

RELATIONSHIP BETWEEN $0\nu\beta\beta$ AND LBL MEASUREMENTS

From parameter determination to natural flavor predictions

NPN2026 (The Hirosawa City Center, Mito)

May/26/2026

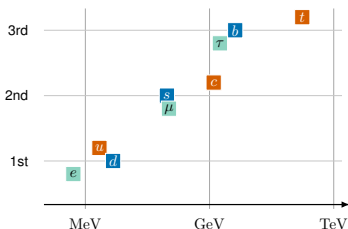
Masahiro Ibe

ICRR, The University of Tokyo

Flavor Structures in the SM + Neutrino Masses

✓ Fermion masses : hierarchical pattern across generations

quark/charged lepton masses

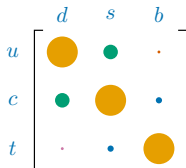


neutrino masses

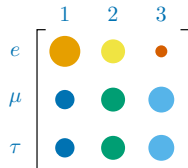
- ✓ $\Delta m_{21}^2 \simeq 7.5 \times 10^{-5} \text{ eV}^2$
($m_2 > m_1$ with $0 < \theta_{12} < \pi/4$)
- ✓ $|\Delta m_{31}^2| \simeq 2.5 \times 10^{-3} \text{ eV}^2$
- ✓ mass ordering : not yet determined
- ✓ absolute mass scale :
 - $\Sigma m_\nu < \mathcal{O}(0.1) \text{ eV}$ (cosmology)
 - $\Sigma m_\nu < \mathcal{O}(1) \text{ eV}$ (β -decay)

✓ Fermions with different flavors mix with each other

CKM: small mixing



PMNS: large mixing



Flavor Puzzle?

The flavor puzzle is the question of the origin of these structures

Many flavor models have been proposed:

flavor symmetries, texture zeros, zero minors . . .

However, the observed flavor structure can simply be **fitted by parameters**

$$\begin{aligned}\mathcal{L}_{\text{flavor}} = & - y_i^u \delta_{ij} H Q_i \bar{u}_j - y_i^d V_{ij}^{\text{CKM}} H^\dagger Q_i \bar{d}_j \\ & - y_i^\ell \delta_{ij} H^\dagger \bar{e}_i L_j - \frac{1}{2v_{\text{EW}}^2} (U^* m_\nu^{\text{diag}} U^\dagger)_{ij} (H L_i)(H L_j) + \text{h.c.}\end{aligned}$$

[All fermions are written as left-handed Weyl fermions]

In frequentist statistics:

✓ parameters = fixed constants

⇒ **flavor structure just tells us that the world is built that way.**

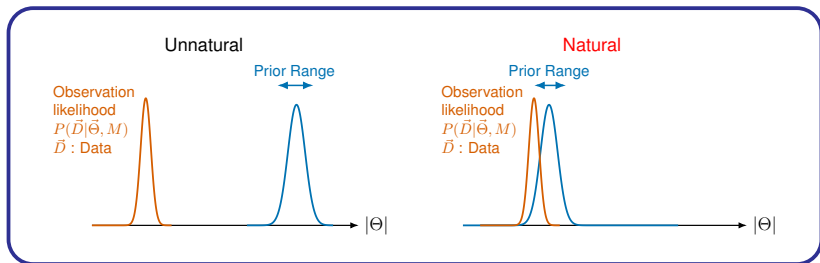
Naturalness as Typicality

To go beyond fitting, we need a notion of naturalness.

A flavor model M with $\mathcal{O}(1)$ parameters $\vec{\Theta}$

In Bayesian statistics :

- ✓ parameters = random variables
drawn from the prior distribution $\pi(\vec{\Theta}|M)$
- ✓ **Naturalness = whether the observed data lie in a typical region predicted by the prior**



Naturalness is quantified as the overlap between likelihood and prior distribution.

$\mathcal{O}(1)$ Yukawa Coupling Constants

The prior distribution of the complex-valued 3×3 Yukawa couplings,

$$P(Y^F) = \left(\frac{1}{\sqrt{2\pi}\sigma} \right)^{18} \exp \left[-\frac{\text{tr}[Y^\dagger Y]}{2\sigma^2} \right], \quad \sigma \sim 1$$

[integration measure $d\text{Re}(Y_{ij}^F)d\text{Im}(Y_{ij}^F)$]

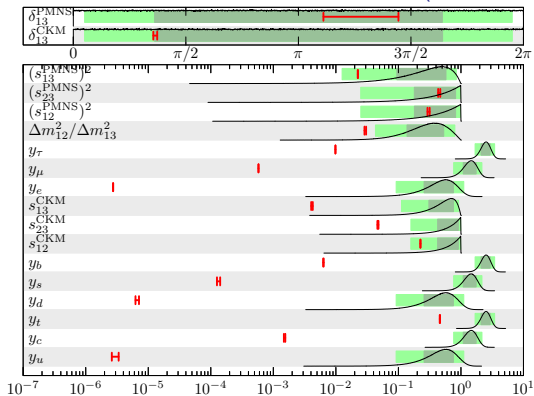
Why Gaussian?

- ✓ Zero mean and common variance: no preferred flavor entry.
- ✓ Maximum entropy ($S[p] = \int dx p(x) \log p(x)$) for fixed mean and variance: least biased choice.
- ✓ The width $\sigma \sim 1$ implements generic $\mathcal{O}(1)$ couplings.

$\mathcal{O}(1)$ Yukawa Coupling Constants

What happens without flavor structure?

Distributions of dimensionless observables (normal ordering)



Red bars : Observed Likelihood

Curves and green shaded region :

Posterior distribution (1σ and 2σ ranges)

Without flavor structure, fitting is quite unnatural ...

Example : Froggatt–Nielsen Mechanism

In the FN mechanism, hierarchies arise from powers of a small parameter:

$$Y_{ij} \sim c_{ij} \epsilon^{|f_i+f_j|}, \quad c_{ij} = \mathcal{O}(1).$$

- ✓ f 's : U(1) charge of quarks/leptons
- ✓ ϵ : order parameter of U(1) symmetry
- ✓ c_{ij} random $\mathcal{O}(1)$ variables

$$\text{Ex. } f_Q = (3, 2, 0), \quad f_{\bar{u}} = (5, 2, 0), \quad f_{\bar{d}} = (4, 4, 3), \\ f_L = (4, 3, 3), \quad f_{\bar{e}} = (4, 1, 0),$$

$$y_u : y_c : y_t = \epsilon^8 : \epsilon^4 : \epsilon^0, \quad y_d : y_s : y_b = \epsilon^7 : \epsilon^6 : \epsilon^3,$$

$$y_e : y_\mu : y_\tau = \epsilon^8 : \epsilon^4 : \epsilon^3,$$

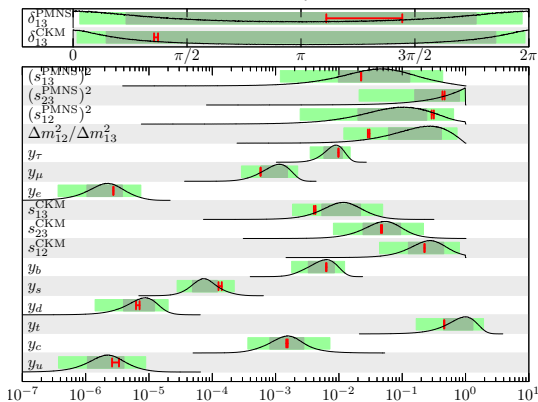
$$|V_{\text{CKM}}| \sim \begin{pmatrix} \epsilon^0 & \epsilon^1 & \epsilon^3 \\ \epsilon^1 & \epsilon^0 & \epsilon^2 \\ \epsilon^3 & \epsilon^2 & \epsilon^0 \end{pmatrix}, \quad |U_{\text{PMNS}}| \sim \begin{pmatrix} \epsilon^0 & \epsilon^1 & \epsilon^1 \\ \epsilon^1 & \epsilon^0 & \epsilon^0 \\ \epsilon^1 & \epsilon^0 & \epsilon^0 \end{pmatrix}.$$

$$\Rightarrow \epsilon \sim 0.2 \text{ (Cabbibo angle)}$$

Example : Froggatt-Nielsen Mechanism

Distributions of dimensionless observables (normal ordering)

$$\epsilon = 0.185, f_{\bar{u}} = (5, 2, 0), f_{\bar{d}} = (4, 4, 3), f_Q = (3, 2, 0), f_{\bar{e}} = (4, 1, 0), f_L = (4, 3, 3)$$



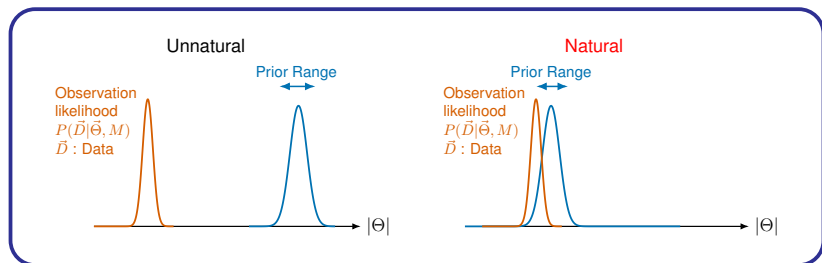
Red bars : Observed Likelihood

Curves and green shaded region :

Posterior distribution (1σ and 2σ ranges)

Flavor structure is naturally explained !

Bayes Factor as Relative Naturalness



- ✓ Evidence = overlap between $P(\vec{D}|\vec{\Theta}, M)$ and $\pi(\vec{\Theta}|M)$:

$$P(\vec{D}|M_i) = \int d\Theta P(\vec{D}|\vec{\Theta}, M)\pi(\vec{\Theta}|M_i)$$

- ✓ Model comparison is quantified by the Bayes factor:

$$B_{ij} = P(\vec{D}|M_i)/P(\vec{D}|M_j)$$

The larger the Bayes factor, the more “natural” the model is

What is a Good Flavor Model?

Flavor model building = flavor origin \times typicality (reasonable prior)

- ✓ **Origin:** the model should explain why flavor structures appear:

mass hierarchy, small CKM mixing, large PMNS mixing.

- ✓ **Typicality:** the observed data should lie in a typical region of the model prediction:

$$P(\vec{D}|M) \text{ large.}$$

- ✓ **Caveat:** The prior distributions are not unique.
They should be chosen and interpreted carefully.

**A good flavor model should not only fit the data,
but should predict them naturally.**

Long-Base-Line Experiments

- ✓ Hyper-K/DUNE ($\nu_\mu/\bar{\nu}_\mu \rightarrow \nu_{e,\mu}/\bar{\nu}_{e,\mu}$)
⇒ sensitive to δ_{CP} , θ_{23} , $|\Delta m_{31}^2|$ and mass ordering
- ✓ JUNO (reactor $\bar{\nu}_e \rightarrow \bar{\nu}_e$)
⇒ sensitive to θ_{12} , Δm_{21}^2 , $|\Delta m_{31}^2|$ and mass ordering
- ✓ Combination with atmospheric-neutrino data and other oscillation experiments
⇒ important for reducing degeneracies and improving precision

These measurements will provide new insight into flavor models !

Long-Baseline Experiments: Special vs Typical?

- ✓ Suppose future data find

$$\theta_{23}^{\text{obs}} = \frac{\pi}{4} \pm \sigma_{23}, \quad \sigma_{23} \ll 1.$$

- ✓ **Sharp model** M_{sharp} : $\theta_{23} = \pi/4$ is predicted by a symmetry.

$$\Rightarrow P(D|M_{\text{sharp}}) \simeq \frac{1}{\sqrt{2\pi}\sigma_{23}}.$$

- ✓ **Broad model** M_{broad} : θ_{23} is broadly distributed (e.g., FN model).

$$\Rightarrow P(D|M_{\text{broad}}) \simeq \pi \left(\frac{\pi}{4} \middle| M_{\text{broad}} \right).$$

If $\theta_{23} = \pi/4$ is confirmed with high precision,

\Rightarrow a sharp prediction can beat many broad predictions!

The same argument applies to other special values, such as

$$\delta_{\text{CP}} = 0, \pm\pi/2.$$

Future observations will teach us about the “structure” part of flavor models!

Absolute Neutrino Mass

- ✓ Neutrino masses in the seesaw mechanism

$$(m_\nu)_{ij} = -v_{\text{EW}}^2 (y_D)_{i\alpha} (M_R^{-1})_{\alpha\beta} (y_D^T)_{\beta j},$$

$$(y_D)_{i\alpha} = \epsilon^{|f_{L_i} + f_{\bar{N}_\alpha}|} \circ c_{i\alpha}^D,$$

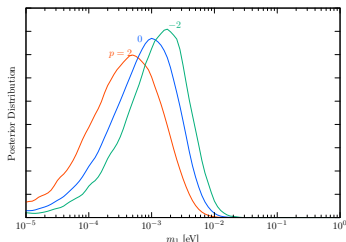
$$(M_R)_{\alpha\beta} = M_0 \epsilon^{|f_{\bar{N}_\alpha} + f_{\bar{N}_\beta}|} \circ c_{\alpha\beta}^R.$$

⇒ The prediction for m_ν depends on the prior for M_0 .

- ✓ For M_0 , the choice of prior distribution is nontrivial.

We adopt a power-law prior: $\pi(M_0) \propto M_0^p$, $-2 \leq p \leq 2$.

- ✓ Posterior distribution in the FN model of $m_1 \Rightarrow m_1 = \mathcal{O}(1)$ meV.



A blue-tilted prior ($p = 2$) tends to predict smaller m_1 .

Absolute Neutrino Mass

- ✓ Effective Majorana mass of ν_e relevant for $0\nu\beta\beta$ (normal ordering)

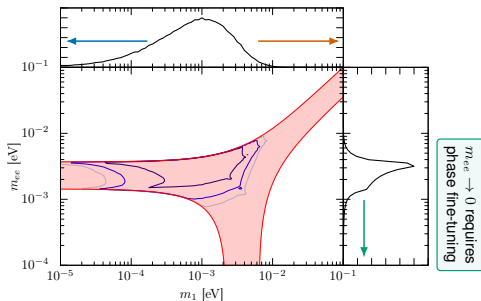
$$m_{ee} := \left| \sum_i U_{ei}^2 m_i \right| = \left| m_1 c_{12}^2 c_{13}^2 + m_2 s_{12}^2 c_{13}^2 e^{2i(\eta_2 - \eta_1)} + m_3 s_{13}^2 e^{-2i(\delta_{CP} + \eta_1)} \right|.$$

$\eta_{1,2}$: Majorana phases in the PMNS matrix.

- ✓ Posterior distribution of m_{ee} in the FN model ($p = 0$)

$P(m_{ee})$ is suppressed
at very small m_1

$P(m_{ee})$ is suppressed
in the quasi-degenerate regime



⇒ **Typicality favors $m_{ee} \gtrsim 1$ meV !**

Summary

- ✓ Typicality provides a quantitative notion of naturalness.
Naturalness \iff large evidence $P(\vec{D}|M)$.
- ✓ **Flavor model = structure \times typicality**
Ex.) FN mechanism gives an overwhelmingly natural fit
- ✓ Future oscillation data will sharpen the test.
 - ✓ IO is challenging for typicality
 - ✓ Special values such as $\theta_{23} = \pi/4$ may also challenge typicality
 \Rightarrow **new insight on the “structure” part of flavor model**
- ✓ Typicality favors the effective majorana mass for $0\nu\beta\beta$ decay

$$m_{ee} \gtrsim 1 \text{ meV.}$$

Future oscillation and $0\nu\beta\beta$ data will provide crucial tests of flavor models.