

Towards unification of quark and lepton flavors in modular invariance

Morimitsu Tanimoto

Niigata University

February 20, 2020
KEK-ph2020, Tsukuba

1 Introduction

We have a big question since the discovery of muon.

What is the principle to control flavors of quarks and leptons ?

Symmetry Approach: $S_3, A_4, S_4, A_5 \dots$

New approach appears by using Modular group

- Flavor symmetry is originated from the modular invariance.
- Flavor symmetry acts non-linearly (Modular forms).
- Quark / Lepton masses and mixing depend on a modulus τ , which is stabilized by some unknown mechanism.

Talk by T. Yoshida, H. Uchida

2 Finite modular groups

$$S : \tau \longrightarrow -\frac{1}{\tau}, \quad \text{Duality}$$

$$T : \tau \longrightarrow \tau + 1. \quad \text{Discrete shift symmetry}$$

$$S^2 = 1, \quad (ST)^3 = 1.$$

Modular group

$$\Gamma \simeq \{S, T \mid S^2 = \mathbb{I}, (ST)^3 = \mathbb{I}\}$$

Modular group has series of subgroups $\Gamma(N)$ level N

Imposing congruence condition $\Gamma(N) = \left\{ \begin{pmatrix} a & b \\ c & d \end{pmatrix} \in SL(2, Z), \begin{pmatrix} a & b \\ c & d \end{pmatrix} \equiv \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} \pmod{N} \right\}$

$$ad - bc = 1$$

called principal congruence subgroups

generate discrete group

$\Gamma_N \equiv \Gamma / \Gamma(N)$ quotient group finite group

$$\Gamma_N \simeq \{S, T \mid S^2 = \mathbb{I}, (ST)^3 = \mathbb{I}, T^N = \mathbb{I}\}$$

$$\Gamma_2 \simeq S_3$$

$$\Gamma_3 \simeq A_4$$

$$\Gamma_4 \simeq S_4$$

$$\Gamma_5 \simeq A_5$$

We can consider effective theories with Γ_N symmetry.

$$\mathcal{L}_{\text{eff}} \in f(\tau) \phi^{(1)} \dots \phi^{(n)} \quad f(\tau), \phi^{(I)}: \text{non-trivial rep. of } \Gamma_N$$

$$f(\gamma\tau) = (c\tau + d)^k f(\tau), \quad \gamma \in \Gamma(N)$$

Modular forms have been explicitly given for some cases with fixed weight k .

On the other hand, chiral superfields are not modular forms and we have no restriction on the possible value of weight k_I , a priori.

Modular transformation of chiral superfields in MSSM

$$\phi^{(I)} \rightarrow (c\tau + d)^{-k_I} \rho^{(I)}(\gamma) \phi^{(I)}$$

Modular weight

Representation matrix

Modular forms with **weight 2** are given by using Dedekind eta-function.

Dedekind eta-function

$$\eta(\tau) = q^{1/24} \prod_{n=1}^{\infty} (1 - q^n) \quad q = e^{2\pi i \tau}$$

$$\eta(-1/\tau) = \sqrt{-i\tau} \eta(\tau), \quad \eta(\tau + 1) = e^{i\pi/12} \eta(\tau)$$

$$Y(\alpha, \beta, \gamma, \delta | \tau) = \frac{d}{d\tau} (\alpha \log \eta(\tau/3) + \beta \log \eta((\tau + 1)/3) + \gamma \log \eta((\tau + 2)/3) + \delta \log \eta(3\tau))$$

$$\alpha + \beta + \gamma + \delta = 0$$

$$S : \tau \longrightarrow -\frac{1}{\tau},$$

$$T : \tau \longrightarrow \tau + 1.$$

$$S : Y(\alpha, \beta, \gamma, \delta | \tau) \longrightarrow \tau^2 Y(\delta, \gamma, \beta, \alpha | \tau),$$

$$T : Y(\alpha, \beta, \gamma, \delta | \tau) \longrightarrow Y(\gamma, \alpha, \beta, \delta | \tau).$$

Consider level $N=3$, $T^3=1$: A_4 group

For triplet

$$\rho(S) = \frac{1}{3} \begin{pmatrix} -1 & 2 & 2 \\ 2 & -1 & 2 \\ 2 & 2 & -1 \end{pmatrix}, \quad \rho(T) = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \omega & 0 \\ 0 & 0 & \omega^2 \end{pmatrix},$$

A_4 triplet of modular forms with weight 2

S transformation

$$\begin{pmatrix} Y_1(-1/\tau) \\ Y_2(-1/\tau) \\ Y_3(-1/\tau) \end{pmatrix} = \tau^2 \rho(S) \begin{pmatrix} Y_1(\tau) \\ Y_2(\tau) \\ Y_3(\tau) \end{pmatrix},$$

T transformation

$$\begin{pmatrix} Y_1(\tau+1) \\ Y_2(\tau+1) \\ Y_3(\tau+1) \end{pmatrix} = \rho(T) \begin{pmatrix} Y_1(\tau) \\ Y_2(\tau) \\ Y_3(\tau) \end{pmatrix}.$$

$$\begin{aligned} Y_1(\tau) &= \frac{i}{2\pi} \left(\frac{\eta'(\tau/3)}{\eta(\tau/3)} + \frac{\eta'((\tau+1)/3)}{\eta((\tau+1)/3)} + \frac{\eta'((\tau+2)/3)}{\eta((\tau+2)/3)} - \frac{27\eta'(3\tau)}{\eta(3\tau)} \right), \\ Y_2(\tau) &= \frac{-i}{\pi} \left(\frac{\eta'(\tau/3)}{\eta(\tau/3)} + \omega^2 \frac{\eta'((\tau+1)/3)}{\eta((\tau+1)/3)} + \omega \frac{\eta'((\tau+2)/3)}{\eta((\tau+2)/3)} \right), \\ Y_3(\tau) &= \frac{-i}{\pi} \left(\frac{\eta'(\tau/3)}{\eta(\tau/3)} + \omega \frac{\eta'((\tau+1)/3)}{\eta((\tau+1)/3)} + \omega^2 \frac{\eta'((\tau+2)/3)}{\eta((\tau+2)/3)} \right), \end{aligned}$$

$$\begin{aligned} |q| \ll 1 \quad Y_1(\tau) &= 1 + 12q + 36q^2 + 12q^3 + \dots, & q &= e^{2\pi i\tau} \\ Y_2(\tau) &= -6q^{1/3}(1 + 7q + 8q^2 + \dots), \\ Y_3(\tau) &= -18q^{2/3}(1 + 2q + 5q^2 + \dots). \end{aligned} \quad \boxed{Y_2^2 + 2Y_1Y_3 = 0}$$

3 Modular A_4 invariance as flavor symmetry

Consider both quarks and leptons in modular symmetry.

A_4 assignments: left-handed doublet $\mathbf{3}$, right-handed singlets $\mathbf{1}$, $\mathbf{1}'$, $\mathbf{1}''$

Quarks

$$\left\{ \begin{array}{l} w_u = \alpha_u u^c H_u Y_3^{(2)} Q + \beta_u c^c H_u Y_3^{(2)} Q + \gamma_u t^c H_u Y_3^{(2)} Q \\ w_d = \alpha_d d^c H_d Y_3^{(2)} Q + \beta_d s^c H_d Y_3^{(2)} Q + \gamma_d b^c H_d Y_3^{(2)} Q \end{array} \right.$$

Charged Lepton

$$w_e = \alpha_e e^c H_d Y_3^{(2)} L + \beta_e \mu^c H_d Y_3^{(2)} L + \gamma_e \tau^c H_d Y_3^{(2)} L$$

$$M_q = \begin{pmatrix} \alpha_q & 0 & 0 \\ 0 & \beta_q & 0 \\ 0 & 0 & \gamma_q \end{pmatrix} \begin{pmatrix} Y_1 & Y_3 & Y_2 \\ Y_2 & Y_1 & Y_3 \\ Y_3 & Y_2 & Y_1 \end{pmatrix}_{RL}$$

for both up- and down-quarks

Simple mass matrix

Can Modular forms reproduce
CKM by fixing modulus τ ?

Real parameters α_q , β_q , γ_q are responsible for quark mass hierarchy.

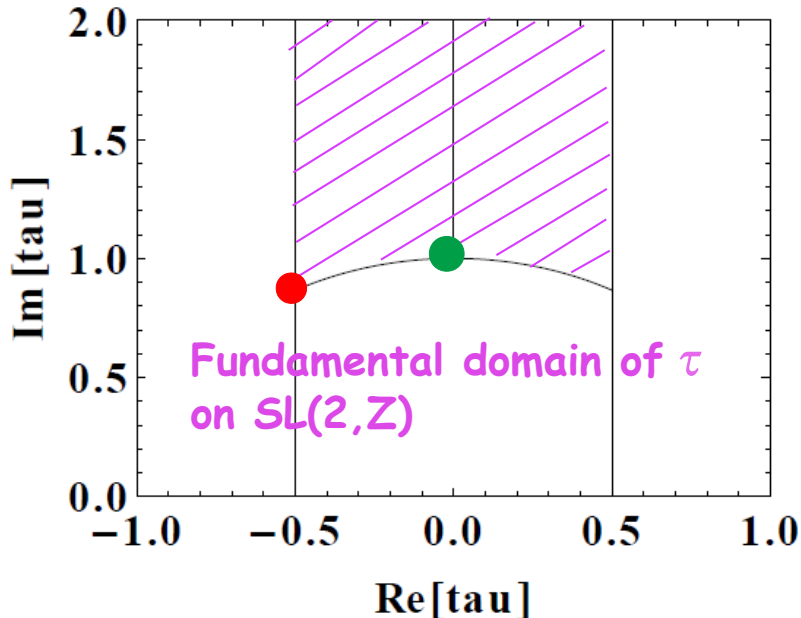
● What is a Principle of fixing modulus τ ?

★ Modular stabilization (non-perturbative effect, model dependent)

Some models indicate the potential minimum at the boundary of fundamental domain.

Talk by H. Uchida

★ Fixed point of τ Residual symmetry



● $Z_2: \tau = i$ ● $Z_3: \tau = -1/2 + \sqrt{3}/2i$
S symmetry *ST* symmetry

P.P.Novichkov, J.T.Penedo, S.T.Petcov, A.V.Titov,
 JHEP04 (2019) 005, arXiv:1811.04933
 P.P.Novichkov, S.T.Petcov and M.T,
 PLB 793 (2019) 247, arXiv:1812.11289

Fixed points of τ

T ($\tau \rightarrow \tau + 1$) preserved : $\langle \tau \rangle = \infty$ i ($q=0$) (Y_1, Y_2, Y_3) = (1, 0, 0)

$Z_3: \{1, T, T^2\}$

S ($\tau \rightarrow -1/\tau$) preserved : $\langle \tau \rangle = i$ ($q=e^{-2\pi}$) (Y_1, Y_2, Y_3) = $Y_1(i)$ (1, $1-\sqrt{3}$, $-2+\sqrt{3}$)

$Z_2: \{1, S\}$

$$Y_1(\tau) = 1 + 12q + 36q^2 + 12q^3 + \dots,$$

$$Y_2(\tau) = -6q^{1/3}(1 + 7q + 8q^2 + \dots),$$

$$Y_3(\tau) = -18q^{2/3}(1 + 2q + 5q^2 + \dots).$$

$$q = e^{2\pi i \tau}$$

$$Z_3: T^\dagger M_q^\dagger M_q T = M_q^\dagger M_q \Rightarrow [T, M_q^\dagger M_q] = 0$$

$$Z_2: S^\dagger M_q^\dagger M_q S = M_q^\dagger M_q \Rightarrow [S, M_q^\dagger M_q] = 0$$

Mixing matrices diagonalise $M_q^\dagger M_q$ also diagonalize T and S, respectively !

$$\rho(S) = \frac{1}{3} \begin{pmatrix} -1 & 2 & 2 \\ 2 & -1 & 2 \\ 2 & 2 & -1 \end{pmatrix}, \quad \rho(T) = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \omega & 0 \\ 0 & 0 & \omega^2 \end{pmatrix},$$

Eigenvalues (1, -1, -1)

Define the new basis of generators, \hat{S} and \hat{T} by a unitary transformation as:

$$\hat{S} = USU^\dagger, \quad \hat{T} = UTU^\dagger \quad \longrightarrow \quad \hat{M}_{RL} = M_{RL}U^\dagger.$$

Since one can take the diagonal basis of S or T, both $M_q^\dagger M_q$ (q=u,d) could be diagonal if there is a residual symmetry Z_2 or Z_3 of A_4 .

Hierarchical flavor structure is realized around $\tau = i$ or ∞ in general !

$$M_q = \begin{pmatrix} \alpha_q & 0 & 0 \\ 0 & \beta_q & 0 \\ 0 & 0 & \gamma_q \end{pmatrix} \begin{pmatrix} Y_1 & Y_3 & Y_2 \\ Y_2 & Y_1 & Y_3 \\ Y_3 & Y_2 & Y_1 \end{pmatrix}_{RL}$$

at $\langle \tau \rangle = i$ $(Y_1, Y_2, Y_3) = Y_1(i) (1, 1 - \sqrt{3}, -2 + \sqrt{3})$

at $\langle \tau \rangle = i \infty$ (q=0) $(Y_1, Y_2, Y_3) = (1, 0, 0)$

$$M_q^\dagger M_q = \begin{pmatrix} 0 & 0 & 0 \\ 0 & \times & \times \\ 0 & \times & \times \end{pmatrix} \begin{pmatrix} \tilde{\alpha}_q^2 & 0 & 0 \\ 0 & \tilde{\beta}_q^2 & 0 \\ 0 & 0 & \tilde{\gamma}_q^2 \end{pmatrix}$$

Too simple !

Vub is inconsistent with data even if τ is scanned.

Let us consider Modular forms with higher weights $k=4, 6 \dots$

of modular forms is $k+1$

Weight 2
3 Modular forms

$$Y_3^{(2)} = \begin{pmatrix} Y_1 \\ Y_2 \\ Y_3 \end{pmatrix}$$

Modular forms with higher weights are constructed by the tensor product of modular forms of weight 2

$$\begin{aligned} \begin{pmatrix} a_1 \\ a_2 \\ a_3 \end{pmatrix}_3 \otimes \begin{pmatrix} b_1 \\ b_2 \\ b_3 \end{pmatrix}_3 &= (a_1b_1 + a_2b_3 + a_3b_2)_1 \oplus (a_3b_3 + a_1b_2 + a_2b_1)_{1'} \\ &\oplus (a_2b_2 + a_1b_3 + a_3b_1)_{1''} \\ &\oplus \frac{1}{3} \begin{pmatrix} 2a_1b_1 - a_2b_3 - a_3b_2 \\ 2a_3b_3 - a_1b_2 - a_2b_1 \\ 2a_2b_2 - a_1b_3 - a_3b_1 \end{pmatrix}_3 \oplus \frac{1}{2} \begin{pmatrix} a_2b_3 - a_3b_2 \\ a_1b_2 - a_2b_1 \\ a_3b_1 - a_1b_3 \end{pmatrix}_3 \end{aligned}$$

$$1 \otimes 1 = 1, \quad 1' \otimes 1' = 1'', \quad 1'' \otimes 1'' = 1', \quad 1' \otimes 1'' = 1.$$

J.T.Penedo, S.T.Petcov, Nucl.Phys.B939(2019)292

Weight 4
5 Modular forms

$$Y_1^{(4)} = Y_1^2 + 2Y_2Y_3, \quad Y_{1'}^{(4)} = Y_3^2 + 2Y_1Y_2, \quad Y_{1''}^{(4)} = Y_2^2 + 2Y_1Y_3 = 0,$$

$$Y_3^{(4)} = \begin{pmatrix} Y_1^2 - Y_2Y_3 \\ Y_3^2 - Y_1Y_2 \\ Y_2^2 - Y_1Y_3 \end{pmatrix},$$



Weight 6
7 Modular forms

$$Y_1^{(6)} = Y_1^3 + Y_2^3 + Y_3^3 - 3Y_1Y_2Y_3,$$

$$Y_3^{(6)} \equiv \begin{pmatrix} Y_1^{(6)} \\ Y_2^{(6)} \\ Y_3^{(6)} \end{pmatrix} = \begin{pmatrix} Y_1^3 + 2Y_1Y_2Y_3 \\ Y_1^2Y_2 + 2Y_2^2Y_3 \\ Y_1^2Y_3 + 2Y_3^2Y_2 \end{pmatrix}, \quad Y_{3'}^{(6)} \equiv \begin{pmatrix} Y_1'^{(6)} \\ Y_2'^{(6)} \\ Y_3'^{(6)} \end{pmatrix} = \begin{pmatrix} Y_3^3 + 2Y_1Y_2Y_3 \\ Y_3^2Y_1 + 2Y_1^2Y_2 \\ Y_3^2Y_2 + 2Y_2^2Y_1 \end{pmatrix}$$

4 A lesson in Quarks and Leptons

Quark Sector

	Q	(d^c, s^c, b^c)	(u^c, c^c, t^c)	$H_{u,d}$	$Y_3^{(2)}$	$Y_3^{(6)}, Y_{3'}^{(6)}$
$SU(2)$	2	1	1	2	1	1
A_4	3	$(1, 1'', 1')$	$(1, 1'', 1')$	1	3	3, 3'
$-k_I$	-2	$(0, 0, 0)$	$(-4, -4, -4)$	0	$k = 2$	$k = 6$

$$M_d = \begin{pmatrix} \alpha_d & 0 & 0 \\ 0 & \beta_d & 0 \\ 0 & 0 & \gamma_d \end{pmatrix} \begin{pmatrix} Y_1 & Y_3 & Y_2 \\ Y_2 & Y_1 & Y_3 \\ Y_3 & Y_2 & Y_1 \end{pmatrix}_{RL}, \quad \text{Weight 2 modular forms}$$

$$M_u = \begin{pmatrix} \alpha_u & 0 & 0 \\ 0 & \beta_u & 0 \\ 0 & 0 & \gamma_u \end{pmatrix} \left[\begin{pmatrix} Y_1^{(6)} & Y_3^{(6)} & Y_2^{(6)} \\ Y_2^{(6)} & Y_1^{(6)} & Y_3^{(6)} \\ Y_3^{(6)} & Y_2^{(6)} & Y_1^{(6)} \end{pmatrix} + \begin{pmatrix} g_{u1} & 0 & 0 \\ 0 & g_{u2} & 0 \\ 0 & 0 & g_{u3} \end{pmatrix} \begin{pmatrix} Y_1'^{(6)} & Y_3'^{(6)} & Y_2'^{(6)} \\ Y_2'^{(6)} & Y_1'^{(6)} & Y_3'^{(6)} \\ Y_3'^{(6)} & Y_2'^{(6)} & Y_1'^{(6)} \end{pmatrix} \right]_{RL}$$

Weight 6 modular forms

After removing parameters $\alpha_q, \beta_q, \gamma_q$ by inputting quark masses, we have 3 complex parameters in addition to τ (8 real parameters).

We set $g_{u1}=g_{u2}=g_{u3}$. 4 parameters \Leftrightarrow 4 CKM elements are reproduced.

Lepton Sector

Common modulus τ for both quarks and leptons

	L	(e^c, μ^c, τ^c)	H_u	H_d	$Y_r^{(2)}, Y_r^{(4)}$
$SU(2)$	2	1	2	2	1
A_4	3	$(1, 1'', 1')$	1	1	$3, \{3, 1, 1'\}$
$-k_I$	-2	$(0, 0, 0)$ or $(-4, -4, -4)$	0	0	$2, 4$

$$M_E^{(6)} = v_d \begin{pmatrix} \alpha_e & 0 & 0 \\ 0 & \beta_e & 0 \\ 0 & 0 & \gamma_e \end{pmatrix} \left[\begin{pmatrix} Y_1^{(6)} & Y_3^{(6)} & Y_2^{(6)} \\ Y_2^{(6)} & Y_1^{(6)} & Y_3^{(6)} \\ Y_3^{(6)} & Y_2^{(6)} & Y_1^{(6)} \end{pmatrix} + g_e \begin{pmatrix} Y_1'^{(6)} & Y_3'^{(6)} & Y_2'^{(6)} \\ Y_2'^{(6)} & Y_1'^{(6)} & Y_3'^{(6)} \\ Y_3'^{(6)} & Y_2'^{(6)} & Y_1'^{(6)} \end{pmatrix} \right]_{RL}$$

$$w_\nu = -\frac{1}{\Lambda} (H_u H_u L L Y_r^{(k)})_1 \quad \text{Weinberg operator by using weight 4 modular forms}$$

$$M_\nu = \frac{v_u^2}{\Lambda} \left[\begin{pmatrix} 2Y_1^{(4)} & -Y_3^{(4)} & -Y_2^{(4)} \\ -Y_3^{(4)} & 2Y_2^{(4)} & -Y_1^{(4)} \\ -Y_2^{(4)} & -Y_1^{(4)} & 2Y_3^{(4)} \end{pmatrix} + g_{\nu 1} Y_1^{(4)} \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix} + g_{\nu 2} Y_{1'}^{(4)} \begin{pmatrix} 0 & 1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 1 \end{pmatrix} \right]_{LL}$$

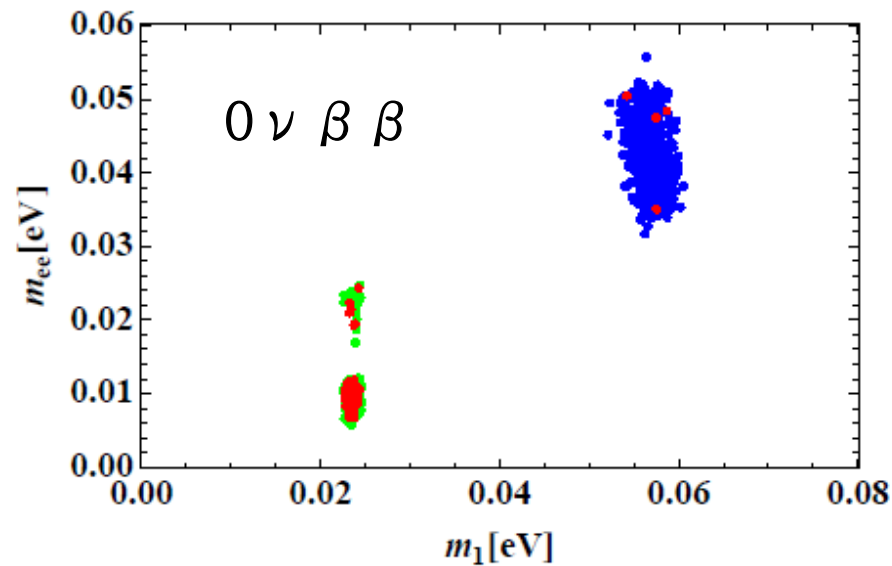
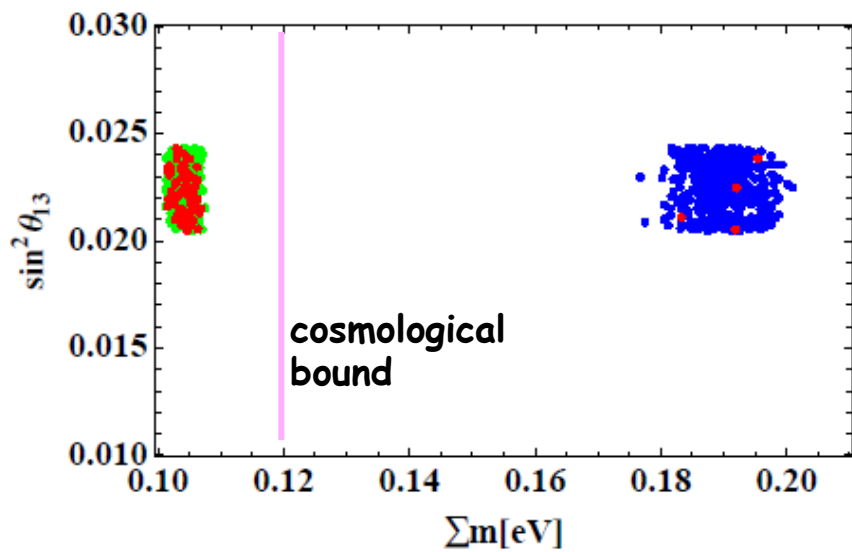
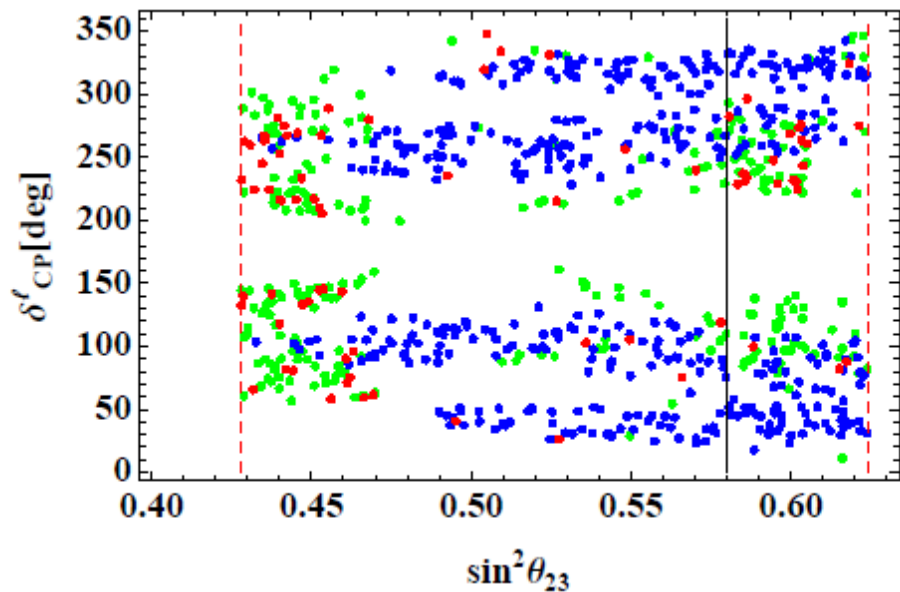
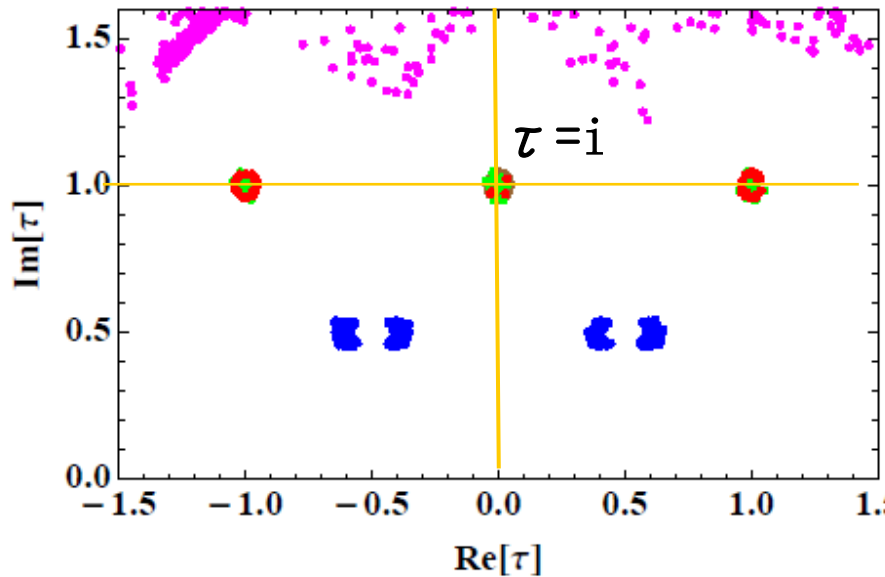
At $\tau = i$ (S symmetric limit) $M_\nu = \begin{pmatrix} \times & 0 & 0 \\ 0 & \times & \times \\ 0 & \times & \times \end{pmatrix}$

Input of 4 observed values: $\theta_{12}, \theta_{23}, \theta_{13}, \Delta m_{sol}^2 / \Delta m_{atm}^2$

output: $\delta_{CP}, \langle m_{ee} \rangle, \sum m_i$

● lepton ● lepton ● quark / lepton
● quark ● quark

NH



5 Summary

We present a A_4 modular invariant model with **common τ** for quarks and leptons by using weight 2, 4 and 6 modular forms.

- Modulus τ is common in both quarks and leptons close to **$\tau = i$** .
- Quark Mass matrices is consistent with observed CKM matrix.
- Lepton mass matrix is consistent with 3 observed mixing angles NH is favored. (IH is not allowed).

By imposing common τ for quarks and leptons, δ_{CP} , $\langle m_{ee} \rangle$ and $\sum m_i$ are predicted.

New approach for Quark-Lepton unification with modular invariance !

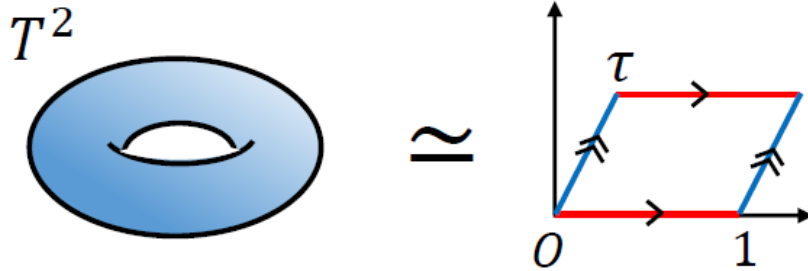
Back up slides

Alternative assignment of weight for quarks

	Q	(q_1^c, q_2^c, q_3^c)	H_q	$Y_3^{(k)}$
$SU(2)$	2	1	2	1
A_4	3	$(1, 1'', 1')$	1	3
$-k_I$	-2	$(-4, -2, 0)$	0	$k = 2, 4, 6$

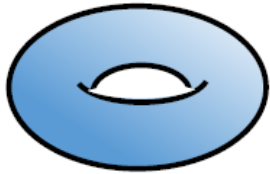
$$M_q = v_q \begin{pmatrix} \alpha_q & 0 & 0 \\ 0 & \beta_q & 0 \\ 0 & 0 & \gamma_q \end{pmatrix} \begin{pmatrix} Y_1^{(6)} + g_q Y_1'^{(6)} & Y_3^{(6)} + g_q Y_3'^{(6)} & Y_2^{(6)} + g_q Y_2'^{(6)} \\ Y_2^{(4)} & Y_1^{(4)} & Y_3^{(4)} \\ Y_3^{(2)} & Y_2^{(2)} & Y_1^{(2)} \end{pmatrix}$$

2D torus (T^2) is equivalent to parallelogram with identification of confronted sides.



Two-dimensional torus T^2 is obtained as $T^2 = \mathbb{R}^2 / \Lambda$
 Λ is two-dimensional lattice

The shape of torus is represented by a modulus $\tau \in \mathbb{C}$.

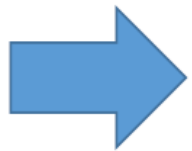


$\tau = \tau_1$



$\tau = \tau_2$

The different value of τ realize the different shape of T^2

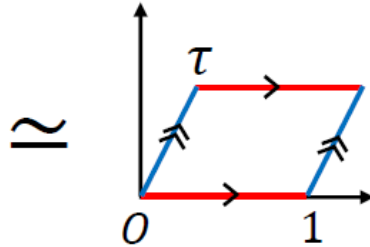
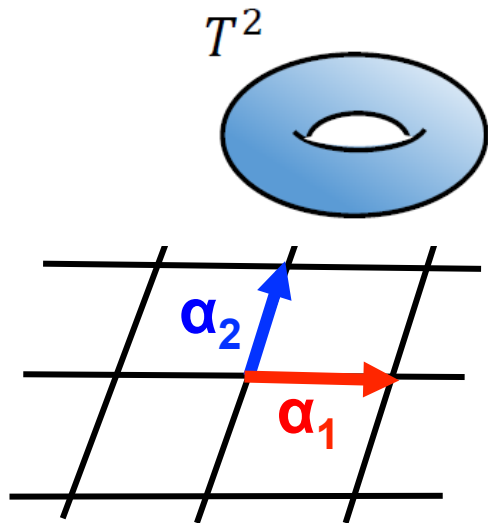


\mathcal{L}_{eff} depends on τ . e.g.) $\mathcal{L}_{\text{eff}} \supset Y(\tau)_{ij} \phi \bar{\psi}_i \psi_j + \dots$

➤ 4D effective theory depends on a modulus τ

Modular transformation

The shape of a torus $T^2 \simeq$ The shape of a lattice on \mathbb{C} -plane



Two-dimensional torus T^2 is obtained as
 $T^2 = \mathbb{R}^2 / \Lambda$

Λ is two-dimensional lattice,
 which is spanned by two lattice vectors

$$\alpha_1 = 2\pi R \quad \text{and} \quad \alpha_2 = 2\pi R \tau$$

$$(x, y) \sim (x, y) + n_1 \alpha_1 + n_2 \alpha_2$$

$\tau = \alpha_2 / \alpha_1$ is a modulus parameter (complex).

The same lattice is spanned by other bases under the transformation

$$\begin{pmatrix} \alpha'_2 \\ \alpha'_1 \end{pmatrix} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} \alpha_2 \\ \alpha_1 \end{pmatrix} \quad \begin{array}{l} ad-bc=1 \\ a, b, c, d \text{ are integer} \end{array} \quad SL(2, \mathbb{Z})$$

$$\begin{pmatrix} \alpha'_2 \\ \alpha'_1 \end{pmatrix} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} \alpha_2 \\ \alpha_1 \end{pmatrix}$$



$$\tau = \alpha_2 / \alpha_1$$

$$\tau \longrightarrow \tau' = \frac{a\tau + b}{c\tau + d}$$

Modular transformation

Modular transf. does not change the lattice (torus)



4D effective theory (depends on τ)
must be invariant under modular transf.

The modular transformation is generated by S and T .

$$\tau \longrightarrow \tau' = \frac{a\tau + b}{c\tau + d}$$

$$S : \tau \longrightarrow -\frac{1}{\tau}$$

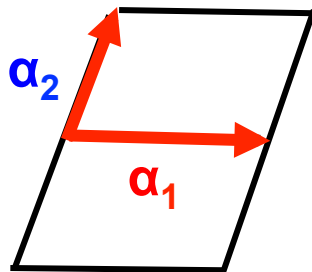
duality

$$T : \tau \longrightarrow \tau + 1$$

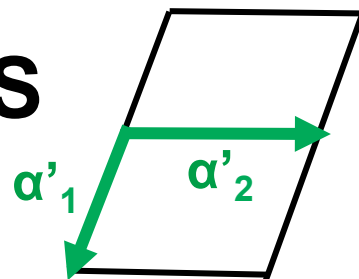
Discrete shift symmetry

$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} = \begin{pmatrix} 0 & 1 \\ -1 & 0 \end{pmatrix}$$

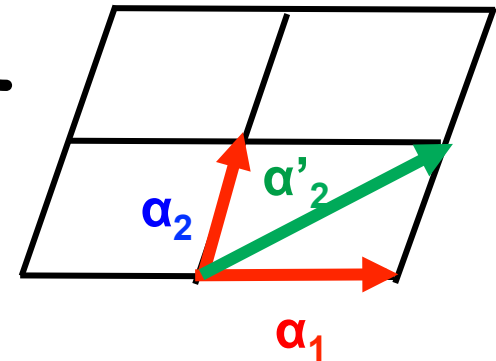
$$\begin{pmatrix} a & b \\ c & d \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 1 & 1 \end{pmatrix}$$



S



T



$$\tau = \alpha_2 / \alpha_1$$

Kinetic Term

Kinetic term of the modulus τ $\frac{|\partial_\mu \tau|^2}{\langle -i\tau + i\bar{\tau} \rangle^2}$

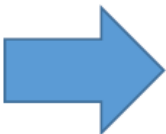
Modular transformation $\tau' = \frac{a\tau + b}{c\tau + d}, ad - bc = 1$

■ numerator

$$\partial_\mu \tau' = \frac{(a\partial_\mu \tau)(c\tau + d) - (a\tau + b)(c\partial_\mu \tau)}{(c\tau + d)^2} = \frac{(ad - bc)\partial_\mu \tau}{(c\tau + d)^2} = \frac{\partial_\mu \tau}{(c\tau + d)^2}$$

■ denominator

$$\tau' - \bar{\tau}' = \frac{(a\tau + b)(c\bar{\tau} + d) - (a\bar{\tau} + b)(c\tau + d)}{|c\tau + d|^2} = \frac{(ad - bc)(\tau - \bar{\tau})}{|c\tau + d|^2} = \frac{\tau - \bar{\tau}}{|c\tau + d|^2}$$

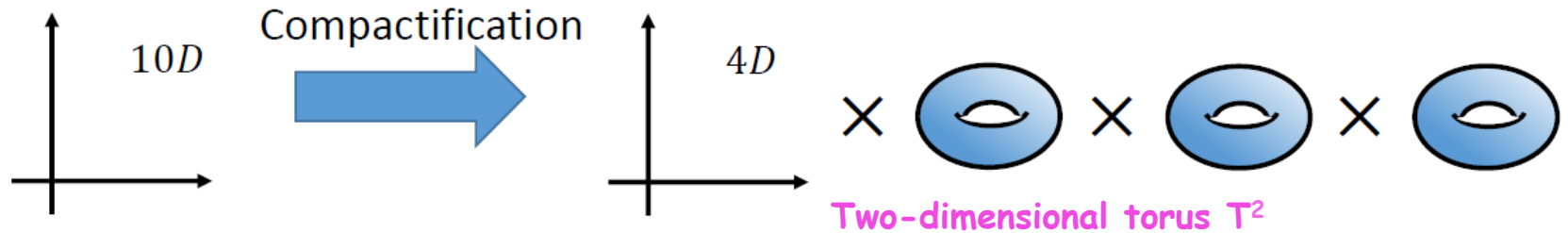
 $\frac{|\partial_\mu \tau'|^2}{\langle -i\tau' + i\bar{\tau}' \rangle^2} = \frac{|\partial_\mu \tau|^2}{\langle -i\tau + i\bar{\tau} \rangle^2}$ Modular invariant

Superstring theory 10D
Our universe is 4D




The extra 6D
should be compactified.

Torus compactification



We get 4D effective Lagrangian by integrating out over 6D.

$$S = \int d^4x d^6y \mathcal{L}_{10D} \rightarrow \int d^4x \mathcal{L}_{\text{eff}}$$

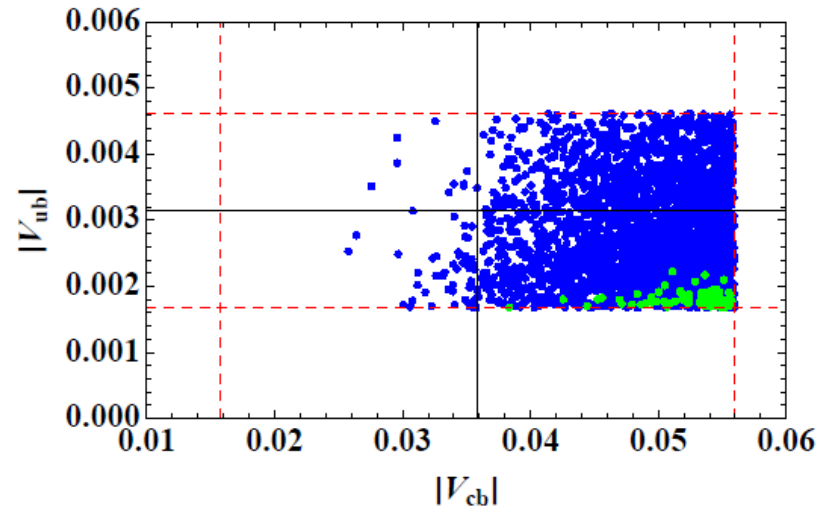
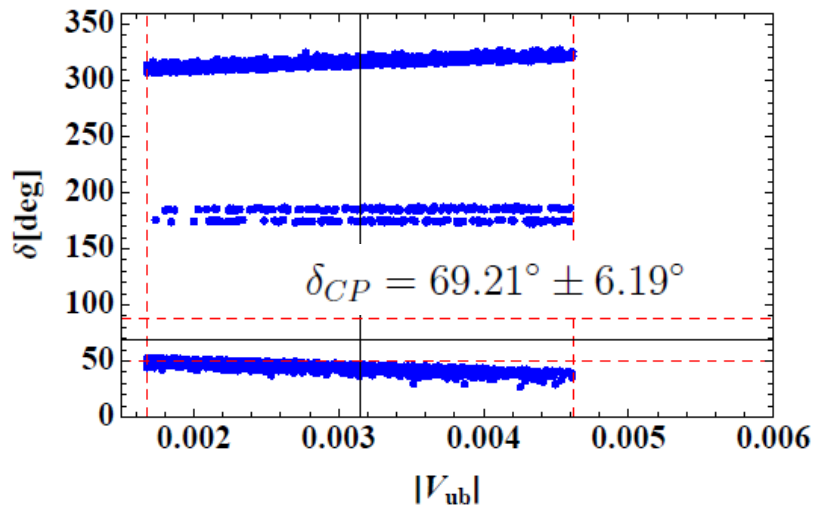
➔ \mathcal{L}_{eff} depends on the structure of 

➤ 4D effective theory depends on internal space

Input : charged lepton masses and three mixing angles at GUT scale

Effect of RGE depends on $\tan \beta$, M_{SUSY} , threshold effect

Output : CP violating phase δ_{CP} (PDG)



($\tan \beta = 5$, $M_{\text{SUSY}} = 1 \text{ TeV}$)

$$|V_d| = \begin{pmatrix} 0.5537 & 0.6135 & 0.5631 \\ 0.8110 & 0.2439 & 0.5317 \\ 0.1889 & 0.7511 & 0.6326 \end{pmatrix}, \quad |V_u| = \begin{pmatrix} 0.4857 & 0.6859 & 0.5419 \\ 0.8198 & 0.2382 & 0.5208 \\ 0.3034 & 0.6876 & 0.6596 \end{pmatrix}$$