

Lepton Flavor Violations from soft SUSY breaking terms in modular flavor models

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based on

“Soft supersymmetry breaking terms and
lepton flavor violations in modular flavor model”

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Introduction

- ▶ The origin of flavor is one of the important questions in particle physics.

- **Non-Abelian flavor symmetries** are interesting approach,

$S_3, A_4, S_4, A_5, \text{ etc}$

A_4 is minimal to embed three families of leptons
in one irreducible rep. [E. Ma and G. Rajasekarn, PRD (2001)]

- ▶ **Modular symmetry** is a new direction of flavor symmetry approach.

[Ferugulio, 1706.08749]

- Modular symmetry arises from the compactification of higher dimensions in superstring theory.

- Modular groups Γ_N are isomorphic to the finite groups,

$$\Gamma_2 \simeq S_3, \Gamma_3 \simeq A_4, \Gamma_4 \simeq S_4, \Gamma_5 \simeq A_5,$$

[Kobayashi et al, PRD (2018)]

[Tanedo, Petcov, NPB (2019)]

[Novichkov et al, JHEP (2019)]

- **The modulus τ** which characterizes the shape of compact space plays an important role to determine flavor structure.

- **The lepton mass matrix can be given by the modular forms.**

(Flavons are not necessary in modular symmetry models)

Introduction

- ▶ Flavor symmetries also control flavor structure of superpartners.

[Ko, et al PRD (2008), Ishimori et al PRD (2008), etc]

- Specific patterns appears in soft SUSY breaking terms.
- Such specific patterns can be observed in LFV processes like $\mu \rightarrow e + \gamma$ if SUSY particles are light.

[See e.g. Kobayashi and Vives, PLB (2001)]

- ▶ In modular symmetry models, the F -term of the modulus τ can be non-vanishing, and lead to SUSY breaking.
 - Such SUSY breaking terms show specific patterns of modular symm.

In this talk,

- ▶ Flavour structure of soft SUSY breaking terms from the modulus F -term in simple modular A_4 models.
- ▶ $\mu \rightarrow e + \gamma$ decay to see flavor structure and parameter dependence.

Modular Symmetry

- ▶ Modular transformations γ acting on the modulus τ

$$\tau \longrightarrow \gamma\tau = \frac{a\tau + b}{c\tau + d} \quad \text{where } a, b, c, d \in \mathbb{Z} \text{ and } ad - bc = 1, \quad \text{Im}[\tau] > 0$$

- ▶ Chiral superfield with modular weight $-k_I$ transforms as

$$\phi^{(I)} \longrightarrow (c\tau + d)^{-k_I} \rho^{(I)}(\gamma) \phi^{(I)}, \quad \rho^{(I)}(\gamma) : \text{unitary rep. matrix}$$

- ▶ Holomorphic functions which **transform under modular trans.**, are called **modular form** with weight k

$$f(\tau) \longrightarrow (c\tau + d)^k f(\tau),$$

- ▶ Superpotential can be formed using chiral superfields and modular forms

$$\mathcal{W} = f(\tau) \phi^{(I_1)} \phi^{(I_2)} \dots \phi^{(I_n)},$$

The superpotential is **invariant under the modular trans.** when

$$\rho^{(I_1)} \times \rho^{(I_2)} \dots \times \rho^{(I_n)} = 1$$

$$k - k_{I_1} - k_{I_2} \dots - k_{I_n} = 0$$

Modular A_4 flavor model

- ▶ We consider two models in which the charged lepton Yukawa is given by A_4 triplet modular form with weight 2.

Model A : Weinberg operator

	L	(e^c, μ^c, τ^c)	$H_{u,d}$	$Y_{\mathbf{r}}^{(2)}, Y_{\mathbf{r}}^{(4)}$
$SU(2)$	2	1	2	1
A_4	3	$(\mathbf{1}, \mathbf{1}'', \mathbf{1}')$	1	3, {3, 1, 1'}
k	-2	0	0	2, 4

Neutrino superpotential

$$W_\nu = -\frac{1}{\Lambda} (H_u H_u L L Y_{\mathbf{r}}^{(4)})_1$$

Model B : type-I Seesaw

L	(e^c, μ^c, τ^c)	ν^c	$H_{u,d}$	$Y_{\mathbf{3}}^{(2)}$
2	1	1	2	1
3	$(\mathbf{1}, \mathbf{1}'', \mathbf{1}')$	3	1	3
-1	-1	-1	0	2

Neutrino superpotential

$$W_\nu = g (L Y_{\mathbf{3}}^{(2)} \nu^c)_1 H_u + \Lambda (Y_{\mathbf{3}}^{(2)} \nu^c \nu^c)_1$$

Charged lepton superpotential

$$W_e = \alpha (L Y_{\mathbf{3}}^{(2)})_1 e_R^c H_d + \beta (L Y_{\mathbf{3}}^{(2)})_{1'} \mu_R^c H_d + \gamma (L Y_{\mathbf{3}}^{(2)})_{1''} \tau_R^c H_d$$

α, β, γ : fixed by charged lepton masses

Charged Lepton Yukawa

- ▶ The triplet modular form with weight 2

$$Y_3^{(2)} = \begin{pmatrix} Y_1 \\ Y_2 \\ Y_3 \end{pmatrix} \quad \text{satisfying} \quad Y_2^2 + 2Y_1Y_3 = 0$$

- ▶ Charged lepton Yukawa

$$W_e = \alpha(LY_3^{(2)})_1 e_R^c H_d + \beta(LY_3^{(2)})_{1'} \mu_R^c H_d + \gamma(LY_3^{(2)})_{1''} \tau_R^c H_d$$

$$\supset l_{Ri} Y_{ijk} H_j l_{Lk}$$

$$\rightarrow Y_{ijk} = \text{diag}[\alpha, \beta, \gamma] \begin{pmatrix} Y_1 & Y_3 & Y_2 \\ Y_2 & Y_1 & Y_3 \\ Y_3 & Y_2 & Y_1 \end{pmatrix}_{RL}, \quad (i, k = e, \mu, \tau, j = H_d)$$

- ▶ Y_i ($i=1,2,3$) is a function of the modulus τ .

$$Y_3^{(2)} = \begin{pmatrix} Y_1 \\ Y_2 \\ Y_3 \end{pmatrix} = \begin{pmatrix} 1 + 12q + 36q^2 + 12q^3 + \dots \\ -6q^{1/3}(1 + 7q + 8q^2 + \dots) \\ -18q^{2/3}(1 + 2q + 5q^2 + \dots) \end{pmatrix} \quad \text{with} \quad q = e^{2\pi i\tau}$$

Once τ is determined, the Yukawa is fixed.

Soft SUSY breaking terms

- ▶ Soft SUSY breaking terms originating from the modulus F -term are given in supergravity theory. ($M_p=1$)

$$\mathcal{L}_{\text{soft}} = \tilde{m}_L^2 \tilde{L}^\dagger \tilde{L} + \tilde{m}_e^2 \tilde{e}_R^\dagger \tilde{e}_R + \underline{(Y A \tilde{L}^\dagger H \tilde{e}_R + h.c.)}$$

charged lepton Yukawa

soft mass

$$\tilde{m}_i^2 = m_{3/2}^2 - k_i \frac{|F^\tau|^2}{(2\text{Im}(\tau))} = m_0^2$$

k_i is common for 3 flavor.

flavor universal

A-term

$$A_{ijk} = A_{ijk}^0 + A'_{ijk}$$

where

$$A_{ijk}^0 = (1 - k_i - k_j - k_k) \frac{F^\tau}{(2\text{Im}(\tau))} = A^0$$

flavor universal

$$A'_{ijk} = \frac{F^\tau}{Y_{ijk}} \frac{dY_{ijk}(\tau)}{d\tau}$$

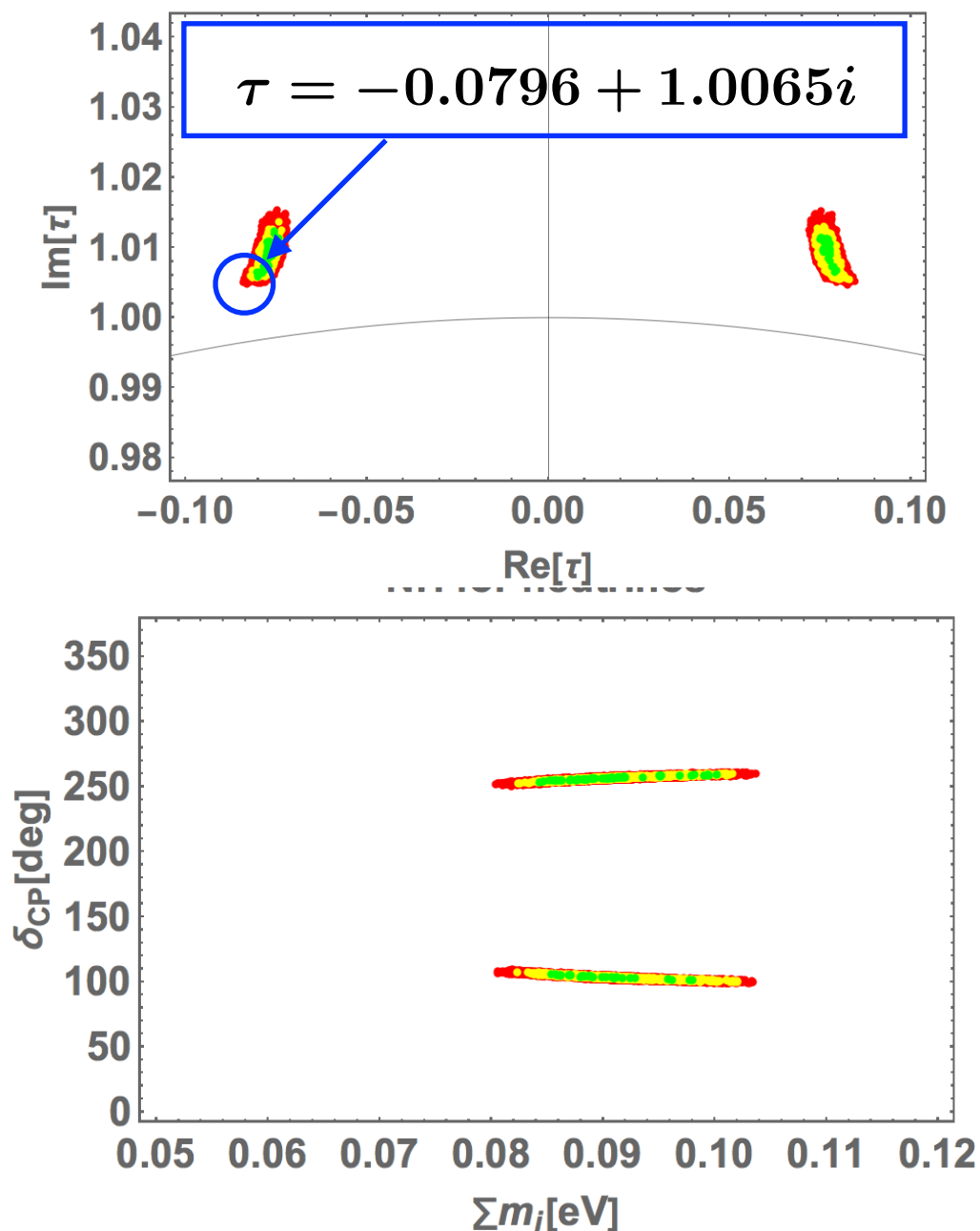
flavor dependent

Modular form (Y) determines the flavor structure of A term

Allowed region of Modulus

- ▶ Neutrino Yukawa is also given by modular forms.
- ▶ The modulus is determined by neutrino oscillation data in each model

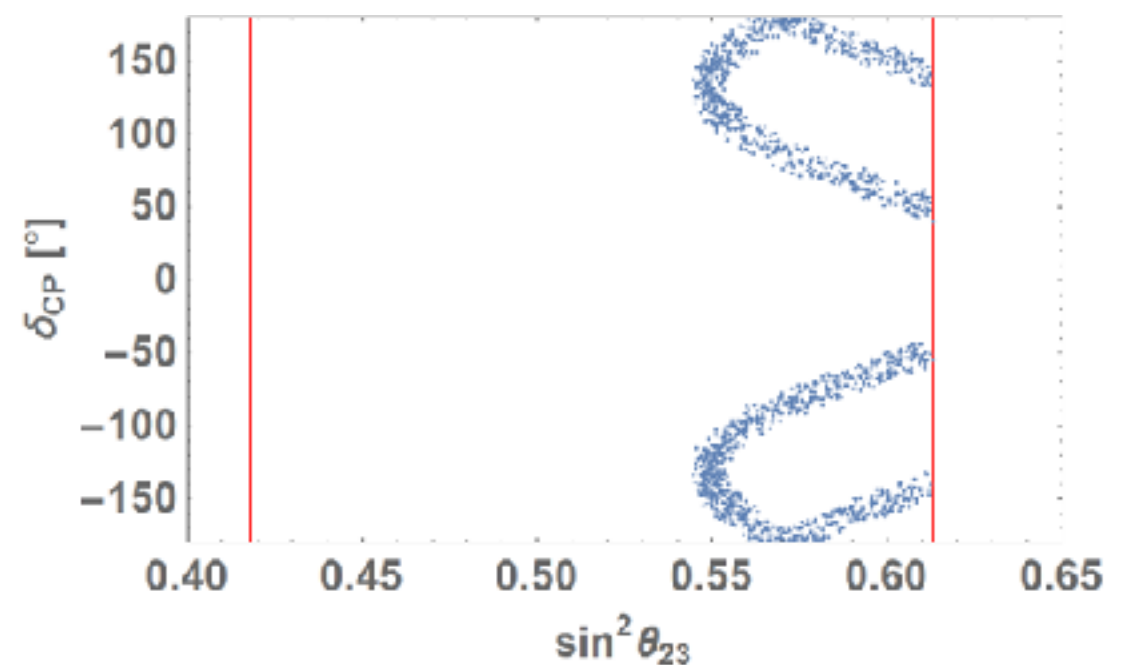
Model A : Weinberg operator
NH for neutrinos



Model B : type-I Seesaw

$\text{Im}[\tau]$	$\text{Re}[\tau]$
$0.66 - 0.73$	$\pm(0.25 - 0.31), \pm(0.46 - 0.54)$
$1.17 - 1.32$	$\pm(0.66 - 0.75), \pm(1.25 - 1.31), \pm(1.46 - 1.50)$

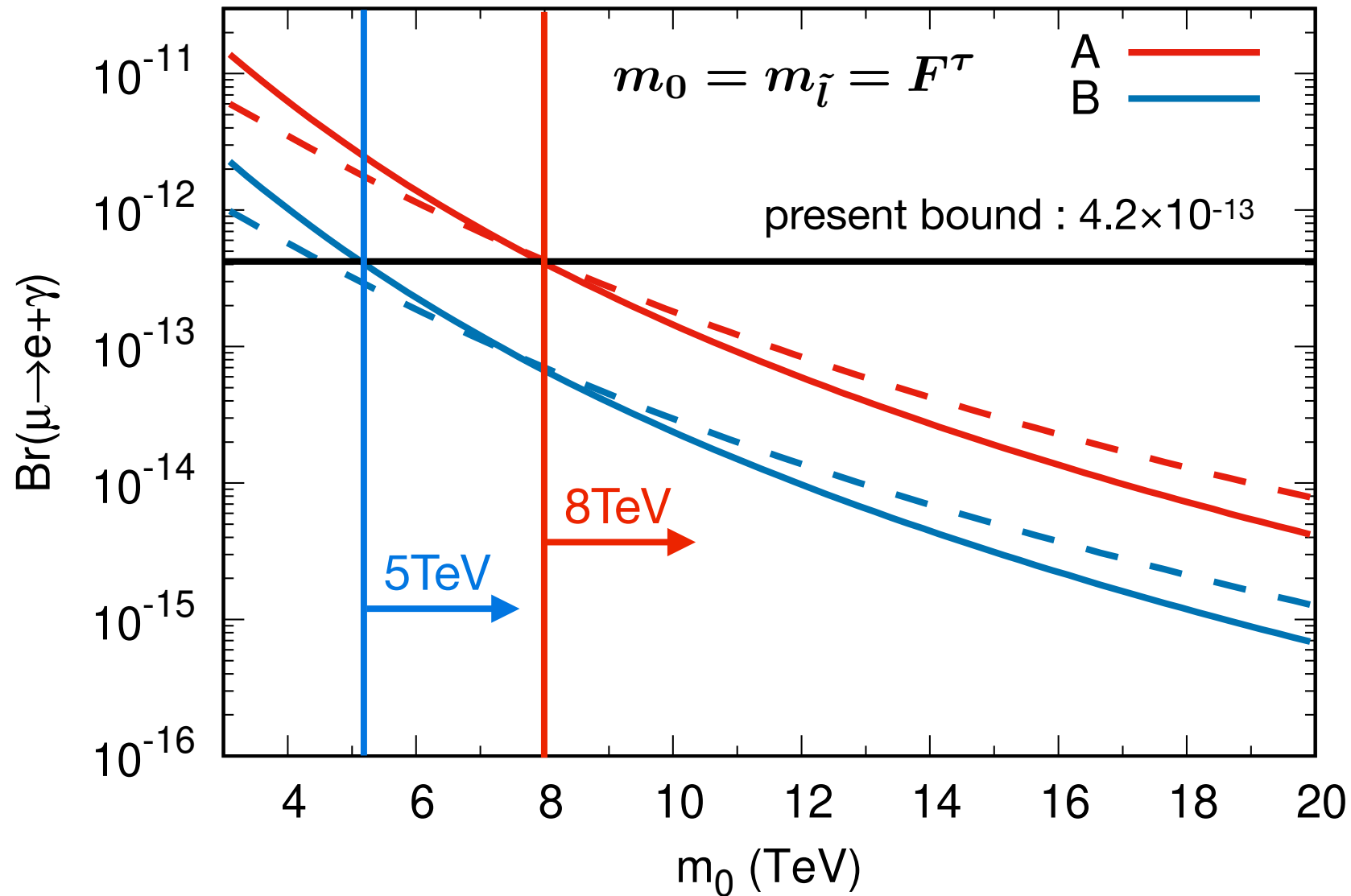
$\tau = 0.48151 + 1.30262i$



SUSY mass scale dependence

A: $\tau = -0.0796 + 1.0065i$
 $\alpha/\gamma = 6.82 \times 10^{-2}$, $\beta/\gamma = 1.02 \times 10^{-3}$

B: $\tau = 0.48151 + 1.30262i$
 $\alpha/\gamma = 2.03 \times 10^2$, $\beta/\gamma = 3.30 \times 10^3$



gaugino mass

solid:

$$M_1 = 3 \text{ TeV}$$

dashed:

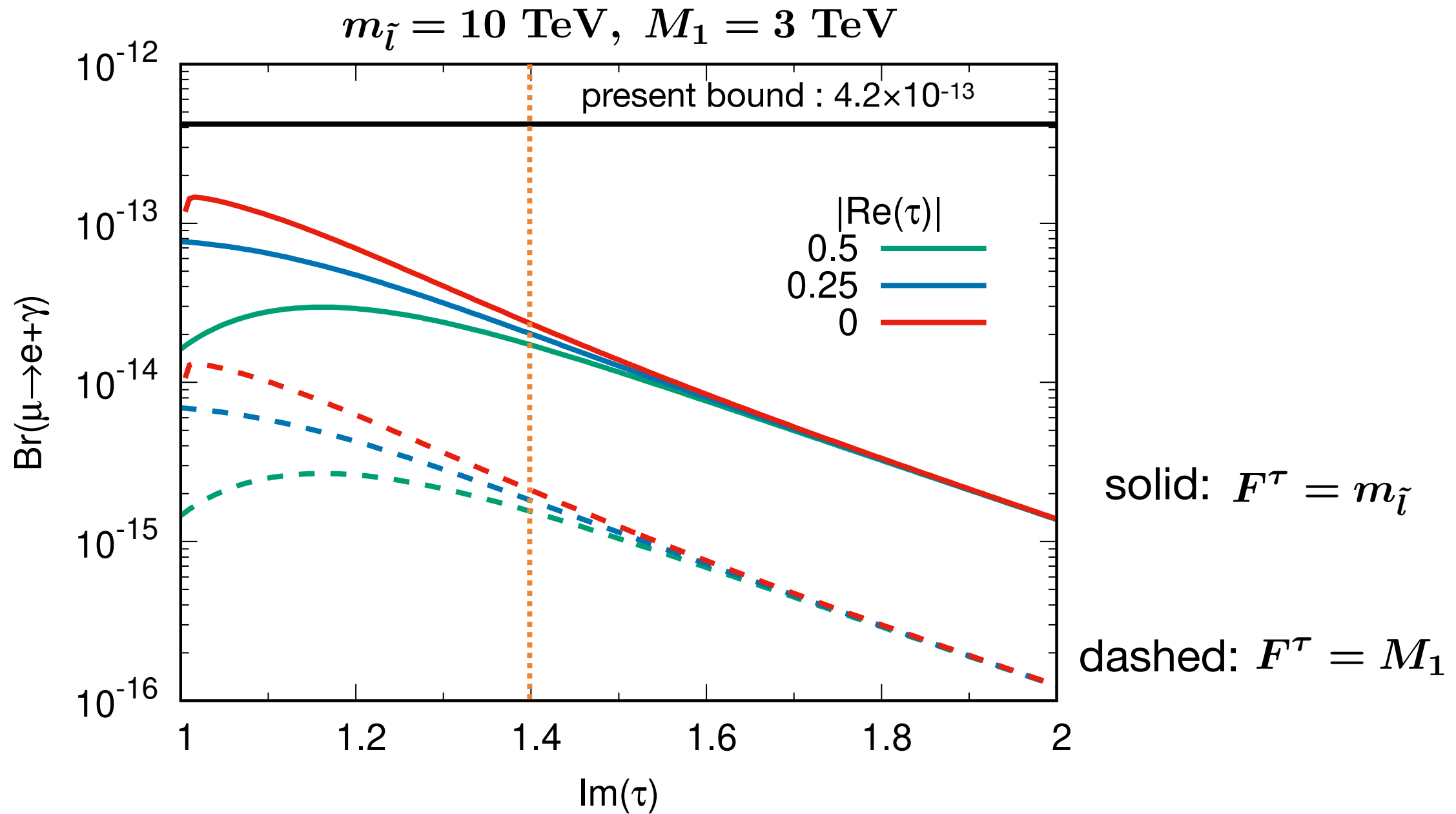
$$M_1 = 5 \text{ TeV}$$

with $\tan \beta = 5$

Mass insertions

A: $|(\delta_\ell^{RL})_{\mu e}| \simeq 2.1 \times 10^{-5} \left(\frac{F^\tau}{10 \text{ TeV}} \right)$ **B:** $|(\delta_\ell^{RL})_{\mu e}| \simeq 8.4 \times 10^{-6} \left(\frac{F^\tau}{10 \text{ TeV}} \right)$
 $|(\delta_\ell^{LR})_{\mu e}| \simeq 9.7 \times 10^{-8} \left(\frac{F^\tau}{10 \text{ TeV}} \right)$ $|(\delta_\ell^{LR})_{\mu e}| \simeq 3.7 \times 10^{-8} \left(\frac{F^\tau}{10 \text{ TeV}} \right)$

Modulus dependence



Depends on both $\text{Re}(\tau)$ and $\text{Im}(\tau)$ below $\text{Im}(\tau) < 1.4$

Summary

We have studied the SUSY breaking and LFV in modular flavor models,

- ▶ Soft SUSY breaking terms from the modulus F-term are obtained.
- ▶ Flavor structure of the SUSY breaking terms are determined by modular form of the charged leptons.

and showed the LFV decay in two A_4 models,

- ▶ The SUSY mass scales are larger than 8 (model A) and 5 (B) TeV.
- ▶ The branching ratio significantly depends on τ for $\text{Im}(\tau) < 1.4$, and is independent of $\text{Re}(\tau)$ for larger $\text{Im}(\tau)$.

Messages are

- ▶ Similar and detailed analyses are important in other flavor models.
- ▶ Specific patterns of soft SUSY br. terms are studied in LFV process.