

QCD Topology to High Temperatures via Reweighting



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UNIVERSITÄT
DARMSTADT

P. Thomas Jahn,^{†,1} Guy D. Moore,¹ and Daniel Robaina²

¹ Institut für Kernphysik, Technische Universität Darmstadt, Germany

[†]email: tjahn@theorie.ikp.physik.tu-darmstadt.de

² Max-Planck Institute of Quantum Optics, Garching, Germany

Introduction

- The QCD axion is a hypothetical particle, predicted in models which solve the strong CP problem via the Peccei-Quinn mechanism [1] and is also a dark matter candidate [2].
- There is currently lots of experimental effort to detect axions and theoretical effort to determine its properties, especially its mass.
- From a theoretical point of view, the axion properties depend on the **QCD topological susceptibility**

$$\chi_{\text{top}} = \frac{1}{\mathcal{V}} \langle Q^2 \rangle, \quad Q = \frac{1}{32\pi^2} \int d^4x F_{\mu\nu}^a \tilde{F}_{\mu\nu}^a \in \mathbb{Z}, \quad \tilde{F}_{\mu\nu}^a \equiv \frac{1}{2} \epsilon_{\mu\nu\alpha\beta} F_{\alpha\beta}^a$$

up to temperatures of $\sim 7 T_c$ [3] with \mathcal{V} the four-dimensional volume, Q the **topological charge**, and $\tilde{F}_{\mu\nu}$ the **dual field-strength tensor**.

- Topologically non-trivial fields are **instantons/calorons** with weight

$$\exp(-S) = \exp\left(-\frac{8\pi^2|Q|}{g^2}\right) \rightarrow 0 \quad \text{as} \quad T \rightarrow \infty (g \rightarrow 0).$$

- Since topological phenomena are inherently non-perturbative, lattice QCD is the method of choice. However, at high temperatures calorons are badly suppressed and it becomes impossible to measure fluctuations of Q with standard lattice techniques.

Reweighting Method

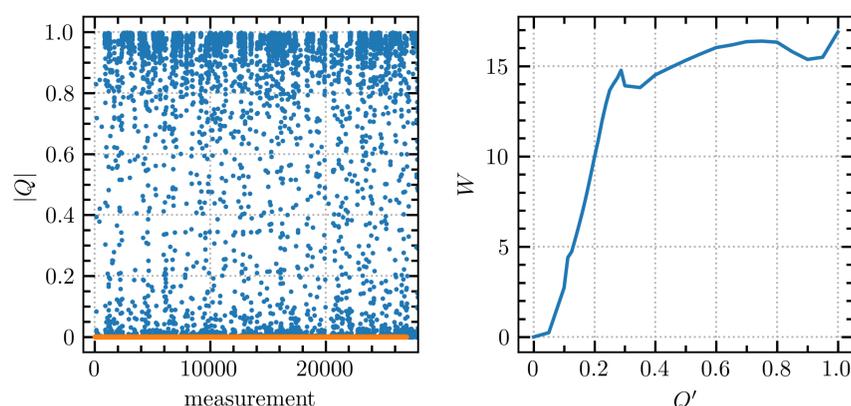
- In order to artificially enhance the number of caloron configurations, we developed a technique based on **reweighting** [4].
- Instead of determining the expectation value via standard importance sampling

$$\langle Q^2 \rangle = \frac{\int \mathcal{D}U Q^2 e^{-S[U]}}{\int \mathcal{D}U e^{-S[U]}} \Rightarrow \langle Q^2 \rangle \simeq \frac{1}{N} \sum_{i=1}^N Q_i^2$$

we apply importance sampling with a modified weight by introducing the **reweighting function** $W(Q')$:

$$\langle Q^2 \rangle = \frac{\int \mathcal{D}U Q^2 e^{-S[U]+W(Q')} e^{-W(Q')}}{\int \mathcal{D}U e^{-S[U]+W(Q')} e^{-W(Q')}} \Rightarrow \langle Q^2 \rangle \simeq \frac{\sum_{i=1}^N Q_i^2 e^{-W(Q'_i)}}{\sum_{i=1}^N e^{-W(Q'_i)}}$$

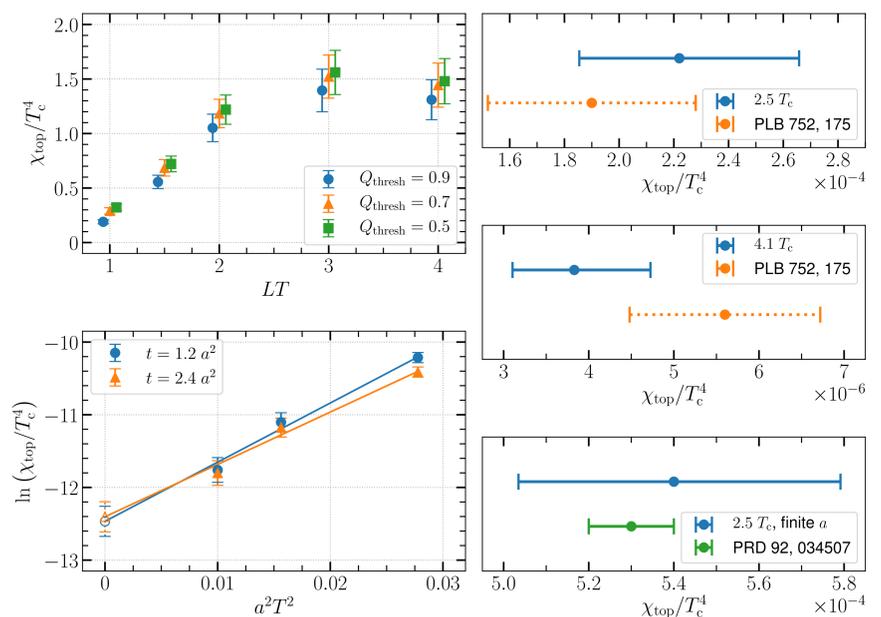
- Since the topological charge is badly contaminated by UV fluctuations, we apply **gradient flow** to both the observable Q (large amount) and the reweighting variable Q' (small amount).
- The reweighting technique is implemented via an additional Metropolis step.
- Using reweighting, the number of caloron configurations can be significantly enhanced if the reweighting function is chosen correctly. We developed an automated way to build the reweighting function:
 - perform a separate Monte Carlo simulation only for building W
 - constrain to the interval $|Q'| \in [0, 1]$ which is sufficient for high T
 - start with a flat function $W(Q') = 0$
 - at each Monte Carlo step measure Q' and lower W at the measured value
 - make the procedure converge by reducing the amount of lowering W after the whole range was explored



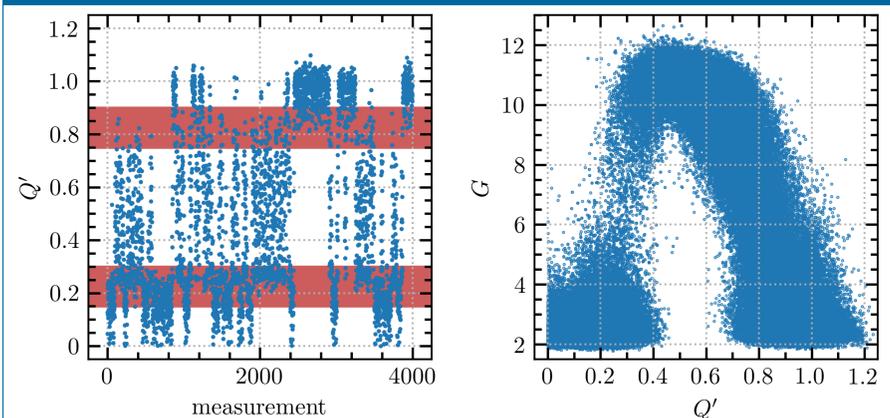
Results

- Continuum extrapolated quenched results at $2.5 T_c$ and $4.1 T_c$
- Thresholded, flowed topological charge as observable:

$$\langle Q^2 \rangle = \frac{\sum_i e^{-W(Q'_i)} \Theta(Q_i^2(t) - Q_{\text{thresh}}^2)}{\sum_i e^{-W(Q'_i)}}$$



Improvement of the Technique



- Reweighting in terms of the topological charge only is incomplete; Q' is not homogeneously distributed but has “barriers”
- Algorithm has problems to move between trivial topology, “dislocations” (very small calorons, lattice artifacts), and genuine calorons
- Peak action density** G is able to distinguish trivial topology and dislocations
→ additional reweighting in terms of G to overcome the lower barrier
- High barrier is mildened by using larger HMC steps and no reweighting at all
- A middle region becomes necessary to connect both regions
- ONGOING:** Continuum extrapolated determination of the topological susceptibility up to $7 T_c$ and extension to full QCD

References

- [1] R.D. Peccei and Helen R. Quinn, *CP Conservation in the Presence of Instantons*, *Phys.Rev.Lett.* **38**, 1440 (1977).
- [2] John Preskill, Mark B. Wise, and Frank Wilczek, *Cosmology of the Invisible Axion*, *Phys. Lett.* **B120**, 127 (1983).
- [3] Guy D. Moore, *Axion dark matter and the Lattice*, *EPJ Web Conf.* **175**, 01009 (2018).
- [4] P Thomas Jahn, Guy D. Moore, and Daniel Robaina, $\chi_{\text{top}}(T \gg T_c)$ in pure-gluon QCD through reweighting, *Phys. Rev.* **D98**, 054512 (2018).

