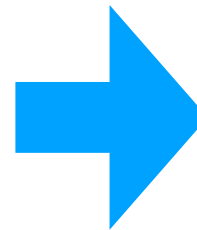
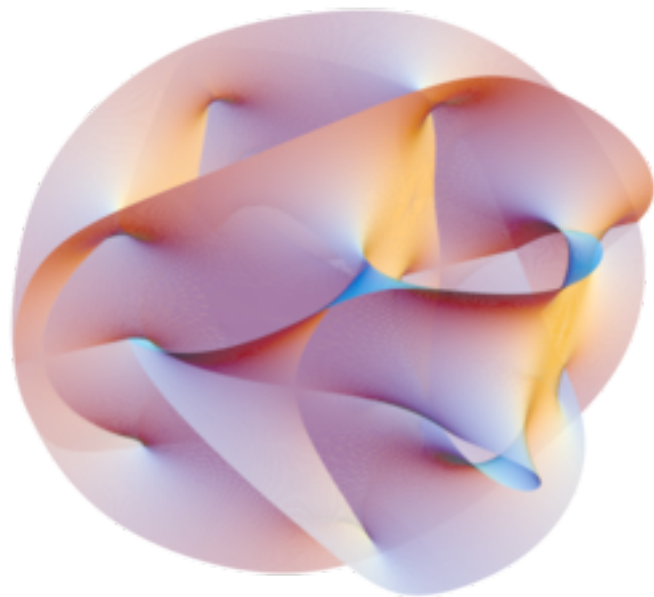

Gravitational wave emission and oscillon formation in string axiverse

Naoya Kitajima

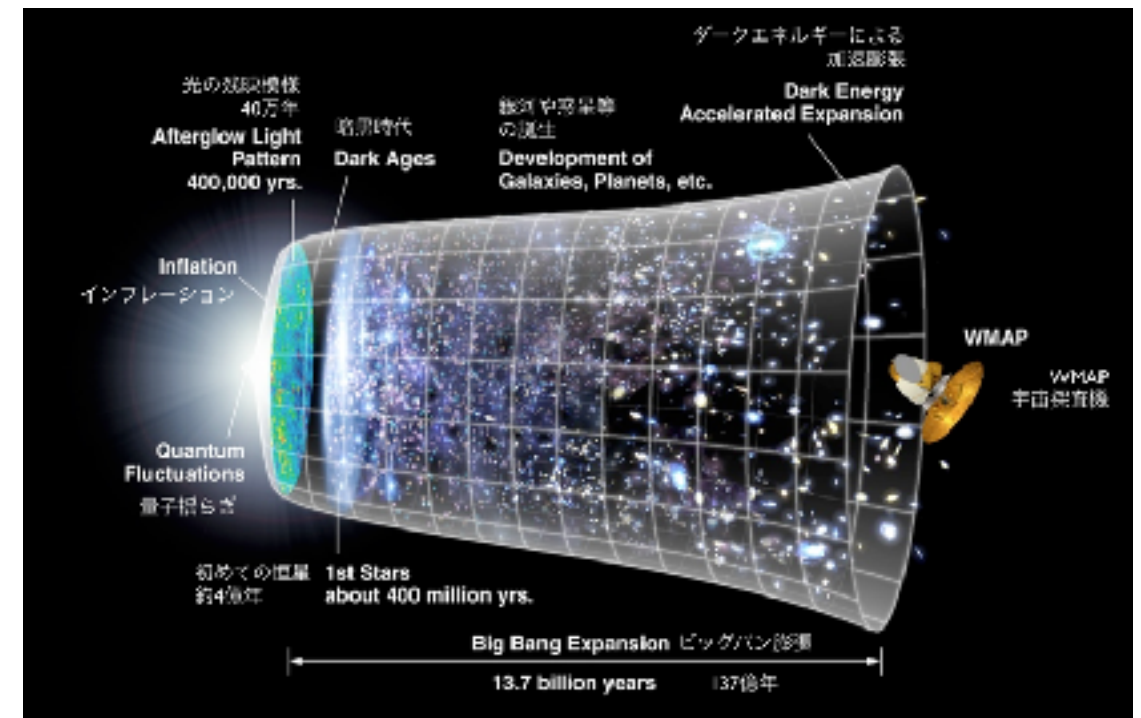


in collaboration with
Jiro Soda (Kobe U.) and Yuko Urakawa (Nagoya U.)

Axiverse



10D string / supergravity theory



4D spacetime

+

Strong CP problem

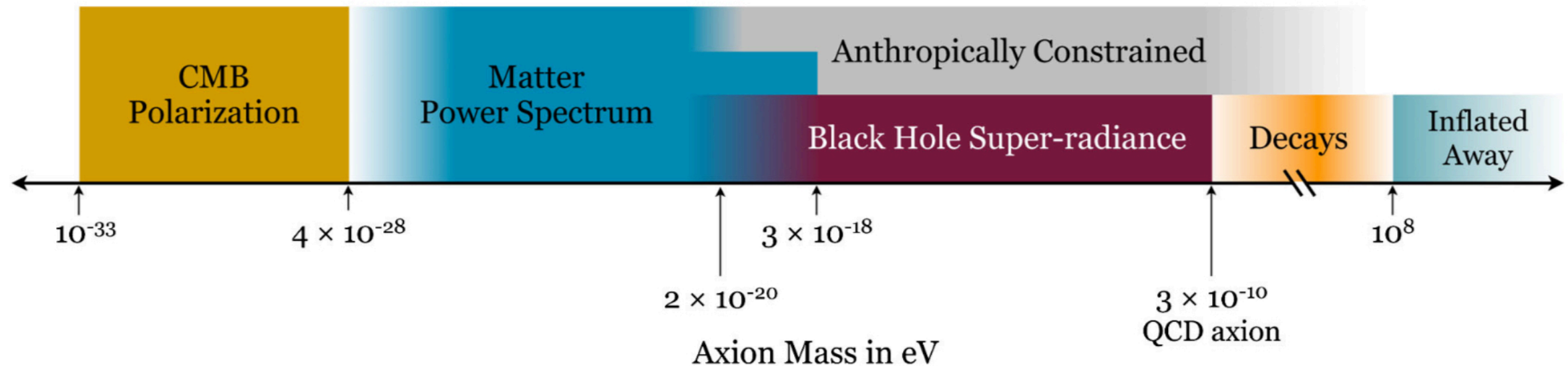
—> QCD axion



String axion

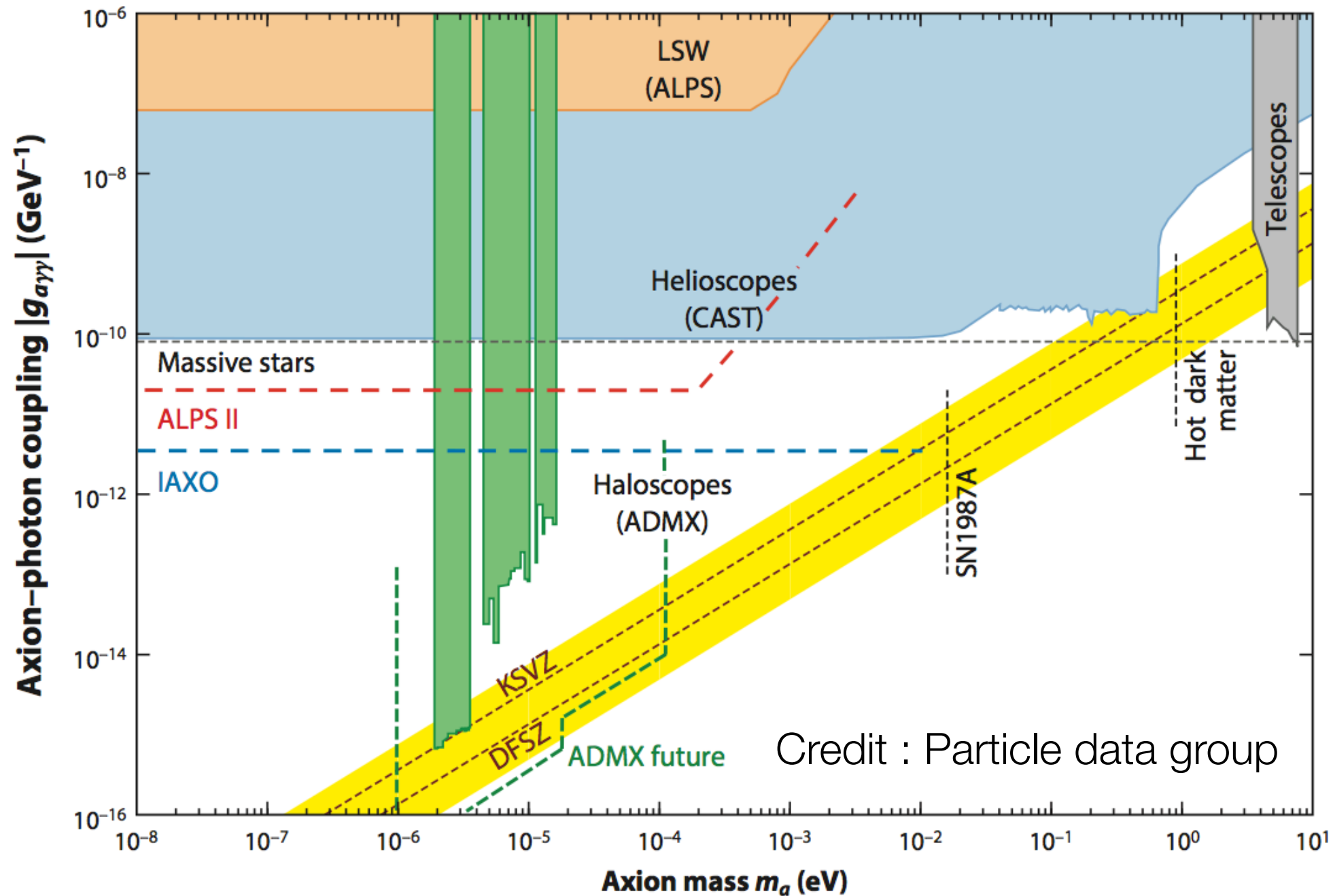
Axiverse

Arvanitaki et al. (2010)



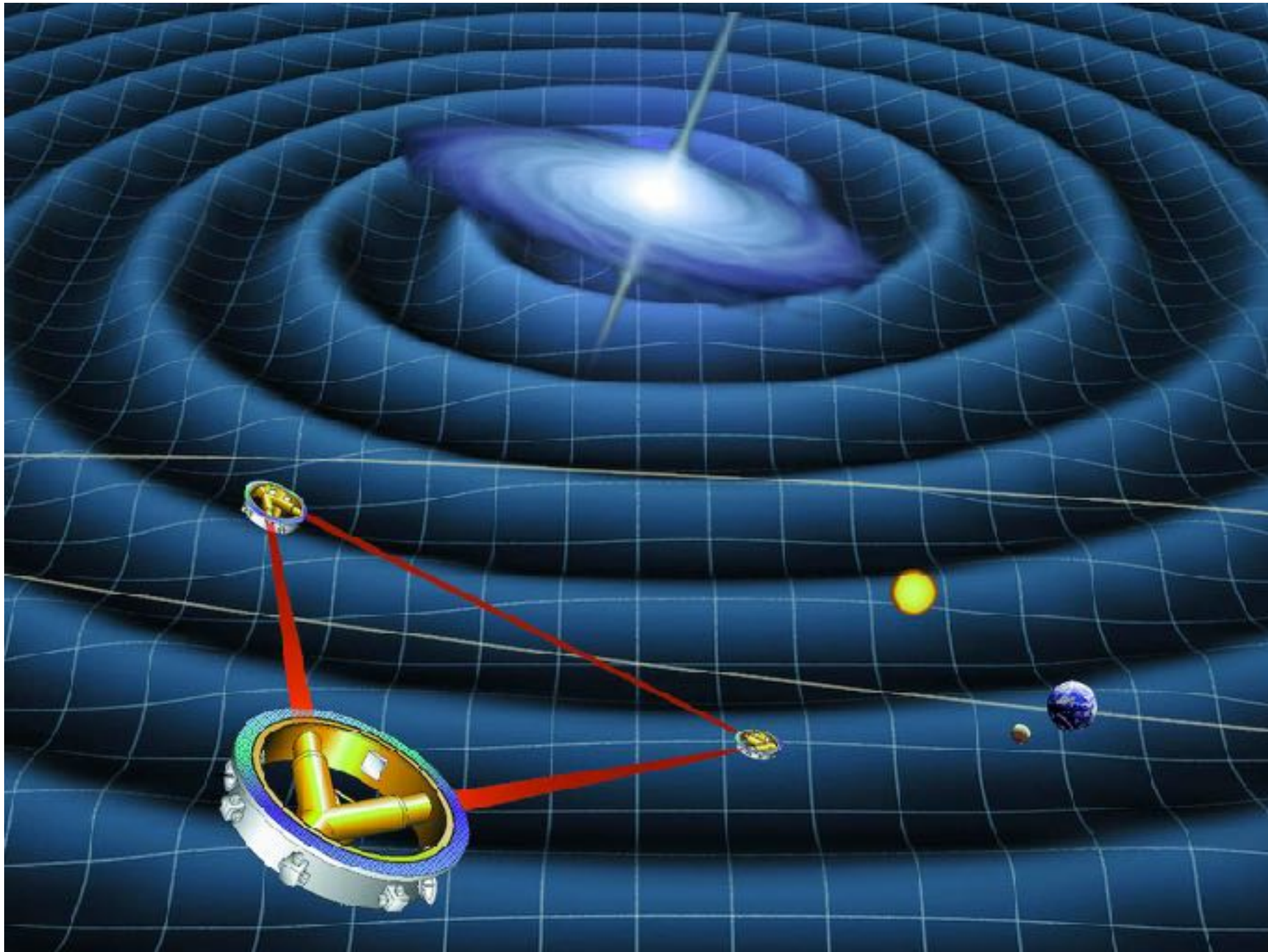
String theory predicts many axions with broad mass range with typical decay constant $f \sim 10^{15-16}$ GeV

Axion detection



Axion-photon interaction is suppressed by $f \sim 10^{16}$ GeV
 —> it is difficult to search string axions by direct detection experiments

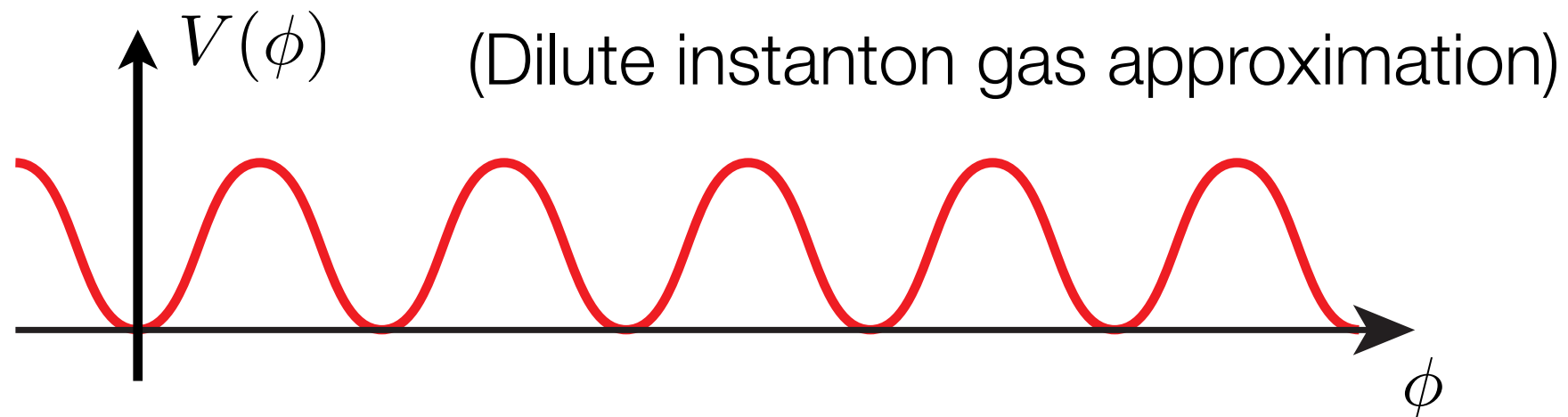
String axion \rightarrow Gravitational wave



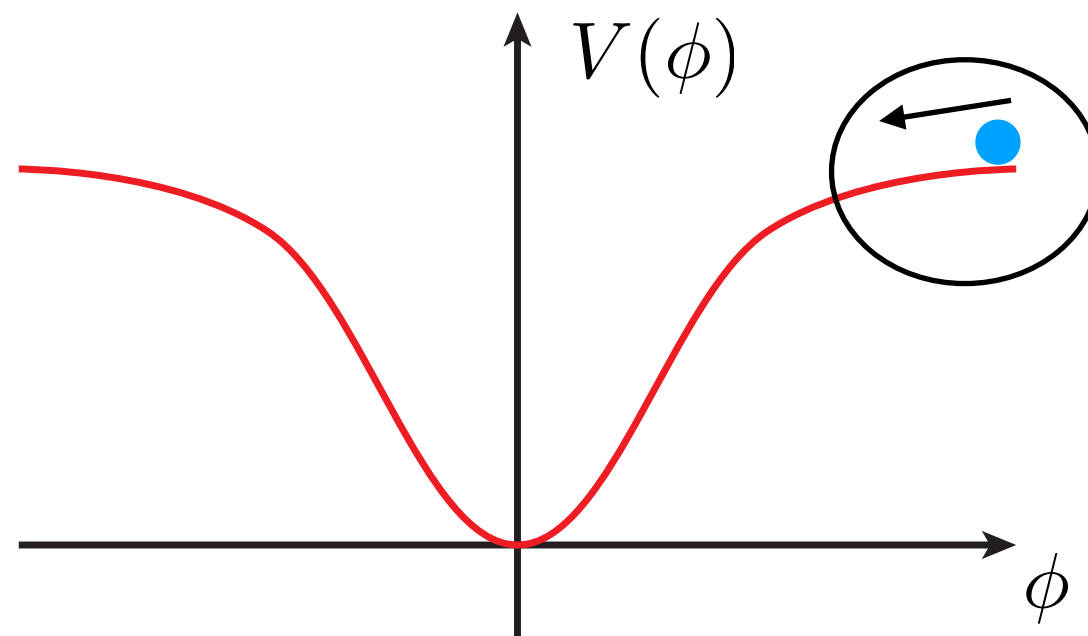
GW signal becomes strong for larger decay constants

Axion potential

Cosine-type potential : $V(\phi) \simeq \Lambda^4 \left(1 - \cos \left(\frac{\phi}{f} \right) \right)$

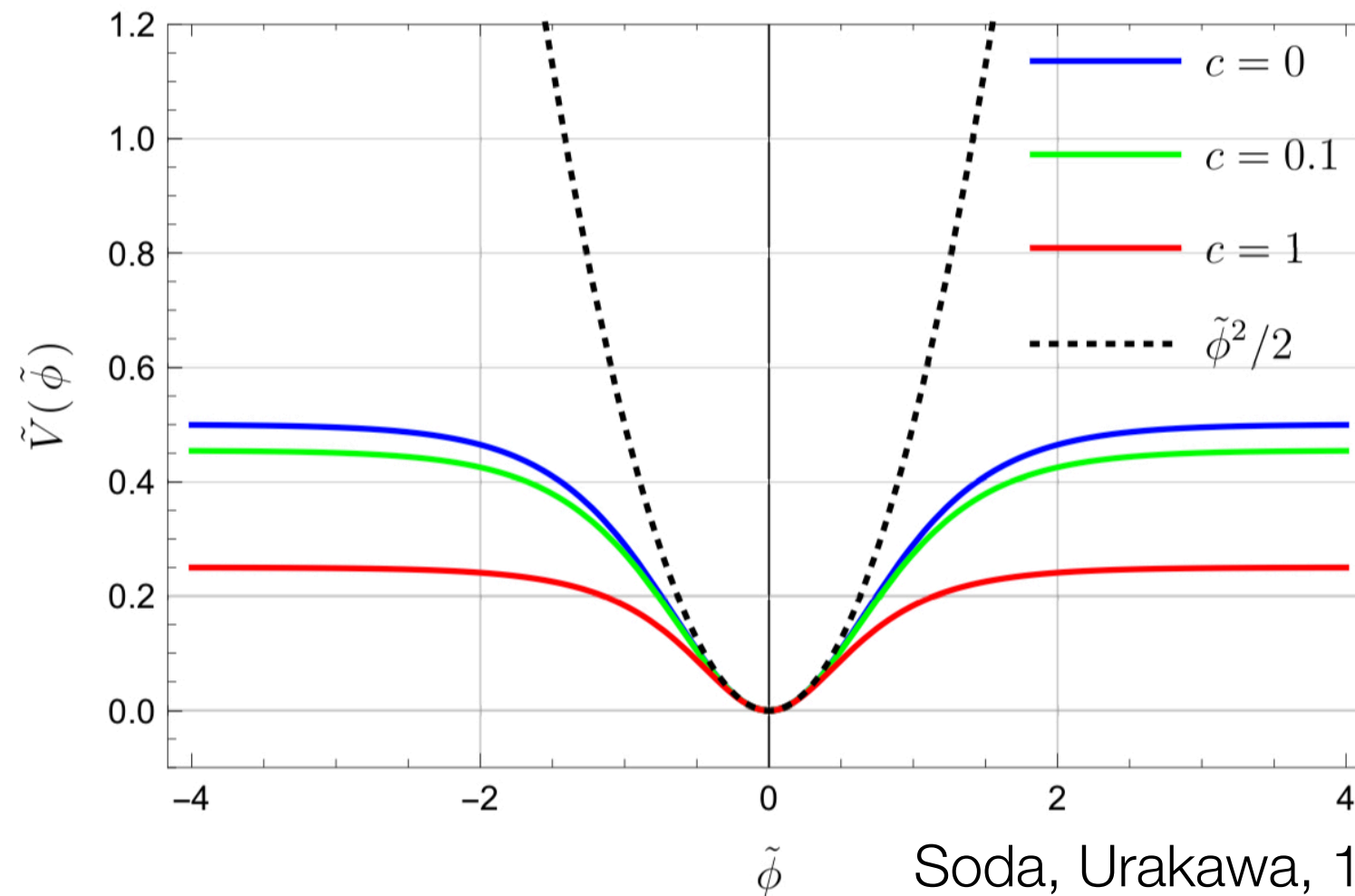


Potential with plateau region



String axion with α -attractor type potential

$$V(\phi) = \frac{m^2 f^2}{2} \frac{\tanh^2(\phi/f)}{1 + c \tanh^2(\phi/f)}$$



Soda, Urakawa, 1710.00305

Self interaction \rightarrow parametric resonance

field fluctuations (with mode $k \sim m$) is amplified exponentially

Kofman, Linde, Starobinsky, 1997

resonance instability

Soda, Urakawa, 1710.00305

$$\ddot{\phi} + 3H\dot{\phi} - \frac{\nabla^2 \phi}{a^2} + \frac{\partial V}{\partial \phi} = 0 \quad \text{with} \quad V(\phi) = \frac{m^2 f^2}{2} \frac{\tanh^2(\phi/f)}{1 + c \tanh^2(\phi/f)}$$

evolution of
nonzero modes

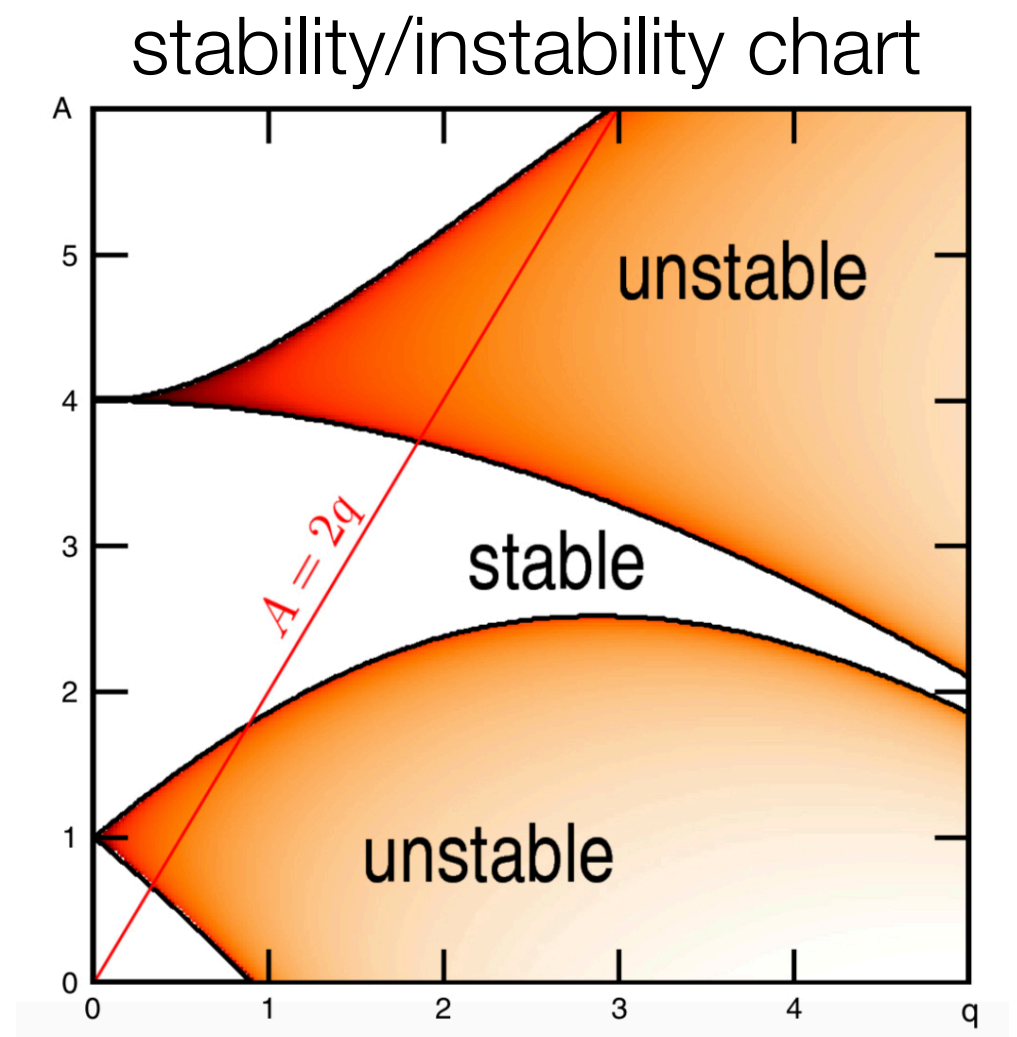
$$\frac{d^2 \varphi_k}{dx^2} + (A - 2q \cos 2x) \varphi_k = 0 \quad (\text{Mathieu equation})$$

$$A = \frac{1}{4} \left[\left(\frac{k}{ma_{\text{osc}}} \right)^2 + 1 - (2 + 3c) \varphi_*^2 \right]$$

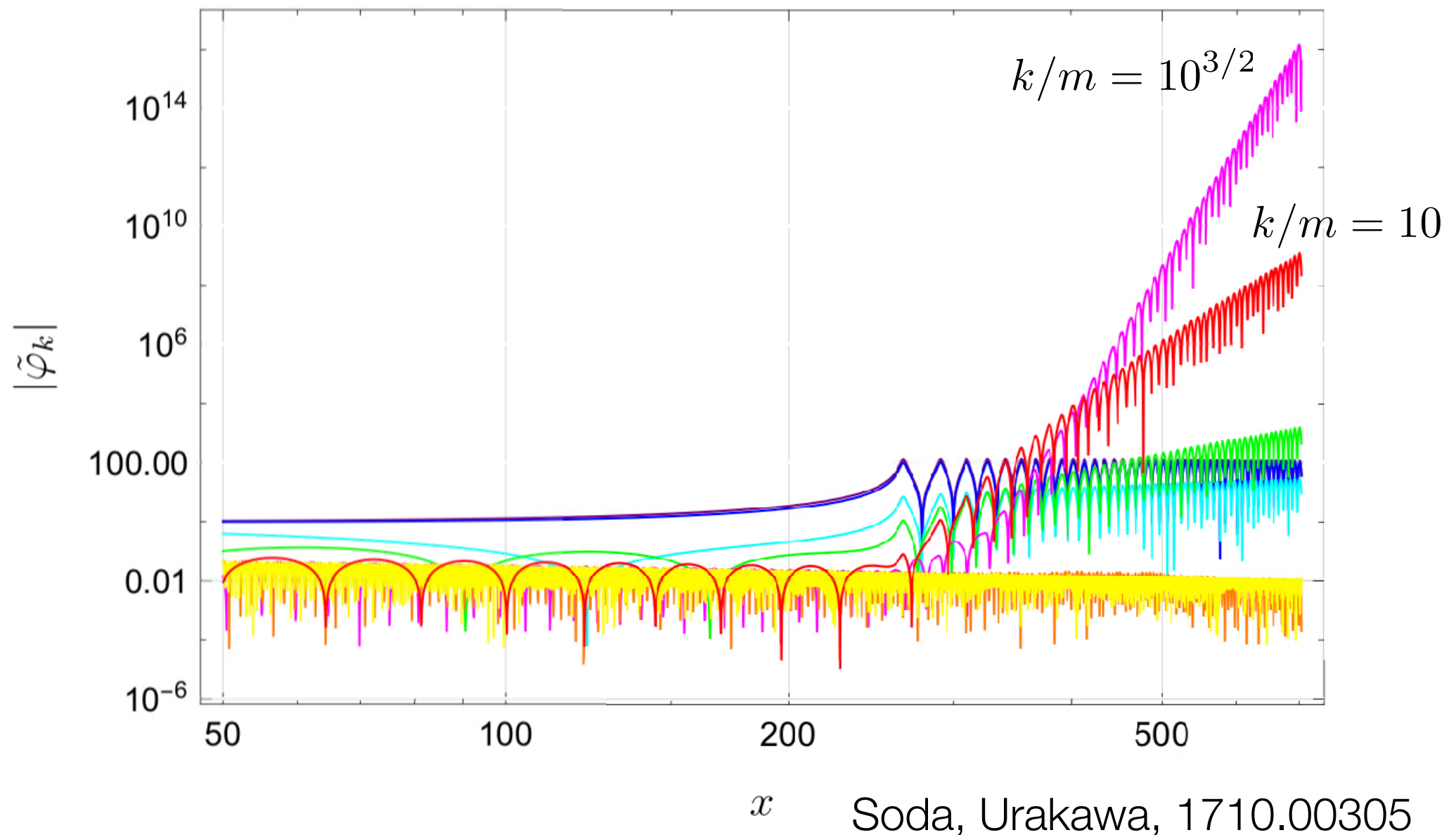
$$q = \frac{2 + 3c}{8} \varphi_*^2, \quad x = m/H$$

amplitude of the zero mode

for k in an instability band $\rightarrow \varphi_k \propto e^{qx/2}$



resonance instability



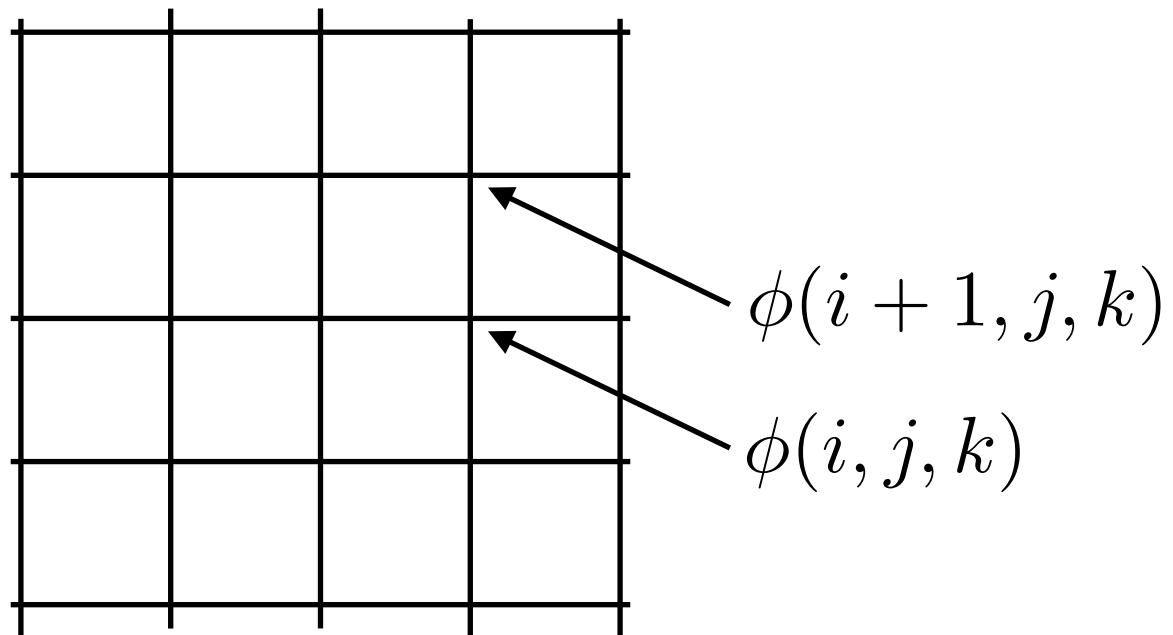
Fluctuations go into the non-linear regime
—> lattice simulation is required

Lattice simulation

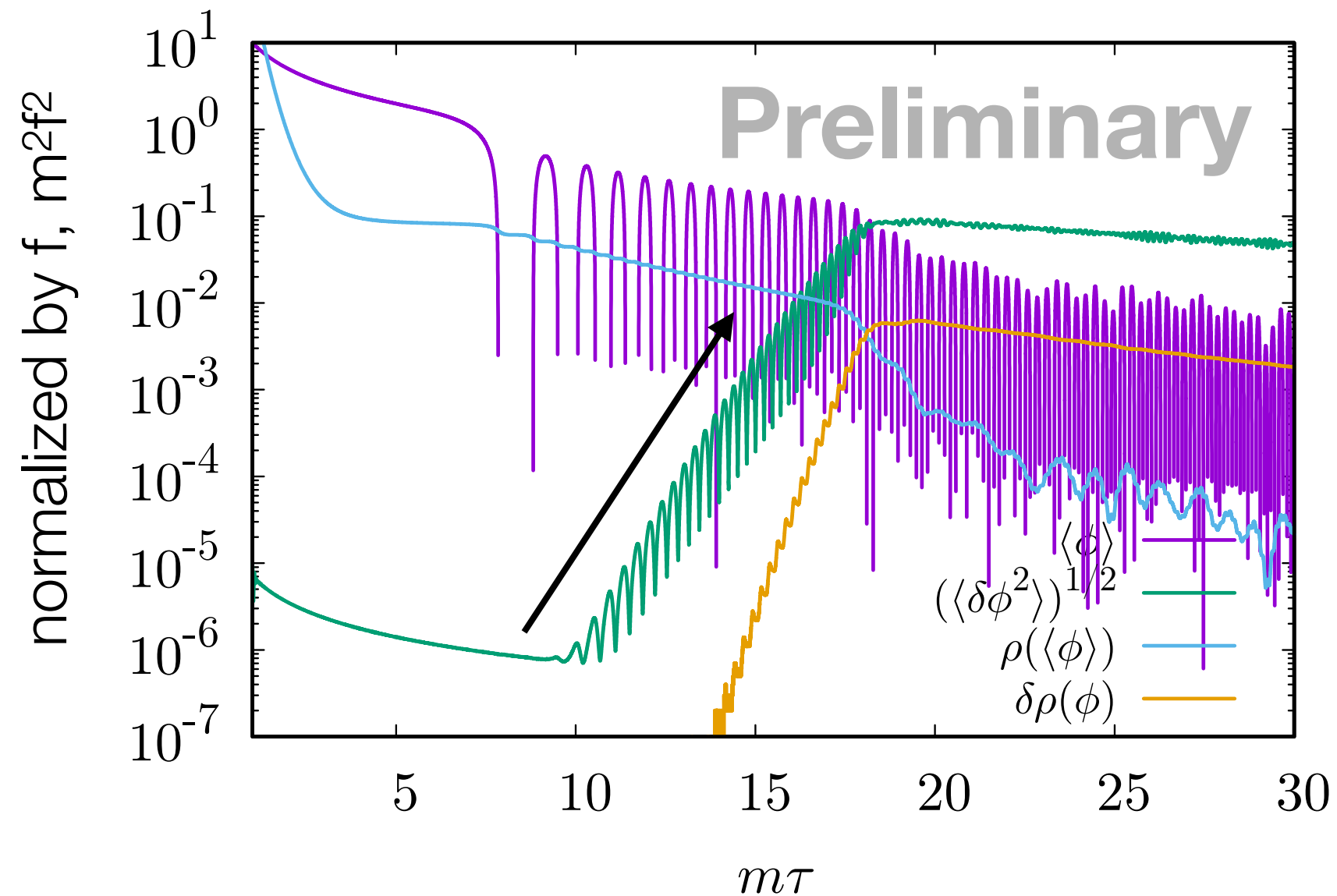
$$\ddot{\phi} + 3H\dot{\phi} - \frac{\nabla^2\phi}{a^2} + \frac{\partial V}{\partial\phi} = 0 \quad \text{with} \quad V(\phi) = \frac{m^2 f^2}{2} \frac{\tanh^2(\phi/f)}{1 + c \tanh^2(\phi/f)}$$

$$\phi(t, \mathbf{x}) \rightarrow \phi(n, i, j, k)$$

$$(t, x, y, z) = \left(t_0 + n\delta t, \frac{iL}{N_{\text{grid}}}, \frac{jL}{N_{\text{grid}}}, \frac{kL}{N_{\text{grid}}} \right)$$

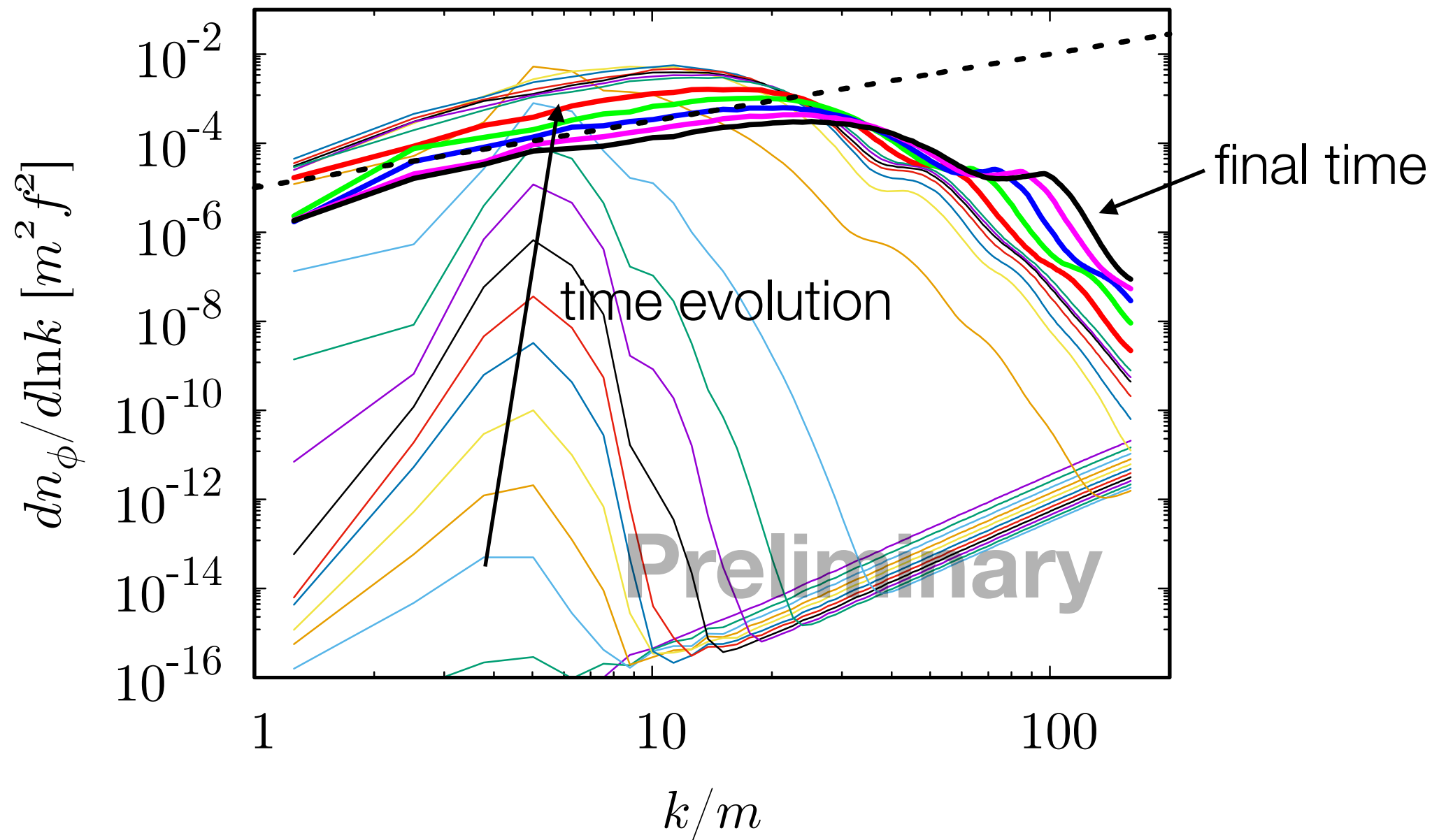


Evolution of the field & the energy density



$$c = 5, \quad \phi_i/f = 10$$

Evolution of the number density spectrum



$$a/a_i = 1 - 50 \quad \text{with} \quad H_i = m$$

Gravitational wave emission

$$ds = -dt^2 + a^2(t)(\delta_{ij} + h_{ij})dx^i dx^j$$



tensor metric perturbation (GW)

evolution equation (Einstein equation) for gravitational waves

$$\ddot{h}_{ij} + 3H\dot{h}_{ij} - \frac{\nabla^2 h_{ij}}{a^2} = 16\pi G \Pi_{ij}^{\text{TT}} \quad \text{with} \quad \Pi_{ij}^{\text{TT}} = \frac{1}{a^2} P_{ij}^{lm} \partial_l \phi \partial_m \phi$$

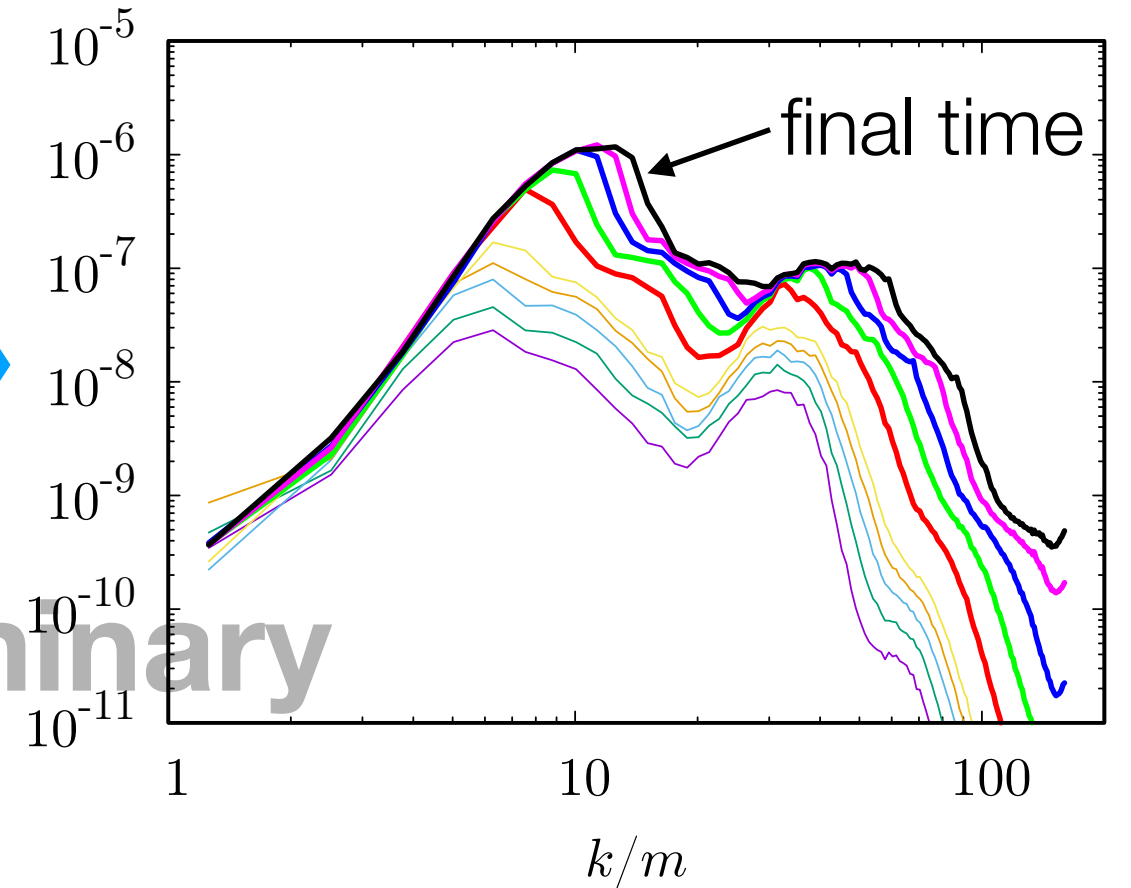
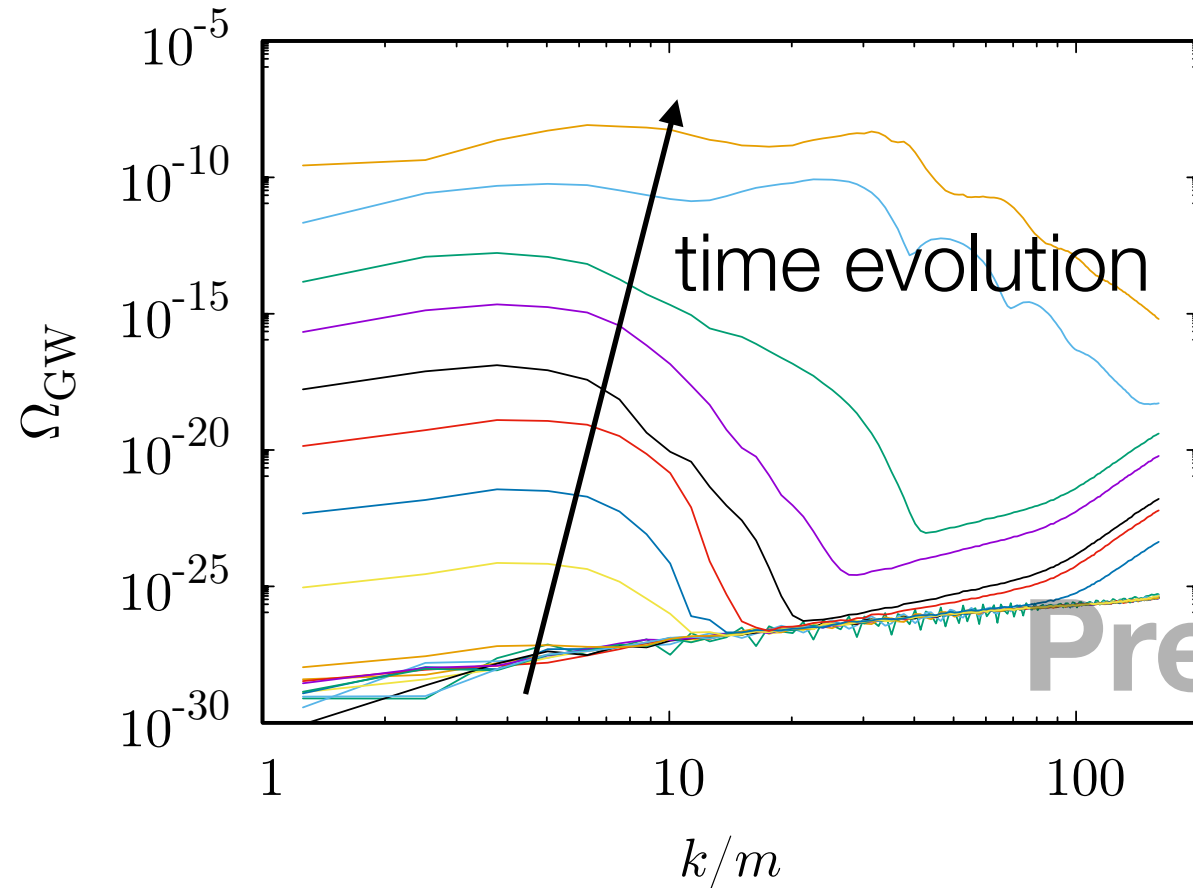


TT projection tensor

Density spectrum of GW

$$\Omega_{\text{GW}}(k) = \frac{1}{\rho_{\text{cr}}} \frac{d\rho_{\text{GW}}}{d \ln k}, \quad \rho_{\text{GW}} = \frac{1}{32\pi G} \langle \dot{h}_{ij} \dot{h}_{ij} \rangle$$

GW density spectrum

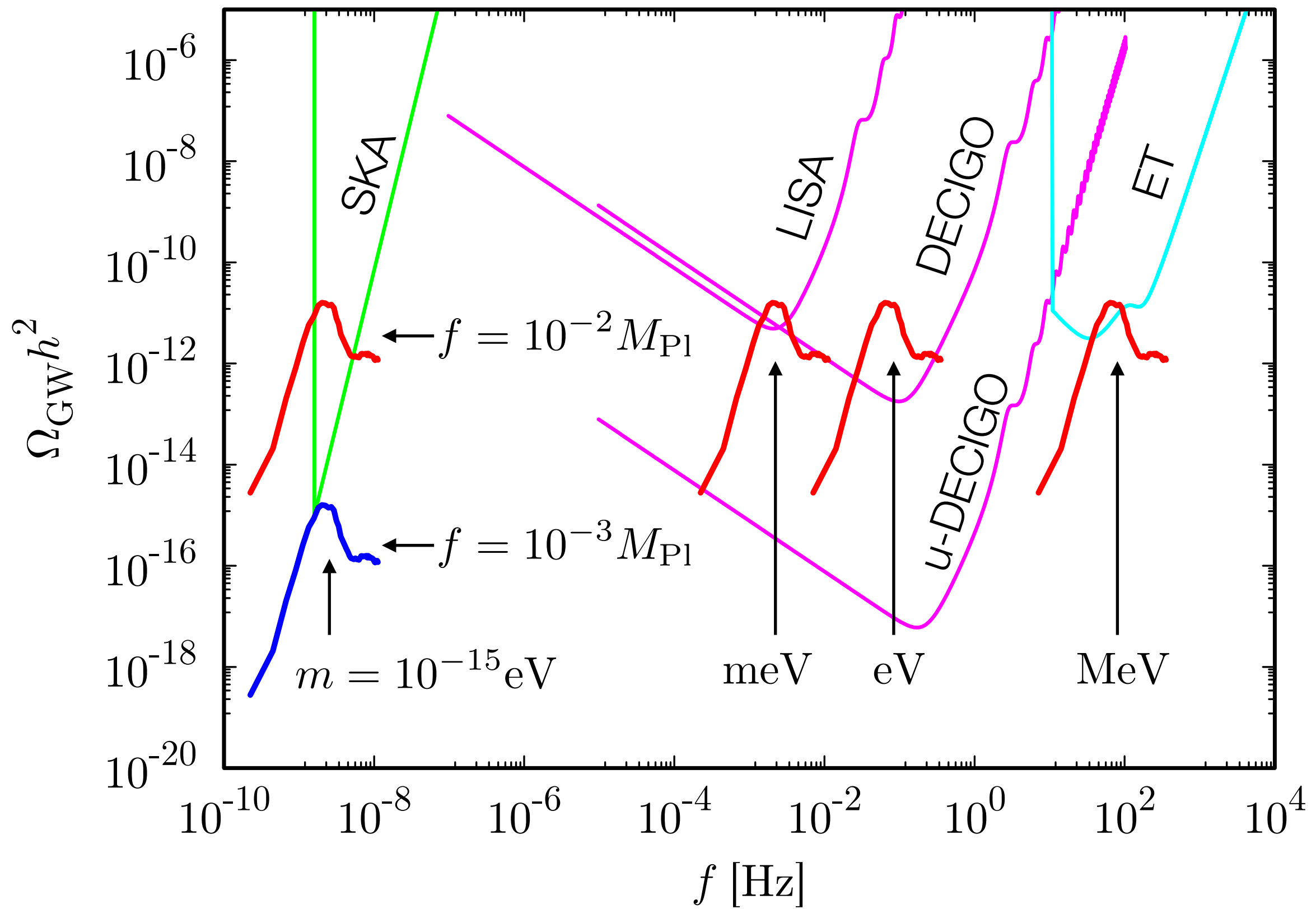


$$c = 5, \quad \phi_i/f = 10, \quad f = 0.01 M_P$$

$$h_{ij} \sim \left(\frac{\phi_{\text{em}}}{M_{\text{Pl}}} \right)^2 \sim \left(\frac{0.1f}{M_{\text{Pl}}} \right)^2 \quad \& \quad \Omega_{\text{GW}} \sim \left(\frac{k}{H_{\text{em}}} \right)^2 h_{ij}^2 \sim 10^2 \left(\frac{f}{M_{\text{Pl}}} \right)^4$$

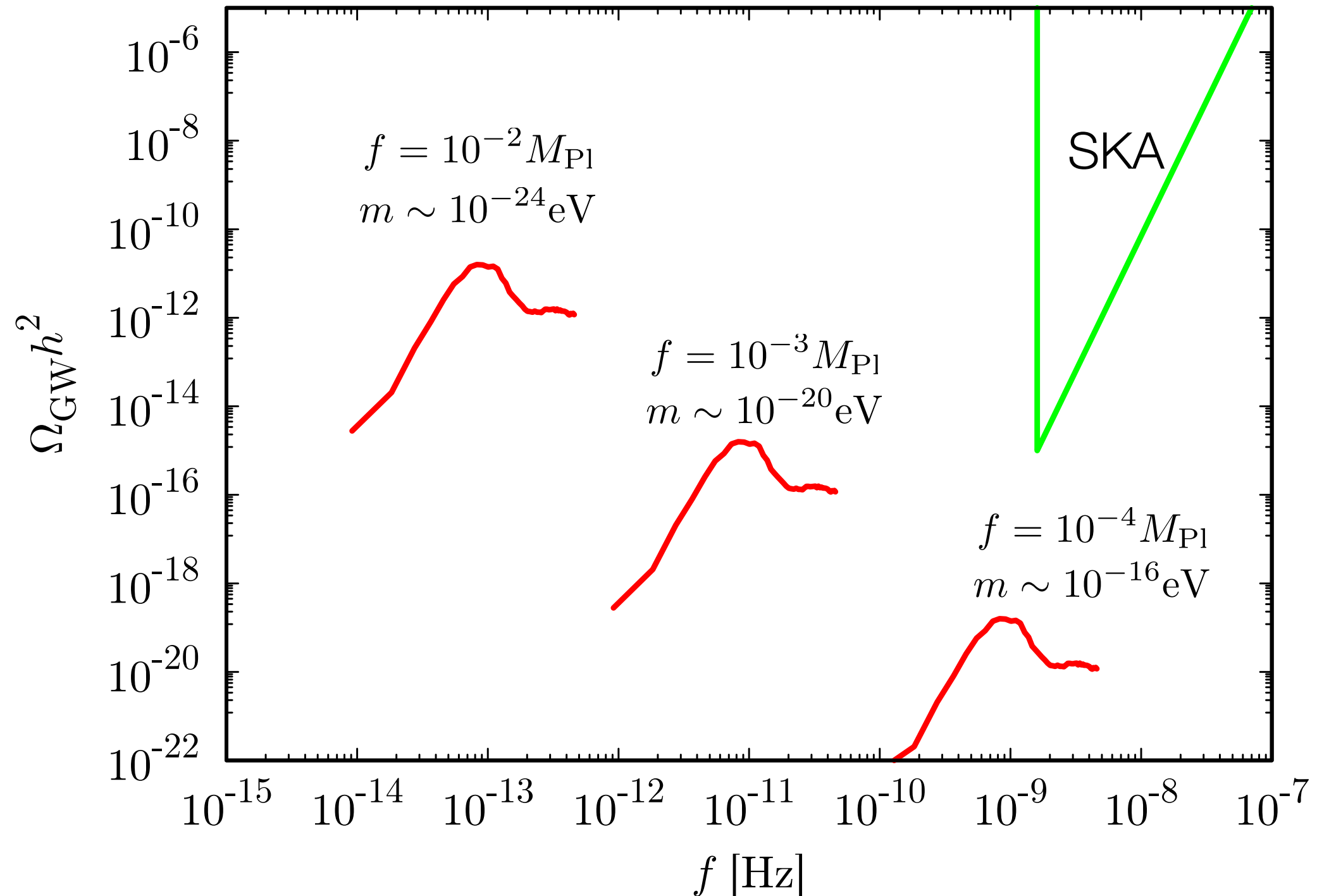
$$\text{with } \phi_{\text{em}} = 0.1f, \quad H_{\text{em}} = 0.01m$$

Detectability

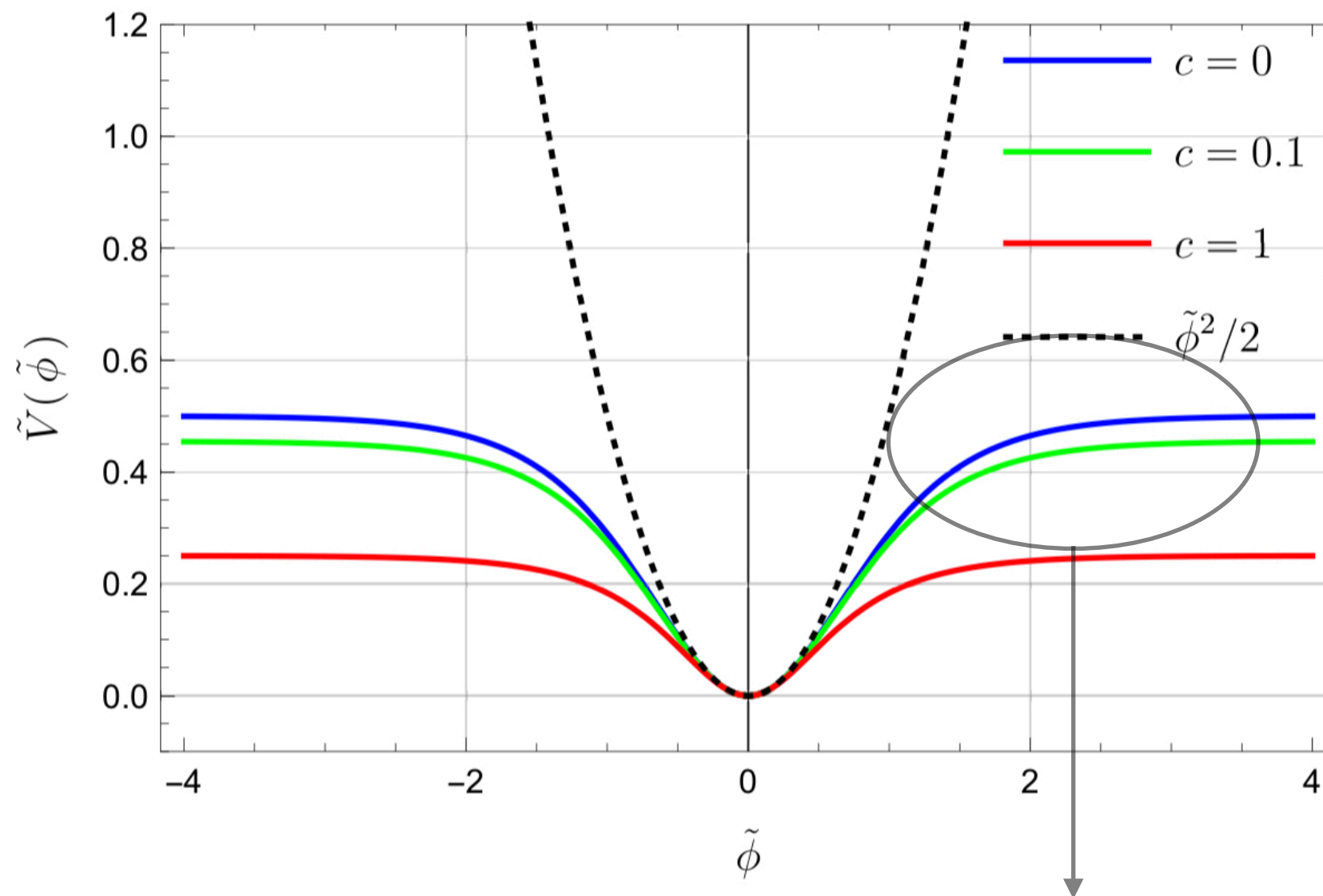


Axion as the present dark matter

$$\Omega_\phi h^2 \sim \left(\frac{g_{*\text{osc}}}{100} \right)^{-1/4} \left(\frac{m}{10^{-24} \text{eV}} \right)^{1/2} \left(\frac{m}{H_{\text{osc}}} \right)^{3/2} \left(\frac{f}{M_{\text{Pl}}} \right)^2$$



Oscillon / I-ball formation

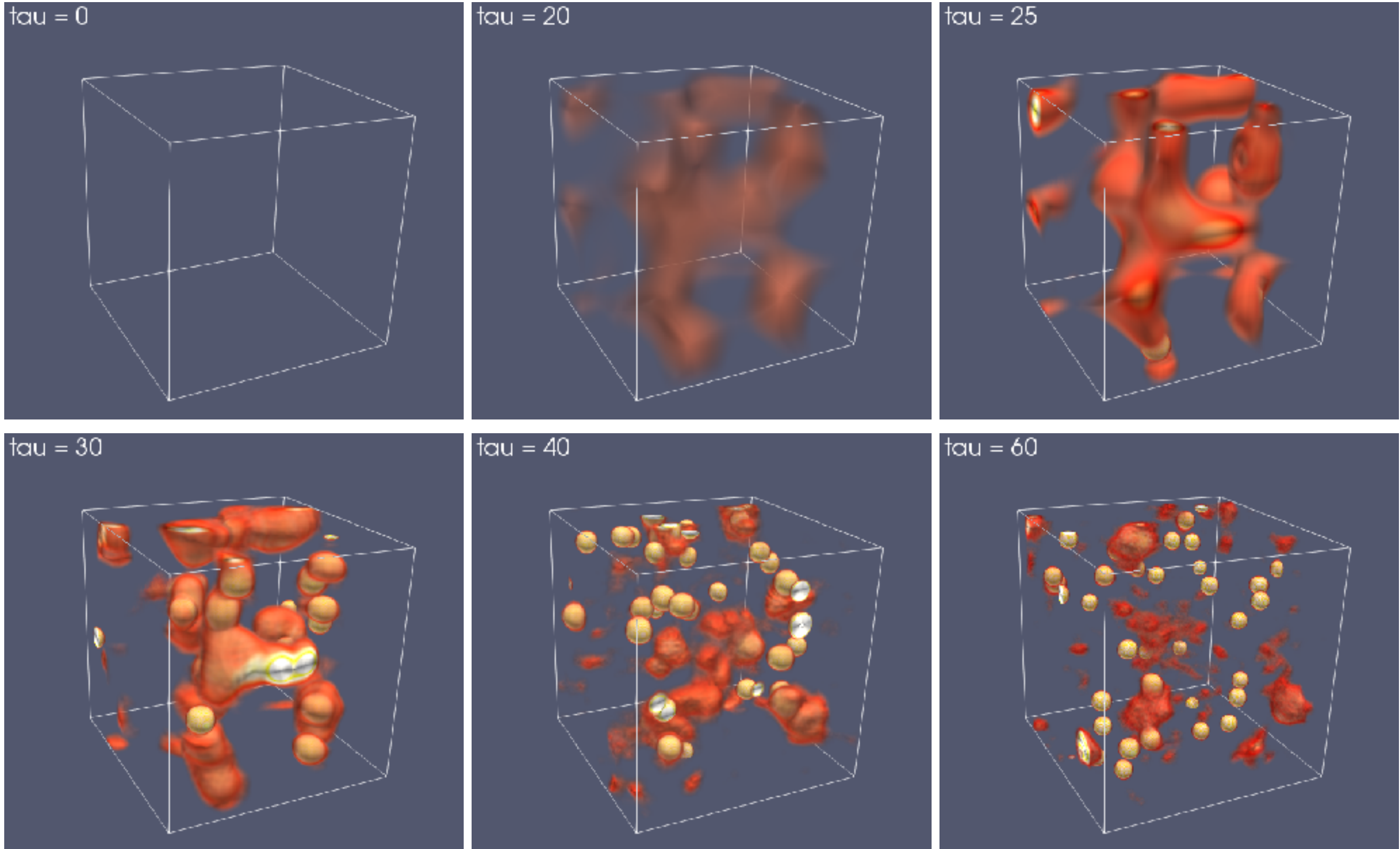


potential becomes flat at large field value
 \rightarrow oscillon formation

Oscillon : quasi-stable, non-topological soliton

Gleiser (1994), Copeland+ (1995), Amin+ (2012) and...

Oscillon / I-ball formation



$$c = 5, \quad \phi_i/f = 10, \quad \rho/\langle\rho\rangle > 2 \quad (\text{red}), \quad 10 \quad (\text{yellow})$$

Summary

- String theory predicts many axions (string axions) with wide mass spectrum
- String axion whose potential has a plateau region (e.g. alpha-attractor model) can show resonance instabilities, leading to GW emission.
- Oscillon formation is followed by the parametric instability.

multi-band GW observation can be a probe of string axions with wide mass range

