



Effects of Attractive KK^{bar} and Repulsive KK Interactions in KKK^{bar} three-body resonance

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Hadron physics

Hadrons as composite systems of hadrons
(Hadron molecules) have been discussed intensively
Physically, hadron number is not a good quantum number
So, the concept of 'compositeness' is important to characterize hadrons

Few-body physics

Resonances in coupled-channel systems
Three-body force
Relativistic effects
Efimov physics



$KK\bar{K}-K\pi\pi-K\pi\eta$ three-body system

This system is very interesting

Three-body boson system

Coupled-channel problem

Strongly interacting : Attractive KK^{bar} and Repulsive KK

Relativistic vs non-relativistic

$K(1460)$ (0^- , $I=1/2$)

Near to (just below) the threshold of $KK\bar{K}$ (1685 MeV)

The same quantum numbers with those of ground state of $KK\bar{K}$

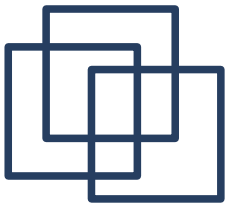
Experiments

G.W. Brandenburg, et al., PRL36(1976)1239

C. Daum, et al., NPB187(1981)1

Theories as a $KK\bar{K}$ resonance

A.M. Torres, D. Jido, and Y. Kanada-Enyo, PRC83, 065205, 2011



Experiment on K(1460) at SLAC

G.W. Brandenburg, et al., PRL36(1976)1239

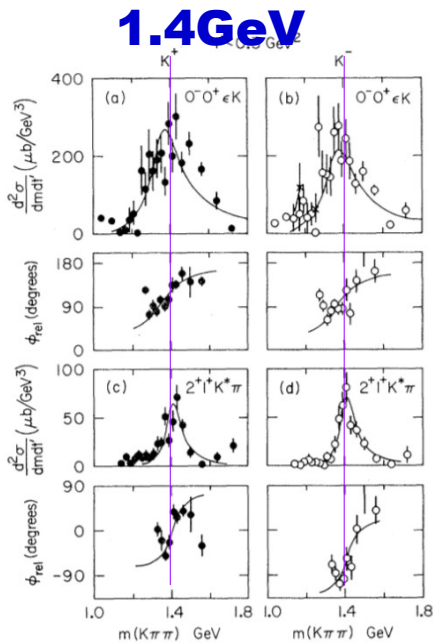


FIG. 1. The mass dependence of the $0^-0^+ \epsilon K$ and the $2^+1^+ K^* \pi$ waves. Crosses correspond to ambiguous solutions, and φ_{rel} is measured with respect to the $1^+0^+ \rho K$ wave. The curves correspond to simple Breit-Wigner parametrizations.

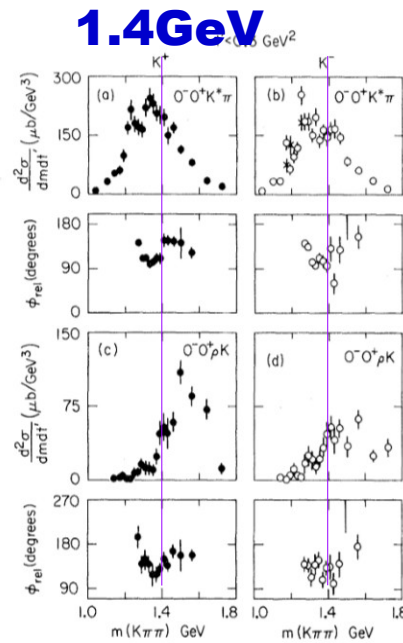


FIG. 2. The mass dependence of the $0^-0^+ K^* \pi$ and the $0^-0^+ \rho K$ waves. Crosses correspond to ambiguous solutions, and φ_{rel} is measured with respect to the $1^+0^+ \rho K$ wave.

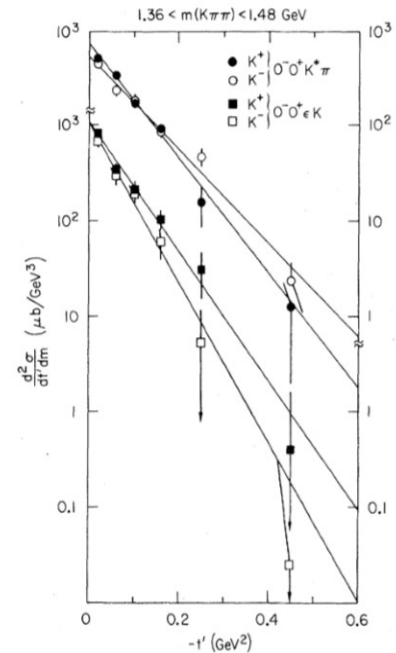


FIG. 3. The momentum transfer dependence of the $0^-0^+ K^* \pi$ and $0^-0^+ \epsilon K$ waves. The lines are exponentials with the parameter values given in Table I.



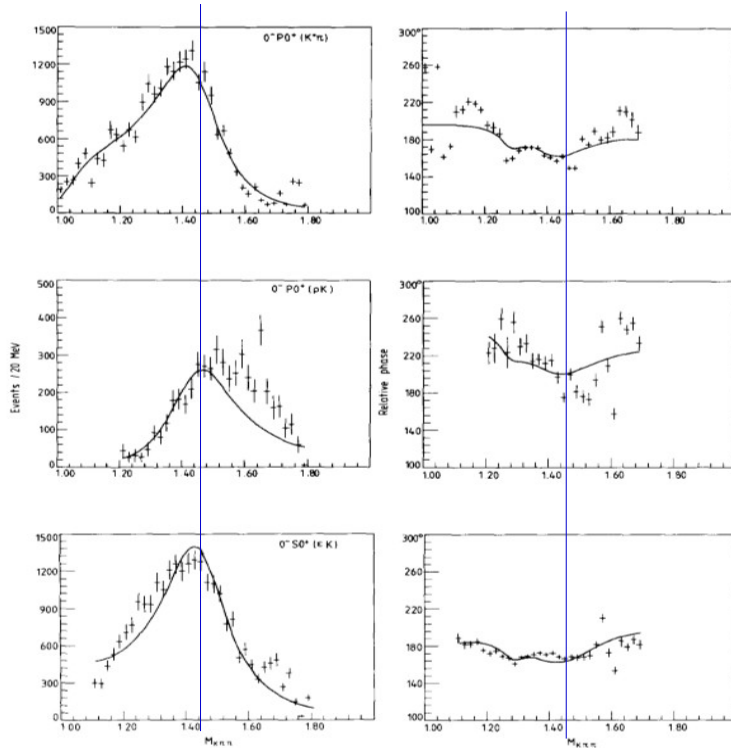
... We have presented evidence for the existence of a $K\pi\pi$ pseudoscalar resonance, the K' of mass $\sim 1400 \text{ MeV}$ and width $\sim 250 \text{ MeV}$. The predominant decay mode is ϵK with weaker couplings to $K^* \pi$ and, possibly, ρK



Experiment on K(1460) at SLAC

C. Daum, et al., NPB187(1981)1

1.46 GeV

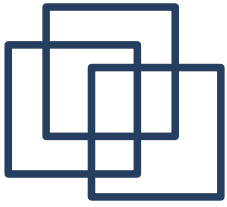


... The fits represent well the qualitative behaviour of the 0^- amplitudes and these data constitute strong evidence for the existence of a diffractively produced 0^- meson with a mass ~ 1.46 GeV and width ~ 260 MeV. The fitted partial widths are

$$\Gamma_{K^*\pi} = 109 \text{ MeV}, \Gamma_{\rho K} = 34 \text{ MeV}, \Gamma_{\varepsilon K} = 117 \text{ MeV}.$$

...

Fig. 18. Intensities and phases of 0^- waves for $1.0 \leq M_{K\pi\pi} \leq 1.8$ GeV, $0 \leq |t'| \leq 0.7$ GeV². The phases are measured relative to $1^+ S_0^+(K^*\pi)$. The curves show the result of fitting with a single resonance and a simple background in all three channels. The fitted phases have been corrected for the variation of the reference phase.



Theoretical works on K(1460)

A.M. Torres, D. Jido, and Y. Kanada-Enyo, PRC83, 065205, 2011

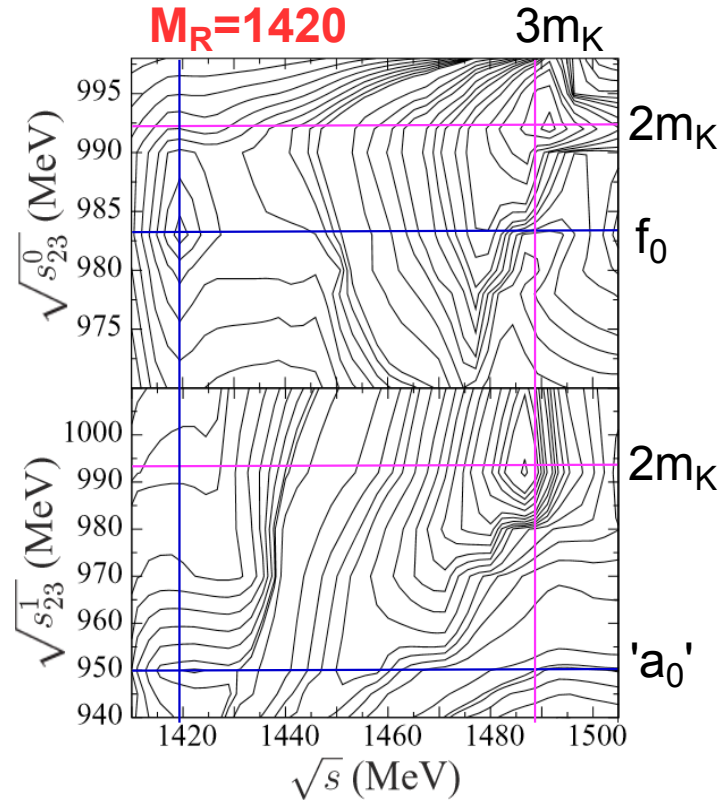


FIG. 2. Contour plots of the three-body squared amplitudes $|T_R^{(1/2,0)}|^2$ and $|T_R^{(1/2,1)}|^2$ for the $KK\bar{K} \rightarrow KK\bar{K}$ transition with total $I = 1/2$ as functions of the total three-body energy, \sqrt{s} , and the invariant mass of the $K\bar{K}$ subsystem with $I_{23} = 0$ (top) or the invariant mass of the $K\bar{K}$ subsystem with $I_{23} = 1$ (bottom).

TABLE I. Comparison of the results of the Faddeev calculation and the potential model. The spatial structure of the $KK\bar{K}$ quasi-bound state obtained with the potential model is also shown in the table.

| Model | Faddeev calculation | Potential Model |
|---|---------------------|-----------------|
| Mass (MeV) | ~ 1420 | 1467 |
| Width (MeV) | ~ 50 | 110 |
| Root mean squared radius (fm) | | 1.6 |
| K - K distance (fm) | | 2.8 |
| (KK) - \bar{K} distance (fm) | | 1.7 |
| K_2 - \bar{K}_3 distance (fm) ^a | | 1.6 |
| K_1 - $(K_2\bar{K}_3)$ distance (fm) ^a | | 2.6 |

^aThe values of the K_2 - \bar{K}_3 and K_1 - $(K_2\bar{K}_3)$ distances are obtained before making the symmetrization of K_1K_2 .

Faddeev calc.

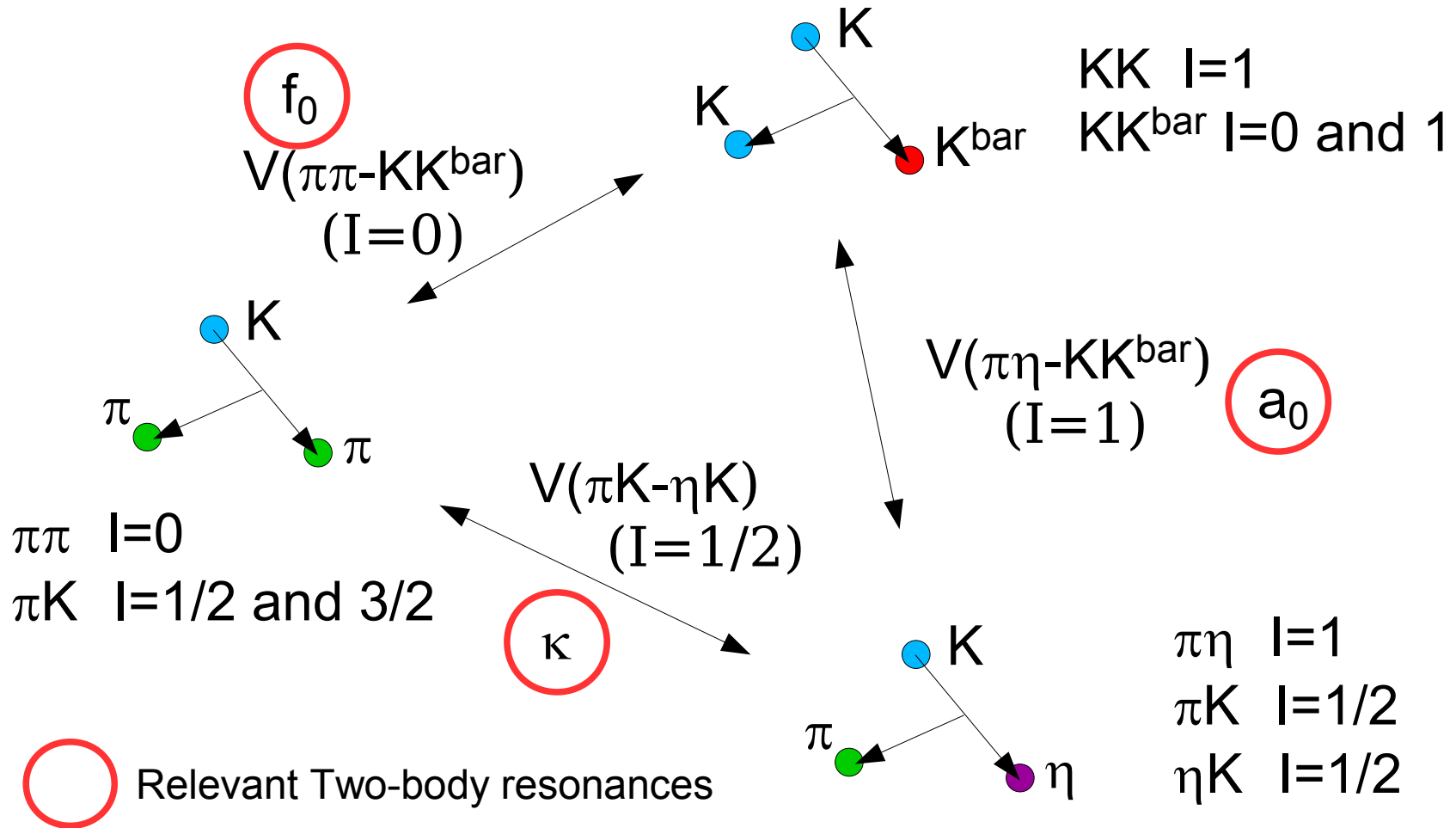
$M \sim 1420 \text{ MeV}$, $\Gamma \sim 50 \text{ MeV}$

Variational calc.

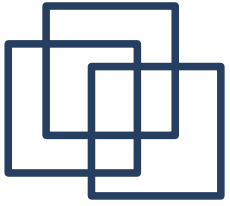
$M = 1467 \text{ MeV}$, $\Gamma = 110 \text{ MeV}$



Our model of Coupled-Channel $KK\bar{K}$ - $K\pi\pi$ - $K\pi\eta$ system



Total : $J=0, I=1/2, L=0$ (S-wave only)



Semi-Relativistic complex-scaled Coupled-Channel Three-Body Equation

A=1:KKK^{bar}, A=2:Kππ, A=3:Kπη

$$\hat{H}^{AA'} = \delta_{AA'} \sum_{i=1}^3 \sqrt{m_i^{A2} + \mathbf{p}_i^2} + \sum_{(i,j)} V_{ij}^{AA'}(r_{ij}), \quad \sum_i \mathbf{p}_i = 0$$



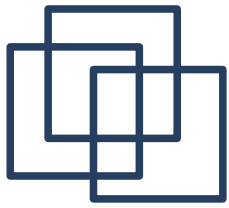
$\mathbf{r}_{ij} \rightarrow \mathbf{r}_{ij} e^{i\theta}$, $\mathbf{p}_i \rightarrow \mathbf{p}_i e^{-i\theta}$: complex rotation

$$\hat{H}_\theta^{AA'} = \delta_{AA'} \sum_{i=1}^3 \sqrt{m_i^{A2} + \mathbf{p}_i^2 e^{-2i\theta}} + \sum_{(i,j)} V_{ij}^{AA'}(r_{ij} e^{i\theta}), \quad \sum_i \mathbf{p}_i = 0$$

$$\sum_{A'} H_\theta^{AA'} \Psi^{A'} |A'\rangle = E_\theta \Psi^A |A\rangle \quad : \text{coupled-channel equation}$$

$$\Psi^A = \Phi_1^A(\boldsymbol{\rho}_1^A, \mathbf{R}_1^A) + \Phi_2^A(\boldsymbol{\rho}_2^A, \mathbf{R}_2^A) + \Phi_3^A(\boldsymbol{\rho}_3^A, \mathbf{R}_3^A)$$

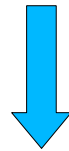
$\boldsymbol{\rho}_i^A, \mathbf{R}_i^A$: Jacobi coordinates in channel A



Three-Body Wave Function in the Gaussian Expansion Method

$$\Psi^A = \Phi_1^A(\rho_1^A, \mathbf{R}_1^A) + \Phi_2^A(\rho_2^A, \mathbf{R}_2^A) + \Phi_3^A(\rho_3^A, \mathbf{R}_3^A)$$

$$\Phi_i^A(\rho_i^A, \mathbf{R}_i^A) = \sum_{\alpha\beta} C_{(i\alpha\beta)}^A N_\alpha N_\beta \exp(-\rho_i^{A2}/\rho_\alpha^2) \exp(-R_i^{A2}/R_\beta^2)$$



Generalized Eigenvalue Problem

$$\sum_{B(j\alpha'\beta')} H_{\theta(i\alpha\beta),(j\alpha'\beta')}^{AB} C_{(j\alpha'\beta')}^B = E_\theta \sum_{(j\alpha'\beta')} N_{(i\alpha\beta),(j\alpha'\beta')}^A C_{(j\alpha'\beta')}^A$$

$H_{\theta(i\alpha\beta),(j\alpha'\beta')}^{AB}$: **complex symmetric** matrix (not Hermitian)

$N_{(i\alpha\beta),(j\alpha'\beta')}^A$: real symmetric matrix

 complex eigenvalues



Two-Body Meson-Meson Interactions in Coupled-Channel $KK\bar{K}-K\pi\pi-K\pi\eta$ Problem

We need meson-meson potentials:

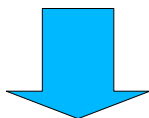
$$V_{KK-KK} (I=1)$$

$$V_{KK\bar{K}-KK\bar{K}}, V_{KK\bar{K}-\pi\pi}, V_{\pi\pi-\pi\pi} (I=0)$$

$$V_{KK\bar{K}-KK\bar{K}}, V_{KK\bar{K}-\pi\eta}, V_{\pi\eta-\pi\eta} (I=1)$$

$$V_{\pi K-\pi K}, V_{\pi K-\eta K}, V_{\eta K-\eta K} (I=1/2)$$

$$V_{\pi K-\pi K} (I=3/2)$$



Coupled-Channel One-range Gaussian potentials

$$V_{mm-mm}(r) = V_{mm-mm} \exp(-(r/r_G)^2)$$

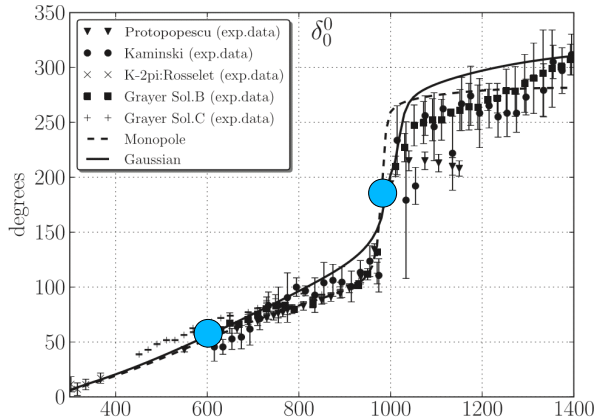
($V_{\pi\eta-\pi\eta} = V_{\pi K-\eta K} = V_{\eta K-\eta K} = 0$ is assumed in present calculation)

We try various ranges : $r_G = 0.3, 0.4, 0.5, 0.6, 0.7$ fm



Determination of $KK^{\text{bar}}-\pi\pi(I=0)$, $KK^{\text{bar}}-\pi\eta(I=1)$ potentials

I=0



$f_0: M=980\text{MeV}, \Gamma=70\text{MeV}$

Strongly attractive

$\delta_{L=0}^{I=0}(\pi\pi)$ at 600 MeV = 55 degree

| $r_G(\text{fm})$ | $V_{KK^{\text{bar}}-KK^{\text{bar}}}$ | $V_{KK^{\text{bar}}-\pi\pi}$ | $V_{\pi\pi-\pi\pi}$ |
|------------------|---------------------------------------|------------------------------|---------------------|
| 0.3 | -2006 | -417 | -1877 |
| 0.4 | -1327 | -383 | -1164 |
| 0.5 | -969 | -350 | -800 |
| 0.6 | -763 | -322 | -590 |
| 0.7 | -635 | -305 | -456 |

I=1

| $r_G(\text{fm})$ | $V_{KK^{\text{bar}}-KK^{\text{bar}}}$ | $V_{KK^{\text{bar}}-\pi\eta}$ | $V_{\pi\eta-\pi\eta}$ |
|------------------|---------------------------------------|-------------------------------|-----------------------|
| 0.3 | -1335 | -1099 | 0 |
| 0.4 | -928 | -686 | 0 |
| 0.5 | -695 | -480 | 0 |
| 0.6 | -550 | -362 | 0 |
| 0.7 | -456 | -288 | 0 |

$a_0: M=980\text{MeV}, \Gamma=70\text{MeV}$

KK^{bar} : strongly attractive

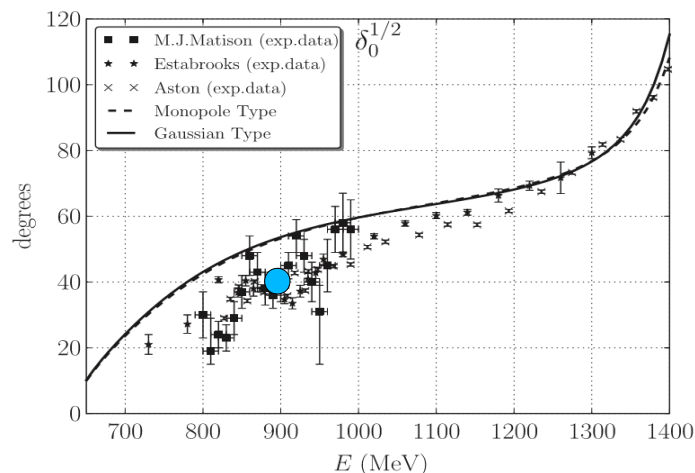
We assume $V_{\pi\eta-\pi\eta}=0$
(tentatively)

*Theoretically, $V_{\pi\eta-\pi\eta}$ may be weak



Determination of $\pi K(I=1/2)$, $\pi K(I=3/2)$ potentials

$I=1/2$

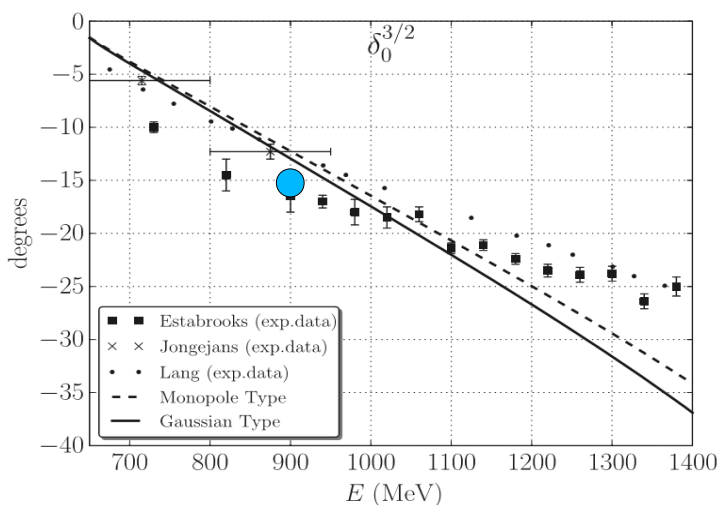


$$\delta_{L=0}^{I=1/2}(\sqrt{s}=900 \text{ MeV}) = 40 \text{ degree}$$

| rG(fm) | $V_{\text{KK}0}$ (MeV) |
|--------|------------------------|
| 0.3 | -1488.7 |
| 0.4 | -895.0 |
| 0.5 | -607.1 |
| 0.6 | -405.1 |
| 0.7 | -356.8 |

attractive

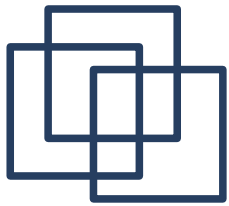
$I=3/2$



$$\delta_{L=0}^{I=3/2}(\sqrt{s}=900 \text{ MeV}) = -15 \text{ degree}$$

| rG(fm) | $V_{\text{KK}0}$ (MeV) |
|--------|------------------------|
| 0.3 | 3694.0 |
| 0.4 | 1174.0 |
| 0.5 | 555.6 |
| 0.6 | 323.8 |
| 0.7 | 215.2 |

repulsive



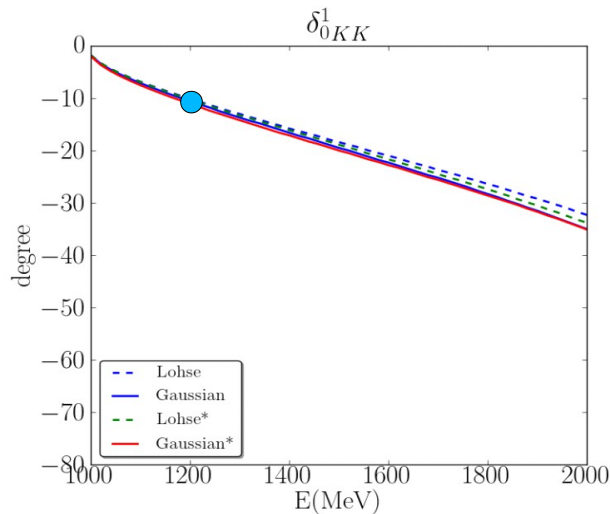
Determination of KK-KK(l=1) potentials

No experimental phase shifts

Theoretical predictions :

Repulsive interaction

$$\delta_{L=0}^{I=1}(\sqrt{s} = 1200 \text{ MeV}) = -10 \text{ degree}$$



| rG(fm) | V _{KK0} (MeV) |
|--------|------------------------|
| 0.3 | 903.0 |
| 0.4 | 366.3 |
| 0.5 | 196.4 |
| 0.6 | 124.5 |
| 0.7 | 88.3 |

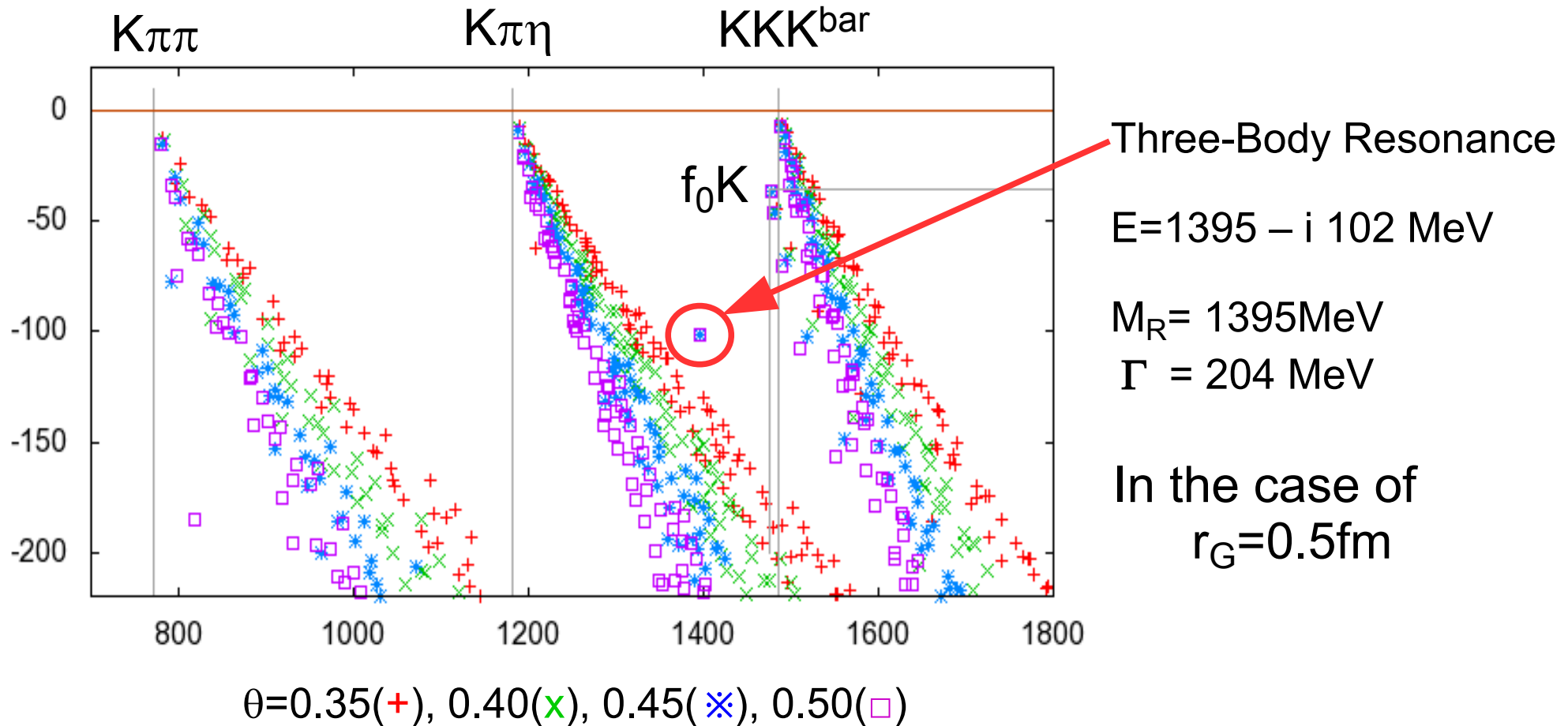
We introduce an artificial factor f

to discuss the effect of repulsive KK interaction

$$V_{KK-KK} = f \times V_{KK0} \exp(-(r/r_G)^2)$$



An example of Eigenvalues in Coupled-Channel $KK\bar{K}$ - $K\pi\pi$ - $K\pi\eta$ system





KKK^{bar}-K $\pi\pi$ -K $\pi\eta$ Three-Body Resonance Poles

$3m_K=1485$ MeV

Two-channel*
KKK^{bar}-K $\pi\pi$

Three-channel
KKK^{bar}-K $\pi\pi$ -K $\pi\eta$

For $r_G=0.7$ fm
no pole

| Range: r_G (fm) | Re(E) (MeV) | Im(E) (MeV) | Re(E) (MeV) | Im(E) (MeV) |
|----------------------|--------------------|-------------------|----------------|----------------|
| 0.3 | -- (1432.7) | -- (-42.8) | 1436.0 | -57.1 |
| 0.4 | 1464.7 (1406.5) | -48.7 (-59.0) | 1406.7 | -74.4 |
| 0.5 | 1457.7 (1400.8) | -72.0 (-86.1) | 1395.2 | -101.9 |
| 0.6 | 1465.3 (1405.6) | -87.3 (-110.2) | 1396.1 | -133.0 |
| 0.7 | -- (1440.2) | -- (-138.8) | -- | -- |

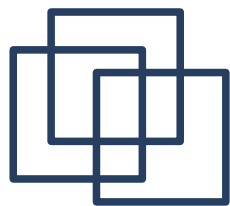
Pole position, especially
Imaginary part, depends
strongly on the potential
range r_G

A good fit to K(1460) for
 $r_G=0.5-0.6$ fm

K $\pi\eta$ channel is important
to explain the width

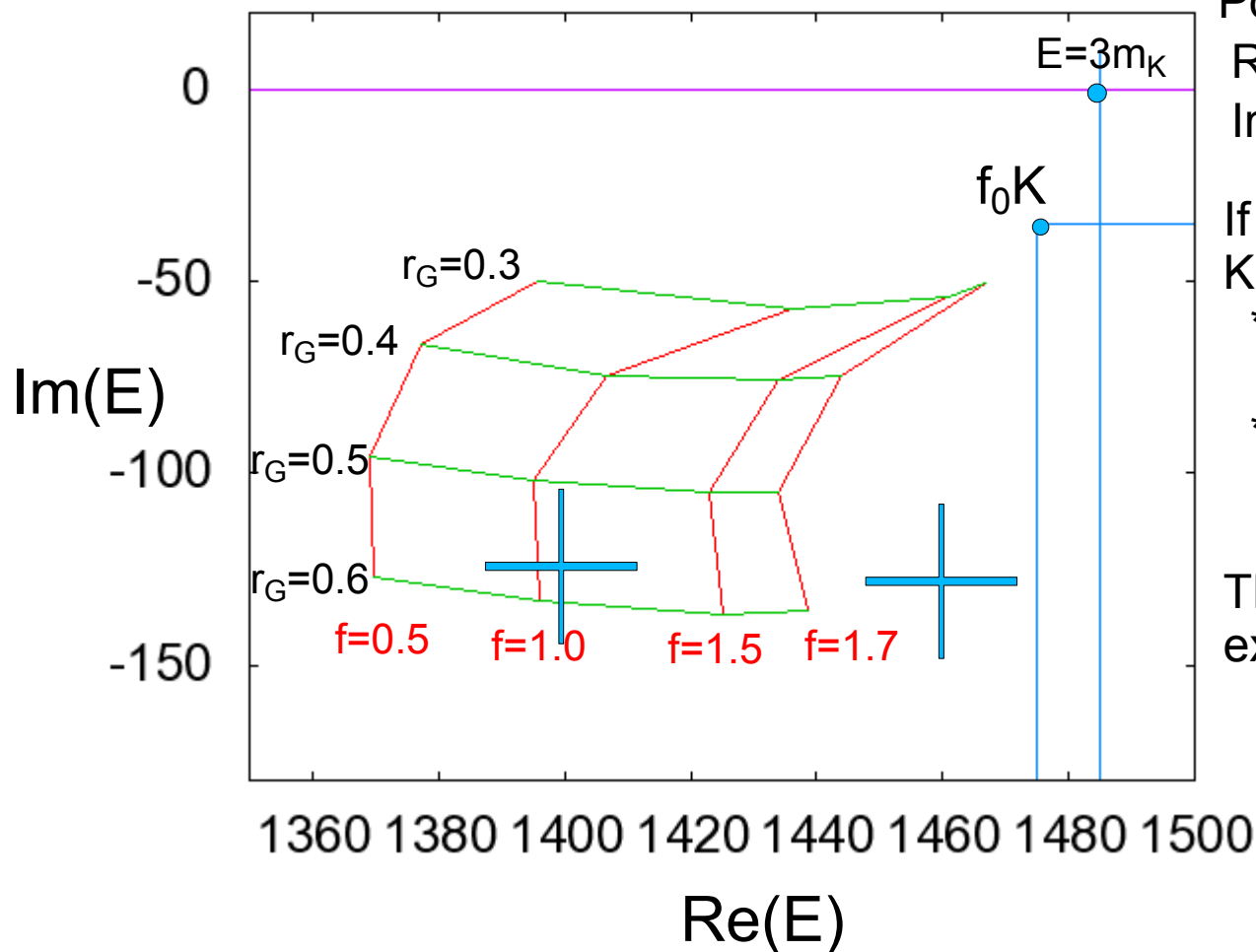
(): $V_{KKbar-KKbar}(l=1)=V_{KKbar-KKbar}(l=0)$ is assumed

[exp K(1460)] 1400 -125
 1460 -130



$KK^{\text{bar}}-K\pi\pi-K\pi\eta$ resonance pole positions Dependence on KK interaction

+ [exp K(1460)]



Pole position E_p

$\text{Re}(E_p)$ depends strongly on f
 $\text{Im}(E_p)$ depends strongly on r_G

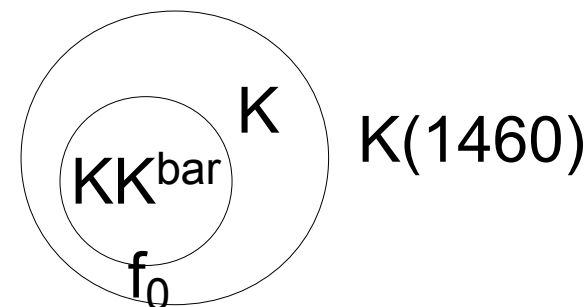
If this pole can be interpreted as K(1460),

*KK Interaction must be repulsive.

*KK Interaction range:
 $r_G \sim 0.5-0.6 \text{ fm}$ is supported.

The three-body resonance is expected to have a f_0K structure.

$$V_{KK-KK} = f \times V_{KK0} \exp(-(r/r_G)^2)$$





Summary

We solved the $KK\bar{K}-K\pi\pi-K\pi\eta$ three-meson problem, using the complex scaling method and found a three-body resonance pole.

Semi-relativistic Hamiltonian
Gaussian Expansion method

The resonance position depends strongly on the range of two-body potentials even if the potentials are adjusted to reproduce the same two-body properties.

This resonance can be interpreted as $K(1460)$ discovered in two SLAC experiments. For this interpretation, repulsive KK interaction is essential and the potential range $r_G=0.5-0.6\text{fm}$ is supported.

Future plans:

To use more realistic two-body interactions

Two- or Three-range gaussian potentials

Inclusion of $V_{\pi\eta-\pi\eta}, V_{\pi K-V\eta K}, V_{\eta K-\eta K}$

To introduce the P -wave components

To calculate the partial decay widths : $\varepsilon K, K^*\pi, \rho K$ modes



Thank you for your attention