

# Form Factors For Heavy $\rightarrow$ Strange Semileptonic Decays

William Parrott

University of Glasgow



Supervisor: Dr. Chris Bouchard

B. Chakraborty: Today, 16:20, this session

D. Hatton: Tomorrow, 17:00, Had. Spec.

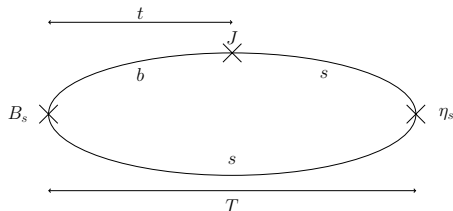


- ▶ Overview of heavy HISQ
- ▶ Results for  $B_s \rightarrow \eta_s$  and what we can learn from it
- ▶ Moving towards  $B \rightarrow K\ell^+\ell^-$
- ▶ Preliminary  $B \rightarrow K\ell^+\ell^-$  and  $D \rightarrow K\ell^-\bar{\nu}$  results



# Overview of heavy HISQ

- ▶ Calculate meson form factors over the full range of  $q^2 = (p_{\text{parent}} - p_{\text{daughter}})^2$  values
- ▶ Interested in  $f_+(q^2)$  and  $f_0(q^2)$  form factors for pseudoscalar to pseudoscalar decays
- ▶ Require three-point correlators with scalar and vector current insertions



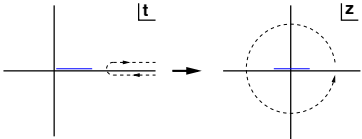
# Overview of heavy HISQ

- ▶ MILC HISQ 2+1+1 ensembles. All valence quarks HISQ
- ▶ 0.09fm 0.06fm and 0.045fm lattices for  $B_s \rightarrow \eta_s$
- ▶ Physical  $b$  is  $am_b \approx 0.9$  on finest lattice
- ▶ Choose several heavy masses and daughter momenta for each ensemble
- ▶ Combine heavy mass fit with continuum extrapolation
- ▶  $D_s \rightarrow \eta_s$  comes ‘for free’
- ▶ Cover whole physical  $q^2$  range



# Overview of heavy HISQ

$$\Lambda_{\text{QCD}} = 0.5 \text{ GeV}$$

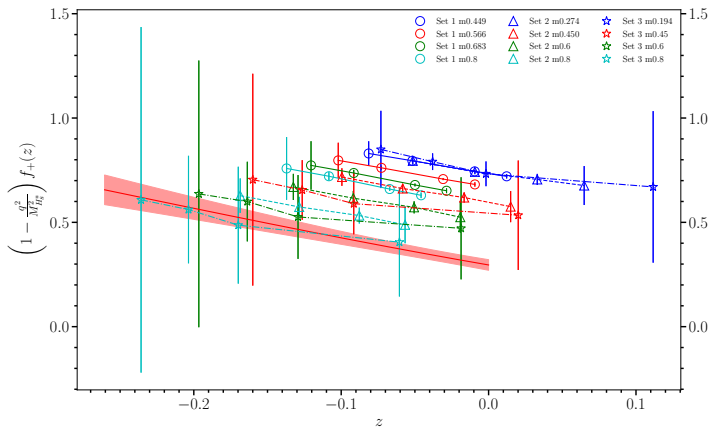
$$f_0(q^2) = \frac{1}{1 - \frac{q^2}{M_{H_s^0}^2}} \sum_{n=0}^{N-1} a_n^0 z^n, \quad (1)$$


$$f_+(q^2) = \frac{1}{1 - \frac{q^2}{M_{H_s^*}^2}} \sum_{n=0}^{N-1} a_n^+ \left( z^n - \frac{n}{N} (-1)^{n-N} z^N \right),$$

$$a_n^{0,+} = \left( 1 + \rho_n^{0,+} \log \left( \frac{M_{H_s}}{M_{D_s}} \right) \right) \times \sum_{i,j,k=0}^{N_{ijk}-1} d_{ijkn}^{0,+} \left( \frac{\Lambda_{\text{QCD}}}{M_{H_s}} \right)^i \left( \frac{am_h^{\text{val}}}{\pi} \right)^{2j} \left( \frac{a\Lambda_{\text{QCD}}}{\pi} \right)^{2k} \times (1 + \mathcal{N}_n^{0,+}). \quad (2)$$



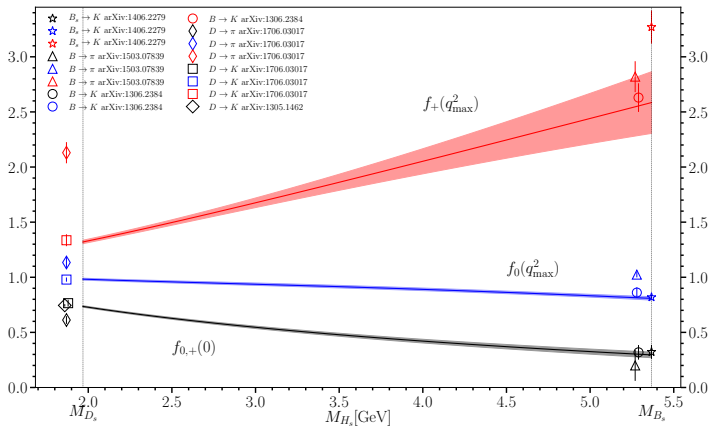
# $B_s \rightarrow \eta_s$ results



Continuum result at the  $b$  mass in red. h-HISQ allows us to evaluate at any mass from  $c$  to  $b$ .



# $B_s \rightarrow \eta_s$ results



Form factors largely independent of spectator quark mass



## Moving to $B \rightarrow K$

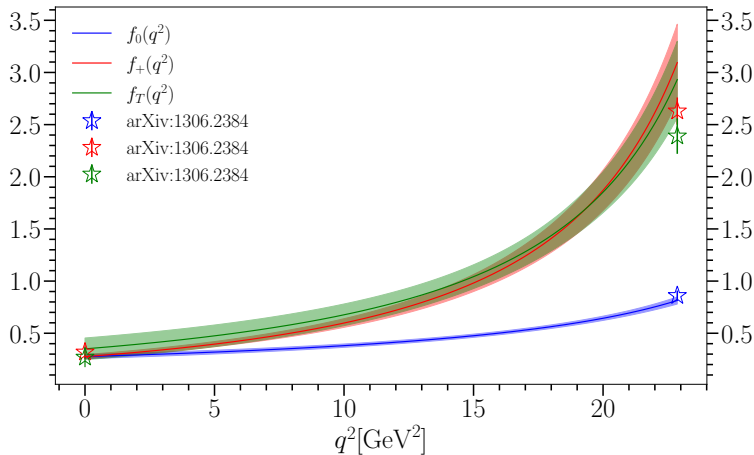
- ▶ Change spectator to light
- ▶ Calculate tensor form factor, using accurate tensor normalisation
- ▶ Include lattices with physical light quarks
- ▶ Include an overall chiral log term:

$$\text{logs} = 1 - \frac{9g^2}{8} \frac{m_l}{10m_s^{\text{tuned}}} \left( \log \left( \frac{m_l}{10m_s^{\text{tuned}}} \right) + \delta_{FV} \right) \quad (3)$$





# $B \rightarrow K$ preliminary results

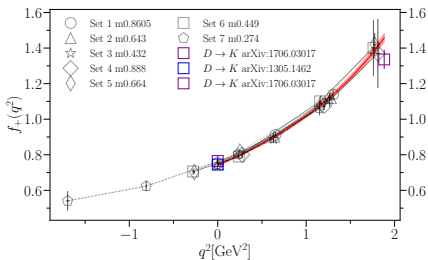
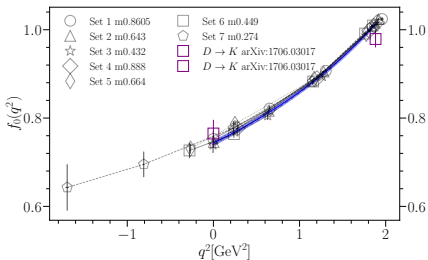


Tensor important for SM  $B \rightarrow K$  due to  $b \rightarrow s$  transition  
Normalisation with  $\mu = 2\text{GeV}$ , matched to  $\overline{\text{MS}}$  at 3 loop at  $b$  mass



# $D \rightarrow K$ preliminary results (*B. Chakraborty, C. T. H. Davies*)

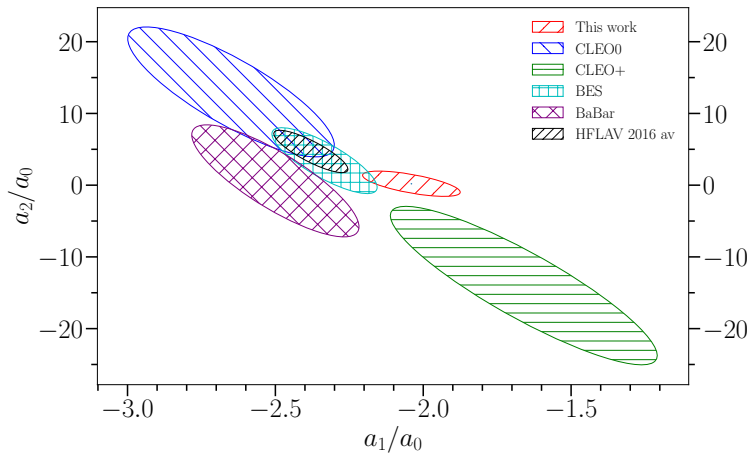
Sets 1-3 physical *v.* coarse to fine Sets 4-7  $m_s/m_l = 5$ , *v.* coarse to superfine



- ▶ Charm mass easy to reach on ensembles
- ▶ Full  $q^2$  range  $\implies$  can compare bin by bin with exp. partial decay rate data
- ▶ Lots of good exp. data available, can compare shape



# $D \rightarrow K$ preliminary results (*B. Chakraborty, C. T. H. Davies*)

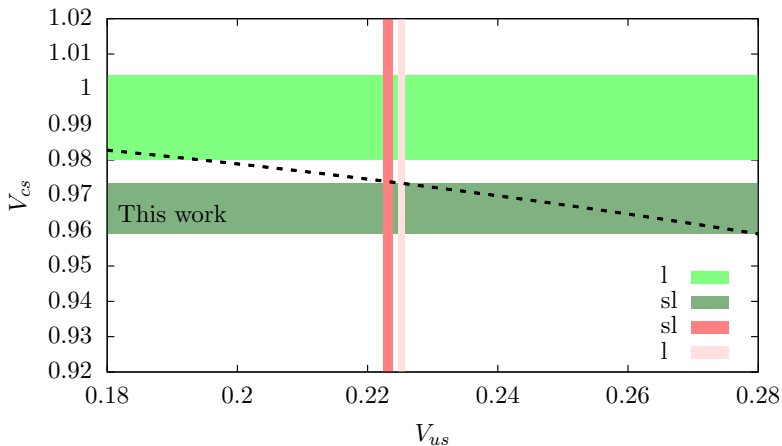


One  $\sigma$  error ellipses. Ratios of  $f_+$   $z$  expansion coefficients  $a_n$ , directly comparable with experiment.



# $D \rightarrow K$ preliminary results (*B. Chakraborty, C. T. H. Davies*)

Preliminary  $V_{cs} = 0.9662(71)$ , improvement on current PDG sl value of  $0.967(25)$



# Conclusions

- ▶ Heavy HISQ an effective method for studying heavy to strange decays and form factors
- ▶ Form factors largely independent of spectator quark mass
- ▶ Can improve upon  $B \rightarrow K\ell^+\ell^-$  and  $D \rightarrow K\ell\bar{\nu}$  results
- ▶ Improvement on  $V_{cs}$  determination from  $D \rightarrow K\ell^-\bar{\nu}$  using bin by bin comparisons with experiment

Thanks for listening. Any questions?



$$\begin{aligned}
 Z_V^0 Z_{\text{disc}} \langle \eta_s | V^0 | \hat{H}_s \rangle = & \\
 f_+^{H_s \rightarrow \eta_s}(q^2) \left( p_{H_s}^0 + p_{\eta_s}^0 - \frac{M_{H_s}^2 - M_{\eta_s}^2}{q^2} q^0 \right) & \quad (4) \\
 + f_0^{H_s \rightarrow \eta_s}(q^2) \frac{M_{H_s}^2 - M_{\eta_s}^2}{q^2} q^0, &
 \end{aligned}$$

$$Z_{\text{disc}} \langle \eta_s | S | H_s \rangle = \frac{M_{H_s}^2 - M_{\eta_s}^2}{m_h - m_s} f_0^{H_s \rightarrow \eta_s}(q^2), \quad (5)$$



$$z(q^2) = \frac{\sqrt{t_+ - q^2} - \sqrt{t_+ - t_0}}{\sqrt{t_+ - q^2} + \sqrt{t_+ - t_0}}. \quad (6)$$

$$\frac{m_l}{m_s} \approx \frac{M_\pi^2}{M_{\eta_s}^2} \quad (7)$$

