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# *Flavor Symmetries and CLFV*

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**CP3**

**SDU** 

- flavor symmetries (and beyond)
  - why flavor symmetries (and beyond)?
  - which ones?
  - what is their predictive power?
- CLFV
  - in models with extra dimension
  - in supersymmetric models
- continuous flavor symmetries and CLFV
- phenomenological imprints beyond CLFV



***Flavor symmetries (and beyond)***

- in the Standard Model of particle physics most free parameters are in the **flavor sector**

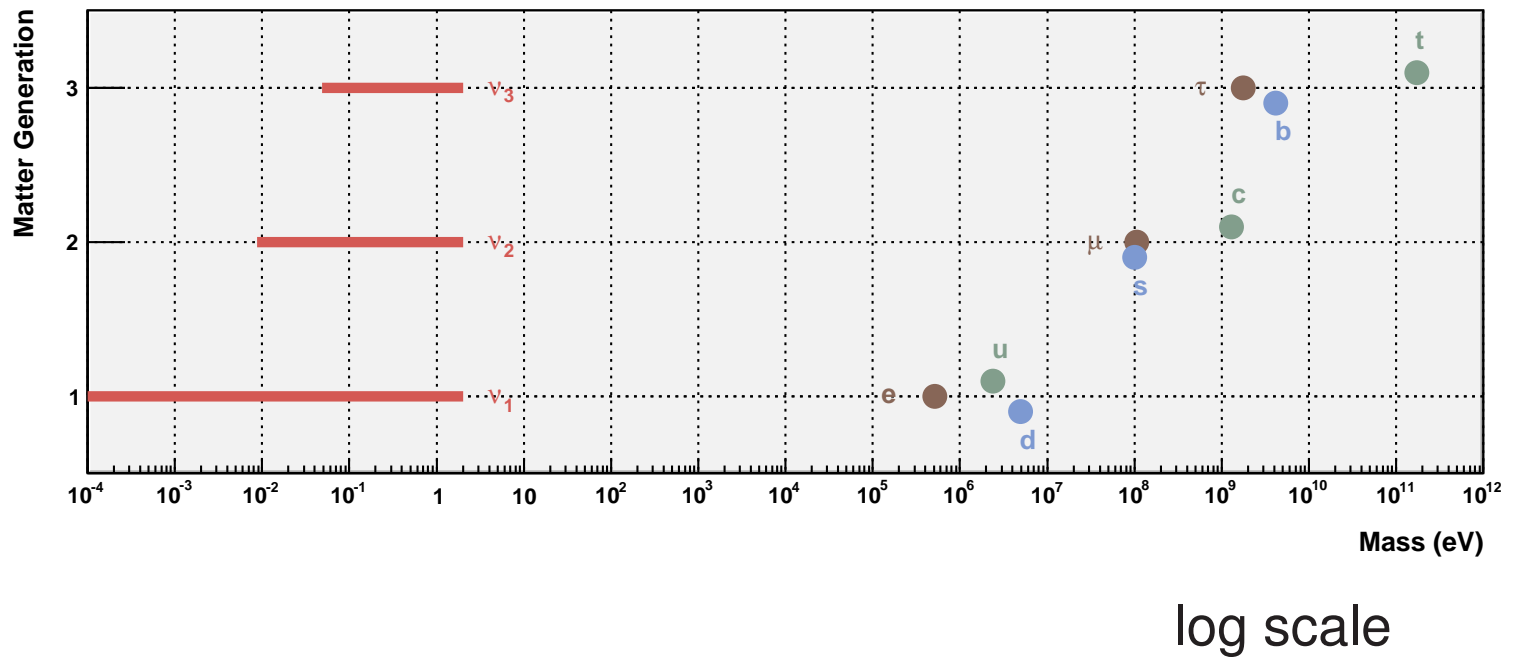
$$-\mathcal{L}_{\text{Yuk}} = y_{ij}^e \bar{L}_i H e_{Rj} + y_{ij}^u \bar{Q}_i H^* u_{Rj} + y_{ij}^d \bar{Q}_i H d_{Rj}$$

and coming from the mechanism which gives mass to neutrinos, e.g. the type-I seesaw mechanism

$$-\mathcal{L}_\nu = y_{ij}^\nu \bar{L}_i H \nu_{Rj} + \frac{1}{2} M_{ab} \overline{\nu_{Ra}^c} \nu_{Rb}$$

- they are needed to accommodate **fermion masses and mixing**

# Why?



- three generations of elementary fermions
- strong mass hierarchy among charged fermions
- (possibly) much milder one among neutrinos

# Why?

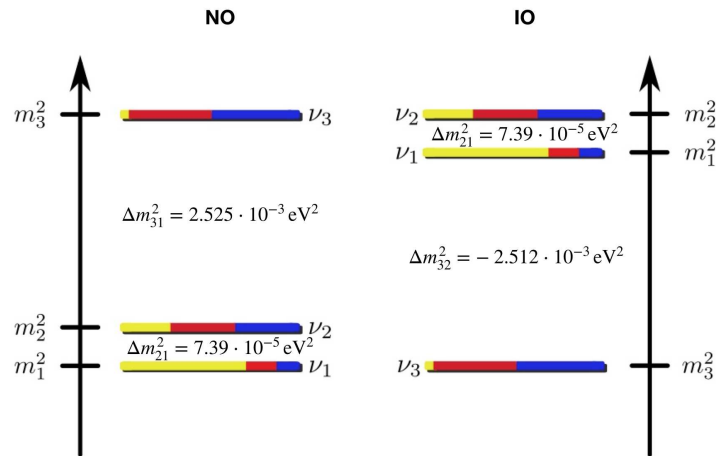
## Summary of charged fermion masses

(Xing et al. ('07))

	mass at $M_Z$	in units of heaviest mass
$e$	$(0.486570161 \pm 0.000000042) \text{ MeV}$	$\lambda^{4\div 5}$
$\mu$	$(102.7181359 \pm 0.0000092) \text{ MeV}$	$\lambda^2$
$\tau$	$1.74624^{+0.00020}_{-0.00019} \text{ GeV}$	1
$u$	$(1.27^{+0.50}_{-0.42}) \text{ MeV}$	$\lambda^8$
$c$	$(0.619 \pm 0.084) \text{ GeV}$	$\lambda^4$
$t$	$(171.7 \pm 3.0) \text{ GeV}$	1
$d$	$(2.90^{+1.24}_{-1.19}) \text{ MeV}$	$\lambda^4$
$s$	$(55^{+16}_{-15}) \text{ MeV}$	$\lambda^2$
$b$	$(2.89 \pm 0.09) \text{ GeV}$	1

# Why?

## Summary of current knowledge about neutrino masses

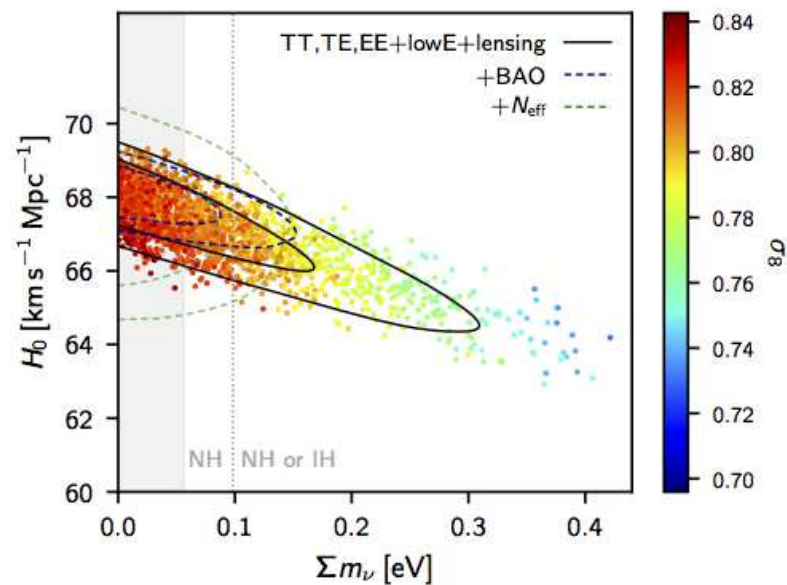


Their ordering is unknown, although NO seems preferred.

(NuFIT ('18))

Their absolute scale is also unknown.

(Planck ('18))



# Why?

- similarly, striking differences in **fermion mixing** among quarks and leptons are observed
- summary of current knowledge about **lepton mixing**

NuFIT 4.0 (2018)

	Normal Ordering (best fit)		Inverted Ordering ( $\Delta\chi^2 = 4.7$ )		
	bfp $\pm 1\sigma$	$3\sigma$ range	bfp $\pm 1\sigma$	$3\sigma$ range	
without SK atmospheric data	$\sin^2 \theta_{12}$	$0.310^{+0.013}_{-0.012}$	0.275 $\rightarrow$ 0.350	$0.310^{+0.013}_{-0.012}$	0.275 $\rightarrow$ 0.350
	$\theta_{12}/^\circ$	$33.82^{+0.78}_{-0.76}$	31.61 $\rightarrow$ 36.27	$33.82^{+0.78}_{-0.76}$	31.61 $\rightarrow$ 36.27
	$\sin^2 \theta_{23}$	$0.580^{+0.017}_{-0.021}$	0.418 $\rightarrow$ 0.627	$0.584^{+0.016}_{-0.020}$	0.423 $\rightarrow$ 0.629
	$\theta_{23}/^\circ$	$49.6^{+1.0}_{-1.2}$	40.3 $\rightarrow$ 52.4	$49.8^{+1.0}_{-1.1}$	40.6 $\rightarrow$ 52.5
	$\sin^2 \theta_{13}$	$0.02241^{+0.00065}_{-0.00065}$	0.02045 $\rightarrow$ 0.02439	$0.02264^{+0.00066}_{-0.00066}$	0.02068 $\rightarrow$ 0.02463
	$\theta_{13}/^\circ$	$8.61^{+0.13}_{-0.13}$	8.22 $\rightarrow$ 8.99	$8.65^{+0.13}_{-0.13}$	8.27 $\rightarrow$ 9.03
	$\delta_{CP}/^\circ$	$215^{+40}_{-29}$	125 $\rightarrow$ 392	$284^{+27}_{-29}$	196 $\rightarrow$ 360



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- summary of current knowledge about **lepton mixing**

(NuFIT ('18))

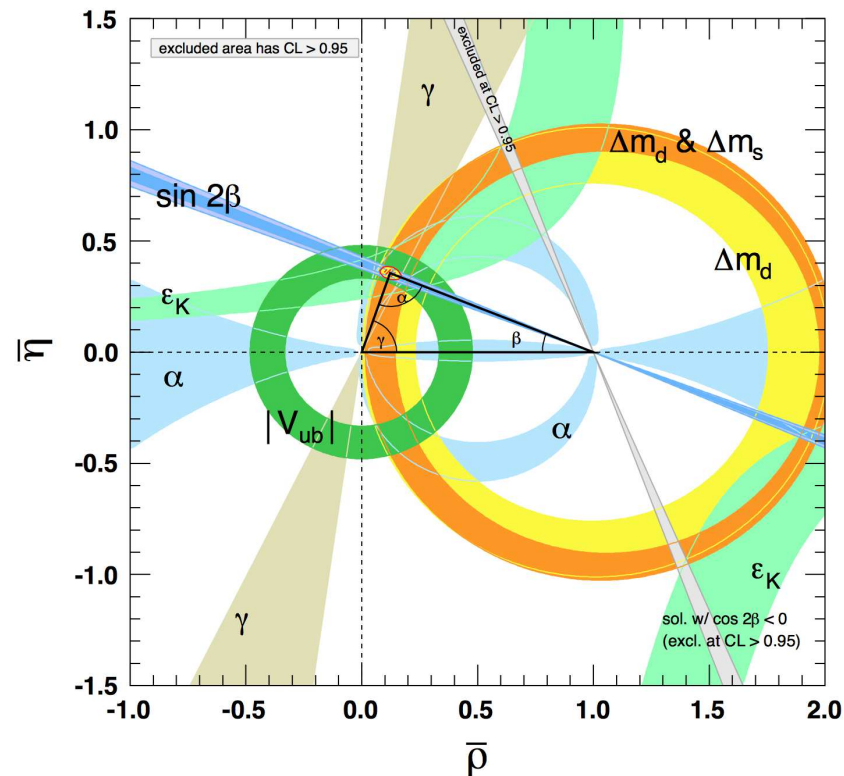
$$||U_{\text{PMNS}}|| \approx \begin{pmatrix} 0.82 & 0.55 & 0.15 \\ 0.29 & 0.59 & 0.75 \\ 0.49 & 0.59 & 0.64 \end{pmatrix} \quad [\text{NO}]$$

and hint for CP violation:  $\delta \approx 215^\circ$  ,  $\alpha = ?$  ,  $\beta = ?$

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(PDG ('18))

$$\|V_{\text{CKM}}\| \approx \begin{pmatrix} 0.97 & 0.22 & 3.7 \cdot 10^{-3} \\ 0.22 & 0.97 & 0.042 \\ 9.0 \cdot 10^{-3} & 0.041 & 0.999 \end{pmatrix}$$

and CP violation:  $J_{\text{CP}} = 3.18 \cdot 10^{-5}$

# Which ones?

$G_f$  could be ...

- ... abelian or non-abelian
- ... continuous or discrete
- ... local or global
- ... spontaneously broken or explicitly
- ... broken arbitrarily or to non-trivial subgroups
- ... broken at low or high energies

Its maximal possible size depends on the gauge group,

e.g. in the Standard Model:  $G_f \subset U(3)^5$ .

in  $SO(10)$ :  $G_f \subset U(3)$ .

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# Which ones?

You have (too) many choices ...

- permutation symmetries:  $S_N$  and  $A_N$  with  $N \in \mathbb{N}$
- dihedral symmetries:  $D_n$  and  $D'_n$  with  $n \in \mathbb{N}$
- further double-valued groups:  $T'$ ,  $O'$ ,  $I'$
- subgroups of  $SU(3)$ : series of  $\Delta(3n^2)$  and  $\Delta(6n^2)$  groups with  $n \in \mathbb{N}$ , as well as finite number of  $\Sigma$  groups
- subgroups of  $U(3)$  such as  $\Sigma(81)$  and subgroups of the listed groups such as  $T_7 \cong Z_7 \rtimes Z_3 \subset \Delta(147)$

NB: several isomorphisms among the groups,

e.g.  $S_3 \cong D_3 \cong \Delta(6)$ .



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# *Which ones?*

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So, you need

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# *Which ones?*

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So, you need

- a further **theoretical guiding principle**: **non-trivial subgroups**
- more **experimental data**: **CLFV**

In the following, ...

- ... I will use **non-trivial subgroups** as additional constraint
- ... I will show results in different types of models for **CLFV**

# *Predictive power*

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I would like to focus on

explanations for lepton mixing

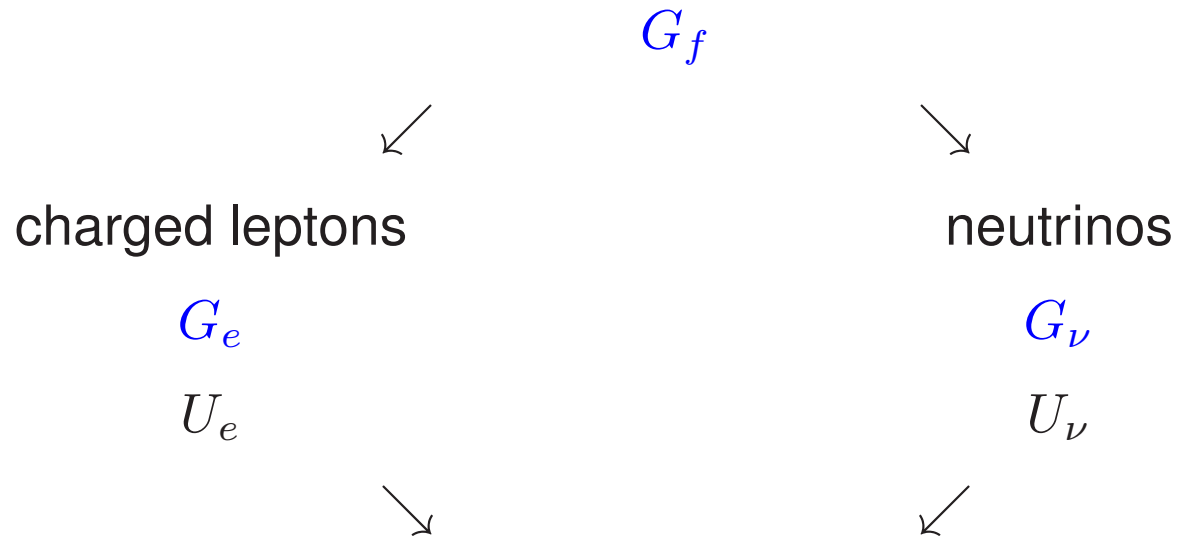
Two very successful approaches are:

- flavor symmetry  $G_f$  and its breaking as explanation
- flavor symmetry  $G_f$  and CP and its breaking as explanation

# Predictive power

Flavor symmetry  $G_f$  and its breaking as explanation

(Lam ('07,'08), Blum/H/Lindner ('07))



$$U_{\text{PMNS}} = U_e^\dagger U_\nu$$

Note: Masses do not play a role in this approach.

# Predictive power

You expect to fix

- the 3 lepton mixing angles
- the Dirac phase  $\delta$

up to permutations of rows and columns of the PMNS mixing matrix, since masses are not fixed.

Example for  $G_f$  for  $\theta_{13} \neq 0$ ,  $\theta_{23} \neq \frac{\pi}{4}$

*(de Adelhart Toorop/Feruglio/H ('11))*

- $G_f = \Delta(384)$
- $G_e = Z_3$
- $G_\nu = Z_2 \times Z_2$

# *Predictive power*

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*(de Adelhart Toorop/Feruglio/H ('11))*

$$\|U_{\text{PMNS}}\| \approx \begin{pmatrix} 0.81 & 0.58 & 0.11 \\ 0.31 & 0.58 & 0.75 \\ 0.50 & 0.58 & 0.65 \end{pmatrix} \quad \text{and} \quad \sin \delta = 0$$

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NB: The value of  $\theta_{13}$  is controlled by the choice of flavor symmetry.  
The result  $\sin \delta = 0$  instead is generic.



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NB: The value of  $\theta_{13}$  is controlled by the choice of flavor symmetry.  
The result  $\sin \delta = 0$  instead is generic.

Non-trivial values of  $\delta$  arise from the combination of flavor and CP symmetries.

With less stringent assumptions on flavor (and residual) symmetries non-trivial CP violation can also be achieved,  
see e.g. *Hernandez/Smirnov ('12)*.

# Predictive power

Flavor symmetry  $G_f$  and CP and its breaking as explanation

(Feruglio/H/Ziegler ('12))

$G_f$  and CP

charged leptons

$$G_e$$

$$U_e$$

neutrinos

$$G_\nu = Z_2 \times \text{CP}$$

$$U_\nu = \Omega_\nu R(\theta) K_\nu$$

$$U_{\text{PMNS}} = U_e^\dagger \Omega_\nu R(\theta) K_\nu$$

Note: Masses do not play a role in this approach.

# Predictive power

You expect to express

- the 3 lepton mixing angles
- **all** CP phases  $\delta$ ,  $\alpha$  and  $\beta$

in terms of **one single real** parameter  $\theta$  and up to permutations of rows and columns of the PMNS mixing matrix, since one has one  $Z_2$  only and masses are not fixed.

Example: study of  $G_f = A_5$  with CP

*(Di Iura/H/Meloni ('15))*

- $G_f = A_5$  and **CP**
- $G_e = Z_2 \times Z_2, Z_3$  or  $Z_5$
- $G_\nu = Z_2 \times$  **CP**

**4 different mixing patterns with different characteristics**

# Predictive power

(Di Iura/H/Meloni ('15))

case ||  $\tan \varphi = 1/\phi, \phi = \frac{1}{2} (1 + \sqrt{5})$  and  $\Phi = \frac{2\pi}{5}$

$$U_{\text{PMNS}} = \frac{1}{\sqrt{2}} \begin{pmatrix} -\sqrt{2} \cos \varphi & -\sqrt{2} \sin \varphi & 0 \\ -e^{-3i\Phi} \sin \varphi & e^{-3i\Phi} \cos \varphi & -e^{-7i\Phi/4} \\ -e^{-2i\Phi} \sin \varphi & e^{-2i\Phi} \cos \varphi & e^{-3i\Phi/4} \end{pmatrix} R_{13}(\theta) K_\nu$$

$$\sin^2 \theta_{13} = \frac{1}{10} (5 + \sqrt{5}) \sin^2 \theta, \quad \sin^2 \theta_{12} = \frac{2}{2 + (3 + \sqrt{5}) \cos^2 \theta}, \quad \sin^2 \theta_{23} = \frac{1}{2}$$

and

$$|\sin \delta| = 1, \quad \sin \alpha = 0, \quad \sin \beta = 0$$

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$$\sin^2 \theta_{13} \approx 0.0219, \quad \sin^2 \theta_{12} \approx 0.283, \quad \sin^2 \theta_{23} = \frac{1}{2}$$

and

$$|\sin \delta| = 1, \quad \sin \alpha = 0, \quad \sin \beta = 0 \quad \text{for } \theta \approx 0.175$$

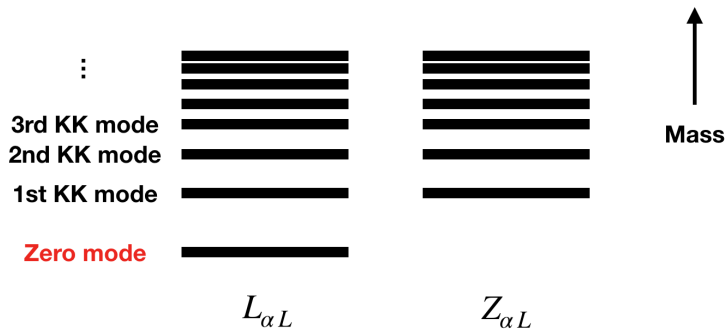


***CLFV***

# Extra dimension

Kaluza-Klein fermions ...

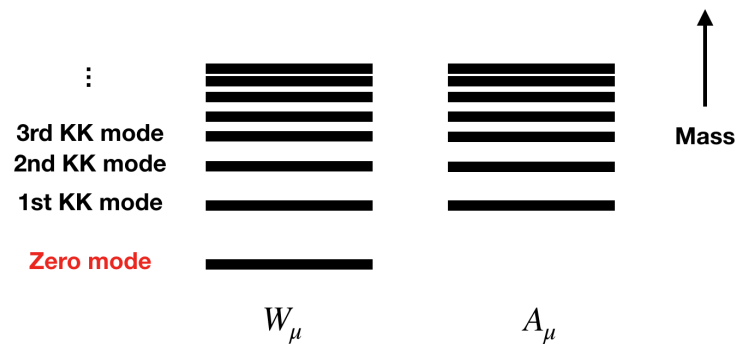
- ... usually also appear in three generations.



# Extra dimension

Kaluza-Klein gauge bosons ...

- ... can be associated with Standard Model gauge group.
- ... can belong to new gauge interactions.

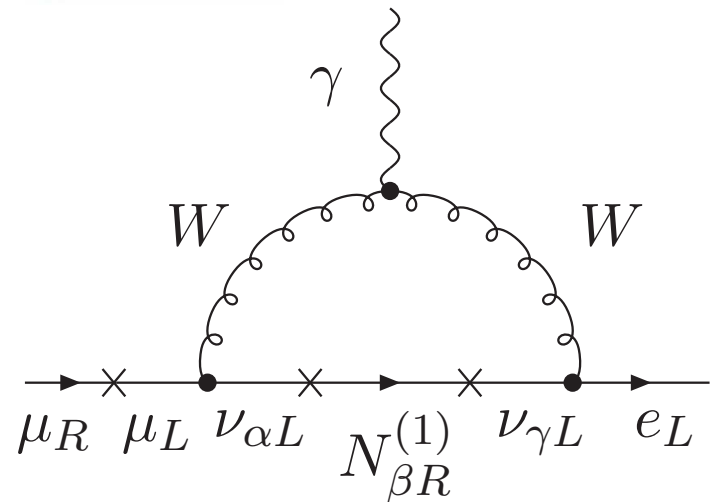
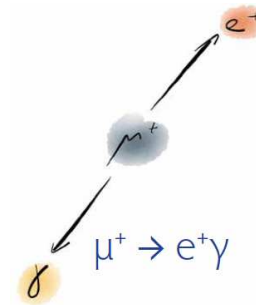
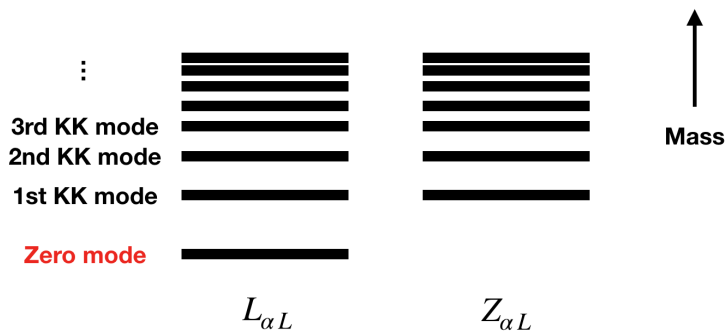




# Extra dimension

Kaluza-Klein particles ...

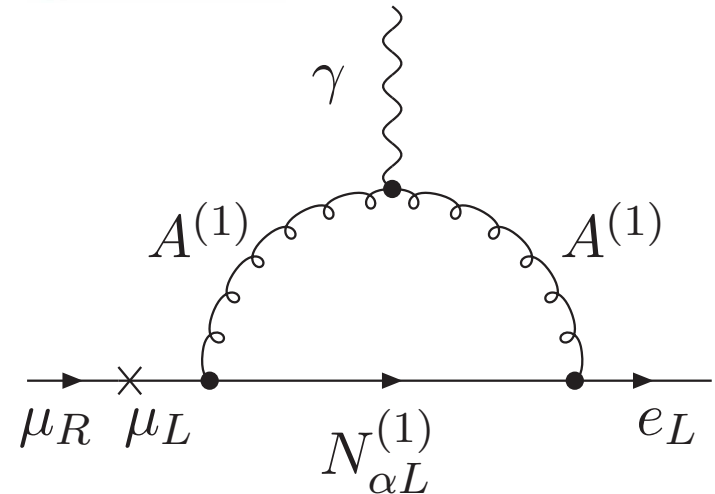
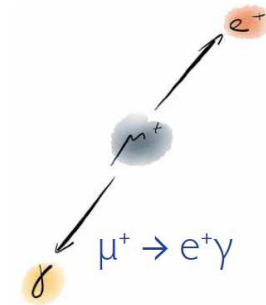
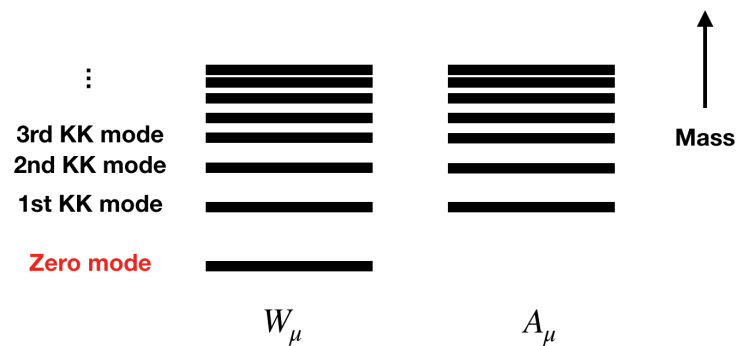
- ... can induce, in particular, CLFV.



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# Extra dimension

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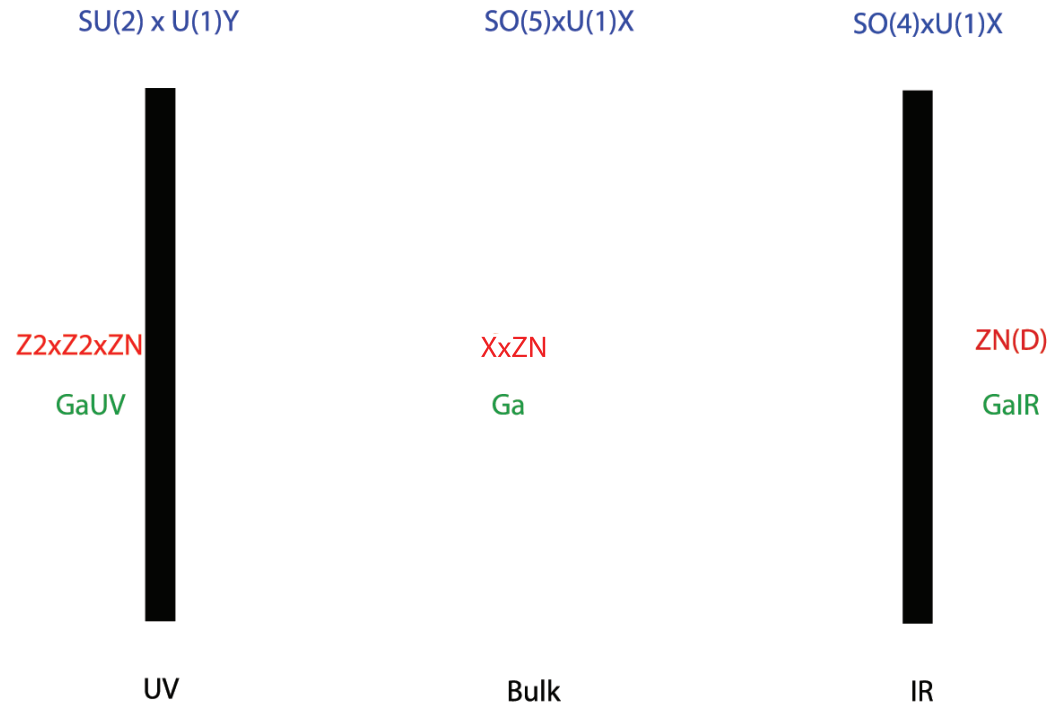
(H/Serone ('11), H ('13))

- study 5D realization of holographic composite Higgs models (gauge-Higgs unified model with warped extra dimension)
- impose flavor symmetry  $G_f = X \times Z_N$  (these correspond to the different colors in the plots)
- chosen symmetries can describe lepton mixing angles

# Extra dimension

(H/Serone ('11), H ('13))

- impose flavor symmetry  $G_f = X \times Z_N$   
(these correspond to the different colors in the plots)
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# Extra dimension

(H/Serone ('11), H ('13))

- assume neutrinos are **Majorana particles**
- particle content ( $X = 0$ )

$$\xi_{l,\alpha} \sim (\mathbf{5}, \mathbf{3}, 1, \omega_3, \omega_3), \quad \xi_{e,\alpha} \sim (\mathbf{10}, \mathbf{1}, \omega_N^{n_\alpha}, \omega_3, \omega_3), \quad \xi_{\nu,\alpha} \sim (\mathbf{1}, \mathbf{3}, 1, \omega_3, 1)$$

with

$$\xi_{l,\alpha} = \begin{pmatrix} \left[ \tilde{L}_{1,\alpha L} (-+), L_{\alpha L} (++) \right] \\ \hat{\nu}_{\alpha L} (-+) \end{pmatrix}, \quad \xi_{e,\alpha} = \begin{pmatrix} x_{\alpha L} (+-) \\ \tilde{\nu}_{\alpha L} (+-) \quad Z_{\alpha L} (+-) \\ e_{\alpha L} (--) \\ \left[ \tilde{L}_{2,\alpha L} (+-), \hat{L}_{\alpha L} (+-) \right] \end{pmatrix}$$

$$\xi_{\nu,\alpha} = \nu_{\alpha L} (--)$$

# Extra dimension

(H/Serone ('11), H ('13))

- assume neutrinos are **Majorana particles**
- mass mixing terms at IR brane

$$-\mathcal{L}_{\text{IR}} = \left(\frac{R}{R'}\right)^4 \sum_{\alpha=e,\mu,\tau} \left( m_{\text{IR},\alpha}^l \left( \widetilde{L}_{1,\alpha L} \widetilde{L}_{2,\alpha R} + \overline{L}_{\alpha L} \hat{L}_{\alpha R} \right) + m_{\text{IR},\alpha}^\nu \overline{\hat{\nu}}_{\alpha L} \nu_{\alpha R} + h.c. \right)$$

and Majorana mass terms for RH neutrinos at UV brane

$$-\mathcal{L}_{\text{UV}} = \frac{1}{2} \overline{\nu_{\alpha R}^c} \mathcal{M}_{\text{UV},\alpha\beta} \nu_{\beta R} + h.c.$$

are constrained by symmetries as well

# Extra dimension

(H/Serone ('11), H ('13))

- assume neutrinos are **Dirac particles**
- crucial change in particle content

$$\xi_{l,\alpha} = \begin{pmatrix} \left[ \tilde{L}_{1,\alpha L} (-+), L_{\alpha L} (++) \right] \\ \hat{\nu}_{\alpha L} (+-) \end{pmatrix}$$

- also additional symmetries are slightly adjusted

# Extra dimension

(H/Serone ('11), H ('13))

- assume neutrinos are **Dirac particles**
- mass mixing terms at IR brane

$$-\mathcal{L}_{\text{IR}} = \left(\frac{R}{R'}\right)^4 \sum_{\alpha=e,\mu,\tau}^3 m_{\text{IR},\alpha}^l \left( \widetilde{L}_{1,\alpha L} \widetilde{L}_{2,\alpha R} + \overline{L}_{\alpha L} \hat{L}_{\alpha R} \right) + h.c.$$

and mass mixing at UV brane

$$-\mathcal{L}_{\text{UV}} = \overline{\hat{\nu}}_{\alpha L} \mathcal{M}_{\text{UV},\alpha\beta} \nu_{\beta R} + h.c.$$

are constrained by symmetries as well



# Extra dimension

(H/Serone ('11), H ('13))

- in addition to mass mixing terms there are  
boundary kinetic terms (BKTs)  
where the most important ones are

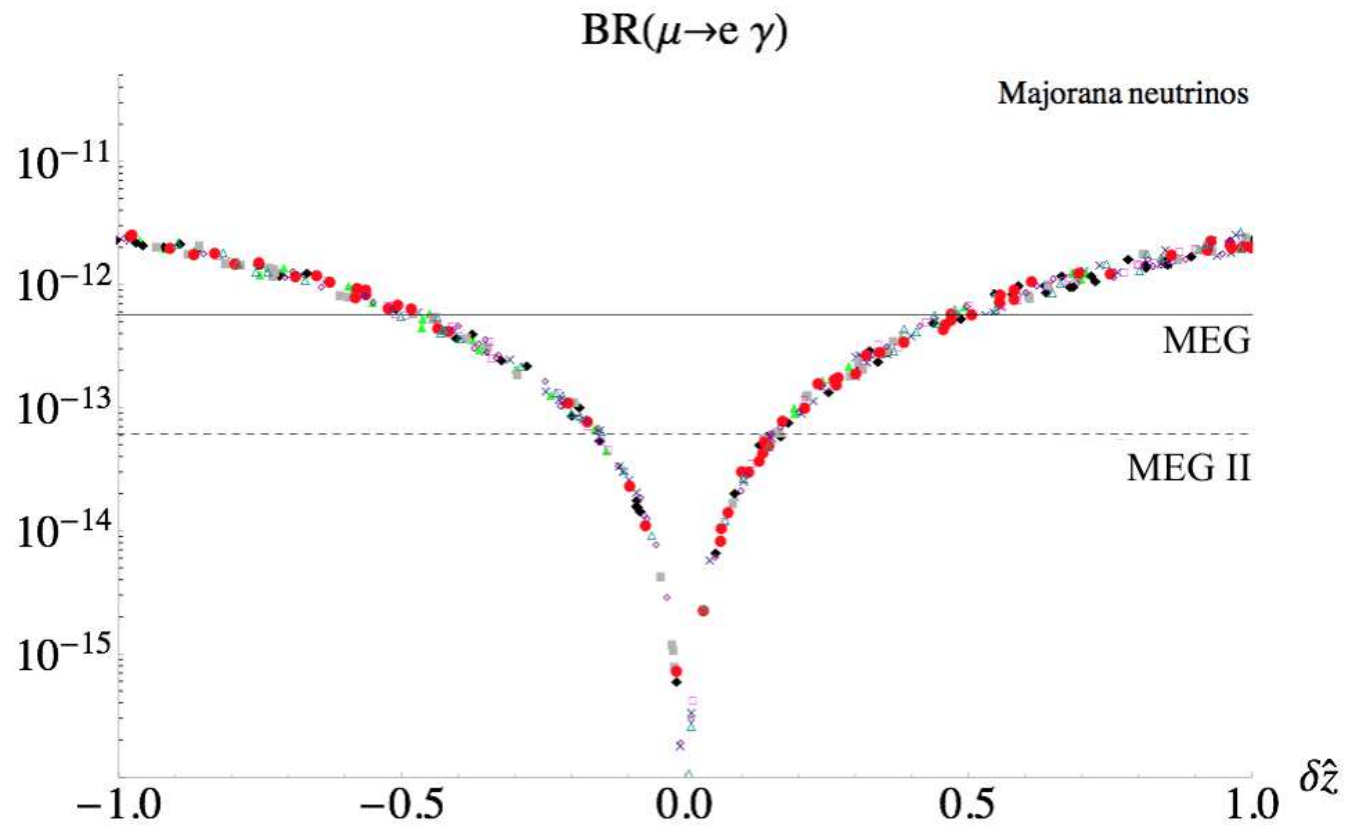
$$\mathcal{L}_{\text{BKT}} = \bar{L}_L(x, R)(R\hat{Z}_l)i\not{D}L_L(x, R).$$

Also, these are constrained by the symmetries.

# Extra dimension

Model with Majorana neutrinos

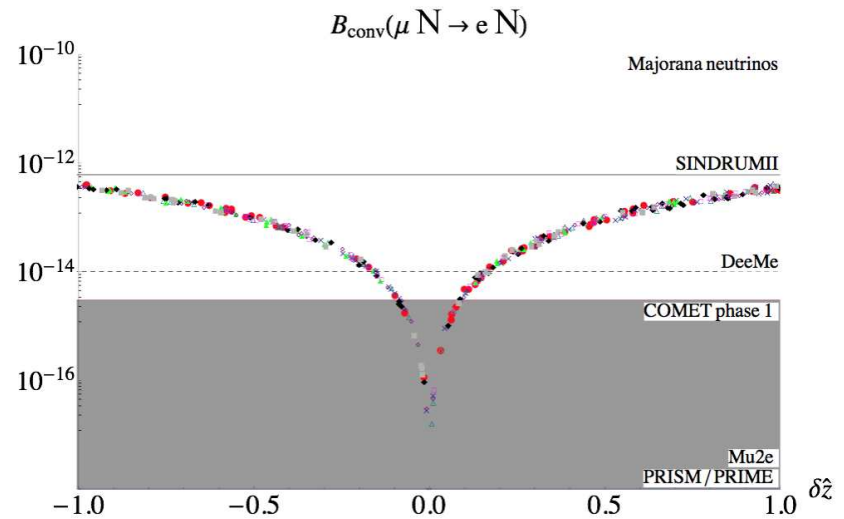
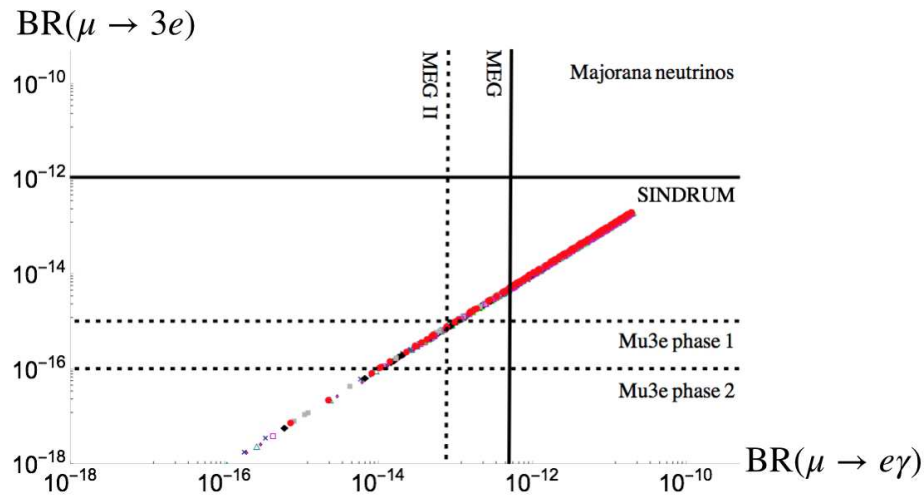
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# Extra dimension

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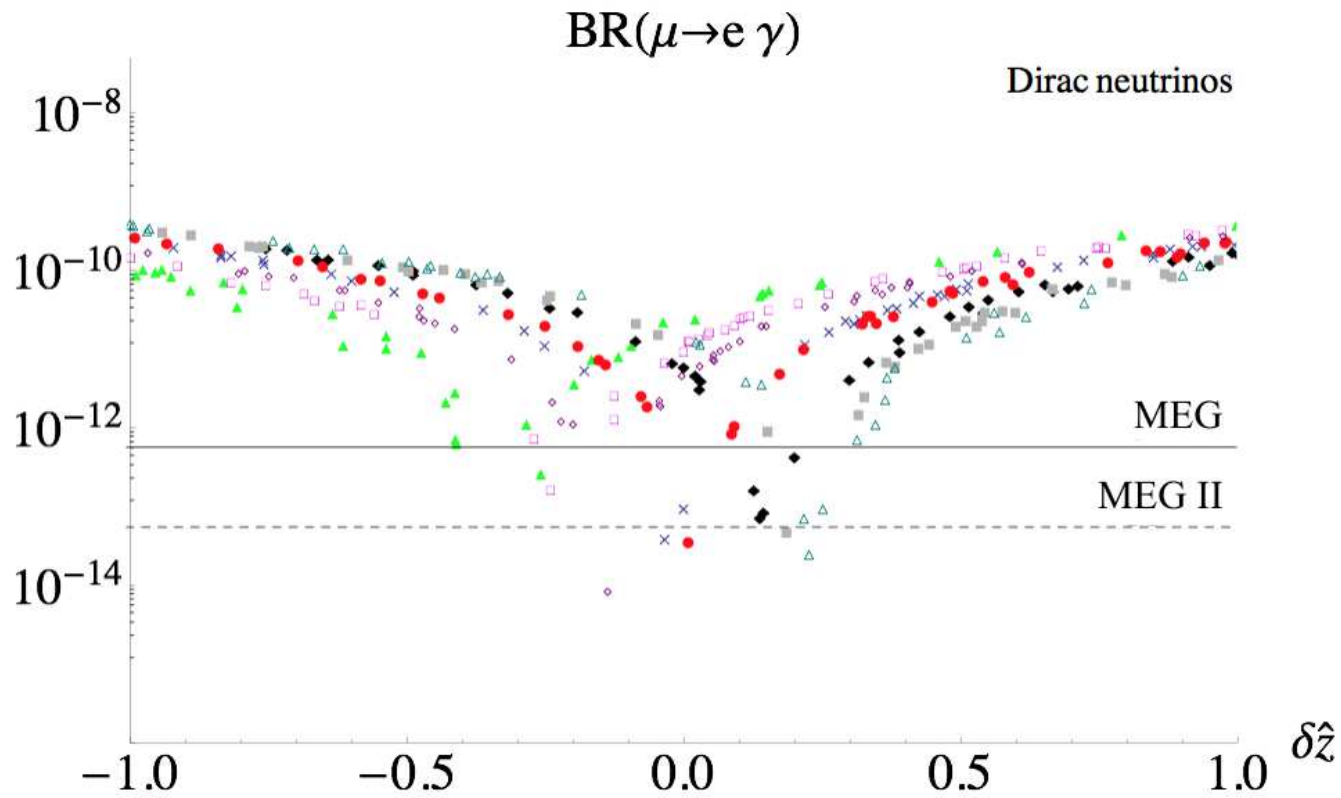
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# Extra dimension

Model with Dirac neutrinos

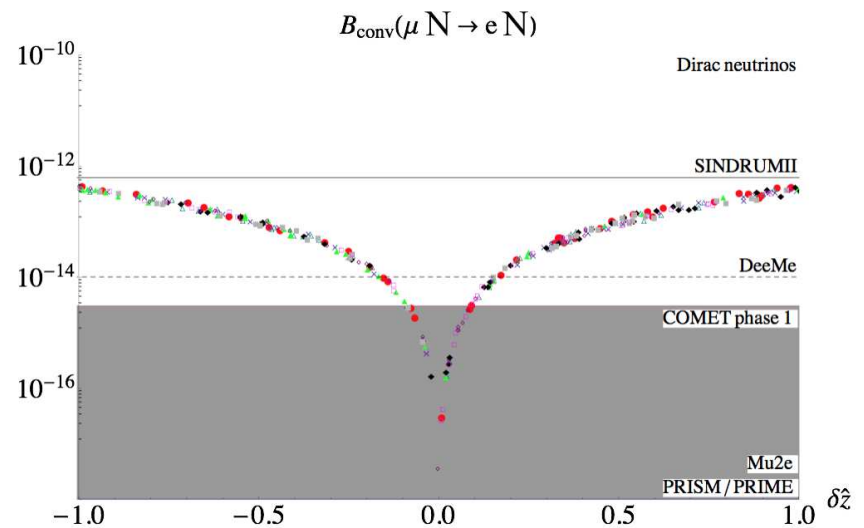
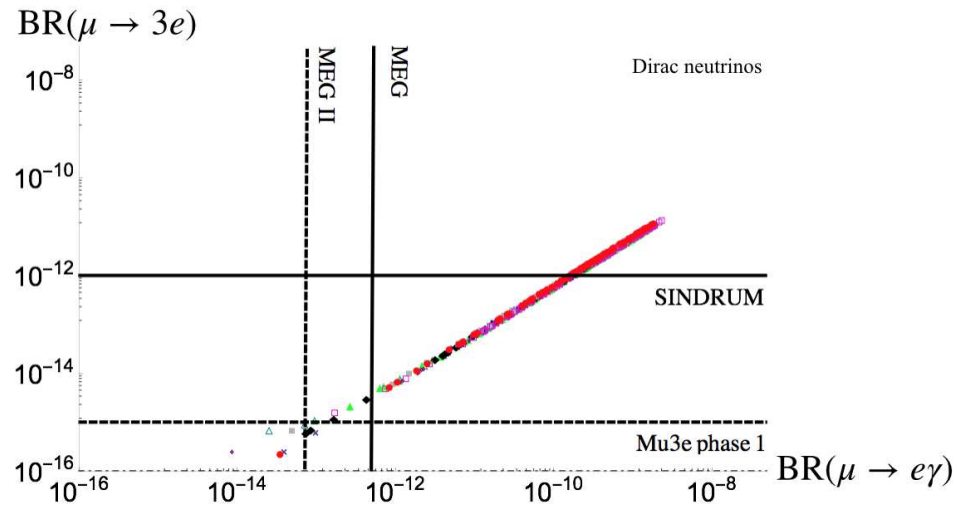
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# Extra dimension

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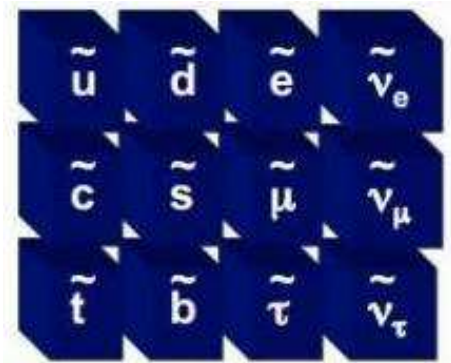
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# *Supersymmetric models*

Supersymmetric particles ...

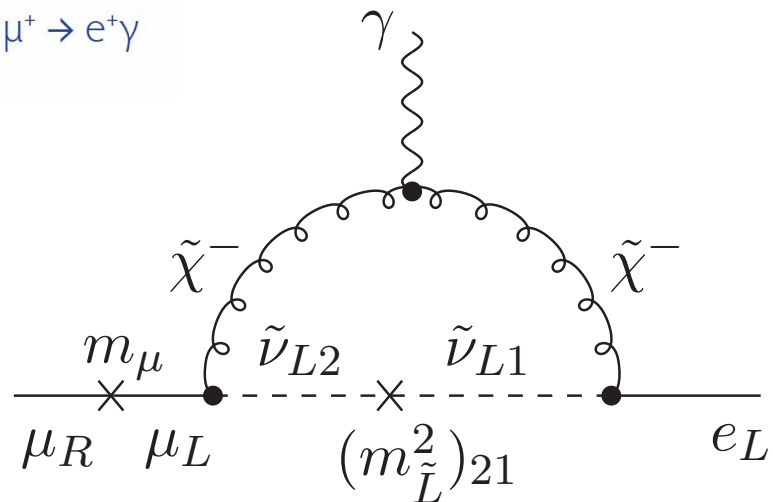
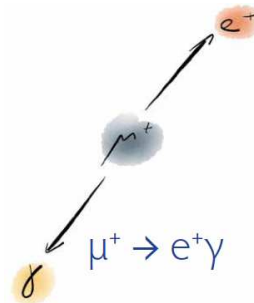
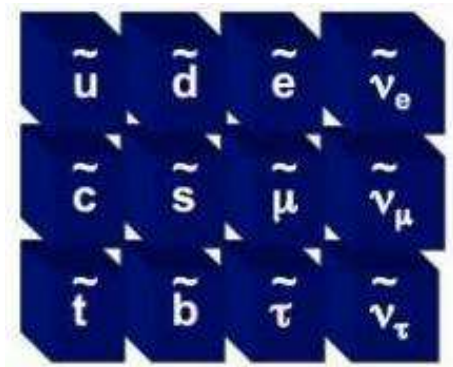
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# Supersymmetric models

Supersymmetric particles ...

- ... usually also appear in three generations.
- ... can induce, in particular, CLFV.



# Supersymmetric models

*(Lopez-Ibanez et al. ('19))*

- impose flavor symmetry  $G_f = A_5$  and CP on MSSM theory
- chosen symmetries can describe lepton mixing well

*(Di Iura/H/Meloni ('15))*

- assume soft SUSY breaking terms

$$\mathcal{L}_{\text{soft}} \supset (m_{\tilde{L}}^2)_{ij} \tilde{L}_i^\dagger \tilde{L}_j$$

are constrained by these symmetries as well

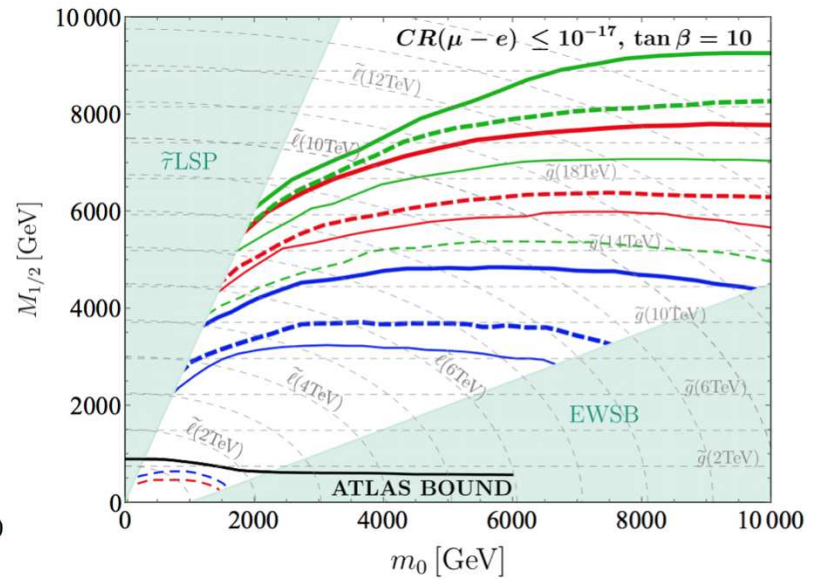
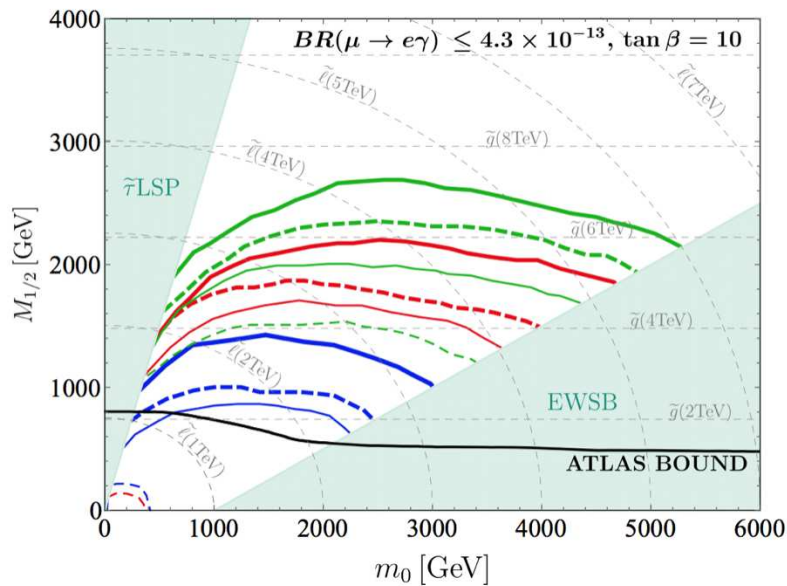
- consider different (minimal) scenarios for neutrino masses  
(these correspond to the differently colored lines in the plots)



# Supersymmetric models

CLFV signals can probe soft masses larger than direct searches

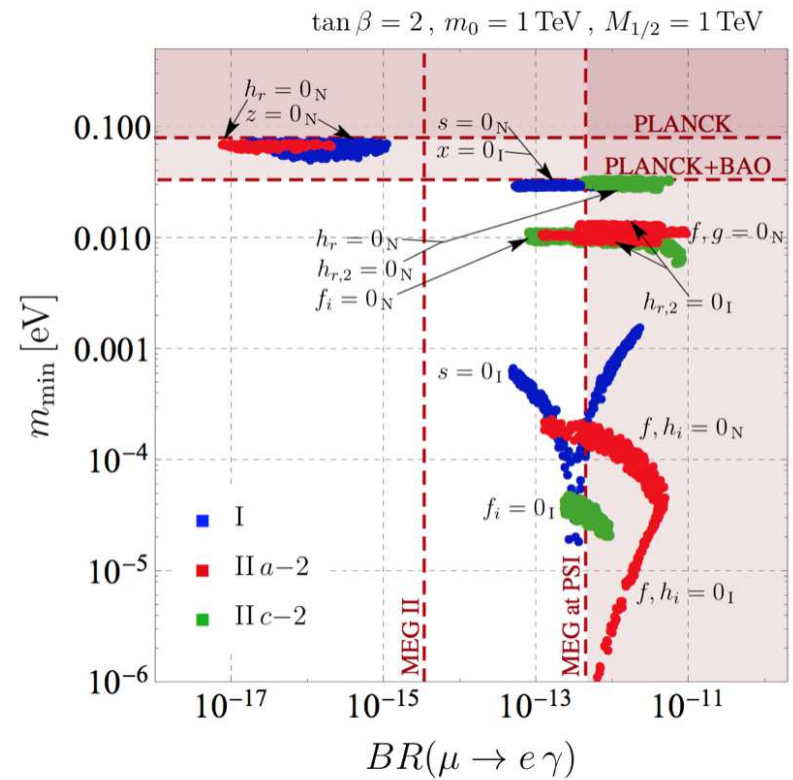
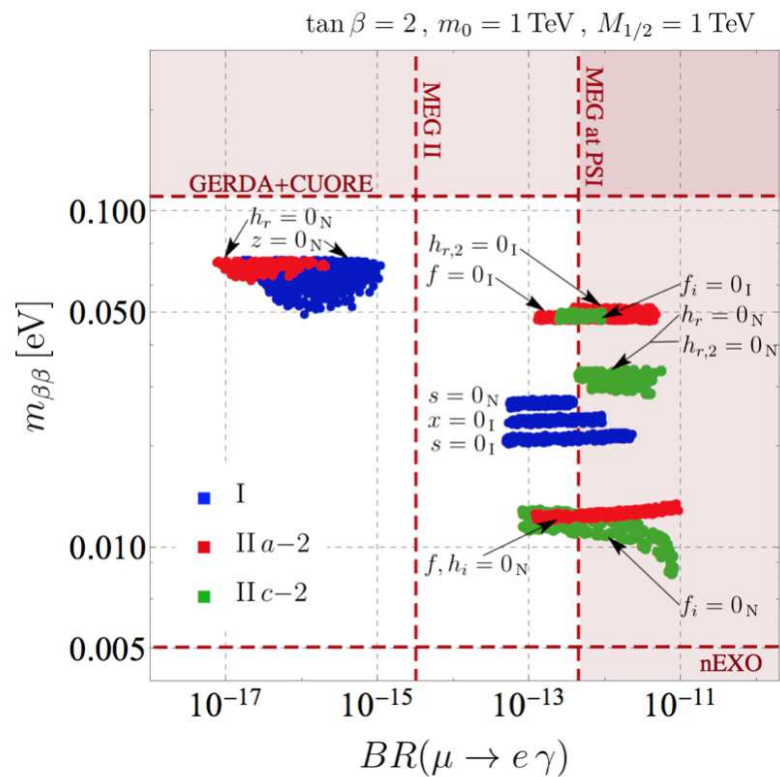
(Lopez-Ibanez et al. ('19))



# Supersymmetric models

Correlations between CLFV signals and neutrino masses arise

(Lopez-Ibanez et al. ('19))





***Continuous flavor symmetries  
and CLFV***

# Some examples

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- models with  $SU(3)$  flavor symmetry,  
see e.g. *Antusch/King/Malinsky ('07)*, *Calibbi/Jones-Perez/Vives ('08)*
- models with so-called minimal flavor violation can also be considered in this category,  
flavor symmetry  $U(3)^5$  or alike is broken by spurions
- flavor symmetry can also be abelian,  
see e.g. *Calibbi et al. ('12)*



***Phenomenological imprints  
beyond CLFV***

# *Examples of studied signals*

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- neutrinoless double beta decay,  
see e.g. *King/Neder ('14)*, *Ding/King/Neder ('14)*, *H/Molinaro ('16)*
- generation of baryon asymmetry through unflavored and flavored leptogenesis in scenarios with flavor and CP symmetries,  
see e.g. *H/Molinaro ('16)*, *Mohapatra/Nishi ('15)*, *Chen/Ding/King ('16)*
- anomalous magnetic moment and electric dipole moment of charged leptons,  
see e.g. *Feruglio et al. ('08)*, *Dimou/King/Luhn ('15)*
- possible explanation of (some) B physics anomalies,  
see e.g. *de Medeiros Varzielas/Hiller ('15)*

# *Conclusions*

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- discrete non-abelian flavor symmetries (and CP) and their breaking can lead to a good description of lepton mixing angles and allow for predictions of the leptonic CP phase(s)
- in theories beyond the Standard Model such symmetries also constrain the signal strength of CLFV processes (models with extra dimension and supersymmetric models)
- such symmetries can leave further phenomenological imprints, e.g. by predicting neutrinoless double beta decay or the electric dipole moment of charged leptons

Thank you for your attention.