

Σ^0 - Λ^0 state mixing from lattice QCD+QED

Zeno Kordov 6/8/2020

Research supervisors: Assoc. Prof. Ross Young and Assoc. Prof. James Zanotti

Table of contents

1. Introduction to mixing

- Mixing on the lattice correlation functions and mixing angle extraction expansions and ensembles
- 3. Fit results and observations

Electromagnetic contribution to Σ - Λ mixing using lattice QCD+QED Phys.Rev.D 101 (2020) 3, 034517; arXiv:1911.02186

- · R. Horsley
- · Y. Nakamura
- · H. Perlt
- · P.E.L. Rakow
- · G. Schierholz
- · H. Stüben
- · R.D. Young
- · J.M. Zanotti
- (CSSM/QCDSF/UKQCD Collaboration)

Introduction to mixing

What:

 In spectroscopy context mixing refers to diagonalisation of (isospin) SU(3)-flavour states to form mass eigenstates

What:

- In spectroscopy context mixing refers to diagonalisation of (isospin) SU(3)-flavour states to form mass eigenstates
- The linear combination of SU(3) states that forms mass eigenstates is parametrised by a single mixing angle for 2-state systems like Σ^0 - Λ^0

What:

- In spectroscopy context mixing refers to diagonalisation of (isospin) SU(3)-flavour states to form mass eigenstates
- The linear combination of SU(3) states that forms mass eigenstates is parametrised by a single mixing angle for 2-state systems like $\Sigma^0-\Lambda^0$

Why:

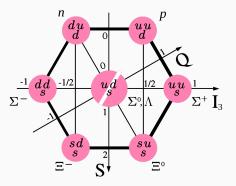
 Mixing is driven by the breaking of flavour symmetry (degeneracy), and for typical isospin eigenstates only occurs once u-d quark degeneracy is broken

What:

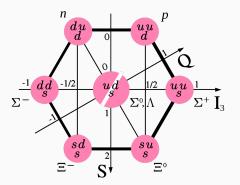
- In spectroscopy context mixing refers to diagonalisation of (isospin) SU(3)-flavour states to form mass eigenstates
- The linear combination of SU(3) states that forms mass eigenstates is parametrised by a single mixing angle for 2-state systems like Σ^0 - Λ^0

Why:

- Mixing is driven by the breaking of flavour symmetry (degeneracy), and for typical isospin eigenstates only occurs once u-d quark degeneracy is broken
- The magnitude of mixing is a measure of isospin symmetry breaking and SU(3) symmetry breaking



 In constructing SU(3) states, an exact isospin 'basis' is usually chosen in favour of U- or V-spin; the analogous SU(2) sub-algebras based upon d-s and u-s symmetry respectively



- In constructing SU(3) states, an exact isospin 'basis' is usually chosen in favour of U- or V-spin; the analogous SU(2) sub-algebras based upon d-s and u-s symmetry respectively
- With QED on the lattice, exact flavour degeneracy cannot be achieved, however d-s (U-spin) symmetry is achieved when quark masses are set equal

Mixing on the lattice

Correlation functions

We employ the general SU(3)-flavour interpolating operators

$$\mathcal{B}_{\Sigma(abc),\alpha}(x) = \frac{1}{\sqrt{2}} \epsilon^{lmn} \left(b_{\alpha}^{l}(x) \left[a^{m}(x)^{\top} C \gamma_{5} c^{n}(x) \right] + a_{\alpha}^{l}(x) \left[b^{m}(x)^{\top} C \gamma_{5} c^{n}(x) \right] \right),$$

$$\mathcal{B}_{\Lambda(abc),\alpha}(x) = \frac{1}{\sqrt{6}} \epsilon^{lmn} \left(2 c_{\alpha}^{l}(x) \left[a^{m}(x)^{\top} C \gamma_{5} b^{n}(x) \right] + b_{\alpha}^{l}(x) \left[a^{m}(x)^{\top} C \gamma_{5} c^{n}(x) \right] \right),$$

$$\dots - a_{\alpha}^{l}(x) \left[b^{m}(x)^{\top} C \gamma_{5} c^{n}(x) \right] \right),$$

Correlation functions

We employ the general SU(3)-flavour interpolating operators

$$\mathcal{B}_{\Sigma(abc),\alpha}(x) = \frac{1}{\sqrt{2}} \epsilon^{lmn} \left(b_{\alpha}^{l}(x) \left[\mathbf{a}^{m}(x)^{\top} C \gamma_{5} c^{n}(x) \right] + \mathbf{a}_{\alpha}^{l}(x) \left[\mathbf{b}^{m}(x)^{\top} C \gamma_{5} c^{n}(x) \right] \right),$$

$$\mathcal{B}_{\Lambda(abc),\alpha}(x) = \frac{1}{\sqrt{6}} \epsilon^{lmn} \left(2c_{\alpha}^{l}(x) \left[\mathbf{a}^{m}(x)^{\top} C \gamma_{5} \mathbf{b}^{n}(x) \right] + b_{\alpha}^{l}(x) \left[\mathbf{a}^{m}(x)^{\top} C \gamma_{5} c^{n}(x) \right] \right),$$

$$\dots - \mathbf{a}_{\alpha}^{l}(x) \left[\mathbf{b}^{m}(x)^{\top} C \gamma_{5} c^{n}(x) \right] \right),$$

Correlation functions

We employ the general SU(3)-flavour interpolating operators

$$\mathcal{B}_{\Sigma(abc),\alpha}(x) = \frac{1}{\sqrt{2}} \epsilon^{lmn} \left(b_{\alpha}^{l}(x) \left[\mathbf{a}^{m}(x)^{\top} C \gamma_{5} c^{n}(x) \right] + \mathbf{a}_{\alpha}^{l}(x) \left[\mathbf{b}^{m}(x)^{\top} C \gamma_{5} c^{n}(x) \right] \right),$$

$$\mathcal{B}_{\Lambda(abc),\alpha}(x) = \frac{1}{\sqrt{6}} \epsilon^{lmn} \left(2c_{\alpha}^{l}(x) \left[\mathbf{a}^{m}(x)^{\top} C \gamma_{5} \mathbf{b}^{n}(x) \right] + \mathbf{b}_{\alpha}^{l}(x) \left[\mathbf{a}^{m}(x)^{\top} C \gamma_{5} c^{n}(x) \right] \right),$$

$$\dots - \mathbf{a}_{\alpha}^{l}(x) \left[\mathbf{b}^{m}(x)^{\top} C \gamma_{5} c^{n}(x) \right] \right),$$

and form the 2x2 correlation matrix

$$C_{ij}(t) \propto \mathrm{Tr}_D \Gamma_{unpol} \left\langle \sum_{\vec{y}} \mathcal{B}_i(\vec{y},t) \bar{\mathcal{B}}_j(\vec{x}_0,0) \right\rangle, \quad i,j = \Sigma(abc), \Lambda(abc)$$

4

Simulation and mixing angle extraction

To extract mixing angles from our simulations we calculate the eigenvectors of the correlation matrices

$$\begin{bmatrix} C_{\Sigma\Sigma,i}(t) & C_{\Sigma\Lambda,i}(t) \\ C_{\Lambda\Sigma,i}(t) & C_{\Lambda\Lambda,i}(t) \end{bmatrix}, \quad 1 \le t \le n_t, \quad i = \text{isospin,u-spin,v-spin}$$

Simulation and mixing angle extraction

To extract mixing angles from our simulations we calculate the eigenvectors of the correlation matrices

$$\begin{bmatrix} C_{\Sigma\Sigma,i}(t) & C_{\Sigma\Lambda,i}(t) \\ C_{\Lambda\Sigma,i}(t) & C_{\Lambda\Lambda,i}(t) \end{bmatrix}, \quad 1 \leq t \leq n_t, \quad i = \text{isospin,u-spin,v-spin}$$

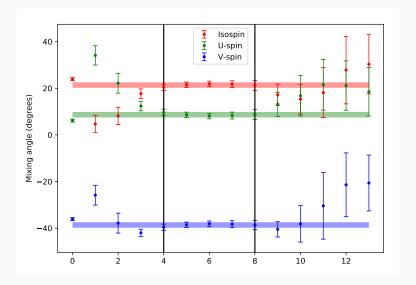
which are parametrised by the mixing angles:

$$\vec{e}_i(t) = \begin{bmatrix} \cos \theta_{\Sigma \Lambda, i}(t) \\ \sin \theta_{\Sigma \Lambda, i}(t) \end{bmatrix}, \begin{bmatrix} -\sin \theta_{\Sigma \Lambda, i}(t) \\ \cos \theta_{\Sigma \Lambda, i}(t) \end{bmatrix}, \quad i = isospin, u-spin, v-spin$$

and fit the mixing angle plateaus

5

Mixing angle extraction

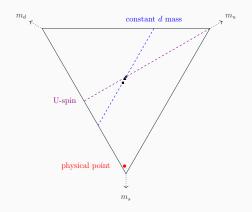


 \cdot Mixing angle extractions from QCDSF dynamical QCD+QED confs

- \cdot Mixing angle extractions from QCDSF dynamical QCD+QED confs
- We use a \approx 10x larger than physical EM coupling to exaggerate QED effects (needs to be corrected for in extrapolations)

- Mixing angle extractions from QCDSF dynamical QCD+QED confs
- We use a \approx 10x larger than physical EM coupling to exaggerate QED effects (needs to be corrected for in extrapolations)
- Simulate along mass trajectory with constant singlet quark mass, ie. $(m_u + m_d + m_s)/3 = {\rm constant} = m_0$

- Mixing angle extractions from QCDSF dynamical QCD+QED confs
- We use a ≈10x larger than physical EM coupling to exaggerate QED effects (needs to be corrected for in extrapolations)
- Simulate along mass trajectory with constant singlet quark mass, ie. $(m_u + m_d + m_s)/3 = {\rm constant} = m_0$



We formulate a LO expansion of the correlation functions themselves, treating them (at a given time) as smooth functions of the quark masses and charges

We formulate a LO expansion of the correlation functions themselves, treating them (at a given time) as smooth functions of the quark masses and charges

$$\rightarrow \tan 2\theta_{\Sigma\Lambda, \text{isospin}}(t) = -\sqrt{3} \left(\frac{D_{\text{QCD}}(t)(\delta m_u - Z\delta m_d) + D_{\text{QED}}(t)}{3D_{\text{QCD}}(t)(\delta m_u + Z\delta m_d) + D_{\text{QED}}(t)} \right),$$

8

We formulate a LO expansion of the correlation functions themselves, treating them (at a given time) as smooth functions of the quark masses and charges

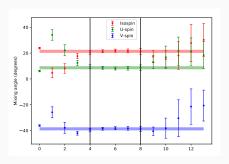
We formulate a LO expansion of the correlation functions themselves, treating them (at a given time) as smooth functions of the quark masses and charges

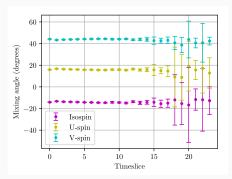
Mixing angle time dependence

$$\tan 2\theta_{\Sigma\Lambda, \text{isospin}}(t) = -\sqrt{3} \left(\frac{D_{\text{QCD}}(t)(\delta m_u - Z\delta m_d) + D_{\text{QED}}(t)}{3D_{\text{QCD}}(t)(\delta m_u + Z\delta m_d) + D_{\text{QED}}(t)} \right),$$

Mixing angle time dependence

$$\tan 2\theta_{\Sigma\Lambda, \mathrm{isospin}}(t) = -\sqrt{3} \left(\frac{D_{\mathrm{QCD}}(t)(\delta m_u - Z\delta m_d) + D_{\mathrm{QED}}(t)}{3D_{\mathrm{QCD}}(t)(\delta m_u + Z\delta m_d) + D_{\mathrm{QED}}(t)} \right),$$





| Lattice Ensembles | | | | | | |
|---------------------|--|-----------|----------|--|-----------------------|--|
| volume | $\kappa_{\sf u},\kappa_{\sf d},\kappa_{\sf s}$ | (valence) | | $\theta_{\Sigma\Lambda, \mathrm{isospin}}$ | M _{uū} (MeV) | |
| $24^{3} \times 48$ | 0.124362 | 0.121713 | 0.121713 | -30° | 442(9) | |
| $24^{3} \times 48$ | 0.124374 | 0.121713 | 0.121701 | | 423(9) | |
| | 0.124387 | 0.121713 | 0.121689 | | 423(10) | |
| | 0.124400 | 0.121740 | 0.121649 | | 378(28) | |
| $24^{3} \times 48$ | 0.124400 | 0.121713 | 0.121677 | | 405(8) | |
| | 0.124420 | 0.121713 | 0.121657 | | 387(8) | |
| | 0.124430 | 0.121760 | 0.121601 | | 377(8) | |
| 48 ³ ×96 | 0.124508 | 0.121821 | 0.121466 | | 284(4) | |
| | 0.124400 | 0.121713 | 0.121677 | | 389(5) | |

| Lattice Ensembles | | | | | | |
|--------------------|--|-----------|----------|--|-----------------------|--|
| volume | $\kappa_{\sf u},\kappa_{\sf d},\kappa_{\sf s}$ | (valence) | | $\theta_{\Sigma\Lambda, \mathrm{isospin}}$ | M _{uū} (MeV) | |
| $24^{3} \times 48$ | 0.124362 | 0.121713 | 0.121713 | -30° | 442(9) | |
| $24^{3} \times 48$ | 0.124374 | 0.121713 | 0.121701 | | 423(9) | |
| | 0.124387 | 0.121713 | 0.121689 | | 423(10) | |
| | 0.124400 | 0.121740 | 0.121649 | | 378(28) | |
| $24^{3} \times 48$ | 0.124400 | 0.121713 | 0.121677 | | 405(8) | |
| | 0.124420 | 0.121713 | 0.121657 | | 387(8) | |
| | 0.124430 | 0.121760 | 0.121601 | | 377(8) | |
| $48^{3} \times 96$ | 0.124508 | 0.121821 | 0.121466 | | 284(4) | |
| | 0.124400 | 0.121713 | 0.121677 | | 389(5) | |

| Lattice Ensembles | | | | | | |
|---------------------|--|-----------|----------|--|-----------------------|--|
| volume | $\kappa_{\sf u},\kappa_{\sf d},\kappa_{\sf s}$ | (valence) | | $\theta_{\Sigma\Lambda, \mathrm{isospin}}$ | M _{uū} (MeV) | |
| $24^{3} \times 48$ | 0.124362 | 0.121713 | 0.121713 | -30° | 442(9) | |
| $24^{3} \times 48$ | 0.124374 | 0.121713 | 0.121701 | | 423(9) | |
| | 0.124387 | 0.121713 | 0.121689 | | 423(10) | |
| | 0.124400 | 0.121740 | 0.121649 | | 378(28) | |
| $24^{3} \times 48$ | 0.124400 | 0.121713 | 0.121677 | | 405(8) | |
| | 0.124420 | 0.121713 | 0.121657 | | 387(8) | |
| | 0.124430 | 0.121760 | 0.121601 | | 377(8) | |
| 48 ³ ×96 | 0.124508 | 0.121821 | 0.121466 | | 284(4) | |
| | 0.124400 | 0.121713 | 0.121677 | | 389(5) | |

red = unitary

| Lattice Ensembles | | | | | | |
|--------------------|--|-----------|----------|--|-----------------------|--|
| volume | $\kappa_{\sf u},\kappa_{\sf d},\kappa_{\sf s}$ | (valence) | | $\theta_{\Sigma\Lambda, \mathrm{isospin}}$ | M _{uū} (MeV) | |
| $24^3 \times 48$ | 0.124362 | 0.121713 | 0.121713 | -30° | 442(9) | |
| $24^{3} \times 48$ | 0.124374 | 0.121713 | 0.121701 | | 423(9) | |
| | 0.124387 | 0.121713 | 0.121689 | | 423(10) | |
| | 0.124400 | 0.121740 | 0.121649 | | 378(28) | |
| $24^{3} \times 48$ | 0.124400 | 0.121713 | 0.121677 | | 405(8) | |
| | 0.124420 | 0.121713 | 0.121657 | | 387(8) | |
| | 0.124430 | 0.121760 | 0.121601 | | 377(8) | |
| $48^{3} \times 96$ | 0.124508 | 0.121821 | 0.121466 | | 284(4) | |
| | 0.124400 | 0.121713 | 0.121677 | | 389(5) | |

red = unitary
blue = partially quenched

Fit results and observations

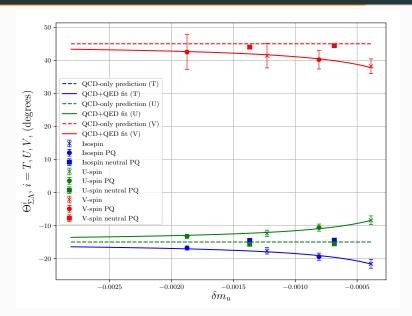
Mixing angle results

| Lattice Ensembles | | | | | | |
|---------------------|--|-----------|----------|--|-----------------------|--|
| volume | $\kappa_{\sf u},\kappa_{\sf d},\kappa_{\sf s}$ | (valence) | | $\theta_{\Sigma\Lambda, \mathrm{isospin}}$ | M _{uū} (MeV) | |
| $24^{3} \times 48$ | 0.124362 | 0.121713 | 0.121713 | -30° | 442(9) | |
| $24^{3} \times 48$ | 0.124374 | 0.121713 | 0.121701 | -21.8(1.1)° | 423(9) | |
| | 0.124387 | 0.121713 | 0.121689 | -19.5(1.2)° | 423(10) | |
| | 0.124400 | 0.121740 | 0.121649 | -6(1)° | 378(28) | |
| $24^{3} \times 48$ | 0.124400 | 0.121713 | 0.121677 | -17.8(7)° | 405(8) | |
| | 0.124420 | 0.121713 | 0.121657 | -16.7(7)° | 387(8) | |
| | 0.124430 | 0.121760 | 0.121601 | -4.8(7)° | 377(8) | |
| 48 ³ ×96 | 0.124508 | 0.121821 | 0.121466 | -3.5(4)° | 284(4) | |
| | 0.124400 | 0.121713 | 0.121677 | -18.5(9)° | 389(5) | |

Mixing angle results

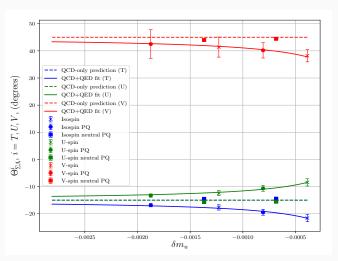
| Lattice Ensembles | | | | | | |
|--------------------|--|-----------|----------|--|-----------------------|--|
| volume | $\kappa_{\sf u},\kappa_{\sf d},\kappa_{\sf s}$ | (valence) | | $\theta_{\Sigma\Lambda, \mathrm{isospin}}$ | M _{uū} (MeV) | |
| $24^{3} \times 48$ | 0.124362 | 0.121713 | 0.121713 | -30° | 442(9) | |
| $24^{3} \times 48$ | 0.124374 | 0.121713 | 0.121701 | -21.8(1.1)° | 423(9) | |
| | 0.124387 | 0.121713 | 0.121689 | -19.5(1.2)° | 423(10) | |
| | 0.124400 | 0.121740 | 0.121649 | -6(1)° | 378(28) | |
| $24^{3} \times 48$ | 0.124400 | 0.121713 | 0.121677 | -17.8(7)° | 405(8) | |
| | 0.124420 | 0.121713 | 0.121657 | -16.7(7)° | 387(8) | |
| | 0.124430 | 0.121760 | 0.121601 | -4.8(7)° | 377(8) | |
| $48^{3} \times 96$ | 0.124508 | 0.121821 | 0.121466 | -3.5(4)° | 284(4) | |
| | 0.124400 | 0.121713 | 0.121677 | -18.5(9)° | 389(5) | |

QED contribution and basis relations



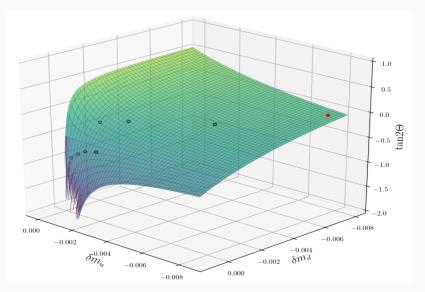
QED contribution and basis relations

$$\tan 2\theta_{\mathbf{\Sigma}\mathbf{\Lambda},\mathrm{isospin}} = -\sqrt{3}\left(\frac{D_{\mathrm{QCD}}(\delta m_u - Z\delta m_d) + D_{\mathrm{QED}}}{3D_{\mathrm{QCD}}(\delta m_u + Z\delta m_d) + D_{\mathrm{QED}}}\right),$$



Isospin extrapolation results

$$D_{\rm QED}/D_{\rm QCD} = -3.8(7) \times 10^{-5}, Z = 0.96(4), \chi^2/{
m DOF} = 0.84$$



- We find a mixing angle $\theta_{\Sigma\Lambda,\mathrm{isospin}} = -1.0(3)^\circ$
- This compares well with other determinations including QED (notably ¹: -0.86(6))

 $^{^{1}}$ R. H. Dalitz and F. Von Hippel, "Electromagnetic $\Lambda-\Sigma^{0}$ mixing and charge symmetry for the Λ -hyperon," Phys. Lett. **10** (1964), 153-157

 $^{^2}$ R. Horsley *et al.*, "Lattice determination of Sigma-Lambda mixing," Phys. Rev. D **91** (2015) no.7, 074512

- We find a mixing angle $\theta_{\Sigma\Lambda,\mathrm{isospin}} = -1.0(3)^\circ$
- This compares well with other determinations including QED (notably ¹: -0.86(6))
- We have not yet investigated systematic errors in this exploratory calculation (finite volume, EM coupling scaling, PQ)

 $^{^1}$ R. H. Dalitz and F. Von Hippel, "Electromagnetic $\Lambda-\Sigma^0$ mixing and charge symmetry for the Λ -hyperon," Phys. Lett. **10** (1964), 153-157

²R. Horsley *et al.*, "Lattice determination of Sigma-Lambda mixing," Phys. Rev. D **91** (2015) no.7, 074512

- . We find a mixing angle $\theta_{\Sigma\Lambda,\mathrm{isospin}} = -1.0(3)^\circ$
- This compares well with other determinations including QED (notably ¹: -0.86(6))
- We have not yet investigated systematic errors in this exploratory calculation (finite volume, EM coupling scaling, PQ)
- Roughly 2x magnitude of our QCD-only determination: $\theta_{\Sigma\Lambda, \rm isospin, QCD-only} = -0.55(3)^{\circ}, \mbox{ and past collaboration result}^2: \\ \theta_{\Sigma\Lambda, \rm isospin, QCD-only} = -0.35(13)(7)^{\circ}$

 $^{^{1}}$ R. H. Dalitz and F. Von Hippel, "Electromagnetic $\Lambda-\Sigma^{0}$ mixing and charge symmetry for the Λ -hyperon," Phys. Lett. **10** (1964), 153-157

²R. Horsley *et al.*, "Lattice determination of Sigma-Lambda mixing," Phys. Rev. D **91** (2015) no.7, 074512

- . We find a mixing angle $\theta_{\Sigma\Lambda,\mathrm{isospin}} = -1.0(3)^\circ$
- This compares well with other determinations including QED (notably ¹: -0.86(6))
- We have not yet investigated systematic errors in this exploratory calculation (finite volume, EM coupling scaling, PQ)
- Roughly 2x magnitude of our QCD-only determination: $\theta_{\Sigma\Lambda, \mathrm{isospin, QCD-only}} = -0.55(3)^{\circ}$, and past collaboration result²: $\theta_{\Sigma\Lambda, \mathrm{isospin, QCD-only}} = -0.35(13)(7)^{\circ}$
- In the future we intend to investigate the π^0 – η – η' mixing in this scheme

 $^{^{1}}$ R. H. Dalitz and F. Von Hippel, "Electromagnetic $\Lambda-\Sigma^{0}$ mixing and charge symmetry for the Λ -hyperon," Phys. Lett. **10** (1964), 153-157

²R. Horsley *et al.*, "Lattice determination of Sigma-Lambda mixing," Phys. Rev. D **91** (2015) no.7, 074512

- . We find a mixing angle $\theta_{\Sigma\Lambda,\mathrm{isospin}} = -1.0(3)^\circ$
- This compares well with other determinations including QED (notably ¹: -0.86(6))
- We have not yet investigated systematic errors in this exploratory calculation (finite volume, EM coupling scaling, PQ)
- Roughly 2x magnitude of our QCD-only determination: $\theta_{\Sigma\Lambda, \mathrm{isospin, QCD-only}} = -0.55(3)^{\circ}$, and past collaboration result²: $\theta_{\Sigma\Lambda, \mathrm{isospin, QCD-only}} = -0.35(13)(7)^{\circ}$
- In the future we intend to investigate the π^0 – η – η' mixing in this scheme
- · Thank you for listening!

 $^{^1}$ R. H. Dalitz and F. Von Hippel, "Electromagnetic $\Lambda-\Sigma^0$ mixing and charge symmetry for the Λ -hyperon," Phys. Lett. 10 (1964), 153-157

²R. Horsley *et al.*, "Lattice determination of Sigma-Lambda mixing," Phys. Rev. D **91** (2015) no.7, 074512