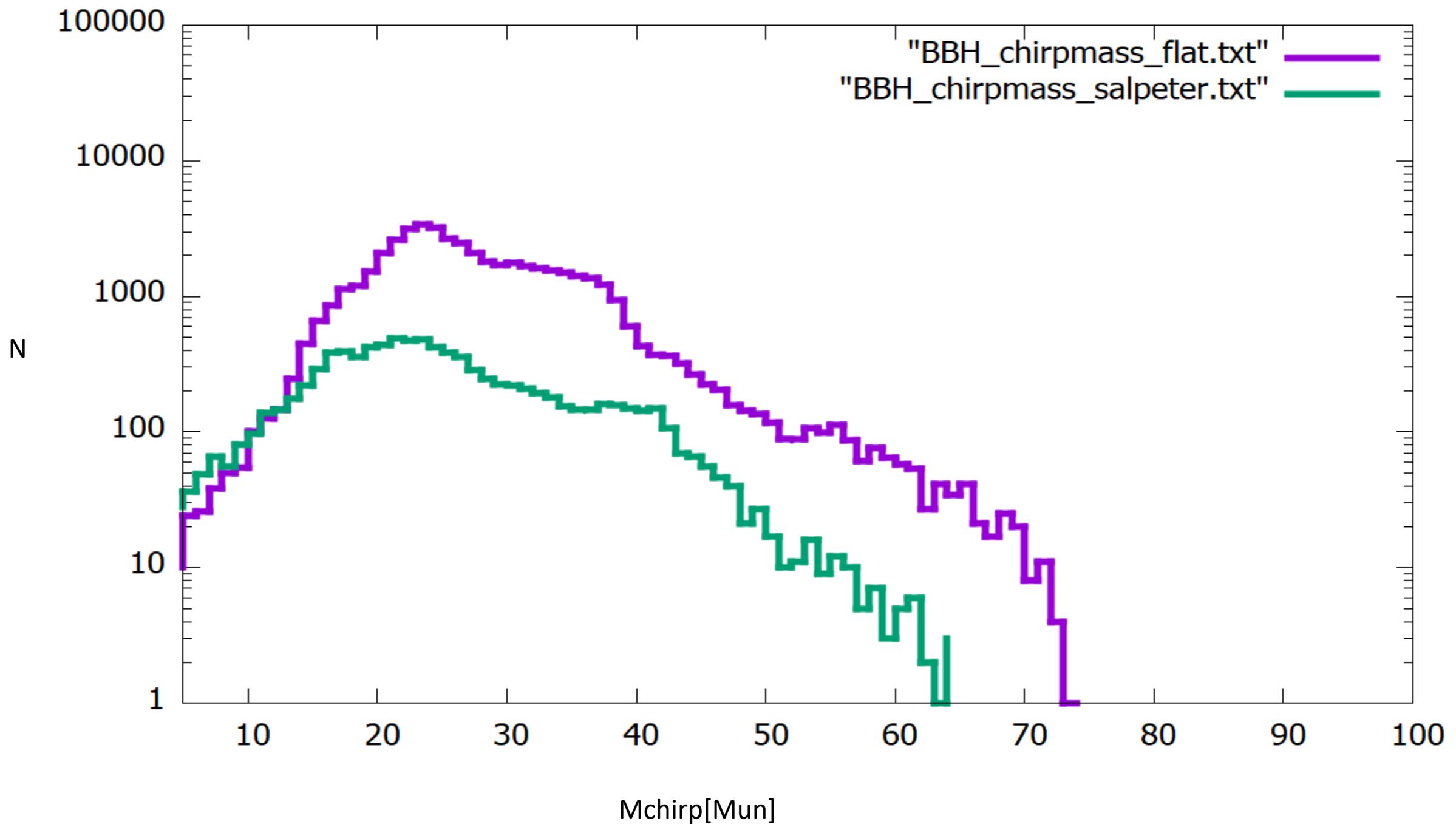
# Event rate of **GW190521** like BH-BH mergers

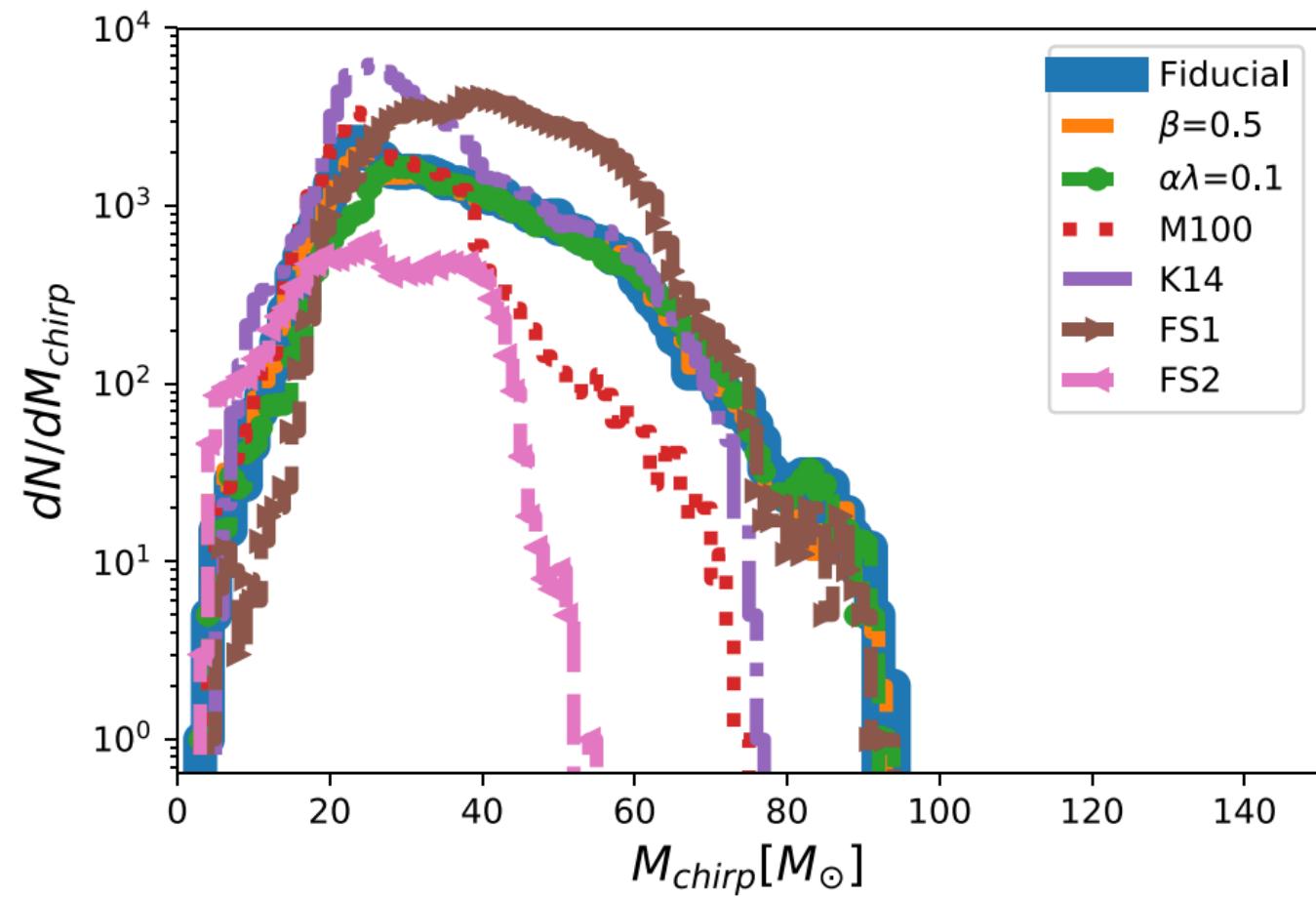
Pop III GW190521 like BBH merger rates at the present day

- $0.13 \text{ /yr/Gpc}^3$  for PPISN model
- $0.66 \text{ /yr/Gpc}^3$  for no PPISN model

Rate of GW190521 by LIGO is

$0.02\text{--}0.43 \text{ /yr/Gpc}^3$





# However...

After GW150914, there are 1 bad news and 1 objection for Pop III BBH scenario

## 1. Bad news

~ decreasing expected Pop III SFR

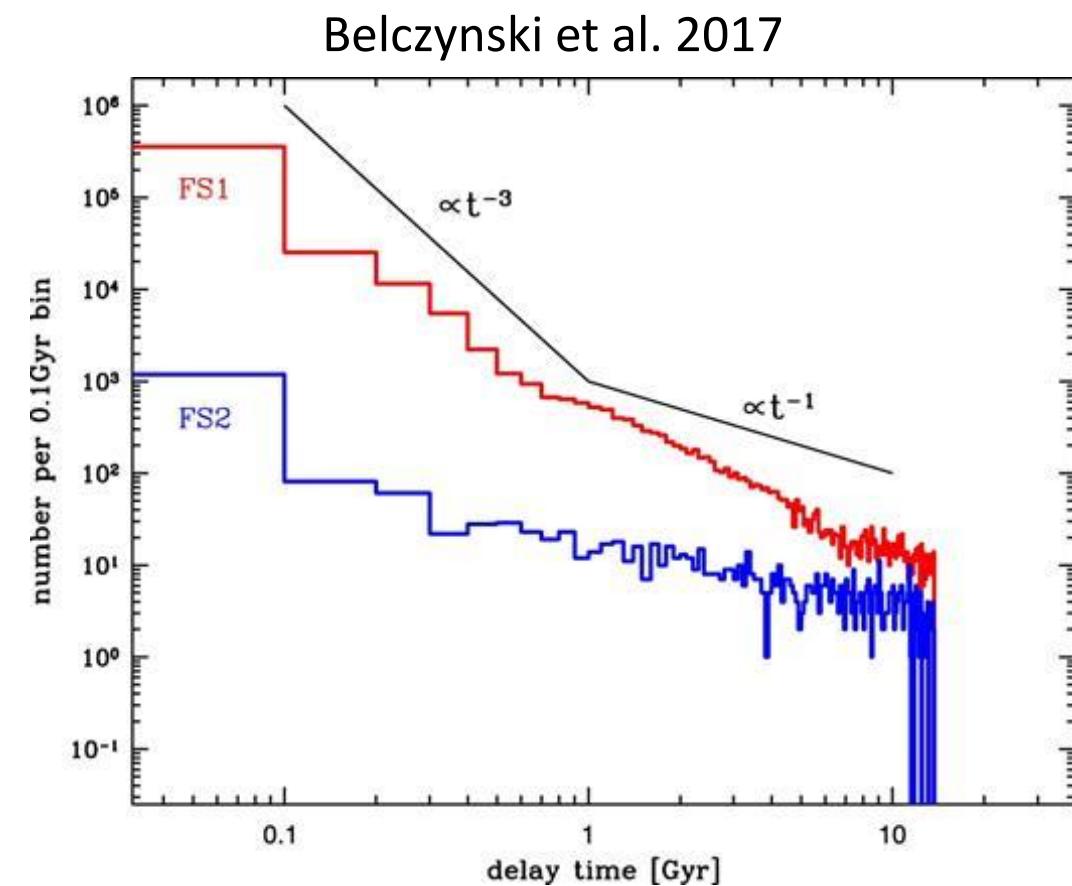
Because of constraints by Planck  $\tau$ e

(Visbal et al. 2015, Hartwig et al. 2016, Inayoshi et al. 2016)

## 2. Objection

Chris Belczynski also tried to calculate Pop III BBH merger rate.

In his calculation, almost all Pop III BBHs *merge at the early universe*



# Pop III star formation constraint by Planck

- The optical depth of the universe to electron scattering was inferred from CMB anisotropies by the *Planck*
- It is lower than previous estimates from *WMAP*
- This makes tight constraints on the star formation history of Pop III

- Before *Planck*

$$\rho = 2 \times 10^6 \text{ Msun/Mpc}^3 \text{ (de Souza et al. 2011)}$$

- After *Planck*

Optimistic constraint  $\rho \leq 6 \times 10^5 \text{ Msun/Mpc}^3$  ← our model uses this value  
(Inayoshi et al. 2016)

Conservative constraint  $\rho \leq 2 \times 10^5 \text{ Msun/Mpc}^3$   
(Visval et al. 2015, Inayoshi et al. 2021)

# However...

After GW150914, there are 1 bad news and 1 objection for Pop III BBH scenario

## 1.Bad news

decreasing expected Pop III SFR

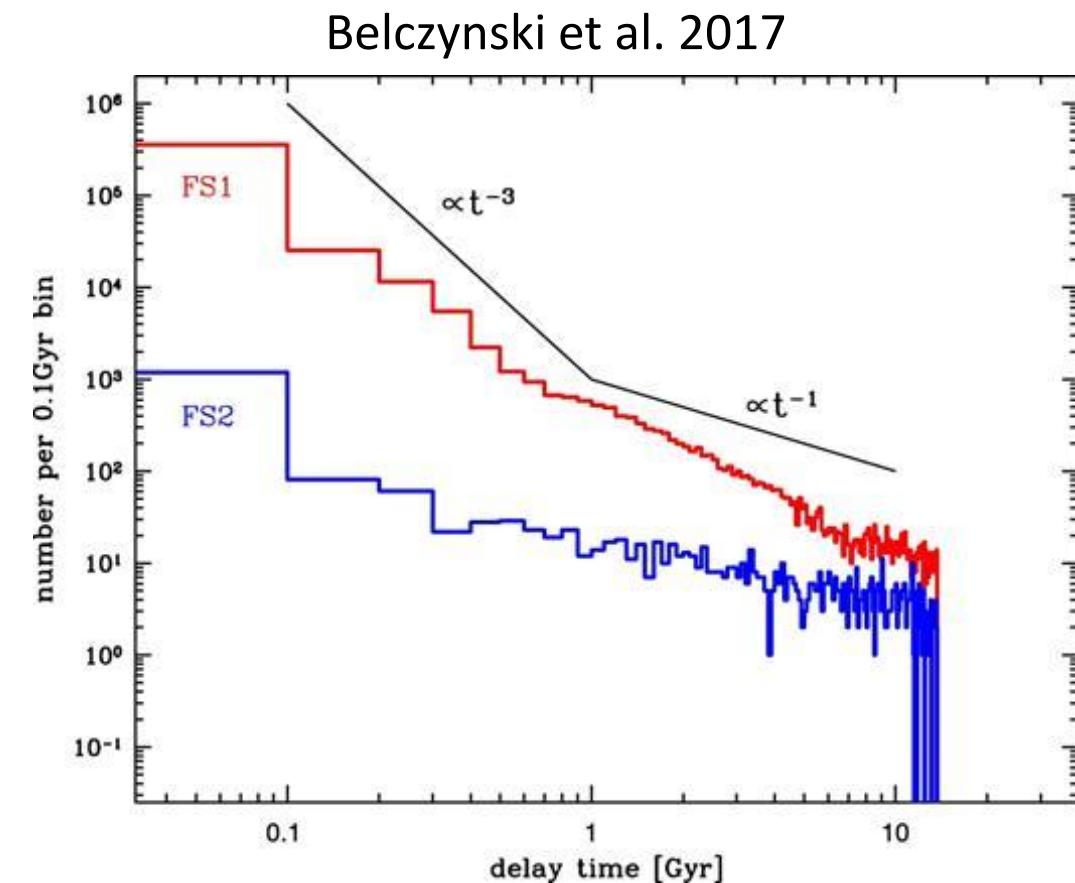
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## 2.Objection

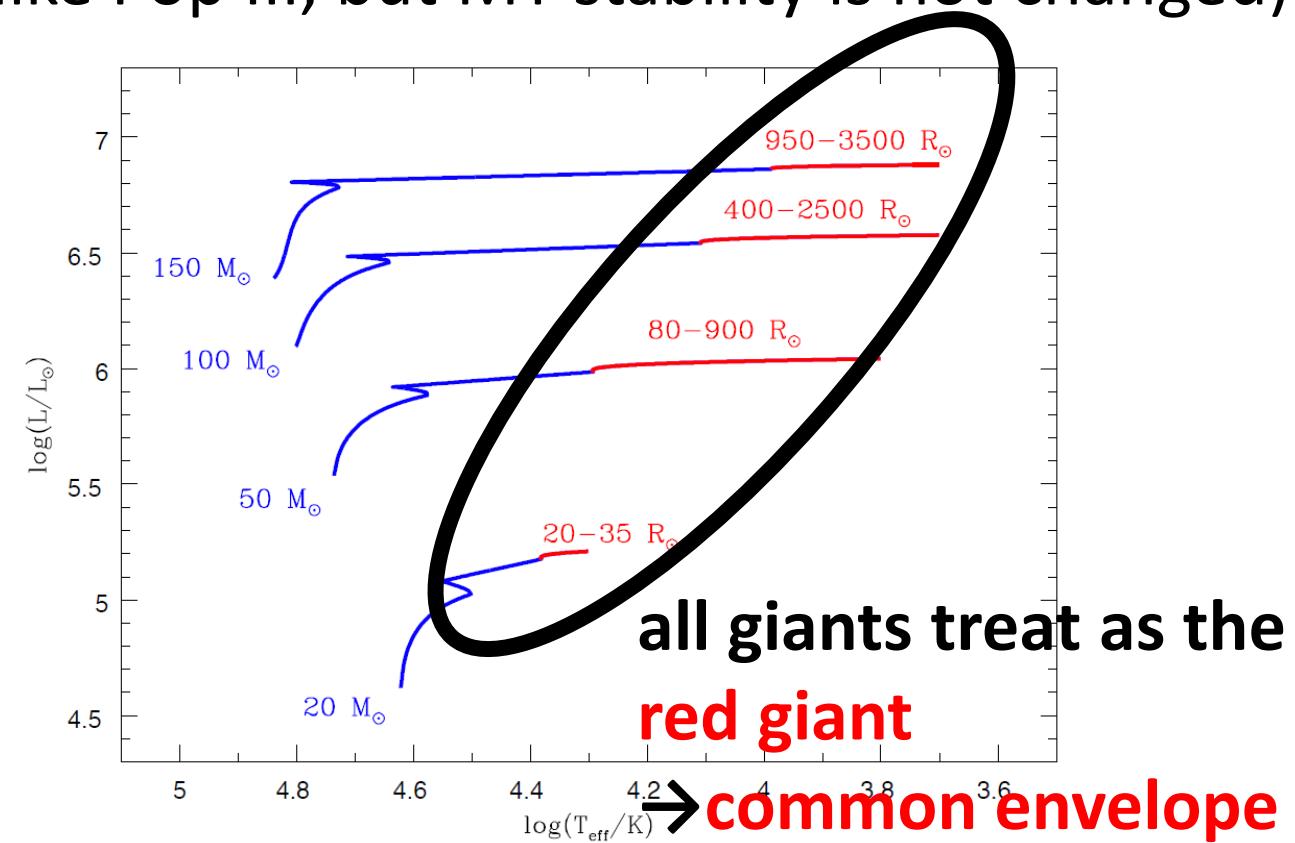
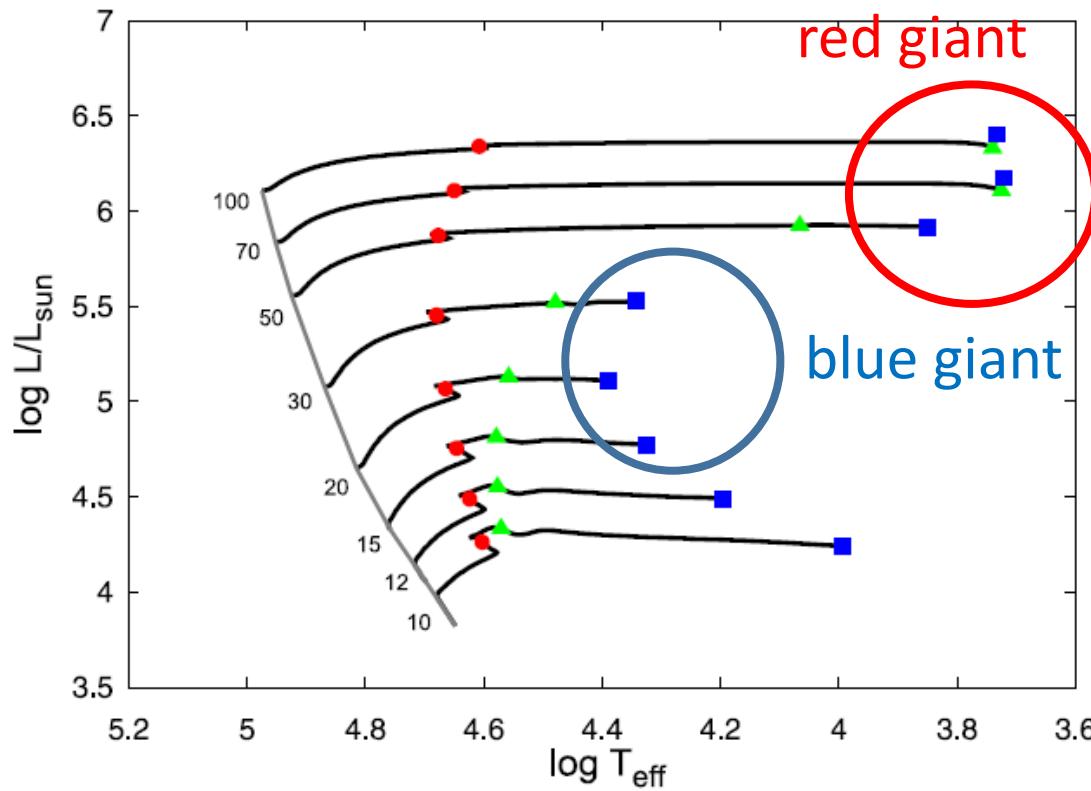
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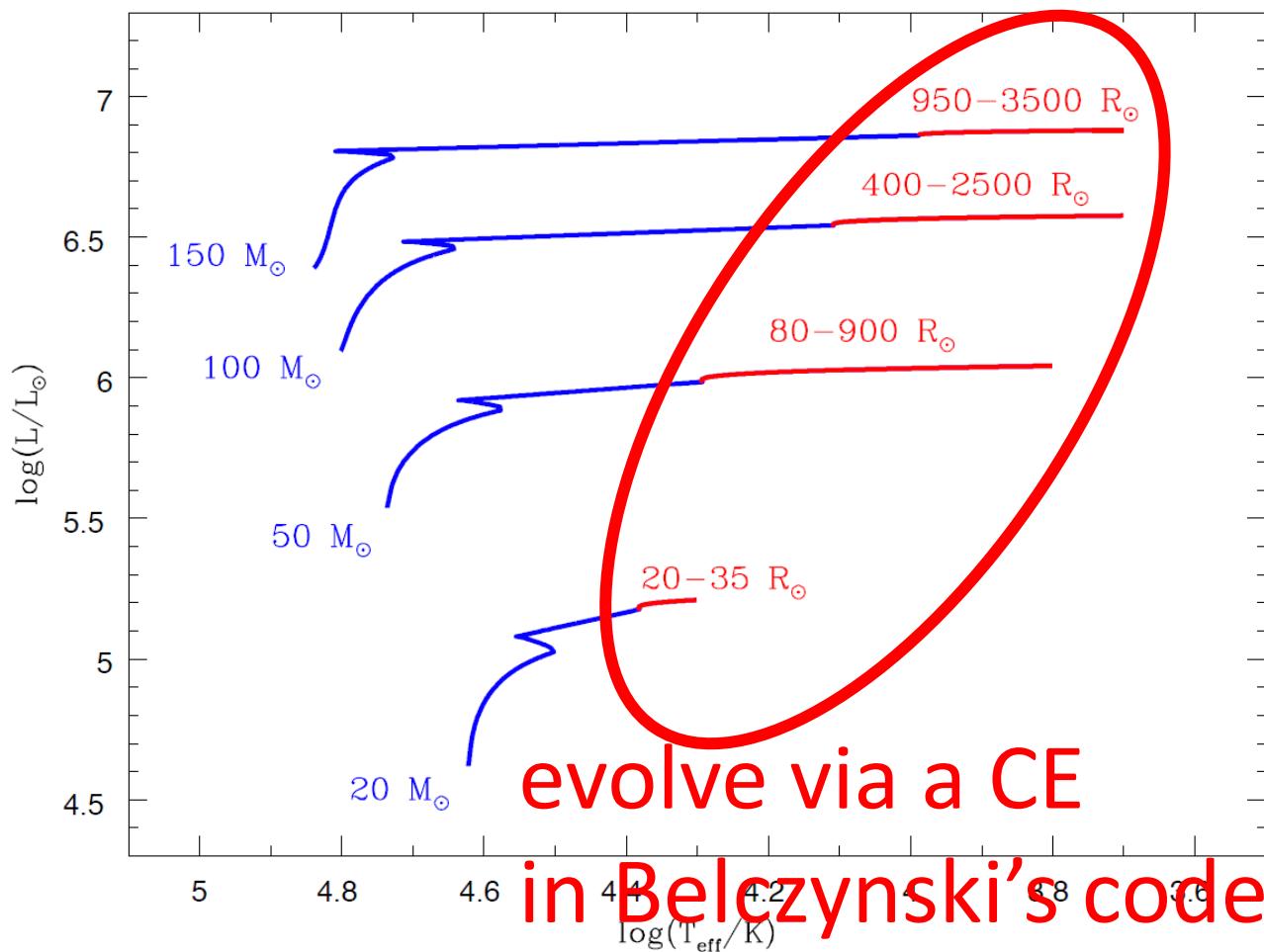


# Difference between K14 and Belczynski's Pop III calc.

- Kinugawa 2014: use Pop III stellar evolution model (Marigo et al.2001)
- Belczynski 2017: use modified  $Z=0.005Z_{\text{sun}}$  model.  
(HR and radius evolution is changed like Pop III, but MT stability is not changed)



# Difference between our code and Belczynski's Pop III

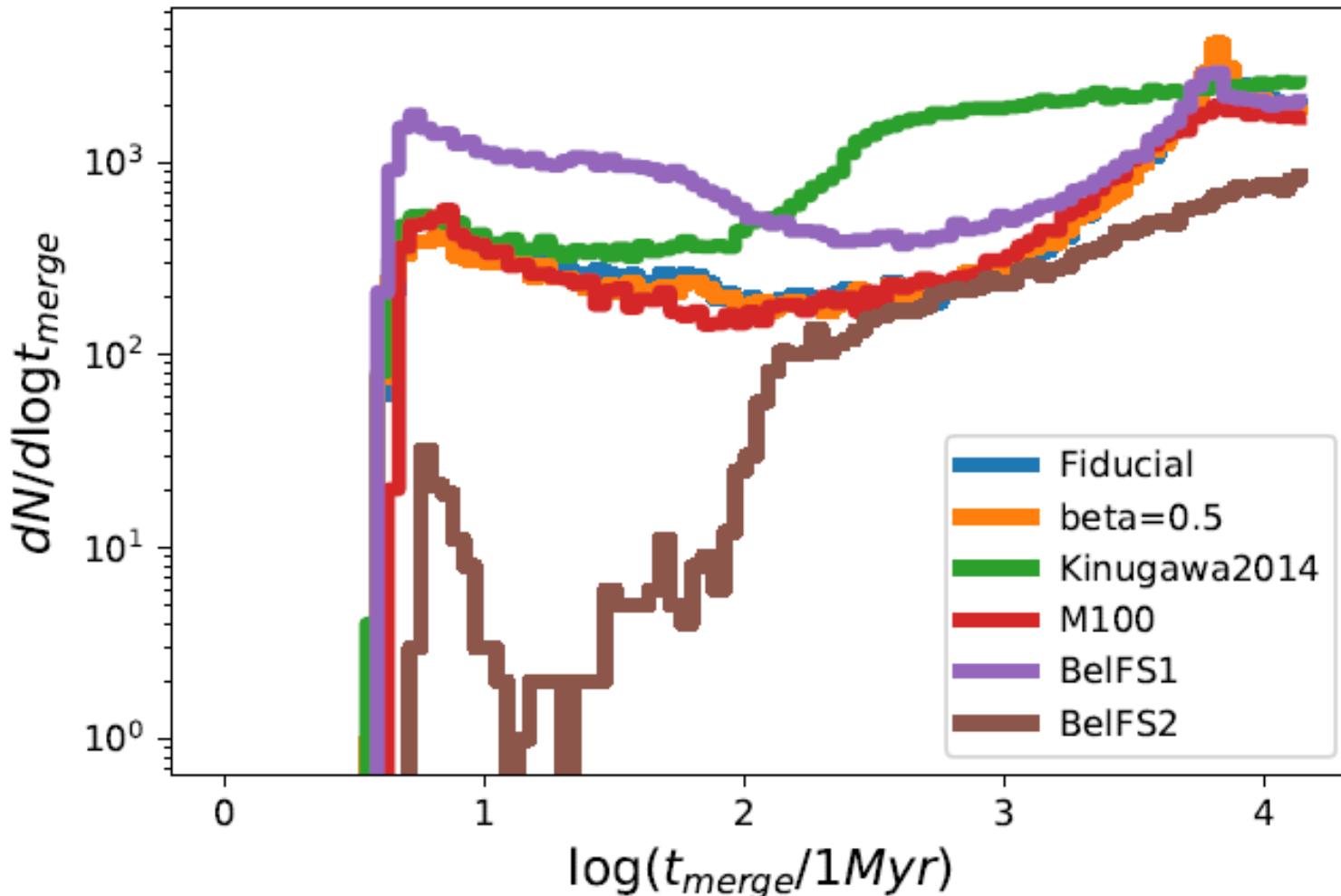


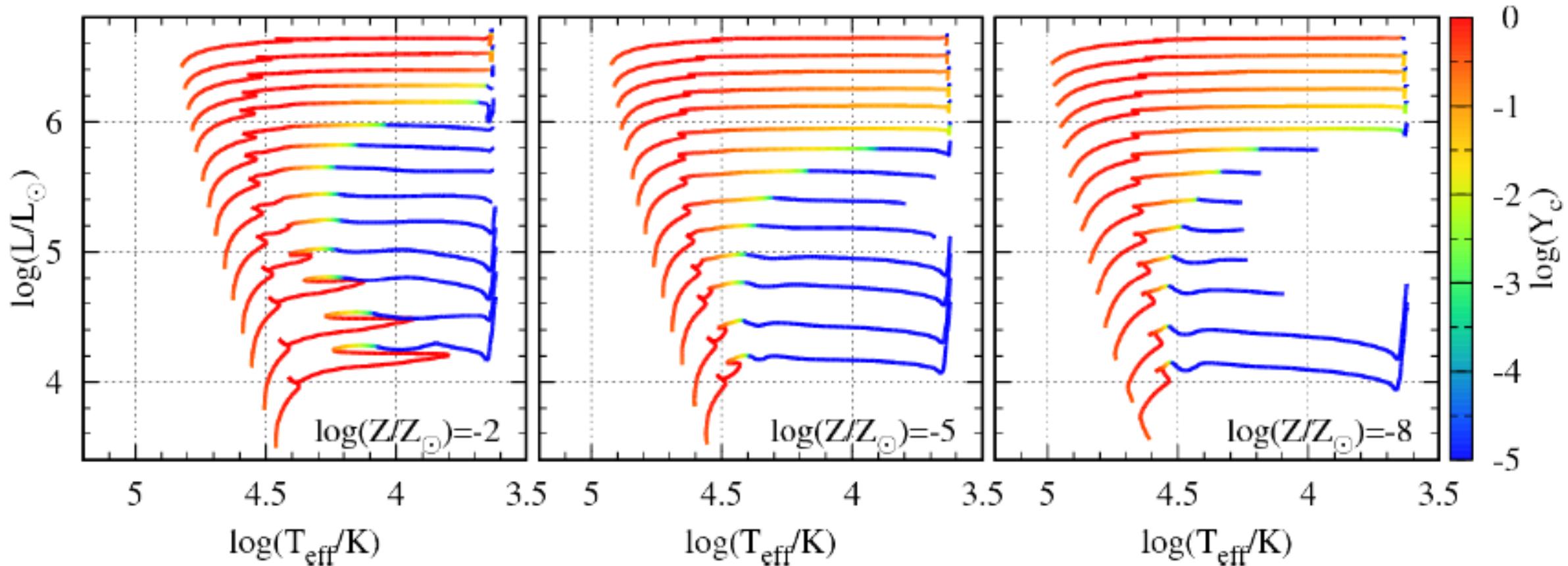
## Belczynski's code

- Modified Pop II ( $Z=10^{-4}$ ) evolution
- The radius evolution is likely.
- But, the mass transfer treatment is same as Pop II

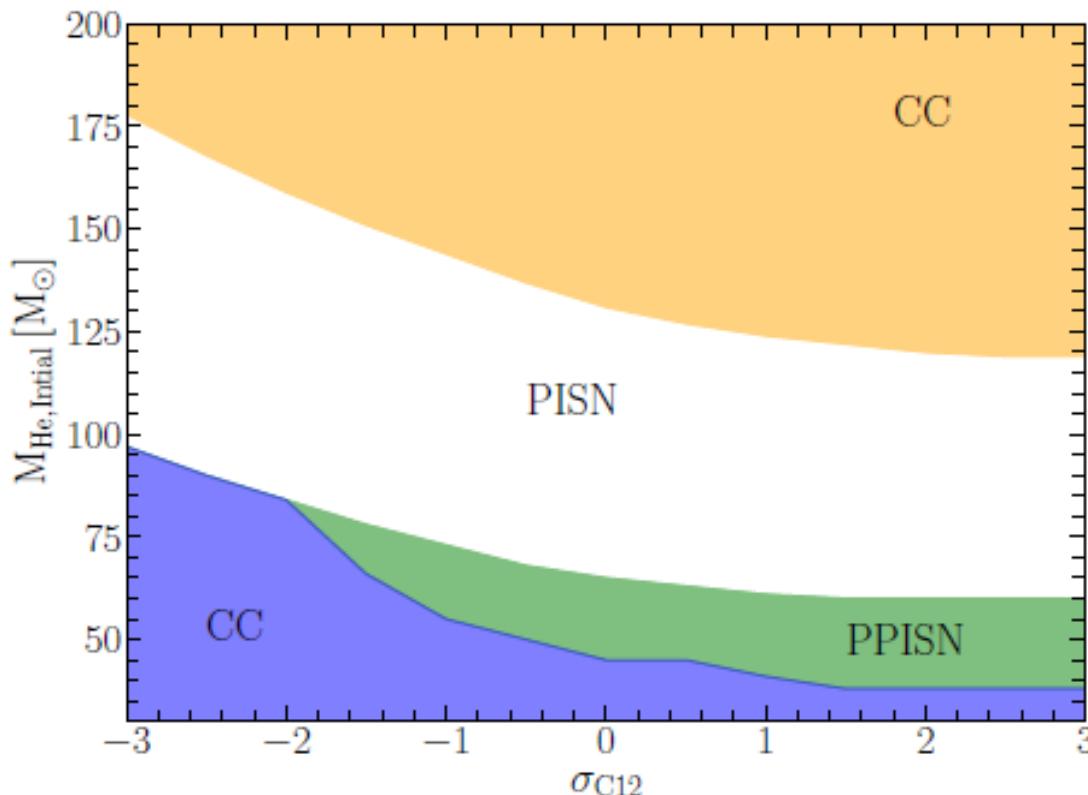
→ all BBH evolved via  
a common envelope (CE)  
many binaries merge during a CE  
Merger Rate of Pop III BBH decrease  
and  $dN/dt$  change

# Merger time distribution of Pop III BBH

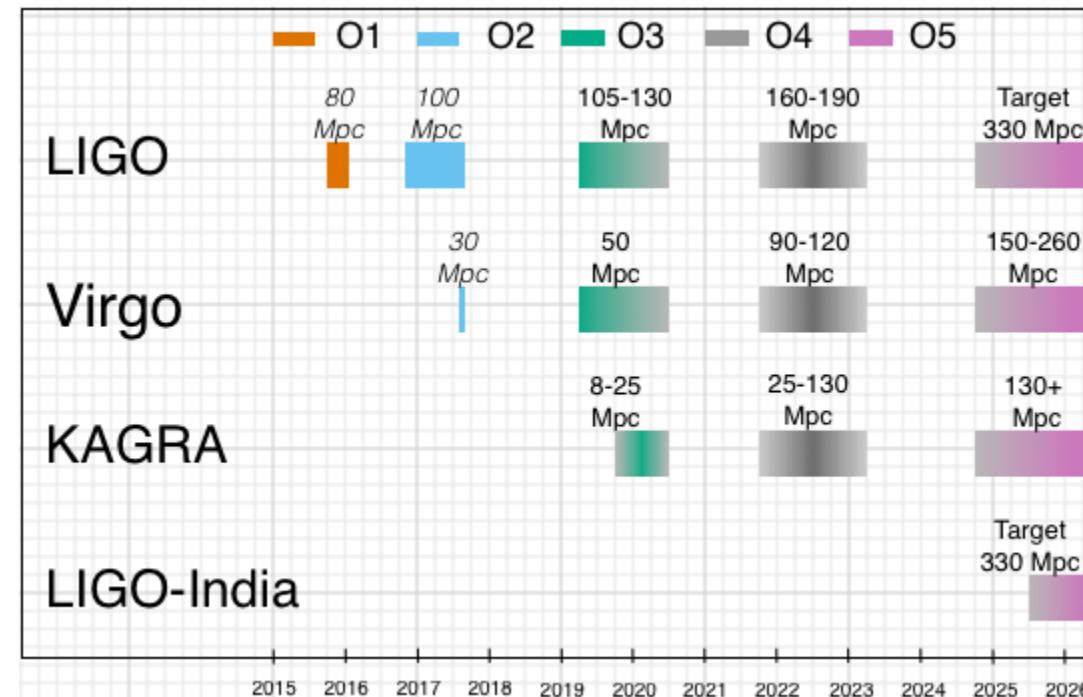




**Figure 1.** HR diagrams for stellar models with  $\log(Z/Z_{\odot}) = -2, -5$ , and  $-8$ . In each panel, curves indicate stellar evolutions with  $M/M_{\odot} = 8, 10, 13, 16, 20, 25, 32, 40, 50, 65, 80, 100, 125$ , and  $160$  from bottom to top. Colors are coded by the helium mass fractions in the stellar cores.



**Figure 1.** Final fate of a star as function of the initial helium core mass and  $^{12}\text{C}(\alpha, \gamma)^{16}\text{O}$  rate.  $\sigma_{\text{C}12}$  denotes how far the  $^{12}\text{C}(\alpha, \gamma)^{16}\text{O}$  is from the median STARLIB rate, measured in standard deviations. Blue regions indicate stars which undergo core collapse (CC) below the pair instability supernovae (PISN) mass gap, green regions form black holes after a pulsational pair instability supernovae (PPISN), while white regions are completely disrupted in a PISN, and models in the orange region form black holes from core collapse for stars above the PISN mass gap. There are 2210 models, in the grid spaced by  $1 M_{\odot}$  and  $0.5\sigma_{\text{C}12}$ .



# Other Pop III SFRs

- simulation  
e.g. Johnson et al. 2013  
 $SFR_p \sim 10^{-3}\text{-}10^{-4} \text{ Msun/yr/Mpc}^3$
- Constraints by Planck  $\tau_e$   
e.g. Visbal et al. 2015, Hartwig et al. 2016,  
Inayoshi et al. 2016  
→ The merger rate might decrease to 1/3-1/10 ?

