

APLAT 2020

Finite-volume effects in $(g - 2)_\mu^{\text{HVP,LO}}$

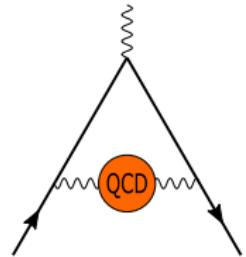
Agostino Patella

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in collaboration with **Maxwell T. Hansen**

MH and AP, Phys. Rev. Lett. 123 (2019) 172001, arXiv:1904.10010 [hep-lat].
MH and AP, accepted on JHEP, arXiv:2004.03935 [hep-lat].

$$a^{\text{HVP}}(T, L) = -\frac{1}{3} \int_0^{T/2} dx_0 K(x_0) \int_0^L d^3x \sum_{\mu=1}^3 \langle j_k(x_0, \mathbf{x}) j_k(0) \rangle_{T,L}$$



D. Bernecker, H. B. Meyer, *Vector Correlators in Lattice QCD: methods and applications*, Eur. Phys. J. A **47** (2011), 148.

Setup:

- ▶ Euclidean $T \times L^3$ torus with (anti)periodic boundary conditions.
- ▶ QCD in isosymmetric limit.
- ▶ Continuum limit has already been taken.

Goal: Understand the finite-volume corrections

$$\Delta a^{\text{HVP}}(T, L) = a^{\text{HVP}}(T, L) - a^{\text{HVP}}(\infty) ,$$

possibly in a model-independent way.

Structure of finite-volume corrections

$$\Delta a^{\text{HVP}}(T, L) = \mathcal{O}\left(e^{-m_\pi L}\right) + \mathcal{O}\left(e^{-\sqrt{2}m_\pi L}\right) + \mathcal{O}\left(e^{-\sqrt{3}m_\pi L}\right) + \mathcal{O}\left(e^{-\sqrt{2+\sqrt{3}}m_\pi L}\right)$$

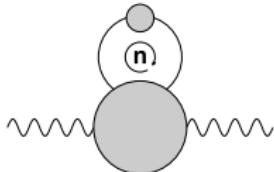
$$+ \mathcal{O}\left(e^{-m_\pi T}\right) + \mathcal{O}\left(e^{-\frac{3}{2}m_\pi T}\right) + \mathcal{O}\left(e^{-m_\pi \sqrt{T^2+L^2}}\right) + \mathcal{O}\left(e^{-m_K L}\right) + \mathcal{O}\left(e^{-m_K T}\right) + \dots$$

Loops can wrap around the spatial torus (in coordinate space!)

M. Lüscher, *Volume Dependence [...] Stable Particle States*, Comm.Math.Phys.104 (1986) 177.

A complete set of independent loops exist such that:

- one pion loop has wrapping number $\mathbf{n} = (n_x, n_y, n_z) \neq \mathbf{0}$;
- all other loops have zero wrapping number.



$$= \mathcal{O}\left(e^{-m_\pi L|\mathbf{n}|}\right)$$

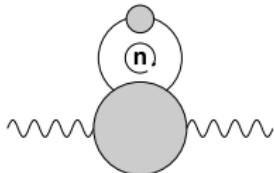
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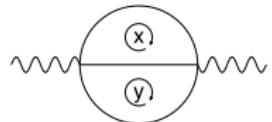
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At least two loops wrap around the spatial torus, e.g.



$$\sqrt{2 + \sqrt{3}} \simeq 1.93$$

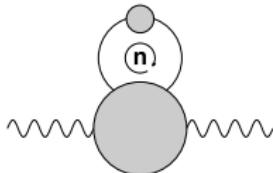
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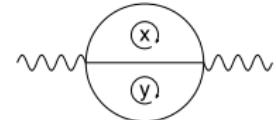
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- Leading finite-temperature corrections. For full analysis see Hansen, AP, arXiv:1904.10010. Subleading (with respect to first line) if $T = 2L$.
- Subleading finite-temperature corrections, and mixed corrections. Subleading (with respect to first line) if $T = 2L$.
- Contributions from kaon loops wrapping around the torus. Subleading (with respect to first line) at the physical point $m_K \simeq 3.5m_\pi$.

Finite- L corrections, $T = \infty$

$$\Delta a^{\text{HVP}}(L) = - \sum_{\mathbf{n} \neq 0} \int \frac{dp}{2\pi} \frac{e^{-|\mathbf{n}|L\sqrt{m_\pi^2 + p^2}}}{24\pi|\mathbf{n}|L} \int_0^\infty dx_0 K(x_0) \times \\ \times \int \frac{dk}{2\pi} \cos(x_0 k) \operatorname{Re} T(-k^2, -kp) + \mathcal{O}(e^{-1.93m_\pi L})$$

Forward Compton scattering amplitude $\pi\gamma^* \rightarrow \pi\gamma^*$

$$T(k^2, kp) = i \lim_{\mathbf{p}' \rightarrow \mathbf{p}} \sum_{q=0, \pm 1} \int d^4x e^{ikx} \langle \mathbf{p}', q | T \mathcal{J}_\rho(x) \mathcal{J}^\rho(0) | \mathbf{p}, q \rangle .$$

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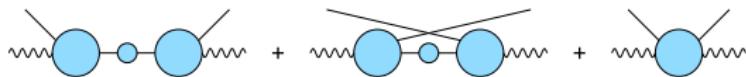
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- ▶ Reference to the effective field theory has disappeared, everything is now defined in terms of the fundamental theory (QCD).
- ▶ We can use this formula for a numerical estimate of finite- L effects!
Use a model for the Compton scattering amplitude, calculate integrals and sum numerically.

Estimates

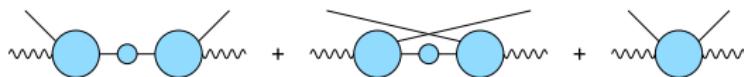
Compton scattering amplitude in the space-like region

$$T(-\mathbf{k}^2, -\mathbf{kp}) = \frac{2(4m_\pi^2 + \mathbf{k}^2) F_\pi^2(\mathbf{k}^2)}{\mathbf{k}^2 + 2\mathbf{p}\mathbf{k} - i\epsilon} + \frac{2(4m_\pi^2 + \mathbf{k}^2) F_\pi^2(\mathbf{k}^2)}{\mathbf{k}^2 - 2\mathbf{p}\mathbf{k} - i\epsilon} + T^{\text{reg}}(-\mathbf{k}^2, -\mathbf{kp})$$



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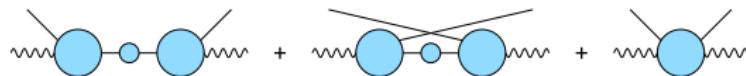


Pole contribution. Space-like electromagnetic form factor of the pion, phenomenological fit
 D. Brömmel *et al.* [QCDSF/UKQCD Collaboration], *The Pion form-factor...*, Eur. Phys. J. C 51 (2007) 335.

$$F_\pi(\mathbf{k}^2) = \left(1 + \frac{\mathbf{k}^2}{M^2}\right)^{-1}, \quad M = 727 \text{ MeV}$$

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Regular contribution. NLO chiral perturbation theory

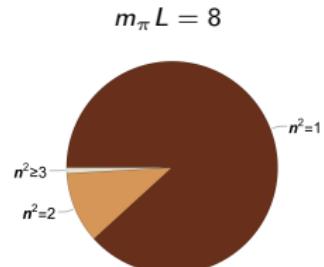
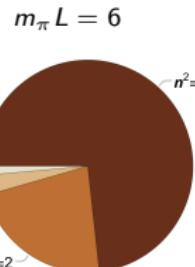
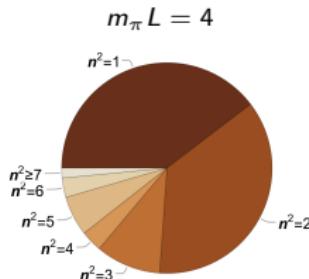
$$T^{\text{reg}}(-\mathbf{k}^2, -\mathbf{kp}) = c_0 + c_1 \mathbf{k}^2 + \frac{7m_\pi^2 + 4\mathbf{k}^2}{6\pi^2 f_\pi^2} \sqrt{1 + \frac{4m_\pi^2}{\mathbf{k}^2}} \cotg^{-1} \sqrt{1 + \frac{4m_\pi^2}{\mathbf{k}^2}}$$

Estimates

$$m_\pi = 137 \text{ Mev} \quad m_\mu = 106 \text{ MeV} \quad M = 727 \text{ MeV} \quad a = 700 \times 10^{-10}$$

$m_\pi L$	$-10^2 \frac{\Delta a(L)}{a}$	pole (form factor)	regular
4	3.19	3.168	0.024
5	1.43	1.424	0.006
6	0.631	0.6300	0.0015
7	0.275	0.2744	0.0004
8	0.118	0.1178	0.0001

Contribution of various $e^{-m_\pi L|\mathbf{n}|}$ to the pole term (i.e. form-factor).



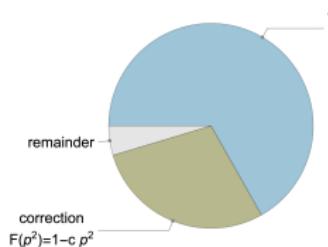
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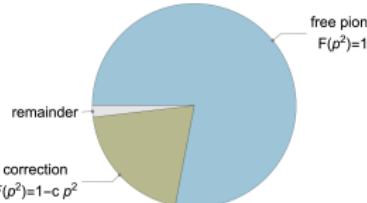
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Contribution of various terms of $F_\pi(k^2) = 1 - k^2/M^2 + \Delta F_\pi(k^2)$.

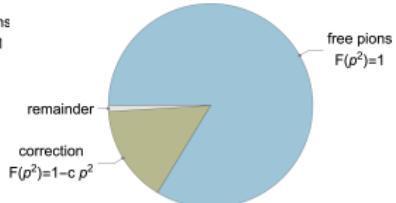
$$m_\pi L = 4$$



$$m_\pi L = 6$$



$$m_\pi L = 8$$



Some comments – Conclusions

- ▶ We studied the finite-volume effects of a particular estimator of a^{HVP} , i.e.

$$a^{\text{HVP}}(T, L) \rightarrow \int_0^{T/2} dx_0 K(x_0) \int_0^L d^3x \sum_{\mu=1}^3 \langle j_\mu(x_0, \mathbf{x}) j_\mu(0) \rangle_{T,L}$$

Different estimators will have different finite-volume effects.

- ▶ Finite- L corrections are exponentially suppressed and the leading exponential accounts only for $\sim 1/3$ of the effect at $m_\pi L = 4$.
- ▶ The space-like form factor is the natural quantity that enters in the finite- L effects, not the time-like one (but one can always use dispersion relation).
- ▶ The pion-structure contribution amounts to $\sim 1/3$ of the finite- L correction at $m_\pi L = 4$.
- ▶ Because of the noise at large distance, one may need to cut the integral over x_0 , i.e.

$$a^{\text{HVP}}(T, L) \rightarrow \int_0^{\tilde{x}_0} dx_0 K(x_0) \int_0^L d^3x \sum_{\mu=1}^3 \langle j_\mu(x_0, \mathbf{x}) j_\mu(0) \rangle_{T,L}$$

The contribution to the finite-volume effects of each time-slice is analyzed in
[arXiv:2004.03935](https://arxiv.org/abs/2004.03935).

- ▶ Finite- T corrections are analyzed in [arXiv:2004.03935](https://arxiv.org/abs/2004.03935), and they are negligible in standard setups $T \geq 2L$.