

Energy-like observables for the chiral phase transition of $2 + 1$ flavor QCD

Mugdha Sarkar
(HotQCD Collaboration)

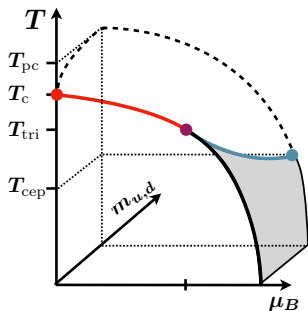
Asia-Pacific Symposium for Lattice Field Theory
(APLAT 2020)



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Chiral phase transition in QCD phase diagram

Chiral phase transition in the limit of vanishing up and down quark masses



[F. Karsch, arxiv:1905.03936]

Expected to belong to the **universality class** of $3d O(4)$ spin model

[R.D. Pisarski, F. Wilczek, PRD 29 338 (1984)]

- ⇒ Determination of the order and nature of the chiral phase transition
- ⇒ Imprint of the criticality on the thermodynamics at physical quark masses

$$T_c = 132^{+3}_{-6} \text{ MeV}$$

Anirban's talk

$$T_{pc} = 156.5(1.5) \text{ MeV}$$

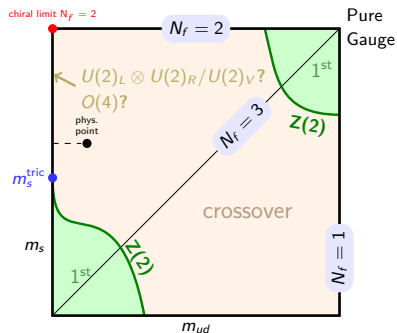
[HotQCD, PRL 123, 062002 (2019)]

[HotQCD, PLB 795 (2019) 15–21]



Exploring the chiral phase transition

- measurements towards the chiral limit with strange mass fixed at physical value
- investigate behavior of conserved charge fluctuations, which are also measurable in HIC, in the chiral limit
- alternative scenario : $3d$ Ising universality transition at non-zero mass, with a first-order chiral limit transition



[O. Philipsen and C. Pinke, PRD93, 114507, 2016]

- ⇒ Gauge ensembles generated with HISQ fermion discretization and Symanzik-improved gauge action, used in chiral T_c determination [HotQCD, arxiv:1905.11610].
- ⇒ Ensembles for smaller-than-physical quark (up, down) masses $m_l = m_s/27, m_s/40, m_s/80, m_s/160$, keeping strange quark mass m_s fixed at physical value.
- ⇒ Corresp. pion masses : 140 MeV, 110 MeV, 80 MeV, 55 MeV
- ⇒ Thermodynamic and continuum limit not yet performed. Measurements done at the largest simulated volumes for each mass at fixed time extent $N_\tau = 8$.
- ⇒ Computing resources : Jülich, Piz Daint, JLAB, Bielefeld and Wuhan supercomputing facilities.

Critical behavior of thermodynamic quantities I

Analysis of QCD thermodynamics in the vicinity of the chiral phase transition, acc. to Wilson's RG approach

infinite
volume

$$\frac{p}{T^4} = \frac{1}{VT^3} \ln Z(T, \vec{\mu}, m_l) = h^{(2-\alpha)/\beta\delta} f_f(z) + f_r(T, \vec{\mu}, m_l), \quad z \equiv t/h^{1/\beta\delta}$$

singular

regular \rightarrow no linear m_l
dependence

$$t = \frac{1}{t_0} \left(\frac{T-T_c}{T_c} + \kappa_2^X \left(\frac{\mu_X}{T} \right)^2 \right)$$

"energy-like" coupling

$$h = \frac{1}{h_0} \frac{m_l}{m_s}$$

X=B,Q,S

"magnetic-like" coupling

critical exponents

Chiral phase transition at
 $m_l \equiv m_u = m_d = 0$ ($h = 0$)
 $T = T_c$ ($t = 0$) at $\mu = 0$

	α	β	δ
$O(4)$	-0.21	0.38	4.82
$O(2)$	-0.017	0.349	4.78
$Z(2)$	+0.109	0.325	4.8

Critical behavior of thermodynamic quantities II

(Anirban's talk)

Conserved charge
fluctuations and

Polyakov loop should
be energy-like w.r.t
chiral phase transition

magnetic-like

$$\frac{\partial^2 \ln Z}{\partial h^2}$$

$$\sim h^{1/\delta-1}$$

$$\sim h^{-0.79}$$

divergence : **strong**

mixed

$$\frac{\partial^2 \ln Z}{\partial h \partial t}$$

$$\sim h^{(\beta-1)/\beta\delta}$$

$$\sim h^{-0.34}$$

moderate

energy-like

$$\frac{\partial^2 \ln Z}{\partial t^2}$$

$$\sim h^{-\alpha/\beta\delta}$$

$$\sim h^{+0.11}$$

vanishes

Conserved charge fluctuations at $\mu = 0$ (Singular part) :

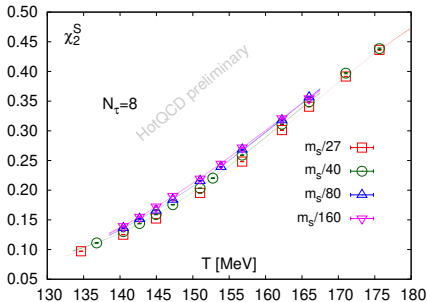
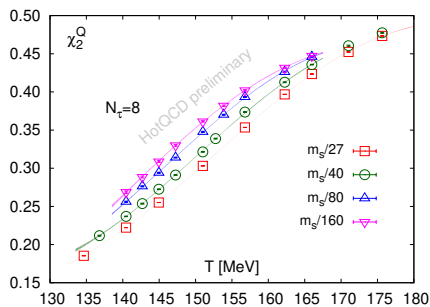
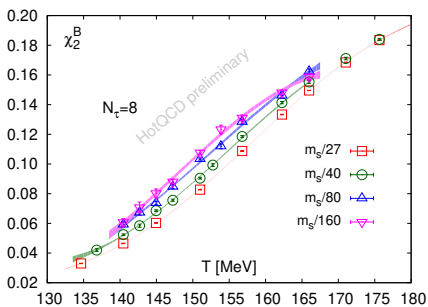
$$\chi_{2n}^X = - \left. \frac{\partial^{2n} f / T^4}{\partial (\mu_X / T)^{2n}} \right|_{\mu_X=0} \sim - (2\kappa_2^X)^n h^{(2-\alpha-n)/\beta\delta} f_f^{(n)}(z)$$

Scaling Expectation : $\frac{\partial}{\partial T} \sim \kappa_2^X \frac{\partial^2}{\partial \mu^2}$

$\Rightarrow \chi_2 \sim$ Energy density, $\chi_4 \sim$ Specific heat

Conserved charge fluctuations as energy-like observables

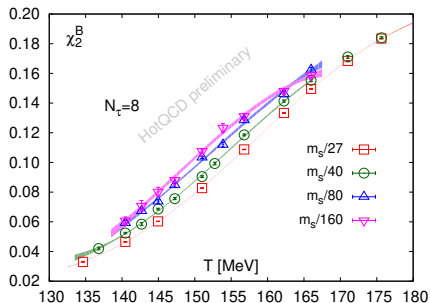
Second order charge fluctuations χ_2



– similar features as energy density

Estimation of singular contribution to χ_2

$$\chi_2^X(T_c, m_l) \sim -\kappa_2^X h^{(1-\alpha)/\beta\delta} f_f^{(1)}(0) + \text{const. reg. term} + \mathcal{O}(m_l^2)$$

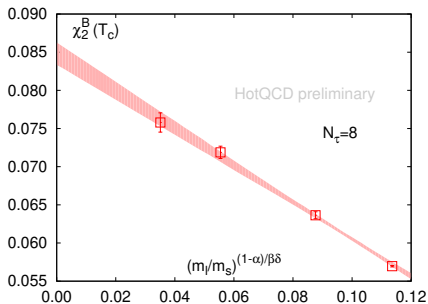


$$h = \frac{1}{h_0} \frac{m_l}{m_s}$$

- expect straight line for $\chi_2(T_c, h)$ vs $h^{0.66}$ if scaling holds ($O(4)$ exponents)
- $T_c \sim 144$ MeV for $N_\tau = 8$

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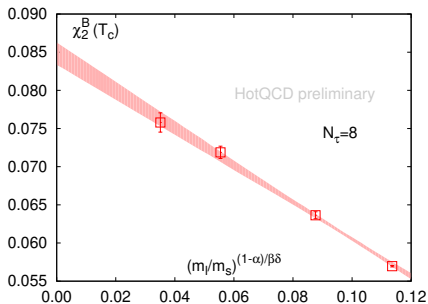
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– $T_c \sim 144$ MeV for $N_\tau = 8$

$$\chi_2^X(T_c, m_l = 0) - \chi_2^X(T_c, m_l = m_s/27) = \text{Singular part of } \chi_2^X(T_c)$$

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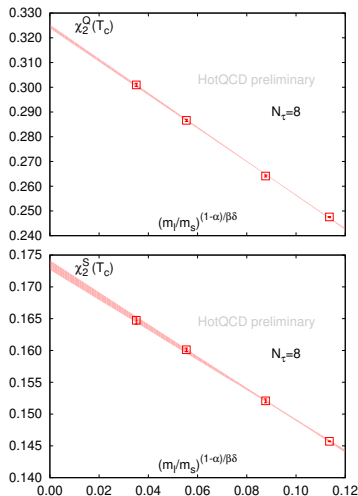
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Singular contribution to χ_2^B at physical masses $\sim 50\%$

Estimation of singular contribution to χ_2

$$\chi_2^X(T_c, m_l) \sim -\kappa_2^X h^{(1-\alpha)/\beta\delta} f_f^{(1)}(0) + \text{const. reg. term} + \mathcal{O}(m_l^2)$$



singular contribution
at physical mass

$$\chi_2^B(T_c) \sim 50\%$$

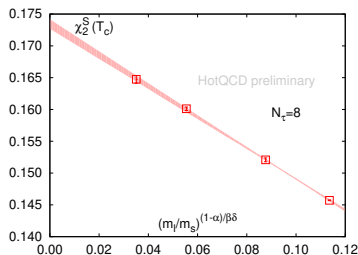
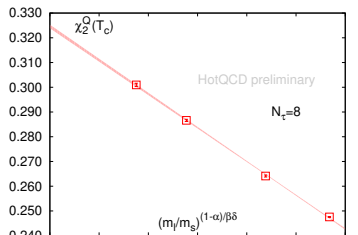
$$\chi_2^Q(T_c) \sim 30\%$$

$$\chi_2^S(T_c) \sim 20\%$$

► Linear in H plot

Estimation of singular contribution to χ_2

$$\chi_2^X(T_c, m_l) \sim -\kappa_2^X h^{(1-\alpha)/\beta\delta} f_f^{(1)}(0) + \text{const. reg. term} + \mathcal{O}(m_l^2)$$



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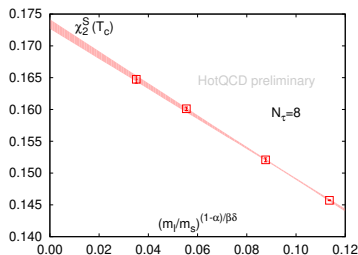
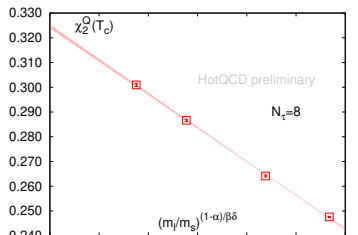
– ratio of singular parts = ratio of κ_2

$$\kappa_2^Q / \kappa_2^B \sim 2.6$$

$$\kappa_2^B / \kappa_2^S \sim 1.0$$

Estimation of singular contribution to χ_2

$$\chi_2^X(T_c, m_l) \sim -\kappa_2^X h^{(1-\alpha)/\beta\delta} f_f^{(1)}(0) + \text{const. reg. term} + \mathcal{O}(m_l^2)$$



singular contribution
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– ratio of singular parts = ratio of κ_2

$$\kappa_2^Q / \kappa_2^B \sim 2.6 \quad 1.8(8)^*$$

$$\kappa_2^B / \kappa_2^S \sim 1.0 \quad 0.9(4)^*$$

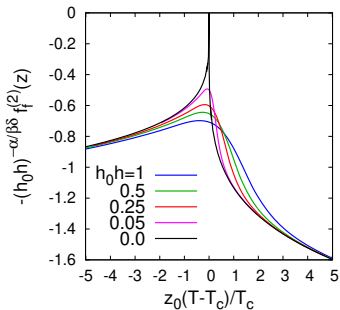
Close to results for physical mass

*[HotQCD, Phys. Lett. B 795 (2019) 15]

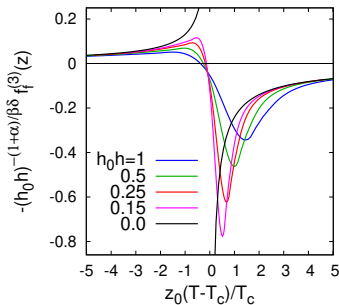
Scaling expectation of higher fluctuations

Derivatives of the singular part for 3d $O(4)$ universality class

$$h^{0.116} f_f^{(2)}$$



$$h^{-0.429} f_f^{(3)}$$



[B. Friman, F. Karsch, K. Redlich, V. Skokov, Eur. Phys. J. C (2011) 71:1694]

$$\sim \chi_4$$

$$\sim \chi_6$$

Scaling behavior of χ_4^Q

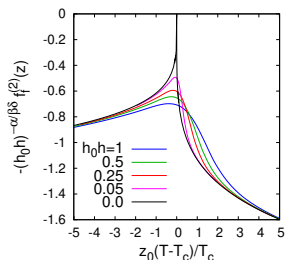
Singular part :

$$\chi_4^Q \sim h^{-\alpha/\beta\delta} f_f^{(2)}(z)$$

$$O(4) : -\alpha/\beta\delta = 0.116$$

Not divergent

– but pronounced spike



Scaling behavior of χ_4^Q

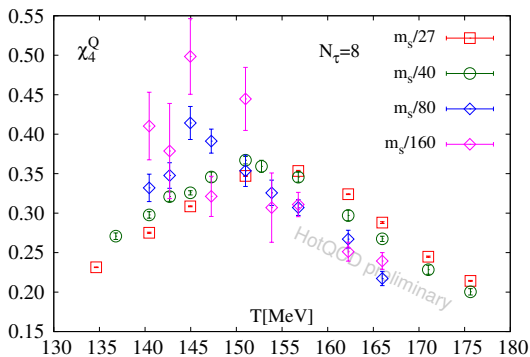
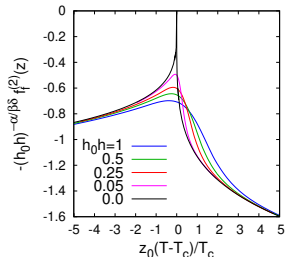
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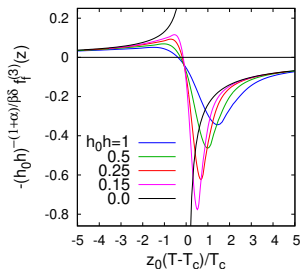


⇒ Expected features are apparent for $m_l = m_s/27, m_s/40$

Singular part :

$$\begin{aligned}\chi_6^Q &\sim h^{-(1+\alpha)/\beta\delta} f_f^{(3)}(z) \\ &\sim h^{-0.429}\end{aligned}$$

Moderate divergence



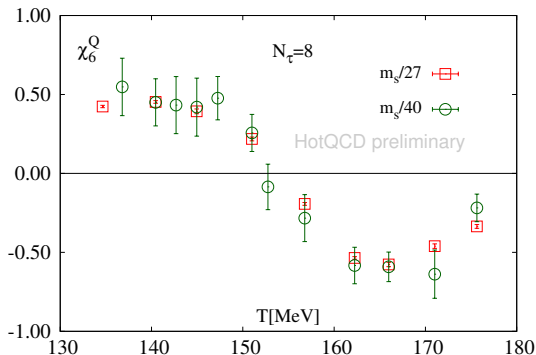
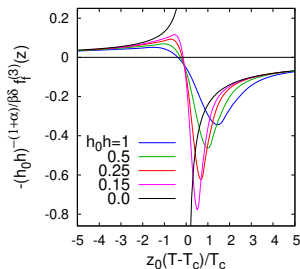
Scaling behavior of χ_6^Q

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$$\chi_6^Q \sim h^{-(1+\alpha)/\beta\delta} f_f^{(3)}(z)$$

$$\sim h^{-0.429}$$

Moderate divergence

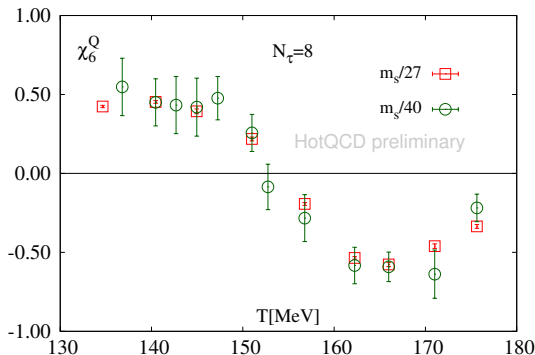
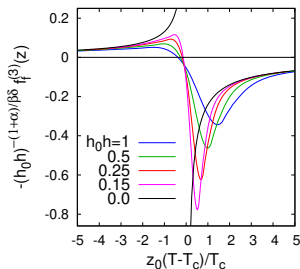


Scaling behavior of χ_6^Q

Singular part :

$$\chi_6^Q \sim h^{-(1+\alpha)/\beta\delta} f_f^{(3)}(z)$$
$$\sim h^{-0.429}$$

Moderate divergence



Ratio of peak heights expected from scaling : $(\chi_6^Q)_{1/40}^{max} / (\chi_6^Q)_{1/27}^{max} \sim 1.18$

Polyakov loop as energy-like observable

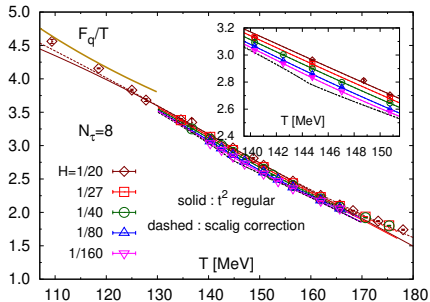
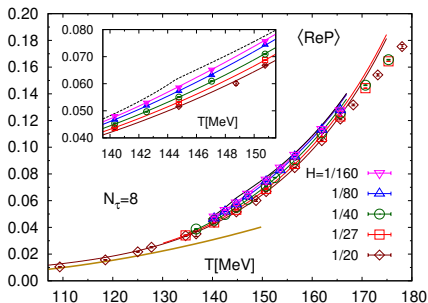
Polyakov loop observables

Polyakov loop $\langle \text{Re} P \rangle$

Heavy quark free energy

$$F_q/T \equiv -\log \langle \text{Re} P \rangle$$

– purely gluonic and invariant
under chiral transformation

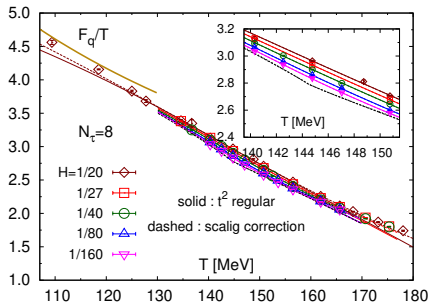
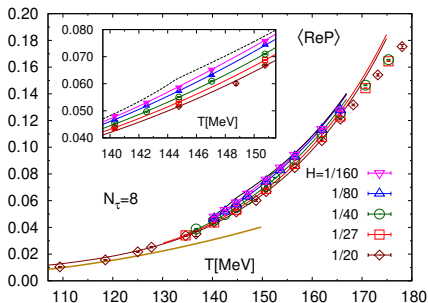


Polyakov loop $\langle \text{Re } P \rangle$

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$$F_q/T \equiv -\log \langle \text{Re } P \rangle$$

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$$F_q/T = AH^{(1-\alpha)/\beta\delta} f'_f(z) + f_{\text{reg}}$$

- fits with singular and regular terms
- singular term with universal $O(2)$ scaling function and critical exponents

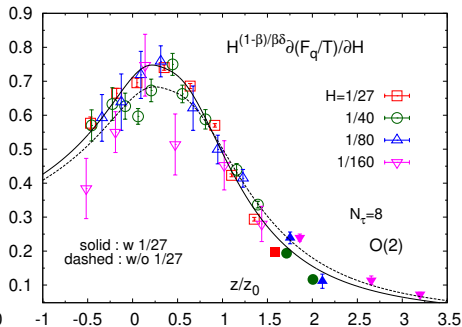
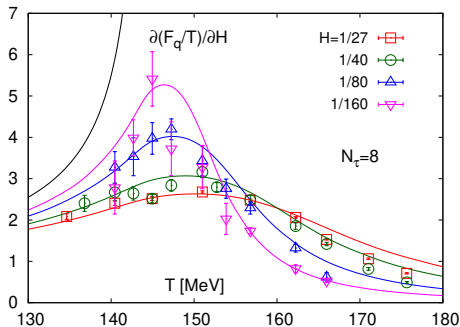
Derivative of Heavy Quark Free Energy

$$\frac{1}{T} \frac{\partial F_q}{\partial H} = -AH^{(\beta-1)/\beta\delta} f'_G(z) + \frac{\partial}{\partial H} f_{\text{reg}}$$

$$H = \frac{m_l}{m_s}$$

Divergent at 2nd order

$$\frac{\beta-1}{\beta\delta} = -0.39, O(2)$$



– consistent with $O(2)$ scaling

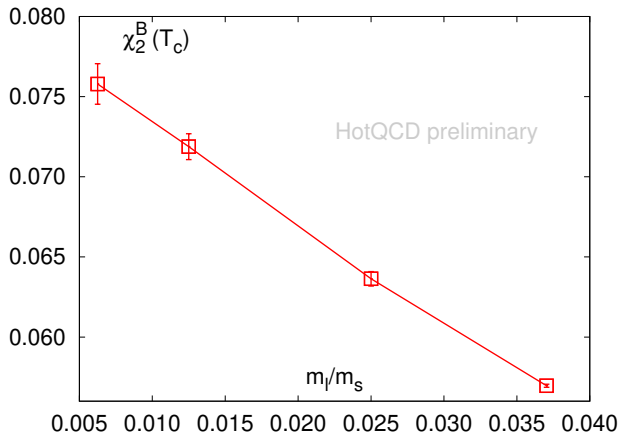
Conclusions and Outlook

- ⇒ Fluctuations of conserved charges are affected by the chiral phase transition belonging to $O(4)$ or $O(2)$ universality class. Expected energy-like behavior of fluctuations w.r.t. chiral phase transition.
- ⇒ Singular part can be extracted from χ_2^X and may be used to determine the curvature coefficients of the chiral critical line. Singular contribution to 2nd order fluctuations roughly $\leq 50\%$
- ⇒ Polyakov loop observables affected by the chiral phase transition at physical mass.
- ⇒ Future comparison with HRG with smaller masses obtained from ChPT.
- ⇒ Requires more statistics at lower masses and proper continuum and thermodynamic limits.

Conclusions and Outlook

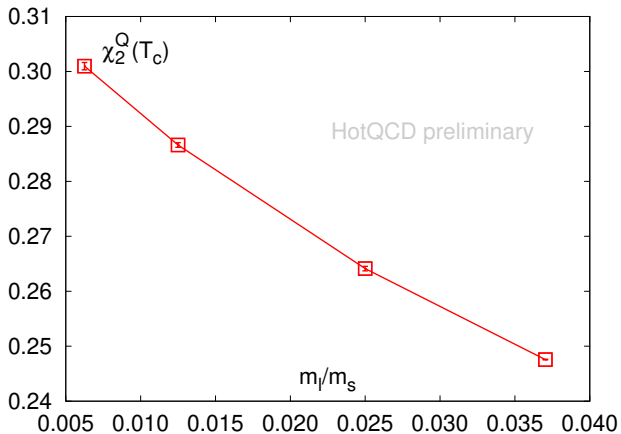
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Thank you for your attention



Points are connected to guide the eye. Lines are NOT fits.

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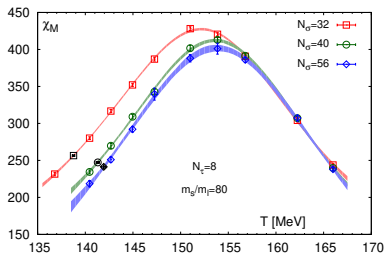
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Symmetry transformations

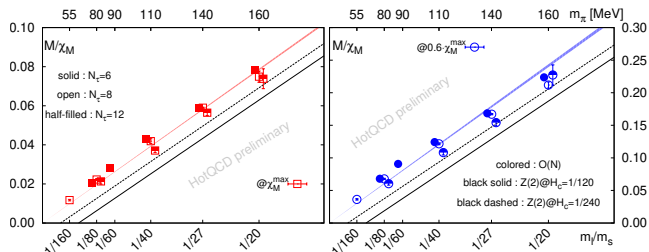
$$\begin{array}{ccccc}
 \chi_{5,\text{con}} & \pi : \bar{\mathbf{q}} \gamma_5 \frac{\tau}{2} \mathbf{q} & \xleftrightarrow{\text{SU}(2)_L \times \text{SU}(2)_R} & \sigma : \bar{\mathbf{q}} \mathbf{q} & \chi_{\text{con}} + \chi_{\text{disc}} \\
 & \updownarrow \text{U}(1)_A & & \updownarrow \text{U}(1)_A & \\
 \chi_{\text{con}} & \delta : \bar{\mathbf{q}} \frac{\tau}{2} \mathbf{q} & \xleftrightarrow{\text{SU}(2)_L \times \text{SU}(2)_R} & \eta : \bar{\mathbf{q}} \gamma_5 \mathbf{q} & \chi_{5,\text{con}} - \chi_{5,\text{disc}}
 \end{array}$$

$O(4)$ criticality : Bounds on first order transition

[HotQCD, arxiv:1905.11610]



- ⇒ No indications of first order volume scaling at $m_\pi = 80$ MeV
- ⇒ Consistent with $O(4)$ volume scaling at non-zero H



Seems no indication of a 1st order transition till $m_\pi \sim 46$ MeV

[H.T. Ding et al, Nucl. Phys. A 982 (2019) 211]

$O(4)$ criticality : Chiral T_c estimation

Chiral phase transition at the chiral limit extracted by $O(4)$ scaling arguments.

$$T_X(H) = T_c \left(1 + \frac{z_X}{z_0} H^{1/\beta\delta} \right)$$

$$X = \delta, 60, pc$$

$$\frac{H\chi_M(T_\delta, H)}{M(T_\delta, H)} = \frac{1}{\delta}$$

$$\chi_M(T_{60}, H) = 0.6\chi_M^{max}$$

$$\mathbf{T_c = 132_{-6}^{+3} \text{ MeV}}$$

[HotQCD, arxiv:1905.11610]

